Last Time

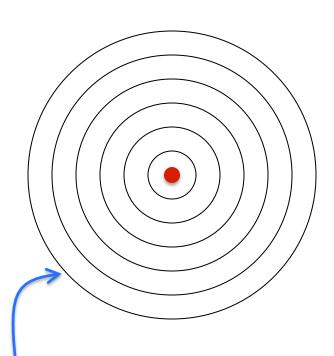
Gauss' Law

Today

Gauss' Law: Examples
"Magnetic Gauss Law"
Ampere's Law

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\text{inside}} \quad \frac{\text{Gauss'}}{\text{Law}}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_o} \frac{q}{r^2} \hat{r} \stackrel{\text{E-field}}{\text{Point Charge}}$$



Gauss' Law for Point Charge:

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{4\pi\epsilon_o} \frac{q}{r^2} (4\pi r^2) = \frac{q}{\epsilon_o}$$

Works for any size sphere because r cancels

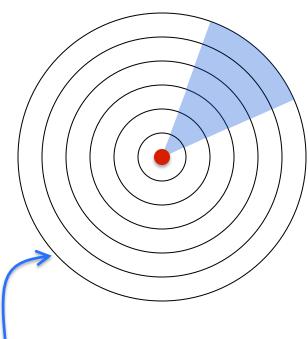
On each sphere:

Field
$$E \propto \frac{1}{r^2}$$

Surface Area $A \propto r^2$



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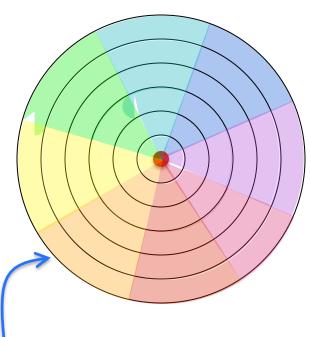
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→ In any segment, the contribution from any "shell" is the **same.** (Like a flashlight.)

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\text{inside}} \quad \text{GAUSS' LAW}$$



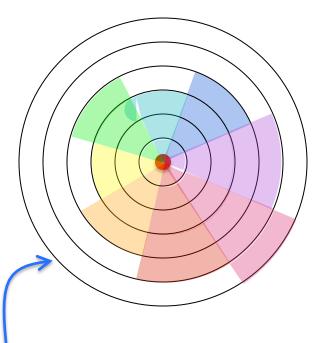
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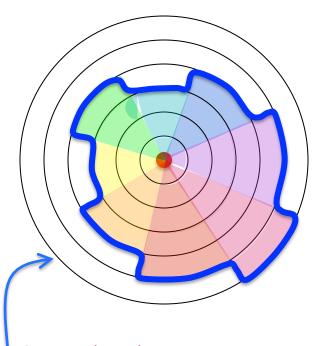
Surround the charge with any shape by following different spheres in different places.

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Surface Area $A \propto r^2$

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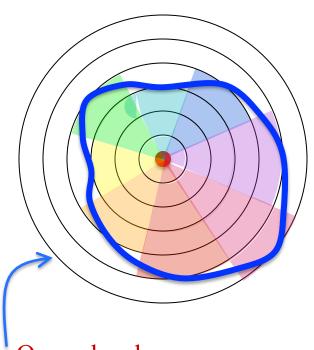
Surface Area $A \propto r^2$

In any segment, the contribution from any "shell" is the **same.** (Like a flashlight.)

Surround the charge with any shape by following different spheres in different places.

Flux through outer surface is always the same.

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\rm inside} \quad \text{GAUSS' LAW}$$



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→ works for any smooth shape.

Example: Gauss' Law for Plane

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\text{inside}} \quad \text{GAUSS' LAW}$$

GIVEN: Infinite plane.

Surface charge density $\sigma = [Q/A]$

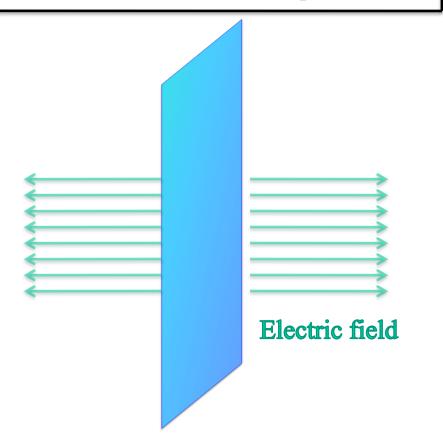
FIND: E-field due to plane

Due to geometry of plane:

E-field is perpendicular to plane

E-field has same magnitude everywhere

→ USE Gauss' Law to find magnitude.



Example: Gauss' Law for Plane

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\text{inside}}$$
 GAUSS' LAW

GIVEN: Infinite plane.

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FIND: E-field due to plane Due to geometry of plane:

E-field is perpendicular to plane

E-field has same magnitude everywhere

→ USE Gauss' Law to find magnitude.

Choose "Gaussian box" wisely! Same area A on ends of cylinders

E is constant and perpendicular

to the surface A of the "Gaussian box":

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} q_{\text{inside}}$$

$$2EA = \frac{1}{\epsilon_o} \sigma A$$

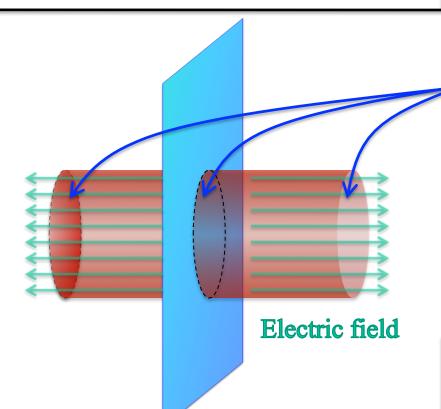
$$q_{\text{inside}} = [Q/A]*A$$

$$= \sigma A$$

$$E = \frac{1}{2\epsilon_o}\sigma$$

 $E = \frac{1}{2\epsilon_o}\sigma \left| \begin{array}{c} \text{ELECTRIC FIELD} \\ \text{of a PLANE} \end{array} \right.$





Gauss' Law for Magnetism?

$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\rm inside} \quad \mbox{GAUSS' LAW for charge}$$

So far, no experiment has found a "magnetic charge" (a.k.a. magnetic monopole)

$$q_{
m magnet} = 0$$
 Big fat ZERO!



→ Gauss' Law for Magnetism is simpler:

$$\oint \vec{B} \cdot \hat{n} dA = 0$$

GAUSS' LAW FOR MAGNETISM

Next Up: Ampere's Law

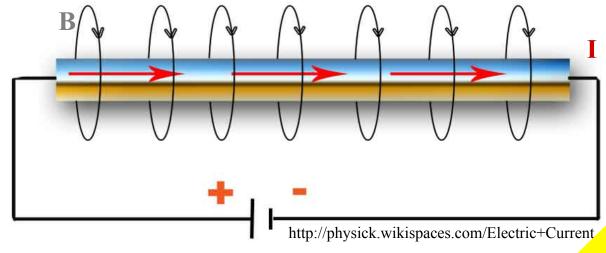
First review Biot-Savart Law

Very Close to the Wire

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r\sqrt{r^2 + (L/2)^2}} \hat{\theta}$$

Very close to the wire: r << L $\sqrt{r^2+(L/2)^2} pprox L/2$

$$\Rightarrow \vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r(L/2)} \hat{\theta} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta} = \vec{B} \quad \begin{array}{c} \text{CLOSE TO} \\ \text{THE WIRE} \end{array}$$



Alast From the 13

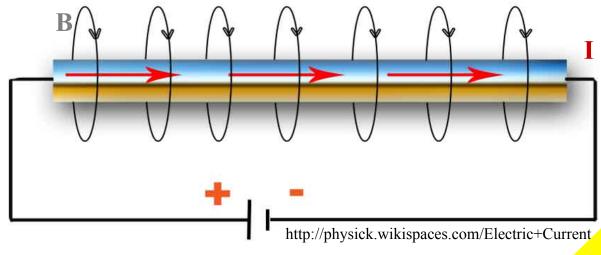
Very Long Wire

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r\sqrt{r^2 + (L/2)^2}} \hat{\theta}$$

Very *Long* wire: L >> r
$$\sqrt{r^2 + (L/2)^2} pprox L/2$$

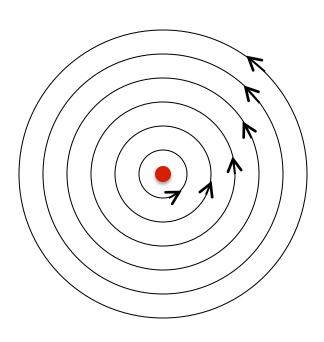
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VERY LONG WIRE



Very Long Wire

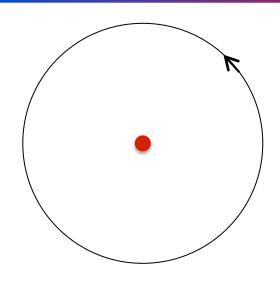
$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta}$$



Viewed from the end Current coming out of board

Very Long Wire

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta}$$



 $d\vec{l}$ is along our circle

$$d\vec{l} = rd\hat{\theta}$$

$$\oint \hat{\theta} \cdot d\vec{l} = r \oint \hat{\theta} d\hat{\theta} = 2\pi r$$

Viewed from the end Current coming out of board

Cylindrical pattern of B-field

→ Let's take a line integral along one circle

$$\oint \vec{B} \cdot d\vec{l} = \oint \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta} \cdot d\vec{l}$$

$$= \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \oint \hat{\theta} \cdot d\vec{l}$$

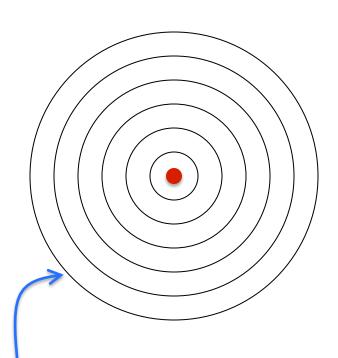
$$= \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} (2\pi r) = \mu_o I$$

$$\oint ec{B} \cdot dec{l} = \mu_o I$$
 AMPERE'S LAW



$$\oint \vec{E} \cdot \hat{n} dA = \frac{1}{\epsilon_o} \sum q_{\text{inside}}$$
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Works for any size sphere because r cancels

On each sphere:

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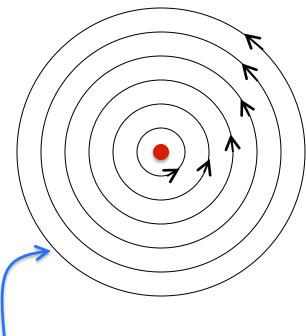
Surface Area $A \propto r^2$

Something similar is going to happen for B of a wire

Remember 183

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta}$$
 Biot-Savart $\oint \vec{B} \cdot d\vec{l} = \mu_o I$ B-field Long Wire

Current coming out of board



Ampere's Law for Long Wire:

$$\oint \vec{B} \cdot d\vec{l} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} (2\pi r) = \mu_o I$$

Works for any size circle because r cancels

On each circle:

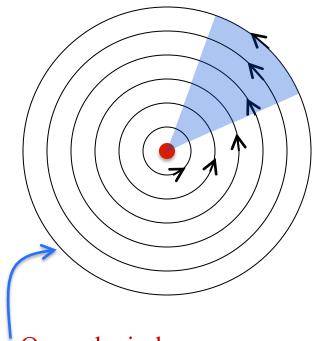
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Circumference $C \propto r$



$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta}$$
 Biot-Savart
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 B-field Long Wire

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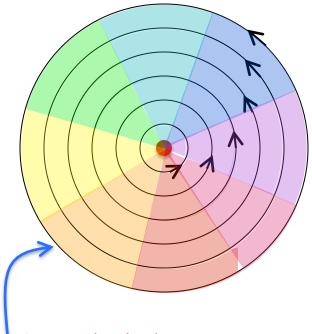
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→ In any segment, the contribution from any circle is the same. (Like a flashlight.)

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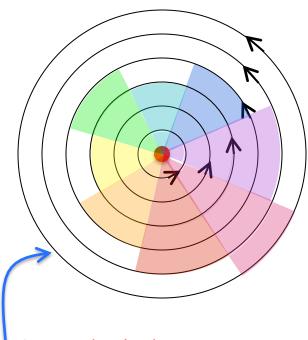
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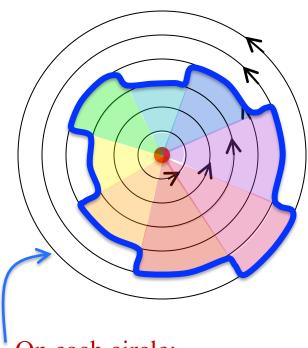
Surround the **wire** with any shape by following different **circles** in different places.

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Biot-Savart Law

$$\oint \vec{B} \cdot d\vec{l} = \mu_o I \qquad \begin{array}{c} \text{B-field} \\ \text{Long Wire} \end{array}$$

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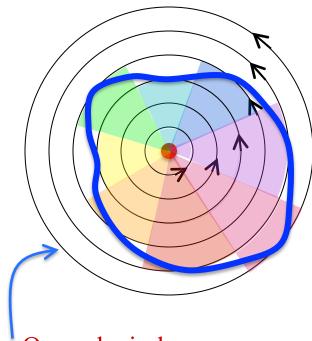
Line integral around outer surface is always **same**.

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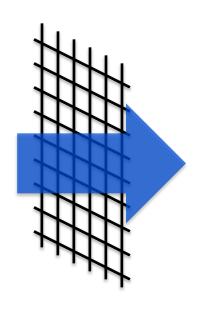
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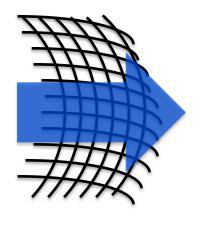
Limit of small segments → works for any smooth path

Electric Current is like Water Flow



Water flows through a net at a certain rate.

Molecules/second through each square



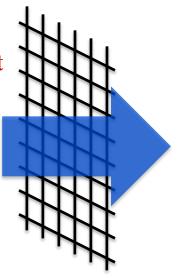
The net can deform, but the flow rate is the same.

Same number of molecules/second through each square

Ampere's Law

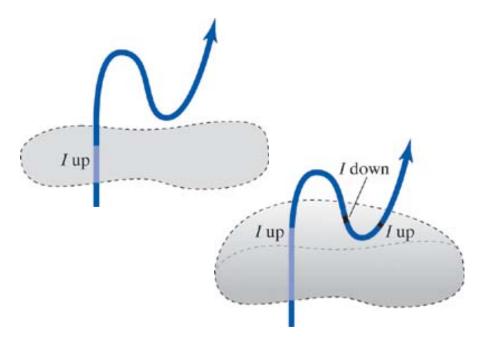
Water flows through a net at a certain rate.

Charges/second through each square



Likewise, Ampere's Law works even when the surface through which current I "pokes" is deformed, even if it means the wire "pokes" through in more places. The net can deform, but the flow rate is the same.

Same number of charges/second through each square



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Ampere's Law