## **Announcements**

EXAM 2 is next Wednesday, Nov. 7

@ 8:00pm-9:30pm in ELLIOT HALL of MUSIC.

#### **Physics Help Center Survey**

#### PHYS Building, Rms. 11-12

#### How often do you use the Help Center on average?

- A. Never
- B. Once or at most twice a semester
- C. Several times in a semester
- D. Around once a week
- E. More than once a week

### **Physics Help Center Survey**

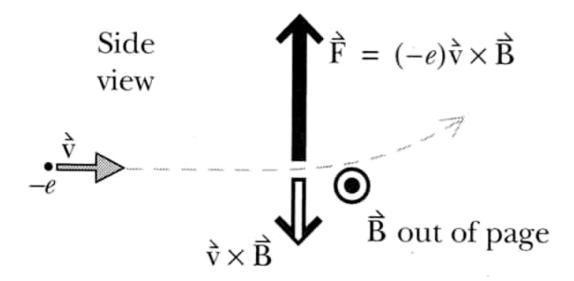
#### PHYS Building, Rms. 11-12

#### How useful was your Help Center visit?

- A. I did not use the Help Center
- B. Useless
- C. Not very useful
- D. Useful
- E. Very useful

# Magnetic Force on a Charge

$$\vec{F}_{magnetic} = q\vec{v} \times \vec{B}$$



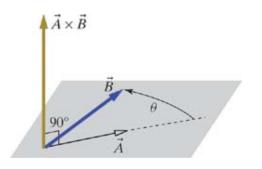
#### Electron charge = -e:

The magnetic force on a moving electron is in opposite direction to the direction of the cross product  $\vec{v} \times \vec{B}$ 

# **Right-Hand Rule**

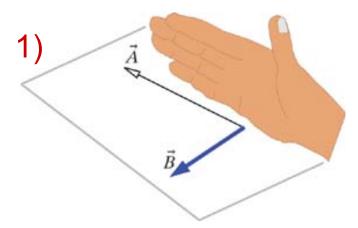
$$\vec{F}_{\text{mag}} = q\vec{v} \times \vec{B}$$

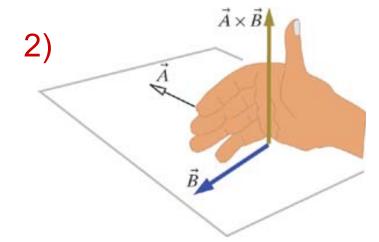
**MAGNETIC FORCE** point charge



Result of Cross Product  $\vec{v} \times \vec{B}$  is Perpendicular to both  $\vec{v}$  and  $\vec{B}$ 

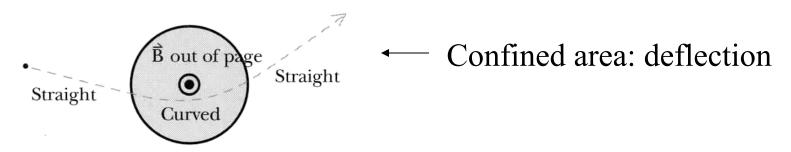
#### Right-Hand Rule:





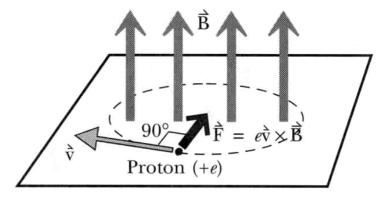
# Motion in a Magnetic Field

$$\vec{F}_{magnetic} = q\vec{v} \times \vec{B}$$



What if we have large (infinite) area with constant  $\overrightarrow{B} \perp \overrightarrow{v}$ 

$$\left| \frac{d\vec{p}}{dt} \right| = qvB$$



# Determining the Momentum of a Particle

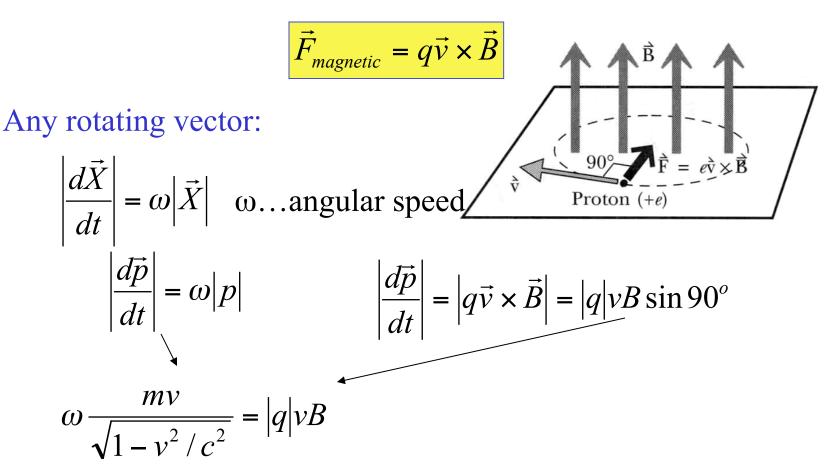
Position vector 
$$r$$
:  $\left| \frac{d\vec{r}}{dt} \right| = v = \omega r \longrightarrow \omega = \frac{v}{r}$  Circular motion

$$\left| \frac{d\vec{p}}{dt} \right| = \omega p = \left( \frac{v}{r} \right) p \qquad \left| \frac{d\vec{p}}{dt} \right| = |q| v B$$

$$\left( \frac{v}{r} \right) p = |q| v B$$

Used to measure momentum in high-energy particle experiments

# Circular Motion at any Speed



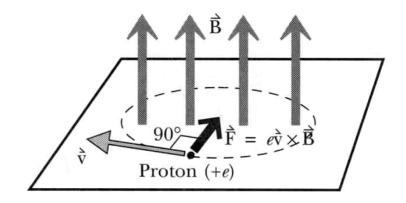
$$\omega = \frac{|q|B}{m}\sqrt{1 - v^2/c^2}$$

Cyclotron Frequency

# **Circular Motion at Low Speed**

$$\omega = \frac{|q|B}{m}\sqrt{1 - v^2/c^2}$$

if 
$$v << c$$
:  $\omega = \frac{|q|B}{m}$ 



#### independent of v!

#### Alternative derivation:

$$F = ma$$

$$|q|vB\sin 90^{\circ} = m\frac{v^{2}}{R}$$

$$|q|B = m\omega \longrightarrow$$

#### Circular motion:

$$a = \frac{v^2}{R} \qquad \omega = \frac{v}{R}$$

$$\omega = \frac{|q|B}{m}$$

Period T: 
$$\omega = \frac{2\pi}{T} \longrightarrow T = 2\pi \frac{m}{|q|B} \longleftarrow$$
 Non-Relativistic

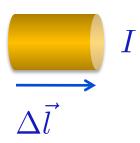
## **Biot-Savart Law**



$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \underbrace{\vec{q\vec{v}} \times \hat{r}}_{|r|^2}$$

**BIOT-SAVART LAW** point charge

We need to understand how these are related

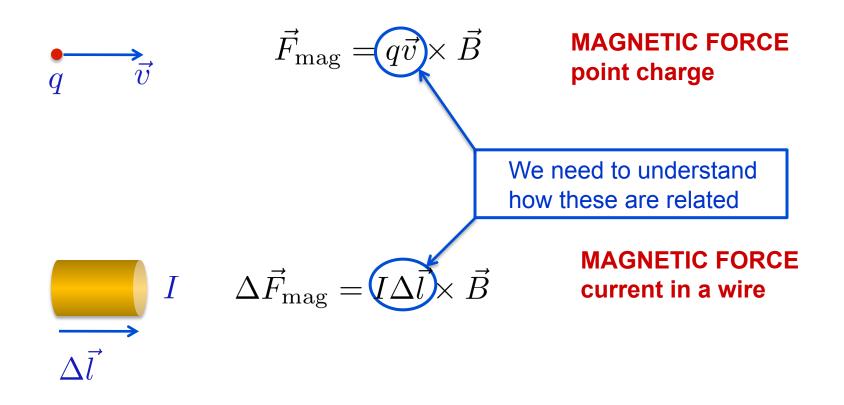


$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{\vec{D} \cdot \vec{l} \cdot \hat{r}}{|r|^2}$$

BIOT-SAVART LAW current in a wire

 $\Delta \vec{l}$  = length of this chunk of wire

## Magnetic Force on a charge or wire



$$\Delta \vec{l}$$
 = length of this chunk of wire

#### GOAL: Show $q\vec{v} \longrightarrow I\Delta\vec{l}$ (point charge) (wire)

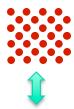
(POSITIVE) **POINT CHARGE** 



 $q\vec{v}$ 

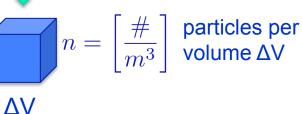
 $q\vec{v}$ 

**MANY POINT CHARGES** 



N particles





 $Nq\vec{v}$  $(n\Delta V)q\vec{v}$ 

**CHUNK OF WIRE** Blast from the Past:



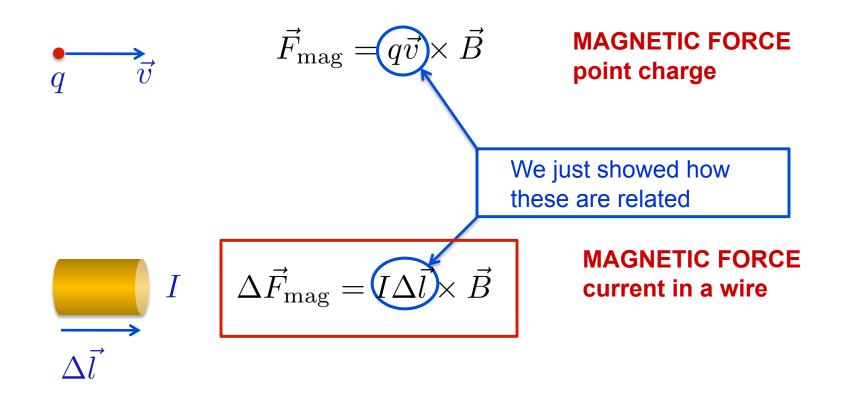
$$\Delta \vec{l}$$
 = length of this chunk of wire

Move the vector symbol

$$q\vec{v}n\Delta V = q\vec{v}nA\Delta l$$

$$=\underline{q|v|nA}\Delta \vec{l}$$
  $I\Delta \vec{l}$   $\equiv I$ 

# Magnetic Force on a charge or wire

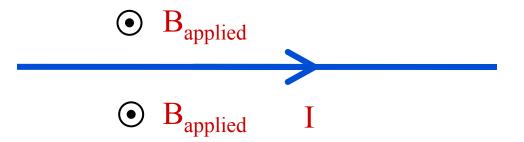


 $\Delta \vec{l}$  = length of this chunk of wire

## Magnetic Force on a Wire

$$\Delta \vec{F}_{\rm mag} = I \Delta \vec{l} \times \vec{B}$$

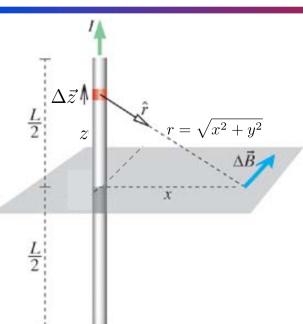
Current-carrying wire in an applied magnetic field B



Which way will the wire move?

# Magnetic Field of a Straight Wire

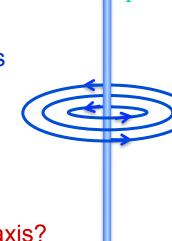
$$|B_z| = \left(rac{\mu_o}{4\pi}
ight) rac{IL}{x\sqrt{x^2+(L/2)^2}}$$
 B in the bisecting plane



Which direction does B point?

→ Always along concentric circles

In cylindrical coordinates, it points in the "  $\hat{\theta}$  " direction



Will the y axis look different from the x axis?

No, so we can trade  $x \rightarrow r$ 

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r\sqrt{r^2 + (L/2)^2}} \hat{\theta}$$

B of a Long Straight Wire (cylindrical coord.)

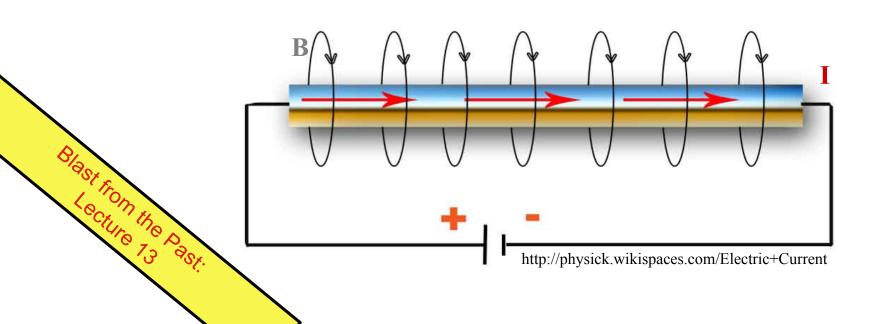
Clast from the Past.

## **Very Close to the Wire**

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r\sqrt{r^2 + (L/2)^2}} \hat{\theta}$$

Very close to the wire: r << L  $\sqrt{r^2+(L/2)^2} pprox L/2$ 

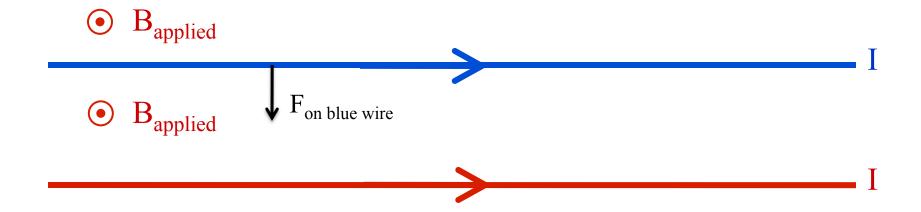
$$\Rightarrow \vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r(L/2)} \hat{\theta} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta} = \vec{B} \quad \begin{array}{c} \text{CLOSE TO} \\ \text{THE WIRE} \end{array}$$



## **Force Between Parallel Wires**

$$\Delta \vec{F}_{\text{mag}} = I \Delta \vec{l} \times \vec{B}$$
  $\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{x} \hat{\theta}$ 

B<sub>applied</sub> = Magnetic field applied by the **red** wire. Blue wire feels a force down.

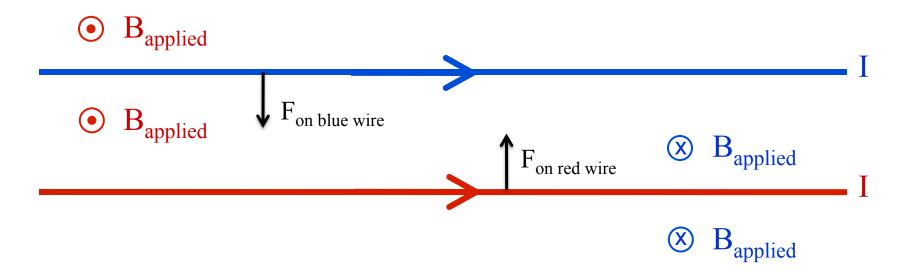


What about reciprocity? (Equal and opposite forces)

### **Force Between Parallel Wires**

$$\Delta \vec{F}_{\text{mag}} = I \Delta \vec{l} \times \vec{B}$$
  $\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{x} \hat{\theta}$ 

B<sub>applied</sub> = Magnetic field applied by the **red** wire. Blue wire feels a force down.

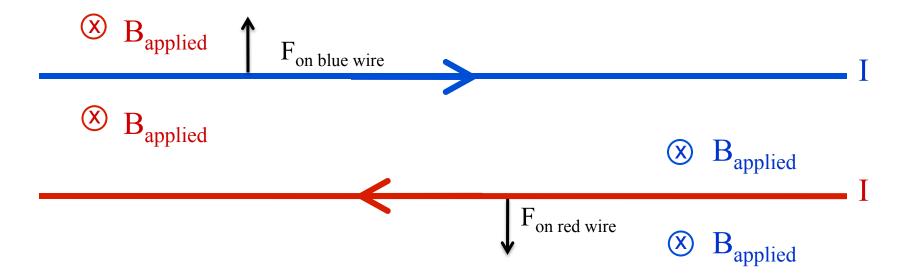


B<sub>applied</sub> = Magnetic field applied by the **blue** wire. **Red** wire feels a force up.

# Force Between (Anti) Parallel Wires

$$\Delta \vec{F}_{\text{mag}} = I \Delta \vec{l} \times \vec{B}$$
  $\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{x} \hat{\theta}$ 

B<sub>applied</sub> = Magnetic field applied by the **red** wire. Blue wire feels a force up.



 $B_{applied}$  = Magnetic field applied by the **blue** wire. Red wire feels a force down.

### **Hall Effect**

By measuring the Hall effect for a particular material, we can determine the sign of the moving particles that make up the current

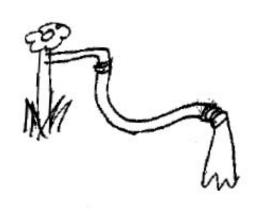
Why would it be anything other than electrons? (Negative charges)

**Semiconductors**: sometimes current is carried by electrons, but sometimes it is carried by the "holes".

In **semiconductors**, "holes" (missing electrons) in the electron sea behave like positive charges.

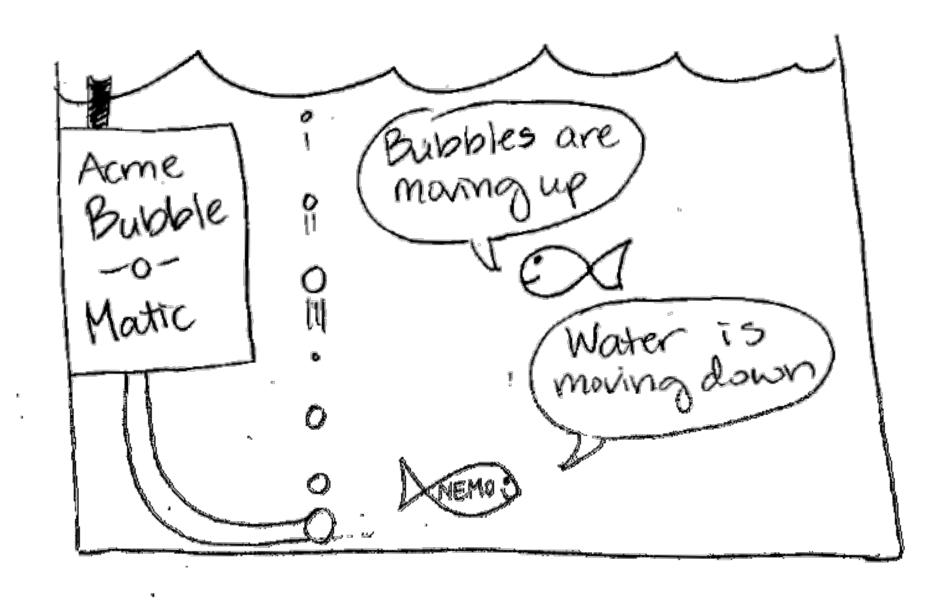
## **ZEN KOAN**

#### Water from a hose



Zen riddle: Is water coming out of the hose, or is the absence of water moving into the hose?

"Holes" not useful here



"Holes" useful here

In a semiconductor, "Holes" in the electron sea act like positive charges.

### **Hall Effect and Electrons**

$$\Rightarrow$$

$$E_y = E_{\rm Hall}$$

 $E_{Hall}$  points down due to buildup of charge.

### **Hall Effect and Holes**

$$F = +|e|(\vec{E} + \vec{v} \times \vec{B})$$

$$\bullet \quad \mathsf{B}$$

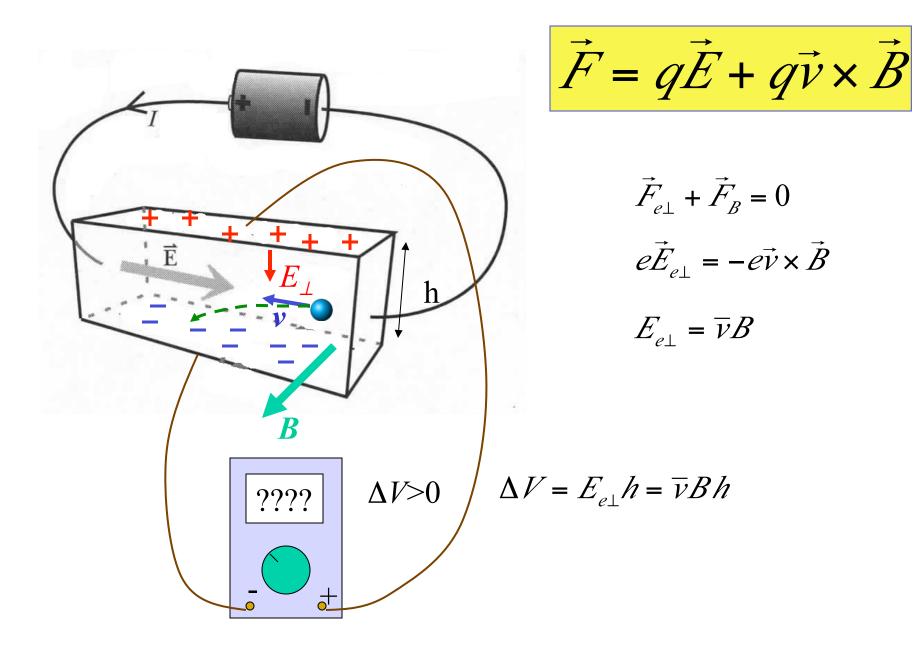
$$+|e|\vec{v} \times \vec{B} \quad +e \quad +++++++$$

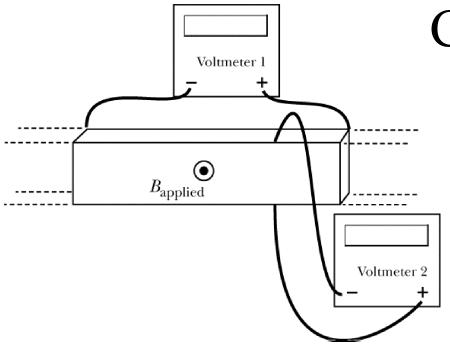
$$+HOLES$$

$$E_y = E_{\mathrm{Hall}}$$

$$E_y = E_{\rm Hall}$$
 E<sub>Hall</sub> points up due to buildup of charge.

## **Hall Effect**





# Clicker Question

Voltmeter 1 reading is POSITIVE Voltmeter 2 reading is POSITIVE

Mobile charges are:

- A) Positive (holes)
- B) Negative (electrons)
- C) Not enough information