

Exam 3 multiple-choice part on CHIP

Please check your scores/answers on CHIP – if you see '0', let us know ASAP.

Finalizing iClicker scores 10-22

Scores for **Lectures 10-22** have been uploaded. Deadline for requesting corrections is **5 PM this Friday** (April 20).

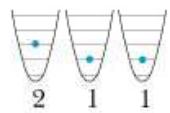
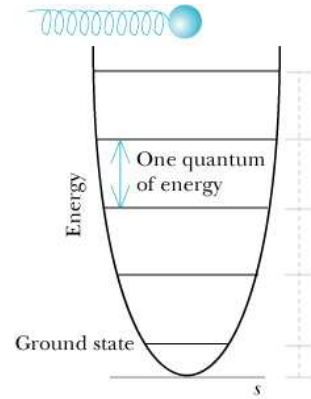
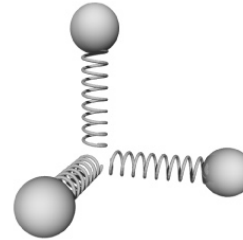
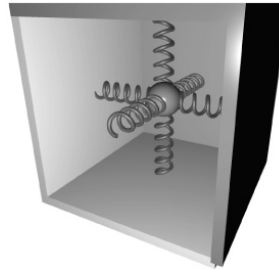
$$\Delta \vec{p} = \vec{F} \Delta t$$

$$\Delta E = W + Q$$

$$\Delta \vec{L} = \vec{\tau} \Delta t$$

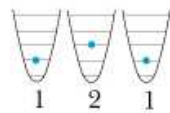
Summary: Foundations

Einstein Model of Solids



= $\bullet \bullet | \bullet | \bullet$ # microstates

$$\equiv \Omega = \frac{(q + N - 1)!}{q!(N - 1)!}$$

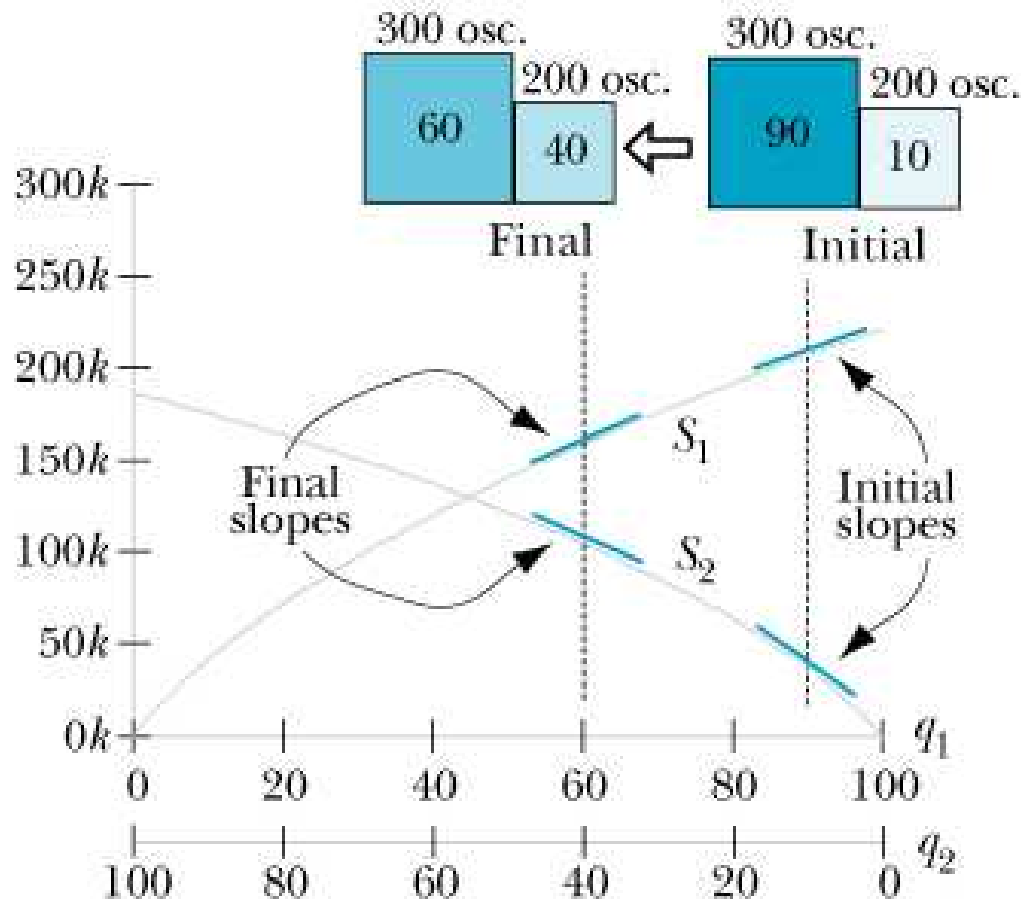


= $\bullet | \bullet \bullet | \bullet$ (N oscillators, q quanta)

Fundamental assumption of statistical mechanics

Over time, an isolated system in a given macrostate (total energy) is equally likely to be found in any of its microstates (microscopic distribution of energy).

Summary: Entropy and Temperature



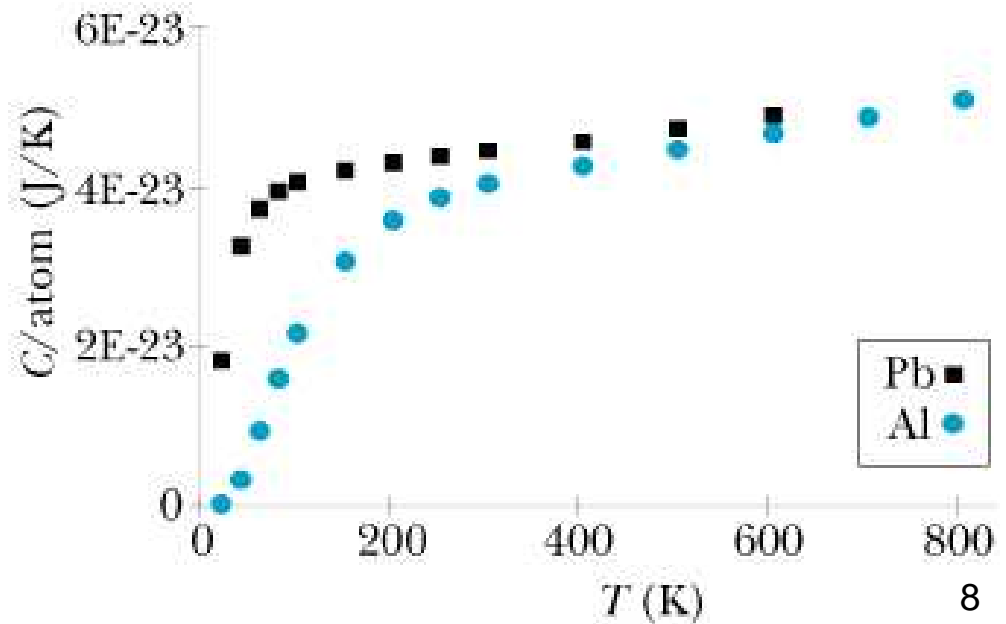
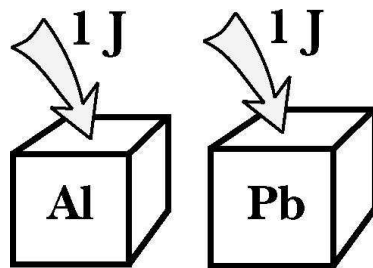
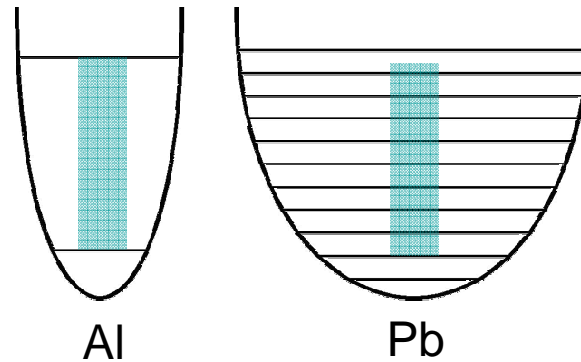
$$S \equiv k \ln \Omega$$

$$\frac{1}{T} \equiv \frac{dS}{dE_{\text{int}}}$$

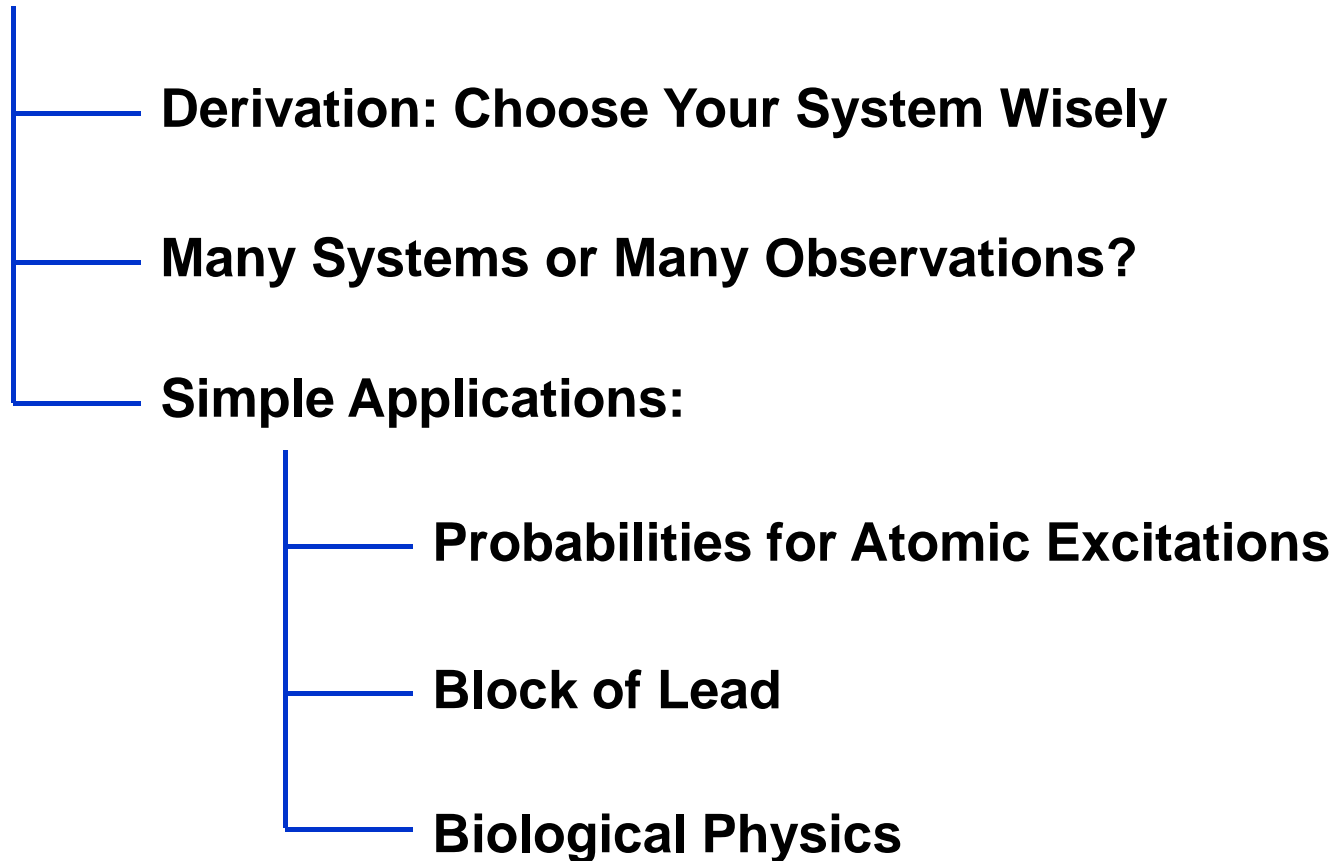
If the initial state is not the most probable, energy is exchanged until the most probable distribution is reached.

Summary: Specific Heat

$$C_{atom} = \frac{\Delta E_{atom}}{\Delta T} \equiv \frac{\Delta E_{system} / N_{atoms}}{\Delta T}$$



Today: The Boltzmann Distribution

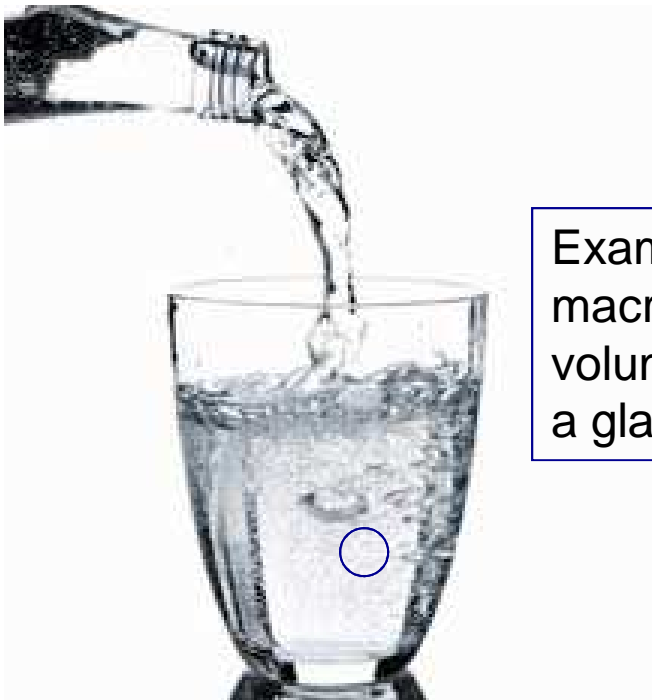


Boltzmann Distribution

The Boltzmann distribution comes about by cleverly picking our system.

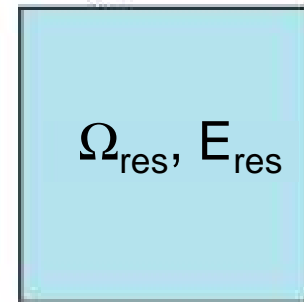
Say that we want to analyze a cup of water.

What system do we pick? A tiny volume of the water!



Example: a
macroscopic
volume of water in
a glass of water

A large "reservoir"

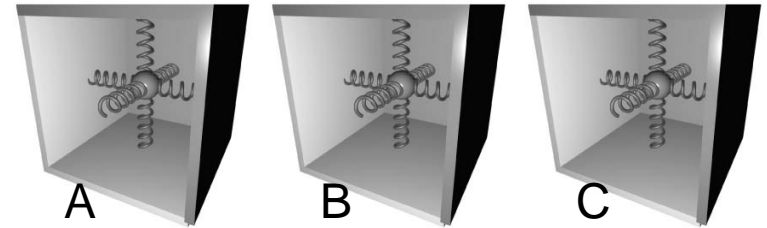


■ Ω, E

A small system

$$E_{\text{tot}} = E_{\text{res}} + E = \text{constant}$$

Recall the Pb Nanoparticle (3 atoms) from Ch 11, p 392-393:



q	Ω	$\ln \Omega$	S
4	465	6.20	$6.20k_B$
5	1287	7.16	$7.16k_B$
6	3003	8.01	$8.01k_B$

q	T(K)
4-5	58.9
5	62.5
5-6	66.6

Temperature: $\frac{1}{T} \equiv \frac{\partial S}{\partial E_{\text{int}}}$

$$T_q \approx \frac{\Delta E}{\Delta S} = \frac{\Delta q \times \hbar \omega_o}{k \Delta(\ln \Omega)}$$

$$C = \frac{1}{3} \frac{\Delta E(3 \text{ atoms})}{\Delta T} = \frac{1}{3} \frac{8 \times 10^{-22}}{(66.6 - 58.9)} = 3.4 \times 10^{-23} \text{ J/K per atom}$$

What if we wanted to know what is happening on **one** of the three atoms?
For example, with $q=6$, how often is atom A found with $q_A=2$?

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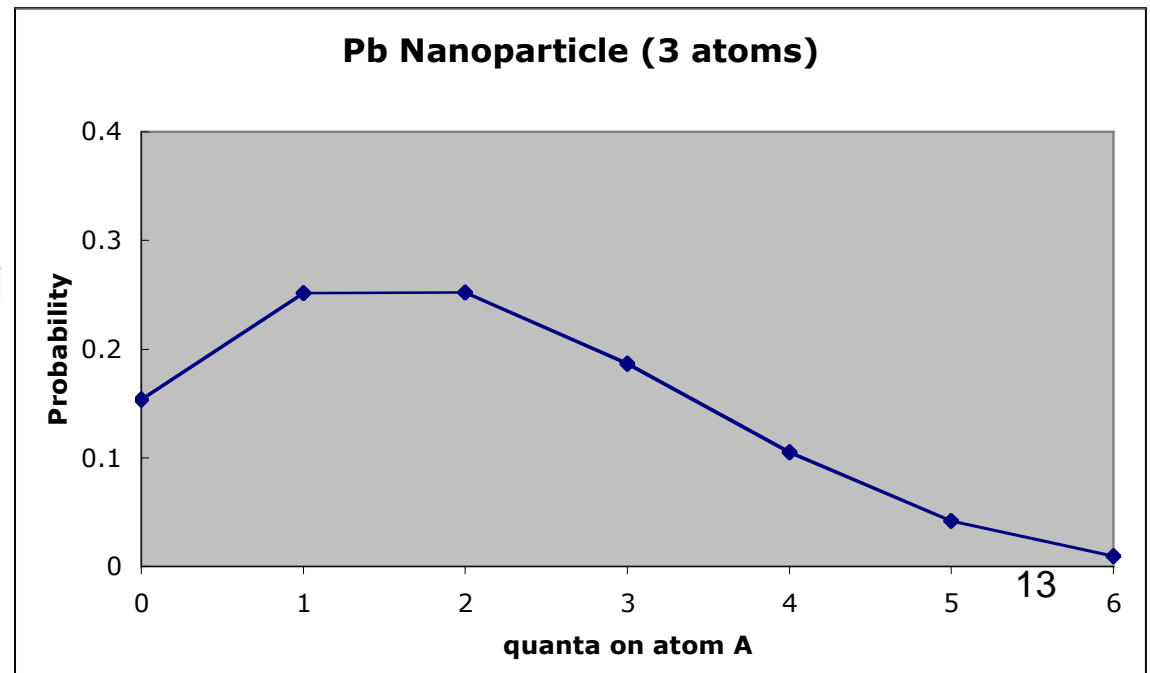
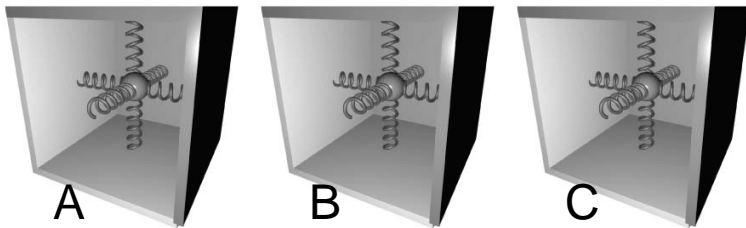
We can simply count the states with 2 quanta on atom A:

$$\Omega_A(q_A = 2) = \frac{(2 + 3 - 1)!}{2!(3 - 1)!} = 6$$

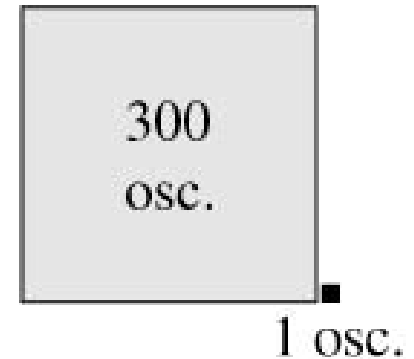
$$\Omega_{BC}(q_{BC} = 6 - 2) = \frac{(4 + 6 - 1)!}{4!(6 - 1)!} = 126$$

$$\Omega_A(2)\Omega_{BC}(4) = 6 \times 126 = 756$$

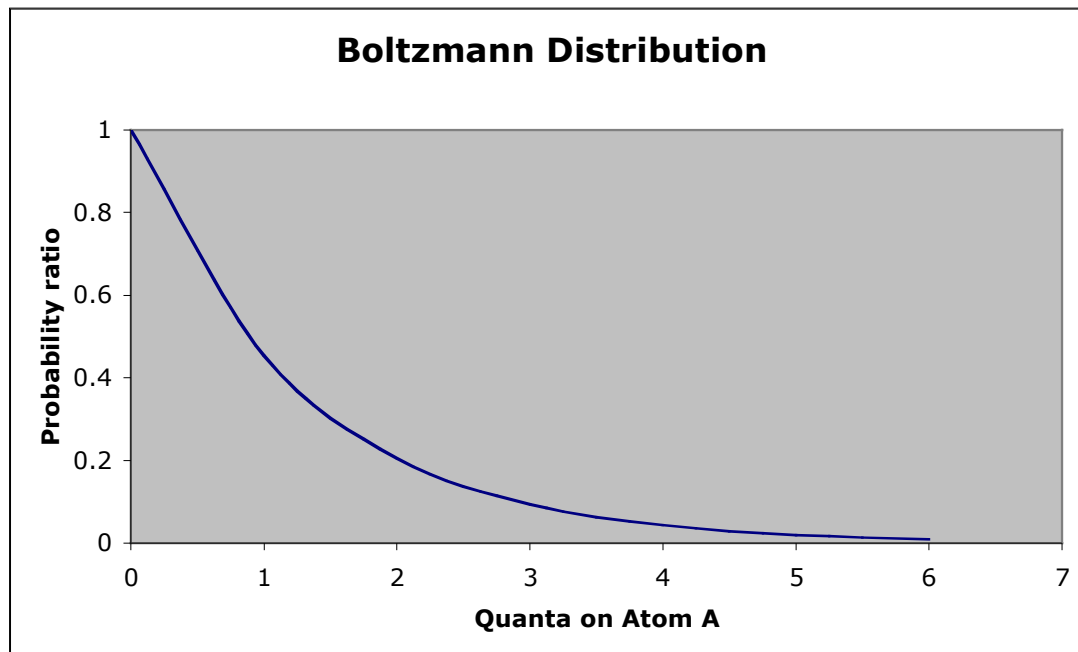
$$P = \frac{\Omega_A(2)\Omega_{BC}(4)}{\Omega_{total}(6)} = \frac{756}{3003} \approx 0.25$$



Now, a more interesting question: What is the probability of finding 2 quanta on atom A if it is in contact with a large Pb block? (Assume same T as before.)



$$\frac{\text{Probability of } \Delta E \text{ above ground state}}{\text{Probability of being in ground state}} \sim \Omega_A(\Delta E) e^{-\frac{\Delta E}{kT}} \sim e^{-\frac{\Delta E}{kT}}$$



This is a property of **any** small system in thermal equilibrium with a large reservoir at a fixed temperature.

Proof (next few slides)

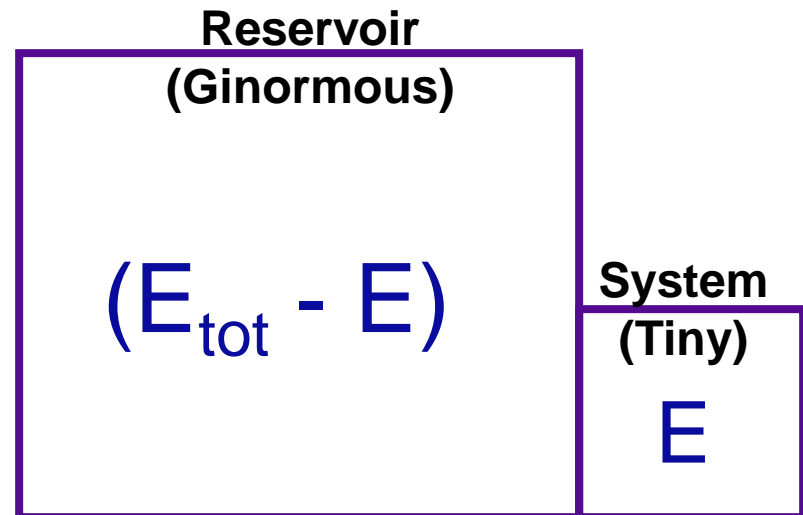
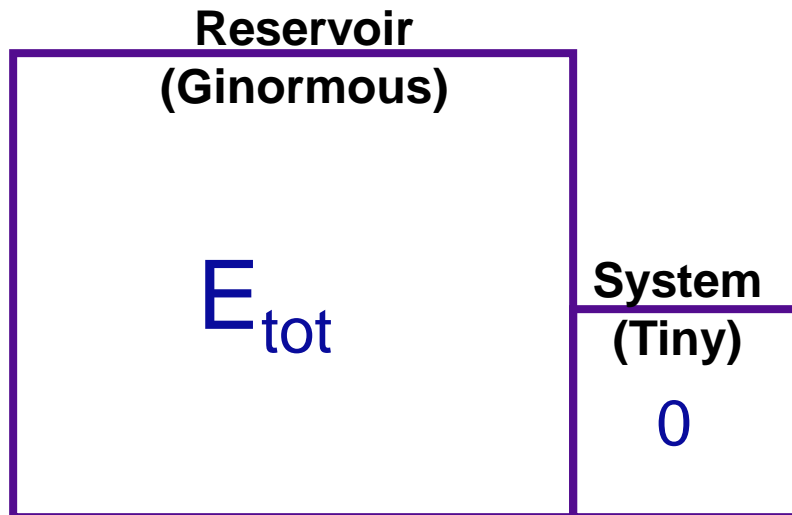
“Reservoir”

*A **ginormous** system
with a constant **temperature***

tiny
system

Probabilities

What is the probability of finding the system in this state?



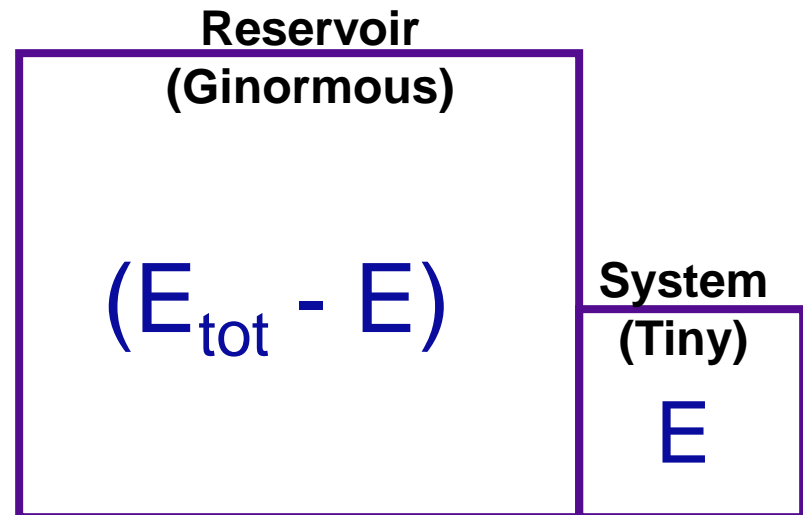
Probabilities

Number of ways of
having $E_{\text{sys}} = E$

We can find the probability
from the number of ways
of arranging the energy.

$$P(E) = \frac{\Omega_{\text{res}}(E_{\text{res}})\Omega(E)}{\Omega_{\text{tot}}(E_{\text{tot}})}$$

Total number of ways
of arranging energy
in combined system:
Tot = Res + Sys



Now Do Some Math...

$$P(E) = \frac{\Omega_{res}(E_{res})\Omega(E)}{\Omega_{tot}(E_{tot})}$$

We can find the probability from the number of ways of arranging the energy.

$$S = k\ln(\Omega)$$

$$\begin{aligned} k\ln P &= k\ln\left(\frac{\Omega_{res}(E_{res})\Omega(E)}{\Omega_{tot}(E_{tot})}\right) \\ &= k\ln(\Omega_{res}(E_{res})) + k\ln(\Omega(E)) - k\ln(\Omega_{tot}(E_{tot})) \\ &= S_{res}(E_{res}) + S(E) - S_{tot}(E_{tot}) \end{aligned}$$

More Math...

$$k \ln P = S_{res}(E_{res}) + S(E) - S_{tot}(E_{tot})$$

$$E_{res} = E_{tot} - E$$

Use a **Taylor Expansion** for the Ginormous Reservoir: ($E \ll E_{res}$)

$$S_{res}(E_{res}) \approx S_{res}(E_{tot}) - \frac{dS_{res}}{dE_{res}} E + \dots = S_{res}(E_{tot}) - \frac{E}{T}$$

$\frac{1}{T}$ **Spot The Temperature!**

$$k \ln P = S_{res}(E_{tot}) - \frac{E}{T} + S(E) - S_{tot}(E_{tot})$$

$$= k \ln(\Omega_{res}(E_{tot})) - \frac{E}{T} + k \ln(\Omega_{sys}(E)) - k \ln(\Omega_{tot}(E_{tot}))$$

More Math...

$$k \ln P = k \ln(\Omega_{res}(E_{tot})) - \frac{E}{T} + k \ln(\Omega_{sys}(E)) - k \ln(\Omega_{tot}(E_{tot}))$$

Exponentiate Everything...

$$P = e^{\ln P} = e^{\text{const.}} e^{-E/kT} \Omega_{sys}(E)$$

$$P \propto e^{-E/kT}$$

BOLTZMANN FACTOR

Very very important!

Boltzmann Distribution

The probability of finding energy E in a small system in contact with a large reservoir is

$$P(E) \propto \Omega(E) \cdot e^{-\frac{E}{kT}}$$

The exponential part, $e^{-E/kT}$, is called the “Boltzmann factor.”

$\Omega(E)$ is the number of microstates in the small system at energy E .

In many circumstances, $\Omega(E)$ changes so slowly compared to $e^{-E/kT}$ that it is essentially constant:

$$P(E) \approx \text{const} \cdot e^{-\frac{E}{kT}}$$

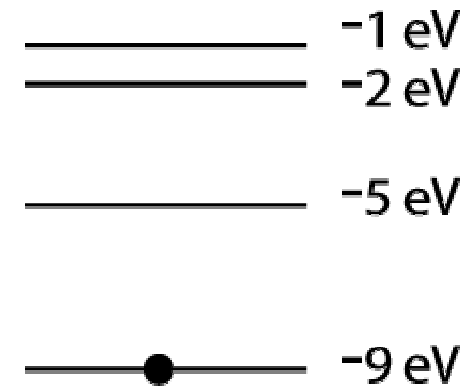
A large "reservoir"



Ω, E
A small system

Application: Atomic Excitations

How likely is an atom to be in 1st excited, compared to odds of being in ground state?



$$\frac{P(E_1)}{P(E_0)} = \frac{e^{-E_1/kT}}{e^{-E_0/kT}} = e^{-\Delta E/kT}$$

NOTE: kT at room temp = 1/40 eV.

For the above atom, odds of being in first excited state are

$$P(1^{\text{st}} \text{ excited state}) = \frac{1}{e^{\Delta E / 1/40 \text{ eV}}} = \frac{1}{e^{40\Delta E}} = \frac{1}{e^{40 \times 4}} = 3 \times 10^{-70}$$

Typical atomic energy gaps are big compared to room temp.

A room-temperature box of neon doesn't glow! (Unless you add energy)₂₃

Application: Block of Lead

$$\frac{\hbar\omega}{k} = \frac{8 \times 10^{-22} J}{k} = 57K$$

T=300K at room temperature

$$P \propto e^{-E/kT}$$

BOLTZMANN FACTOR
Very very important!

Think of adding q quanta of energy to one lead atom...

At room temperature, the Boltzmann factors for exciting vibrations are:

$$q = 0 \quad P \propto e^{-q\hbar\omega/kT} = e^{0/300} = 1 \quad \leftarrow \text{Most probable state}$$

$$q = 1 \quad P \propto e^{-q\hbar\omega/kT} = e^{-57/300} = 0.83$$

$$q = 2 \quad P \propto e^{-q\hbar\omega/kT} = e^{-2*57/300} = 0.68$$

$$q = 3 \quad P \propto e^{-q\hbar\omega/kT} = e^{-3*57/300} = 0.57$$

Watch out!
These are
relative
probabilities

Application: Block of Lead

$$\hbar\omega = 8 \times 10^{-22} \text{ J} = 57 \text{ K}$$

T=300K at room temperature

Think of adding q quanta of vibrational energy to one lead atom...

The actual probabilities use a normalization factor “Z” (partition function).

$$q = 0 \quad P = e^{-q\hbar\omega/kT} / Z = 17\%$$

$$q = 1 \quad P = e^{-q\hbar\omega/kT} / Z = 14\%$$

$$q = 2 \quad P = e^{-q\hbar\omega/kT} / Z = 12\%$$

$$q = 3 \quad P = e^{-q\hbar\omega/kT} / Z = 9.8\%$$

$$q = 4 \quad P = e^{-q\hbar\omega/kT} / Z = 8.1\%$$

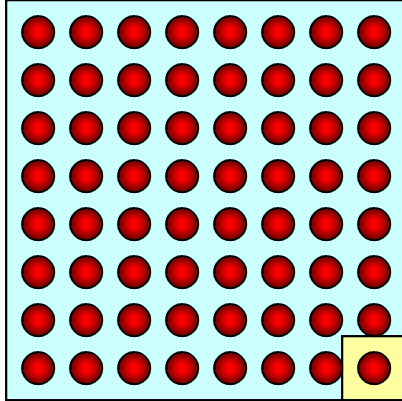
$$q = 5 \quad P = e^{-q\hbar\omega/kT} / Z = 6.7\%$$

$$q = 6 \quad P = e^{-q\hbar\omega/kT} / Z = 5.5\%$$

Most likely to find the vibrations are excited.

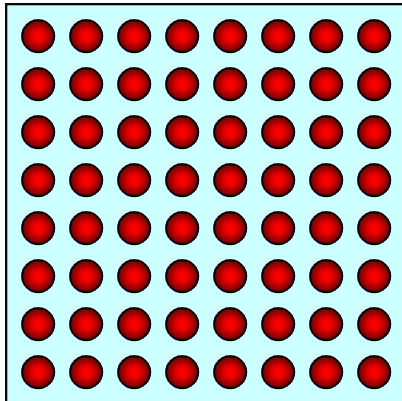
On average, $q=5$.

Boltzmann distribution arises in 2 ways



Measure energy on a single oscillator at many different times. Number of measurements at energy E is proportional to:

$$e^{-\frac{E}{kT}}$$

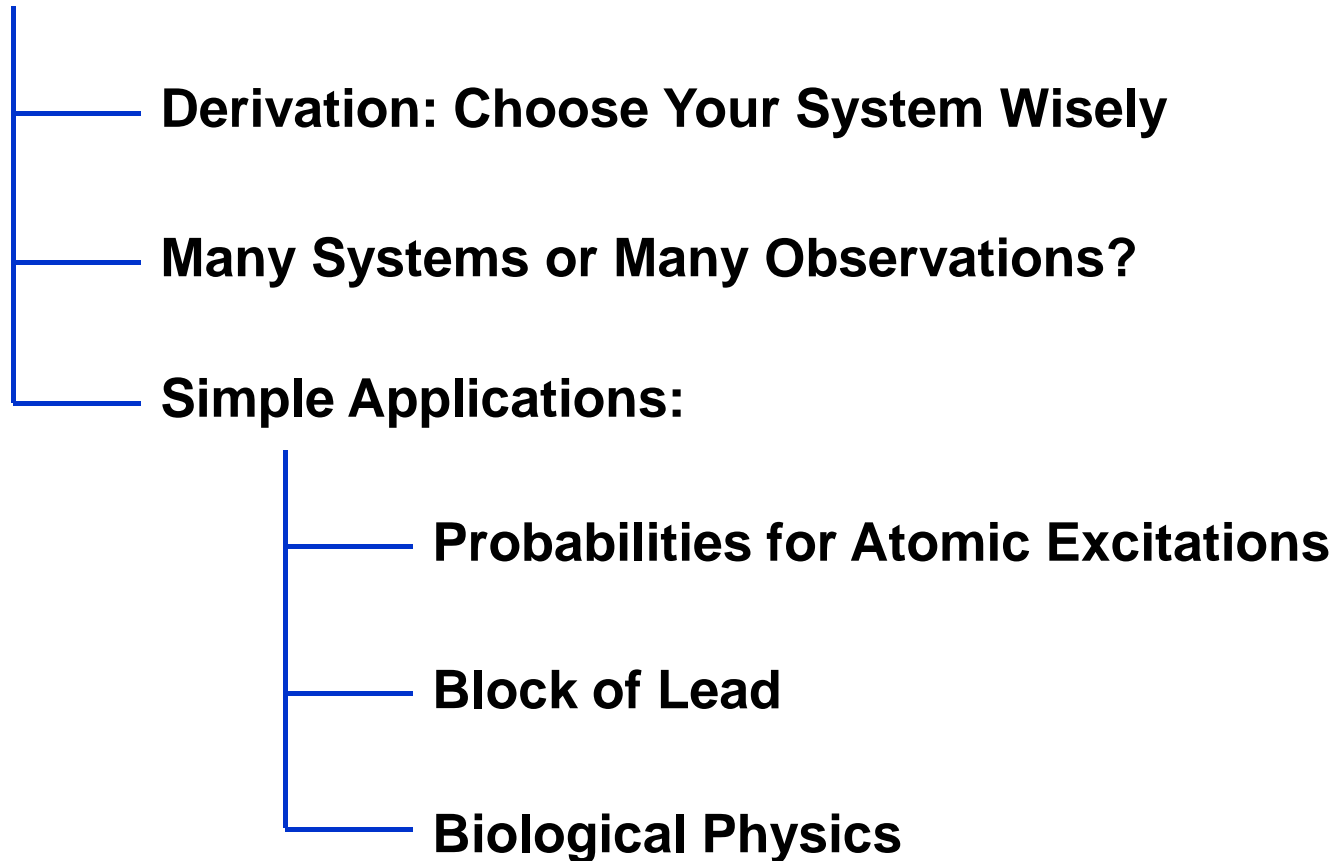


Measure energy on each oscillator in a system at one single time. Number of oscillators at energy E is proportional to:

$$e^{-\frac{E}{kT}}$$

Distribution of oscillator energies in a large system also follows Boltzmann distribution!

Today: The Boltzmann Distribution



Next Time: Boltzmann Applications

- Speed Distribution in a Gas
- Energy Equipartition and Specific Heat
- Pressure and the Ideal Gas Law