Key Ideas in Chapter 18: Magnetic Field

- Moving charged particles make a magnetic field, which is different from an electric field.
- The needle of a magnetic compass aligns with the direction of the net magnetic field at its location.



- Electron current is a number of electrons per second entering a section of a conductor.
- Conventional current (Coulombs/second) is opposite in direction to the electron current, and is assumed to be due to positively charged particles.
- The superposition principle can be applied to calculate the expected magnetic field from current-carrying wires in various configurations.
 - A current-carrying loop is a magnetic dipole.
 - A bar magnet is also a magnetic dipole.
 - Even a single atom can be a magnetic dipole!





Last Time

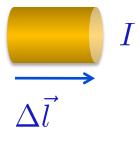
- (Cross Products: Mathematically)
- Electron Current and Conventional Current
- Calculating the Electron Current
- True vs. Useful
- Biot-Savart Law in a Wire
- Relativity??

Biot-Savart Law

$$\overrightarrow{q}$$
 \overrightarrow{v}

$$ec{B} = \left(rac{\mu_o}{4\pi}
ight)rac{qec{v} imes \hat{r}}{|r|^2}$$
 BIOT-SAVAR point charge

BIOT-SAVART LAW



$$\Delta ec{B} = \left(rac{\mu_o}{4\pi}
ight)rac{I\Deltaec{l} imes\hat{r}}{|r|^2}$$
 BIOT-SAVART LA current in a wire

BIOT-SAVART LAW

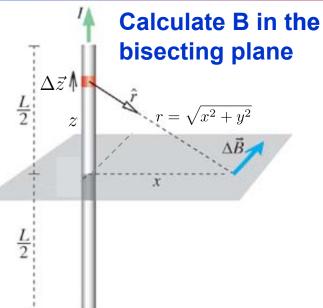
$$\Delta \vec{l}$$
 = length of this chunk of wire

Today

- Magnetic Field of a Straight Wire
- (Magnetic Field of a Current Loop)
- Magnetic Dipole Moment
- Bar Magnet
- Atomic Dipoles

$$I$$
 $\Delta \vec{l}$

$$\vec{I} = \frac{1}{2} \vec{l} \qquad \vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{I\Delta \vec{l} \times \hat{r}}{|r|^2} \qquad \begin{array}{c} \text{BIOT-SAVART LAW} \\ \text{current in a wire} \end{array}$$



$$\vec{r} = \langle x, 0, z \rangle \implies |r|^2 = (x^2 + z^2)$$

$$\hat{r} = \frac{\vec{r}}{|r|} = \frac{\langle x, 0, z \rangle}{\sqrt{x^2 + z^2}}$$

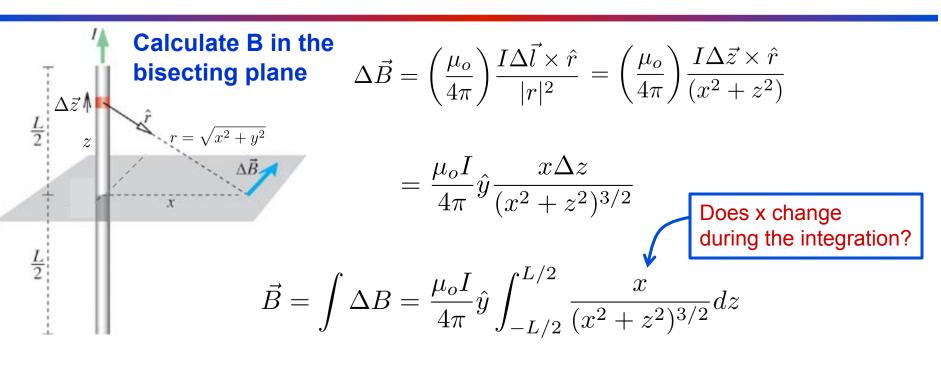
$$\Delta \vec{l} = \Delta \vec{z} = \langle 0, 0, \Delta z \rangle$$

$$\Rightarrow \Delta \vec{z} \times \hat{r} = \frac{1}{|r|} (\Delta \vec{z} \times \vec{r})$$

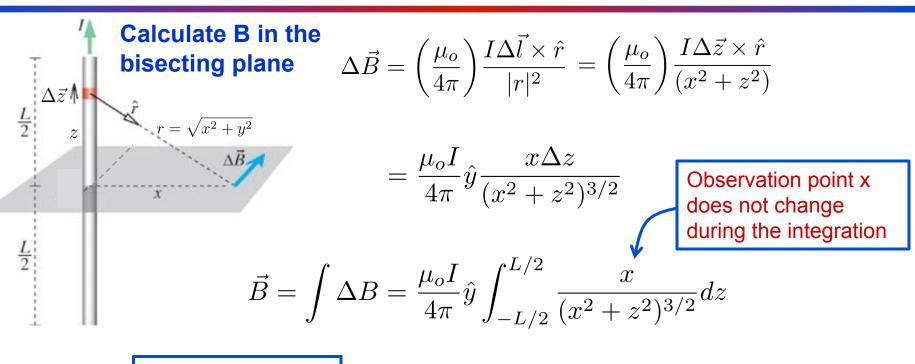
$$\Delta \vec{z} \times \hat{r} = \frac{1}{|r|} \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & \Delta z \\ x & 0 & z \end{vmatrix} = \frac{0\hat{x} + x\Delta z\hat{y} + 0\hat{z}}{|r|}$$

$$\Delta \vec{z} \times \hat{r} = \frac{x\Delta z}{\sqrt{x^2 + z^2}} \hat{y}$$

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{I\Delta \vec{l} \times \hat{r}}{|r|^2} \qquad |r|^2 = (x^2 + z^2) \qquad \Delta \vec{l} = \Delta \vec{z} \qquad \Delta \vec{z} \times \hat{r} = \frac{x\Delta z}{\sqrt{x^2 + z^2}} \hat{y}$$



$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{I\Delta \vec{l} \times \hat{r}}{|r|^2} \qquad |r|^2 = (x^2 + z^2) \qquad \Delta \vec{l} = \Delta \vec{z} \qquad \Delta \vec{z} \times \hat{r} = \frac{x\Delta z}{\sqrt{x^2 + z^2}} \hat{y}$$

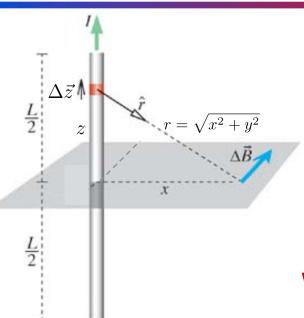


Look up the integral

$$|B_z| = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{x\sqrt{x^2 + (L/2)^2}}$$

B of a Long **Straight Wire**

$$|B_z| = \left(rac{\mu_o}{4\pi}
ight) rac{IL}{x\sqrt{x^2+(L/2)^2}}$$
 B in the bisecting plane



Which direction does B point?

→ Always along concentric circles

In cylindrical coordinates, it points in the " $\hat{\theta}$ " direction



No, so we can trade $x \rightarrow r$

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r\sqrt{r^2 + (L/2)^2}} \hat{\theta}$$

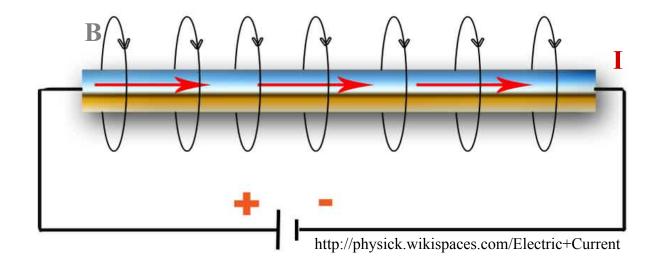
B of a Long Straight Wire (cylindrical coord.)

Very Close to the Wire

$$\vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r\sqrt{r^2 + (L/2)^2}} \hat{\theta}$$

Very close to the wire: r << L $\sqrt{r^2+(L/2)^2} pprox L/2$

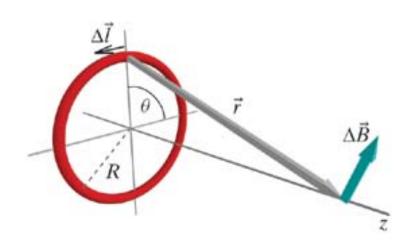
$$\Rightarrow \vec{B} = \left(\frac{\mu_o}{4\pi}\right) \frac{IL}{r(L/2)} \hat{\theta} = \left(\frac{\mu_o}{4\pi}\right) \frac{2I}{r} \hat{\theta} = \vec{B} \quad \begin{array}{c} \text{CLOSE TO} \\ \text{THE WIRE} \end{array}$$



iClicker Question

B along Axis of Circular Loop of Wire

$$\vec{I} = \left(\frac{\mu_o}{4\pi}\right) \frac{I\Delta \vec{l} \times \hat{r}}{|r|^2} \qquad \text{BIOT-SAVART LAW current in a wire}$$



- 1) Draw $\Delta \vec{B}$ for one piece
 - → Notice only z component survives!
- 2) Write $\Delta \vec{B}$ due to one piece

$$\Delta B_z = \frac{\mu_o}{4\pi} \frac{IR^2\Delta\theta}{(R^2+z^2)^{3/2}} \ \ \text{(lots of math!)}$$

3) Integrate $\Delta \vec{B}$ to find total B

$$B_z = \int \Delta B_z = \int_0^{2\pi} \frac{\mu_o}{4\pi} \frac{IR^2 d\theta}{(R^2 + z^2)^{3/2}} = \frac{\mu_o}{4\pi} \frac{IR^2 2\pi}{(R^2 + z^2)^{3/2}} = B_z$$

4) Doulbecheck

$$B_z = \left[\left(rac{T \cdot m}{A}
ight) rac{(m^2 \cdot A)}{(m^2)^{3/2}} \right] \quad \checkmark \quad \text{UNITS} \qquad \text{Right-hand Rule} \quad \checkmark$$

Magnetic Field of a Wire Loop

Special case: far from the loop

$$B_z = \frac{\mu_0}{4\pi} \frac{2\pi R^2 I}{\left(R^2 + z^2\right)^{3/2}}$$

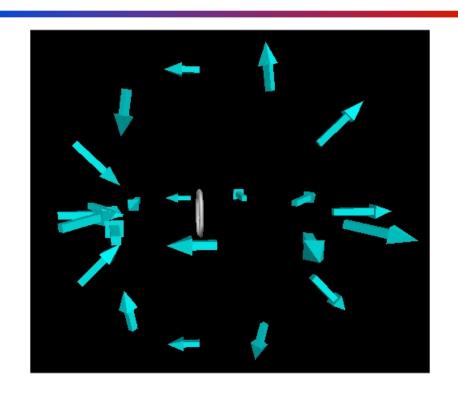
for
$$z >> R$$
: $(z^2 + R^2)^{3/2} \approx (z^2)^{3/2} = z^3$

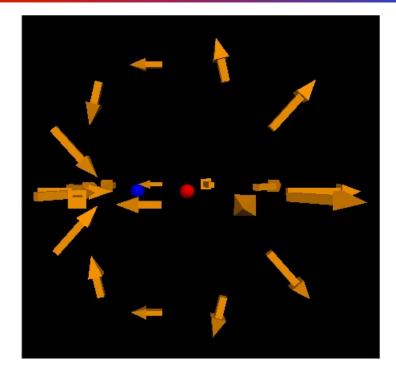
$$B_{z} = \frac{\mu_{0}}{4\pi} \frac{2\pi R^{2} I}{\left(z^{2}\right)^{3/2}}$$

$$B_z = \frac{\mu_0}{4\pi} \frac{2\pi R^2 I}{z^3}$$

The magnetic field of a circular loop falls off like $1/z^3$

Magnetic Dipole Moment





far from coil:
$$B_z = \frac{\mu_0}{4\pi} \frac{2\pi R^2 I}{z^3}$$

$$B_z = \frac{\mu_0}{4\pi} \frac{2\mu}{z^3}$$

far from dipole:
$$E_z = \frac{1}{4\pi\varepsilon_0} \frac{2p}{z^3}$$

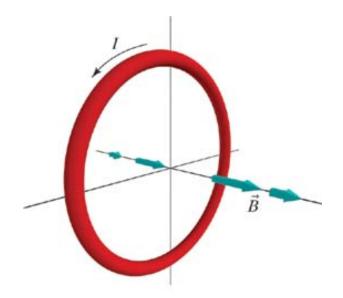
$$p = sq$$

magnetic

dipole moment: $\mu = \pi R^2 I = AI$ $\overrightarrow{\mu}$ - vector in the direction of \overrightarrow{B}

Magnetic Dipole in a B Field

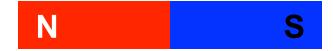
The magnetic dipole moment μ acts like a compass needle!

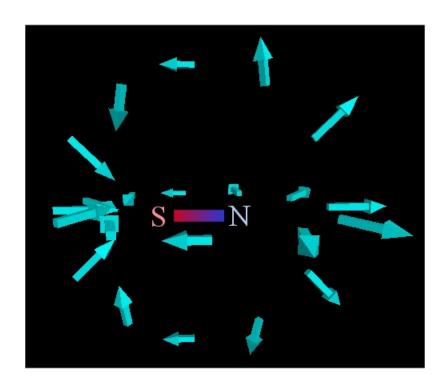


In the presence of external magnetic field a current-carrying loop rotates to align the magnetic dipole moment $\overline{\mu}$ along the field \overline{B} .

The Magnetic Field of a Bar Magnet

How does the magnetic field around a bar magnet look?





Atomic Structure of Magnets

Assume: Electrons make circular current loops

$$B_{atom} = \frac{\mu_0}{4\pi} \frac{2\mu}{z^3}, \quad \mu = \pi R^2 I$$

What is the direction?

One loop:

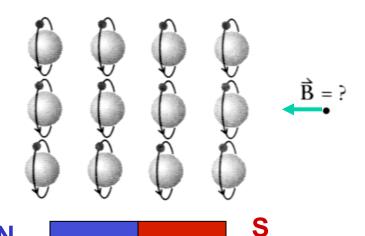
What is the average current *I*?

$$current = charge/second: I = \frac{e}{t}$$

$$T = \frac{2\pi R}{v} \longrightarrow I = \frac{ev}{2\pi R}$$

$$\mu = \pi R^2 \frac{ev}{2\pi R} = \frac{1}{2}eRv$$

Electrons



Atomic Structure of Magnets

Assume: Electrons make circular current loops

$$\mu = \frac{1}{2}eRv$$
 Dipole Moment of 1 Atom

Angular Momentum is Quantized:

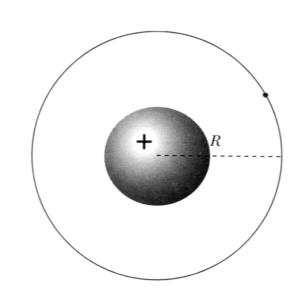
Orbital angular momentum: L = Rmv

$$\mu = \frac{1}{2}eRv = \frac{1}{2}\frac{e}{m}Rmv = \frac{1}{2}\frac{e}{m}L$$



$$L = n\hbar$$
, $\hbar = 1.05 \times 10^{-34} \text{ J} \cdot \text{s}$

If
$$n=1$$
: $\mu = \frac{1}{2} \frac{e}{m} L = 0.9 \times 10^{-23} \text{ A} \cdot \text{m}^2 \text{ per atom}$



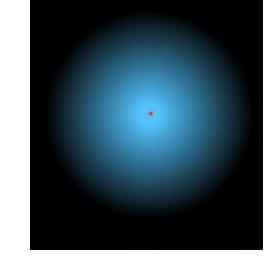
Quantum Magnetism

Magnetic Moment = Orbital Motion + "Spin"

ORBITAL MOTION:

Fuzzy electron cloud.

Shapes are set by "spherical harmonics"



Spherically symmetric cloud (s-orbital) has no μ

Only non spherically symmetric orbitals (p, d, f) contribute to μ

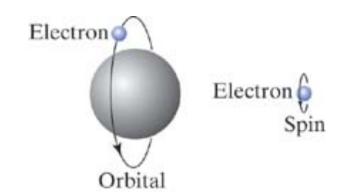
There is more than 1 electron in an atom

Quantum Magnetism

Magnetic Moment = Orbital Motion + "Spin"

SPIN:

Electron acts like spinning charge - contributes to μ



Electron spin contribution to μ is of the same order as one due to orbital momentum

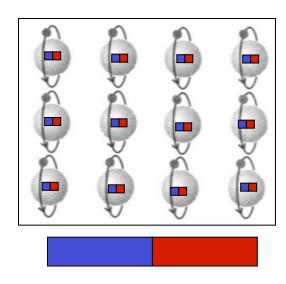
Neutrons and proton in nucleus also have spin but their μ 's are much smaller than for electron

same angular momentum:
$$\mu \approx \frac{1}{2} \frac{e}{m} \hbar$$

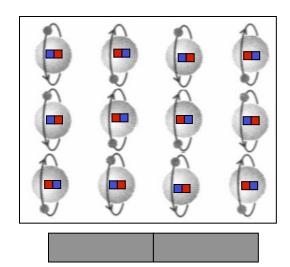
NMR, MRI – use nuclear μ

Refrigerator Magnets

Alignment of atomic dipole moments:

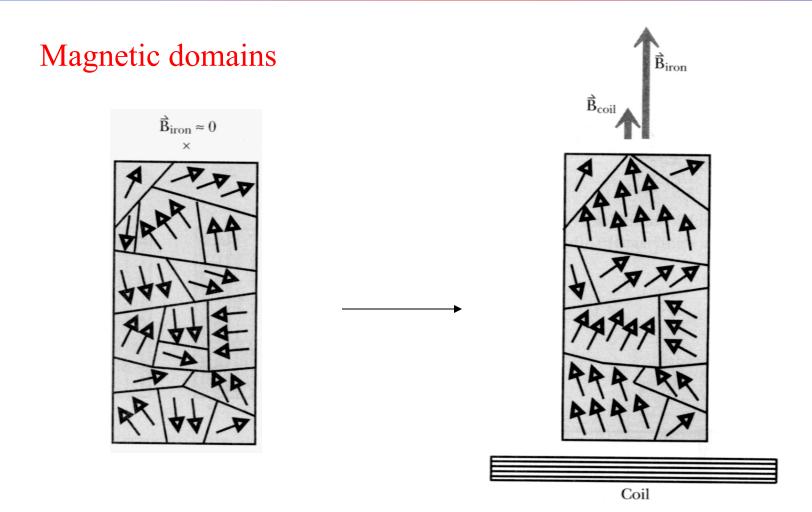


ferromagnetic materials: iron, cobalt, nickel



most materials

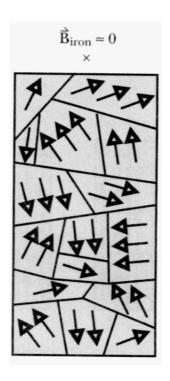
Reality Physics - Domains

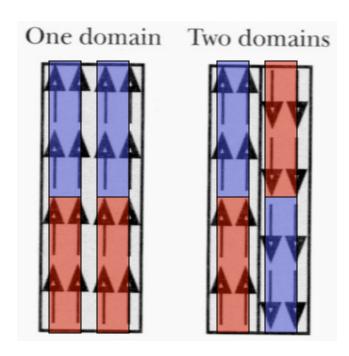


Hitting or heating can also demagnetize

Why are there Multiple Domains?

Magnetic domains





Today

- Magnetic Field of a Straight Wire
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