CONTRACTIBILITY OF SPACE OF STABILITY CONDITIONS ON THE PROJECTIVE PLANE VIA GLOBAL DIMENSION FUNCTION

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ABSTRACT. We compute the global dimension function gldim on the principal component $\mathrm{Stab}^\dagger(\mathbb{P}^2)$ of the space of Bridgeland stability conditions on $\mathbb{P}^2$. It admits 2 as the minimum value and the preimage $\mathrm{gldim}^{-1}(2)$ is contained in the closure $\overline{\mathrm{Stab}^{\mathrm{Geo}}(\mathbb{P}^2)}$ of the subspace consisting of geometric stability conditions. We show that $\mathrm{gldim}^{-1}[2, x)$ contracts to $\mathrm{gldim}^{-1}(2)$ for any real number $x \geq 2$ and that $\mathrm{gldim}^{-1}(2)$ is contractible.

Key words. Bridgeland stability conditions, Coherent sheaves, Contractibility, Global dimension function.

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1. INTRODUCTION

1.1. **Stability conditions.** The notion of stability conditions on triangulated categories was introduced by Bridgeland [Bri07], with motivation coming from string theory and mirror symmetry. Let $\mathcal{D}$ be a triangulated category and $K_{\mathrm{num}}(\mathcal{D})$ be its numerical Grothendieck group. A stability condition $\sigma = (Z, \mathcal{P})$ consists of a central charge $Z \in \mathrm{Hom}(K(\mathcal{D}), \mathbb{C})$ and a slicing $\mathcal{P}$, which is an $\mathbb{R}$-collection of t-structures on $\mathcal{D}$. In this paper, we denote $\mathrm{Stab}(\mathcal{D})$ as the stability manifold of stability conditions with support property with respect to $K_{\mathrm{num}}(\mathcal{D})$. By the seminal result in [Bri07], when $K_{\mathrm{num}}(\mathcal{D})$ is of finite rank, the space $\mathrm{Stab}(\mathcal{D})$ is a complex manifold with local coordinate given by the central charge. The original conjecture [Bri08, Conjecture 1.2] in the K3 surface case is that $\mathrm{Stab}(\mathcal{D})$ has a connected component $\mathrm{Stab}^\dagger(\mathcal{D})$ which is simply-connected and preserved by the automorphism group of $\mathcal{D}$. A more ambitious conjecture expects that the stability manifold $\mathrm{Stab}(\mathcal{D})$ is contractible in general. The contractibility is confirmed in couple of examples at least for the principal component of the space, namely:

- The smooth curves case in [Oka06, Bri07, Mac07].
- The K3 surfaces with Picard rank one in [Bri08, BB17].
- The local $\mathbb{P}^1$ in [IUU10]; the local $\mathbb{P}^2$ in [BM11].
- The projective plane $\mathbb{P}^2$ in [Li17].
- The Abelian surfaces in [Bri08] and Abelian threefolds with Picard rank one in [BMS16].
- The finite type (connected) component Stab_0 in [QW18], where the heart of any stability conditions in Stab_0 is a length category with finite many torsion pairs. The key examples are (Calabi–Yau) ADE Dynkin quiver case and new classes of examples are studied in [AW19].
- The Calabi–Yau-3 affine type A case in [Qiu16].

- The acyclic triangular quiver case in [DK16].
- The wild Kronecker quiver case in [DK19].

The proofs in each case are quite different.

1.2. **Global dimension functions.** Recently, Ikeda and the fourth-named author [Qiu18, IQ18a] introduce the global dimension function gldim on $\mathrm{Stab}(\mathcal{D})$, namely:

$$\mathrm{gldim}\colon \mathrm{Stab}(\mathcal{D}) \to \mathbb{R}_{\geq 0} \cup \{+\infty\}, \tag{1.1}$$

which is given by

$$\mathrm{gldim}\,\sigma = \mathrm{gldim}\,\mathcal{P} \coloneqq \sup\{\phi_2 - \phi_1 \mid \mathrm{Hom}(\mathcal{P}(\phi_1), \mathcal{P}(\phi_2)) \neq 0\}. \tag{1.2}$$

Such a function is continuous and invariant under the natural left action by $\mathrm{Aut}(\mathcal{D})$ and the right action of $\mathbb{C}$, and thus descends to a continuous function

$$\mathrm{gldim}\colon \mathrm{Aut}(\mathcal{D})\backslash \mathrm{Stab}(\mathcal{D})/\mathbb{C} \to \mathbb{R}_{\geq 0} \cup \{+\infty\}. \tag{1.3}$$

The philosophy in [Qiu18] is as follows:

(i) The infimum of gldim on $\mathrm{Stab}(\mathcal{D})$ (or the principal component of it) should be considered as the global dimension $\mathrm{gd}\,\mathcal{D}$ of the category $\mathcal{D}$.
(ii) If the subspace $\mathrm{gldim}^{-1}(\mathrm{gd}\,\mathcal{D})$ is non-empty, then it is contractible. Moreover, $\mathrm{gldim}^{-1}([\mathrm{gd}\,\mathcal{D}, x))$ contracts to $\mathrm{gldim}^{-1}(\mathrm{gd}\,\mathcal{D})$ for any real number $\mathrm{gd}\,\mathcal{D} < x$.
(iii) If $\mathrm{gldim}^{-1}(\mathrm{gd}\,\mathcal{D})$ is empty, then $\mathrm{gldim}^{-1}(\mathrm{gd}\,\mathcal{D}, x)$ contracts to $\mathrm{gldim}^{-1}(\mathrm{gd}\,\mathcal{D}, y)$ for any real number $\mathrm{gd}\,\mathcal{D} < y < x$.

Note that for a Calabi–Yau category, the global dimension function is constant. If the global dimension function gldim is not constant, it sheds some lights on why $\mathrm{Stab}(\mathcal{D})$ should be contractible.

The theme in [IQ18a] is to q-deform stability conditions. More precisely, given a Calabi–Yau-∞ category $\mathcal{D}_\infty$ (e.g. bounded derived category of $\mathbb{P}^2$), the corresponding Calabi–Yau-N category $\mathcal{D}_N$ (e.g. local $\mathbb{P}^2$ for $\mathbb{P}^2$ and $N = 3$) can be obtained by Calabi–Yau-$\mathbb{X}$ completing $\mathcal{D}_\infty$ to $\mathcal{D}_\mathbb{X}$ and specializing $\mathbb{X}$ to be N, in other words, taking the orbit category $\mathcal{D}_N = \mathcal{D}_\mathbb{X} /\!\!/ [\mathbb{X} - N]$. Under this procedure, a stability condition σ on $\mathcal{D}_\infty$ such that

$$\mathrm{gldim}\,\sigma \leq N - 1 \tag{1.4}$$

induces a stability condition on $\mathcal{D}_N$ via q-stability conditions on $\mathcal{D}_\mathbb{X}$. We will discuss such inducing in Section 7 for the example from $\mathbb{P}^2$ to local $\mathbb{P}^2$ where $N = 3$.

1.3. **The projective plane case.** In this paper, we study the case of the projective plane $\mathbb{P}^2$ for the above conjectures/philosophy. The main result is a computation of the global dimension function for the principal component $\mathrm{Stab}^\dagger(\mathbb{P}^2)$ (i.e. the connected component which contains geometric stability conditions, where a stability condition $\sigma \in \mathrm{Stab}(\mathbb{P}^2)$ is called geometric if all skyscraper sheaves are σ-stable of the same phase). Details are in Propositions 3.4 and 5.1. Based on the computation of gldim, we prove the following theorem.

Theorem 1.1 (Corollary 5.10 and Theorem 6.1)**.** *Consider the function*

$$\mathrm{gldim}\colon \mathrm{Stab}^\dagger(\mathbb{P}^2) \to \mathbb{R}_{\geq 0}$$

on the principal component $\mathrm{Stab}^{\dagger}(\mathbb{P}^2)$ of the space of stability conditions on the bounded derived category $\mathcal{D} = \mathcal{D}^b(\mathrm{Coh}\,\mathbb{P}^2)$ of coherent sheaves on $\mathbb{P}^2$. Then

- gd $\mathcal{D} = 2$ and gldim $\mathrm{Stab}^{\dagger}(\mathbb{P}^2) = [2, \infty)$,
- the subspace gldim$^{-1}[2, x)$ contracts to gldim$^{-1}(2)$, for any $x \geq 2$,
- the subspace gldim$^{-1}(2)$ is contractible and is contained in $\overline{\mathrm{Stab}^{\mathrm{Geo}}(\mathbb{P}^2)}$, where $\mathrm{Stab}^{\mathrm{Geo}}(\mathbb{P}^2)$ consists of geometric stability conditions.

The contractibility of $\mathrm{Stab}^{\dagger}(\mathbb{P}^2)$ is already proved by the second-named author [Li17]. The new approach here shows how this stability manifold contracts along the values of the global dimension function.

1.4. **Topological Fukaya case.** In the parallel work [Qiu20], we use the same philosophy to study the contractibility of the space of stability conditions on the topological Fukaya category of a graded marked surface. We prove a slightly weaker version of the corresponding Theorem 1.1, that gldim induces the contractible flow except for certain possible critical values.

We hope that these works will shed lights on how this philosophy would apply to other cases.

2. Preliminaries

2.1. **The category.** In this paper, we let $\mathbb{P}^2$ be the projective plane over the complex number field. We write

$$\mathcal{D}_{\infty}(\mathbb{P}^2) := \mathcal{D}^b(\mathbb{P}^2) = \mathcal{D}^b(\mathrm{Coh}\,\mathbb{P}^2) \tag{2.1}$$

for the bounded derived category of coherent sheaves on $\mathbb{P}^2$. Due to the well-known result by Beĭlinson [Bei83], we have the equivalent description $\mathcal{D}_{\infty}(\mathbb{P}^2) \cong \mathcal{D}^b(\mathbf{k}Q/R)$, where (Q, R) is the quiver

$$1 \underset{}{\overset{x_1,\,y_1,\,z_1}{\Longrightarrow}} 2 \underset{}{\overset{x_2,\,y_2,\,z_2}{\Longrightarrow}} 3$$

with commutative relations

$$a_1 b_2 = b_1 a_2, \quad a, b \in \{x, y, z\}.$$

The Serre functor on $\mathcal{D}_{\infty}(\mathbb{P}^2)$ is given by (see [BK90] or [Huy06])

$$\mathbb{S} = \mathbb{S}_{\mathbb{P}^2} := (-) \otimes \omega_{\mathbb{P}^2}[2] = (-) \otimes \mathcal{O}_{\mathbb{P}^2}(-3)[2].$$

An object $E \in \mathcal{D}_{\infty}(\mathbb{P}^2)$ is called exceptional if $\mathrm{Hom}(E, E[i]) = 0$ for $i \neq 0$ and $\mathrm{Hom}(E, E) = \mathbb{C}$. The right and left mutations of an object F with respect to an

exceptional object E are defined by

$$\mathsf{R}_E(F) := \text{Cone}\left(F \xrightarrow{\text{ev}} E \otimes \text{Hom}(F, E)^*\right)[-1], \qquad (2.2)$$

$$\mathsf{L}_E(F) := \text{Cone}\left(E \otimes \text{Hom}(E, F) \xrightarrow{\text{ev}} F\right). \qquad (2.3)$$

2.2. **An affine plane.** Let $\mathcal{D} = \mathcal{D}_\infty(\mathbb{P}^2)$. Let H be the hyperplane divisor of $\mathbb{P}^2$. For $E \in \mathcal{D}$, we identity the Chern character $\text{ch}(E)$ with the triple of numbers

$$\tilde{\text{v}}(E) = (\text{ch}_0(E), \text{ch}_1(E).H, \text{ch}_2(E)).$$

When we say the *point E* (or the *point $\tilde{\text{v}}(E)$*), we mean the point in the real projective space $\mathbb{P}(\mathbb{R}^3)$ with homogeneous coordinate $[\text{ch}_0(E), \text{ch}_1(E).H, \text{ch}_2(E)]$. We call the locus $\text{ch}_0 = 0$ as the line at infinity and its complement as the affine $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. Moreover, we always assume that the $\frac{\text{ch}_1}{\text{ch}_0}$-axis is horizontal and the $\frac{\text{ch}_2}{\text{ch}_0}$-axis is vertical. If $\text{ch}_0(E) \neq 0$, the *reduced character* of E corresponds to the point

$$\text{v}(E) := (1, s(E), q(E)), \quad \text{with } s(E) := \frac{\text{ch}_1(E).H}{\text{ch}_0(E)}, \quad q(E) := \frac{\text{ch}_2(E)}{\text{ch}_0(E)}, \qquad (2.4)$$

in the affine $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. In particular $\text{v}(E) = \text{v}(E[n])$, i.e. E and its any shift $E[n]$ will be the same point in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane.

The $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane provides a playground for studying both geometric and algebraic stability conditions in the following part of the paper.

2.3. **Stability conditions.** A *stability condition* $\sigma = (Z, \mathcal{P})$ on $\mathcal{D}$ consists of a group homomorphism $Z: K(\mathcal{D}) \to \mathbb{C}$ called the *central charge* and a family of full additive subcategories $\mathcal{P}(\phi) \subset \mathcal{D}$ for $\phi \in \mathbb{R}$ called the *slicing* satisfying certain conditions. We refer the lecture notes [MS17, Definition 5.8] for the details. Nonzero objects in $\mathcal{P}(\phi)$ are called *semistable of phase ϕ* and simple objects in $\mathcal{P}(\phi)$ are called *stable of phase ϕ*. For semistable object $E \in \mathcal{P}(\phi)$, denote by $\phi_\sigma(E) = \phi$ its *phase*.

Let $\mathcal{D} = \mathcal{D}_\infty(\mathbb{P}^2)$ and

$$\text{Stab}(\mathbb{P}^2) := \text{Stab}(\mathcal{D}_\infty(\mathbb{P}^2))$$

be the space of stability conditions on $\mathcal{D}_\infty(\mathbb{P}^2)$.

A stability condition $\sigma \in \text{Stab}(\mathbb{P}^2)$ is called *geometric* if all skyscraper sheaves are σ-stable of the same phase. We denote the set of all geometric stability conditions by $\text{Stab}^{\text{Geo}}(\mathbb{P}^2)$.

Let us briefly recall the construction of geometric stability conditions. There is a fractal curve C_{LP}, the so called Le Potier curve, and a region Geo_{LP} in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane, as in Definition 4.5. For each $(1, s, q) \in \text{Geo}_{\text{LP}}$, one can associate a geometric stability condition $\sigma_{s,q} = (Z_{s,q}, \mathcal{P}_{s,q})$ as follows. The central charge $Z_{s,q}$ is given by

$$Z_{s,q}(E) := (-\text{ch}_2(E) + q \cdot \text{ch}_0(E)) + i(\text{ch}_1(E).H - s \cdot \text{ch}_0(E)), \quad \text{for } E \in \mathcal{D}. \qquad (2.5)$$

Denote H-slope of coherent sheaves by $\frac{\text{ch}_1(-).H}{\text{ch}_0(-)}$. We make a convention that H-slope of a torsion sheaf is $+\infty$. The heart $\mathcal{P}_{s,q}((0,1])$ is the tilting

$$\text{Coh}_{\#s} := \langle \text{Coh}_{\leq s}[1], \text{Coh}_{>s} \rangle,$$

where $\text{Coh}_{\leq s}$ (resp. $\text{Coh}_{>s}$) is the subcategory of $\text{Coh}(\mathbb{P}^2)$ generated by H-slope semistable sheaves of slope $\leq s$ (resp. $> s$) by extension. The slicing for $\phi \in (0, 1]$ is defined by

$$\mathcal{P}_{s,q}(\phi) = \{E \in \text{Coh}_{\#s} \mid E \text{ is } \sigma_{s,q} \text{-semistable of phase } \phi\} \cup \{0\}.$$

For general $\phi \in \mathbb{R}$, we have $\mathcal{P}_{s,q}(\phi + 1) = \mathcal{P}_{s,q}(\phi)[1]$.

The $\widetilde{\text{GL}^+(2, \mathbb{R})}$ acts freely on $\text{Stab}^{\text{Geo}}(\mathbb{P}^2)$ ([Li17, Definition 1.4, Corollary 1.15]) with quotient

$$\text{Stab}^{\text{Geo}}(\mathbb{P}^2)/\widetilde{\text{GL}^+(2, \mathbb{R})} \cong \text{Geo}_{\text{LP}}.$$

We refer to Section 4.1 for the definition of algebraic stability conditions $\text{Stab}^{\text{Alg}}(\mathbb{P}^2)$. Let $\text{Stab}^{\dagger}(\mathbb{P}^2)$ be the connected component in $\text{Stab}(\mathbb{P}^2)$ which contains the geometric stability conditions. It is still a *conjecture* that $\text{Stab}(\mathbb{P}^2) = \text{Stab}^{\dagger}(\mathbb{P}^2)$. The second-named author [Li17] shows that

$$\text{Stab}^{\dagger}(\mathbb{P}^2) = \text{Stab}^{\text{Geo}}(\mathbb{P}^2) \bigcup \text{Stab}^{\text{Alg}}(\mathbb{P}^2) \tag{2.6}$$

and it is contractible. In the following Sections, we will compute the global dimension function gldim on $\text{Stab}^{\dagger}(\mathbb{P}^2)$ and show that the contraction is along the value of gldim.

3. Geometric stability conditions in the parabolic region

Let $\mathcal{D} = \mathcal{D}_{\infty}(\mathbb{P}^2)$. For $a \in \mathbb{R}$, denote by Δ_a the parabola in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane:

$$\Delta_a := \left\{ (1, s, q) \in \{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}\text{-plane} \mid \frac{1}{2}s^2 - q = a \right\}.$$

Similarly we have the notation $\Delta_{<a}$ or $\Delta_{\geq a}$. We study geometric stability conditions in the parabolic region $\Delta_{<0}$. Denote by L_{PE} the line passing through the two points P and E in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. Recall a lemma due to Bayer [LZ19b, Lemma 3].

Lemma 3.1. *Let P and Q be two points in the region $\Delta_{<0}$ in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. Let F be a σ_P-stable object in Coh_P with $\text{ch}(F) \neq (0, 0, 1)$. Let C and D be the intersection points*

$$\{C, D\} := L_{PF} \cap \Delta_0,$$

of the line L_{PF} and the parabola Δ_0. Denote the σ_Q-HN semistable factors of F by F_i. Then for each factor, the phase $\phi_Q(F_i)$ lies in between $\phi_Q(C)$ and $\phi_Q(D)$.

Proof. The case that $\text{ch}_0(F) \neq 0$ is proved in [LZ19b, Lemma 3].

So we assume that $\text{ch}(F) = (0, \text{ch}_1(F), \text{ch}_2(F))$ with $\text{ch}_1(F).H > 0$. Now the point F is in the ∞-line outside the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. But the line L_{PF} still makes sense: it is the line passing through the point P with slope $\frac{\text{ch}_2(F)}{\text{ch}_1(F).H}$ in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane, see [Liu18, Corollary 2.8]. So we still have the notation l^+_{PF}, which is the ray starting at the point P on the line L_{PF} with $s \geq s(P)$. Note that L_{QF} is parallel to L_{PF}. Then the proof follows by Li-Zhao's original argument. $\square$

Let K be the canonical divisor of $\mathbb{P}^2$, and $\omega_{\mathbb{P}^2}$ be the dualizing sheaf. Let $\sigma = \sigma_{s,q}$ with $(1, s, q) \in \text{Geo}_{\text{LP}}$. We identify σ with the point $(1, s, q)$ in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. Then $\sigma(-3) := \sigma \otimes \omega_{\mathbb{P}^2} = (1, s-3, q-3s+\frac{9}{2})$ is the point of moving σ along the parabola to the

left by $-H.K = 3$, where the parabola is of the form $q - \frac{1}{2}s^2 = q(\sigma) - \frac{1}{2}(s(\sigma))^2$ passing through the point σ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. Similarly for $F \in \mathcal{D}$ as a point in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane, if $\mathrm{ch}_0(F) \neq 0$, then $F(-3) := F \otimes \omega_{\mathbb{P}^2} = (1, s(F) - 3, q(F) - 3s(F) + \frac{9}{2})$ is the point of moving the point F along the parabola to the left by 3, where the parabola is of the form $q - \frac{1}{2}s^2 = q(F) - \frac{1}{2}(s(F))^2$ passing through the point F in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane.

Let A, B, $\tilde{A}$, $\tilde{B}$ be the corresponding intersection points

$$\{A, B\} := L_{F\sigma} \cap \Delta_0, \quad \{\tilde{A}, \tilde{B}\} := L_{F(-3)\sigma(-3)} \cap \Delta_0,$$

with $s(B) > s(A)$ and $s(\tilde{B}) > s(\tilde{A})$. We have the following observation.

Lemma 3.2.
$$s(B) - s(A) = s(\tilde{B}) - s(\tilde{A}). \tag{3.1}$$

Proof. This is an elementary calculation. □

We prove a lemma, which is the key calculation for proving $\mathrm{gldim}\,\sigma_{s,q} = 2$ in the region $\Delta_{<0}$.

Lemma 3.3. *Let $\sigma = \sigma_{s,q}$ be a geometric stability condition in the region $\Delta_{<0}$ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. Let F, G be two σ stable objects in a same heart satisfying: $0 < \phi_\sigma(F) < \phi_\sigma(G) \leq 1$, $\mathrm{ch}_0(F) \neq 0$ and $s < s(F)$. Then $\mathrm{Hom}(F, G[2]) = 0$.*

Proof. Let $P := L_{F\sigma} \cap L_{F(-3)\sigma(-3)}$. We have two cases.

Case A. P is in the region $\Delta_{\geq 0}$. Then by Lemma 3.2, we must have $s(\tilde{B}) \leq s(A)$ (i.e. left above of Figure 1) instead of $s(\tilde{B}) > s(A)$ (i.e. left below or right below of Figure 1). So $l^+_{\sigma F}$ is above or equal to $l^+_{\sigma \tilde{A}}$ and $l^+_{\sigma \tilde{B}}$.

By [LLM19, Lemma A.3], $F(-3)$ is $\sigma(-3)$-stable. By Lemma 3.1, the σ-HN factor $F(-3)_i$ of $F(-3)$ lies between $\phi_\sigma(\tilde{A})$ and $\phi_\sigma(\tilde{B})$. By [LZ19b, Lemma 2], $\phi_\sigma(F(-3)_i) \leq \phi_\sigma(F)$. So $\phi^+_\sigma(F(-3)) \leq \phi_\sigma(F) < \phi_\sigma(G)$ and $\mathrm{Hom}(G, F(-3)) = 0$. By the Serre duality, we have

$$\mathrm{Hom}(F, G[2]) \cong (\mathrm{Hom}(G[2], \mathbb{S}(F)))^* = (\mathrm{Hom}(G, F(-3)))^* = 0.$$

Case B. P is in the region $\Delta_{<0}$. So both F and $F(-3)$ are σ_P-stable with

$$\phi_P(F) > \phi_P(F(-3)). \tag{3.2}$$

Let $Q := L_{G\sigma} \cap L_{F(-3)\sigma(-3)}$. We have three subcases.

Case B.(i) Q is in the region $\Delta_{>0}$. Then Q is to the right of $\tilde{B}$ since $\tilde{B}$ is on the Δ_0. Now l^+_{QG} is above $l^+_{\sigma \tilde{A}}$ and $l^+_{\sigma \tilde{B}}$. We must have $l^+_{\sigma G}$ is above $l^+_{\sigma \tilde{A}}$ and $l^+_{\sigma \tilde{B}}$. By Lemma 3.1 again, we have $\mathrm{Hom}(G, F(-3)) = 0$. By the Serre duality, we have $\mathrm{Hom}(F, G[2]) = 0$.

Case B.(ii) Q is in the region $\Delta_{<0}$. We illustrate the picture in right above of Figure 1. Since G is σ-stable, it is also σ_Q-stable. Since $F(-3)$ is $\sigma(-3)$-stable, it is also σ_Q-stable. We then compare their phases at Q and have

$$\phi_Q(G) = \phi_\sigma(G) > \phi_\sigma(F) = \phi_P(F) > \phi_P(F(-3)) = \phi_Q(F(-3)),$$

where each equality is because of colinear condition, and the first inequality is given by the assumption of the Lemma and the second inequality is given by (3.2). So $\mathrm{Hom}(G, F(-3)) = 0$. By the Serre duality, we have $\mathrm{Hom}(F, G[2]) = 0$.
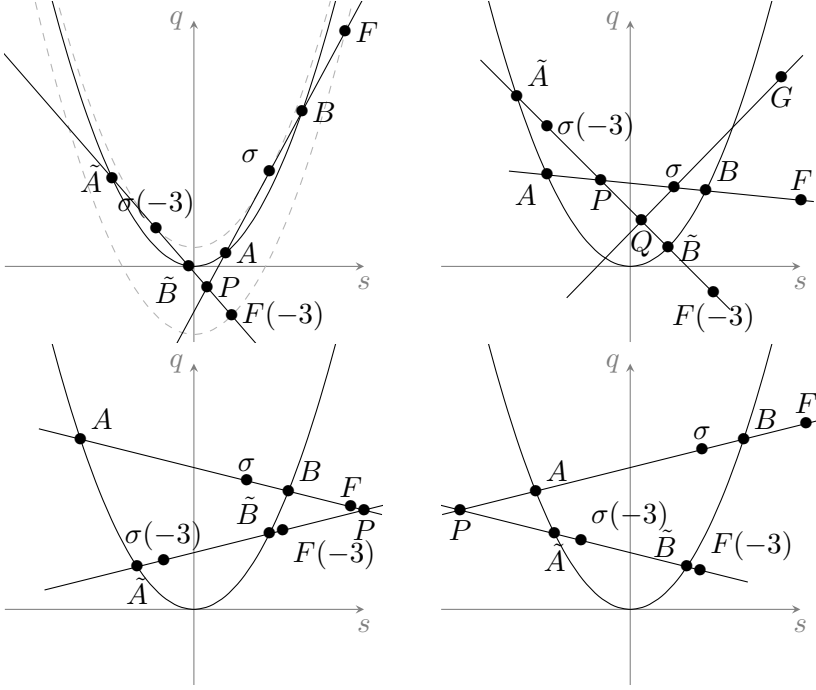
FIGURE 1. Relative positions of $L_{F\sigma}$ and $L_{F(-3)\sigma(-3)}$: $P \in \Delta_{\geq 0}$ (left above); $P \in \Delta_{<0}$ (right above). The below pictures are impossible by Lemma 3.2.

Case B.(iii) Q is on the parabola Δ_0. Since F is σ-stable, we may perturb σ a little bit and reduce to the previous cases. □

Proposition 3.4. *Let $\sigma_{s,q}$ be a geometric stability condition in the region $\Delta_{<0}$ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. Then*

$$\mathrm{gldim}\, \sigma_{s,q} = 2. \tag{3.3}$$

Proof. Denote $\sigma_{s,q}$ by σ. Let F and G be two σ-semistable objects such that

$$\mathrm{Hom}(F, G[2]) \neq 0.$$

Then by the Serre duality,

$$\mathrm{Hom}(F, G[2]) \cong (\mathrm{Hom}(G, F(-3)))^* \neq 0. \tag{3.4}$$

The object $F(-3)$ may not be σ-semistable. We consider its σ-HN factors. Thus by [Bri07, Lemma 3.4] we have $\phi_\sigma(F) \leq \phi_\sigma(G[2]) \leq \phi_\sigma^+(F(-3)) + 2$. So

$$0 \leq \phi_\sigma(G[2]) - \phi_\sigma(F) \leq \phi_\sigma^+(F(-3)) - \phi_\sigma(F) + 2. \tag{3.5}$$

We need to show that

$$\phi_\sigma(G[2]) - \phi_\sigma(F) \leq 2. \tag{3.6}$$

The idea is to give an estimation of $\phi_\sigma^+(F(-3)) - \phi_\sigma(F)$. We could assume that $F \in \mathrm{Coh}_{\#s}$, i.e. $\phi_\sigma(F) \in (0, 1]$. We could also assume that F is σ-stable since we can take its Jordan-Hölder factors. So by [LLM19, Lemma A.3], $F(-3)$ is $\sigma(-3)$-stable.

We have the following three cases according to the Chern characters of F.

Case 1. Assume $\mathrm{ch}_0(F) = 0$, $\mathrm{ch}_1(F) = 0$ and $\mathrm{ch}_2(F) > 0$. Then F is supported at point(s) and $\phi_\sigma(F(-3)) = \phi_\sigma(F)$. So (3.6) holds. On the other hand, for any closed point $x \in \mathbb{P}^2$, we have $\mathrm{Hom}(\mathcal{O}_x, \mathcal{O}_x[2]) \neq 0$ and

$$\phi_\sigma(\mathcal{O}_x[2]) - \phi_\sigma(\mathcal{O}_x) = 2. \tag{3.7}$$

Case 2. Assume that $\mathrm{ch}_0(F) \neq 0$. We have the following three subcases.

(i) σ is to the left of F. This is precisely Lemma 3.3.

(ii) If σ is to the right of F, by applying a shifted derived dual functor, we reduce to case (i).

(iii) If the H-slope of F is s, by local finiteness of walls, we could replace σ by σ' in a small open neighbourhood of σ so that F is σ'-stable. So we reduce to case (i) or (ii).

Case 3. Assume $\mathrm{ch}_0(F) = 0$ and $\mathrm{ch}_1(F).H > 0$. Now $\mathrm{ch}(F(-3)) = (0, \mathrm{ch}_1(F), \mathrm{ch}_2(F) + \mathrm{ch}_1(F).K)$.

The line $L_{F\sigma}$ is the line passing through σ of the slope $\frac{\mathrm{ch}_2(F)}{\mathrm{ch}_1(F).H}$. The line $L_{F(-3)\sigma(-3)}$ is the line passing through $\sigma(-3)$ of the slope $\frac{\mathrm{ch}_2(F)}{\mathrm{ch}_1(F).H} + H.K$. By Lemma 3.1, the phase of $\phi_\sigma((F(-3))_i)$ lies between $\phi_\sigma(\tilde{A})$ and $\phi_\sigma(\tilde{B})$. We have similar analysis as the **Case 2** and still have (3.6).

Therefore for geometric stability conditions $\sigma_\sigma \in \Delta_{<0}$ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane, we have $\mathrm{gldim}(\sigma_\sigma) = 2$. Moreover, it is gldim-reachable by (3.7). This finishes the proof. $\qquad\square$

4. Algebraic stability conditions

4.1. **Reviews.** We first recall the construction of algebraic stability conditions with respect to exceptional triples from [Li17].

Definition 4.1. We call an ordered set $\mathcal{E} = \{E_1, E_2, E_3\}$ *exceptional triple* on $\mathcal{D}^b(\mathbb{P}^2)$ if $\mathcal{E}$ is a full strong exceptional collection of coherent sheaves on $\mathcal{D}^b(\mathbb{P}^2)$.

There is a one-to-one correspondence between the dyadic integers $\frac{p}{2^m}$ and exceptional bundles $E(\frac{p}{2^m})$:

$$\frac{p}{2^m} \Longleftrightarrow E(\frac{p}{2^m}), \text{ for } p \in \mathbb{Z} \text{ and } m \in \mathbb{Z}_{\geq 0}.$$

The exceptional triples have been classified by Gorodentsev and Rudakov [GR87]. Up to a cohomological shift, the exceptional triples are labeled by the following three cases,

$$\left\{\frac{p-1}{2^m}, \frac{p}{2^m}, \frac{p+1}{2^m}\right\}, \quad \left\{\frac{p}{2^m}, \frac{p+1}{2^m}, \frac{p-1}{2^m} + 3\right\}, \quad \left\{\frac{p+1}{2^m} - 3, \frac{p-1}{2^m}, \frac{p}{2^m}\right\},$$

for $p \in \mathbb{Z}$ and $m \in \mathbb{Z}_{\geq 0}$. Note that the last two cases are mutations of the first case.

Proposition 4.2 ([Mac07, Section 3]). *Let $\mathcal{E}$ be an exceptional triple on $\mathcal{D}^b(\mathbb{P}^2)$. For any positive real numbers m_1, m_2, m_3 and real numbers ϕ_1, ϕ_2, ϕ_3 such that:*

$$\phi_1 < \phi_2 < \phi_3, \text{ and } \phi_1 + 1 < \phi_3,$$

there is a unique stability condition $\sigma = (Z, \mathcal{P})$ such that

1°. *each E_j is stable with phase ϕ_j;*
2°. *$Z(E_j) = m_j e^{i\pi\phi_j}$.*

Definition 4.3. Given an exceptional triple $\mathcal{E} = \{E_1, E_2, E_3\}$ on $\mathcal{D}^b(\mathbb{P}^2)$, we write $\Theta_{\mathcal{E}}$ as the space of all stability conditions in Proposition 4.2, which is parametrized by

$$S := \{(m_1, m_2, m_3, \phi_1, \phi_2, \phi_3) \in (\mathbb{R}_{>0})^3 \times \mathbb{R}^3 \mid \phi_1 < \phi_2 < \phi_3, \phi_1 + 1 < \phi_3\}.$$

We make the following notations for some subsets of $\Theta_{\mathcal{E}}$.

$$\Theta_{\mathcal{E}}(A) := \{\sigma \in \Theta_{\mathcal{E}} \mid \sigma \in A\}, \text{ where } A \text{ is a subset of } S;$$
$$\Theta_{\mathcal{E}}^{\mathrm{Pure}} := \{\sigma \in \Theta_{\mathcal{E}} \mid \phi_2 - \phi_1 \geq 1 \text{ and } \phi_3 - \phi_2 \geq 1\};$$
$$\Theta_{\mathcal{E}, E_3}^{\mathrm{left}} := \{\sigma \in \Theta_{\mathcal{E}} \mid \phi_2 - \phi_1 < 1 \text{ and } E_3(3) \text{ is not } \sigma\text{-stable}\};$$
$$\Theta_{\mathcal{E}, E_1}^{\mathrm{right}} := \{\sigma \in \Theta_{\mathcal{E}} \mid \phi_3 - \phi_2 < 1 \text{ and } E_1(-3) \text{ is not } \sigma\text{-stable}\};$$
$$\Theta_{\mathcal{E}}^{\mathrm{Geo}} := \Theta_{\mathcal{E}} \cap \mathrm{Stab}^{\mathrm{Geo}}(\mathbb{P}^2);$$
$$\Theta_{\mathcal{E}, E_3}^{-} := \Theta_{\mathcal{E}}(\phi_2 - \phi_1 < 1) \setminus \Theta_{\mathcal{E}}^{\mathrm{Geo}};$$
$$\Theta_{\mathcal{E}, E_1}^{+} := \Theta_{\mathcal{E}}(\phi_3 - \phi_2 < 1) \setminus \Theta_{\mathcal{E}}^{\mathrm{Geo}}.$$

We denote

$$\mathrm{Stab}^{\mathrm{Alg}}(\mathbb{P}^2) := \bigcup_{\mathcal{E} \text{ exceptional triples}} \Theta_{\mathcal{E}}$$

and call the elements of it as the *algebraic* stability conditions.

Lemma 4.4 ([Li17, Lemma 2.4]). *Let $\mathcal{E} = \{E_1, E_2, E_3\}$ be an exceptional triple, and σ be a stability condition in $\Theta_{\mathcal{E}}^{\mathrm{Pure}}$. The only σ-stable objects are $E_i[n]$ for $i = 1, 2, 3$ and $n \in \mathbb{Z}$.*

4.2. **Five points associated to an exceptional bundle.** For an object $A \in \mathcal{D}$ with $\mathrm{ch}_0(A) \neq 0$, by abusing of notations, we write A for $\mathrm{v}(A) = (1, s(A), q(A))$ in (2.4) as the associated point in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane, and call it the *point A* in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. Moreover, by the Riemann–Roch formula, we have

$$\chi(A, A) = \mathrm{ch}_0^2(A)(1 - s(A)^2 + 2q(A)). \tag{4.1}$$

In particular, for an exceptional bundle E, we have $\mathrm{ch}_0(E) \neq 0$, $\chi(E, E) = 1$, and

$$\frac{1}{2}s(E)^2 - q(E) = \frac{1}{2} - \frac{1}{2\mathrm{ch}_0^2(E)}. \tag{4.2}$$

So the point E is in the region $\Delta_{[0, \frac{1}{2})}$.

For each exceptional bundle E, we define the following two pairs of parallel lines in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane:

$$\{\chi(E, -) = 0\}, \quad \{\chi(E, -) = \frac{\mathrm{ch}_0(-)}{\mathrm{ch}_0(E)}\}; \tag{4.3}$$

$$\{\chi(-, E) = 0\}, \quad \{\chi(-, E) = \frac{\mathrm{ch}_0(-)}{\mathrm{ch}_0(E)}\}. \tag{4.4}$$

We now give a geometric description of above lines. By the Riemann–Roch formula, one can check that the line $\{\chi(E, -) = \frac{\mathrm{ch}_0(-)}{\mathrm{ch}_0(E)}\}$ is the line $L_{E(-3)E}$ passing through the points $E(-3)$ and E. Similarly, the line $\{\chi(-, E) = \frac{\mathrm{ch}_0(-)}{\mathrm{ch}_0(E)}\}$ is the line $L_{E(E(3))}$ passing through the points E and $E(3)$.

The line $\{\chi(E, -) = 0\}$ is the line passing through points E_1 and E_2 for any choice of exceptional triple $\{E_1, E_2, E\}$ ending with E. It is clearly that this line is independent of the choice of representatives E_1 and E_2. Similarly, the line $\{\chi(-, E) = 0\}$ is the line passing through points E_2 and E_3 for any choice of exceptional triple $\{E, E_2, E_3\}$ starting with E. This line is independent of the choice of representatives E_2 and E_3.

For an exceptional bundle E, we define five points in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane as intersection points of the following lines,

$$E^l := L_{E(E(3))} \cap \{\chi(E, -) = 0\},$$
$$E^r := L_{E(-3)E} \cap \{\chi(-, E) = 0\},$$
$$E^+ := \{\chi(E, -) = 0\} \cap \{\chi(-, E) = 0\},$$
$$e^l := \Delta_{\frac{1}{2}} \cap \{\chi(E, -) = 0\} \text{ as the first intersection point staring from } E^+,$$
$$e^r := \Delta_{\frac{1}{2}} \cap \{\chi(-, E) = 0\} \text{ as the first intersection point staring from } E^+.$$

We now give a geometric description of above points. One can also refer to Figure 2 and Figure 3. By the Riemann–Roch formula, (4.2) and (2.4), we have

$$s(E^+) = s(E), \quad q(E^+) = q(E) - \frac{1}{(\mathrm{ch}_0(E))^2}. \tag{4.5}$$

So E^+ is the point of moving E downward of length $\frac{1}{(\mathrm{ch}_0(E))^2}$. By (4.2) and (4.5), we have

$$\frac{1}{2}s(E^+)^2 - q(E^+) = \frac{1}{2} + \frac{1}{2\mathrm{ch}_0^2(E)}. \tag{4.6}$$

So the point E^+ is in the region $\Delta_{(\frac{1}{2}, 1]}$.

We observe that the point E^l stands for $\mathrm{v}(\mathsf{L}_E(E(3)))$, i.e. the reduced character of $\mathsf{L}_E(E(3))$. This is because by the definition (2.3), the point $\mathsf{L}_E(E(3))$ is on the line $L_{EE(3)}$. Also, the object $\mathsf{L}_E(E(3)) \in E^\perp = \langle E_1, E_2 \rangle$ has a resolution (5.2) (by taking $E_3 = E$). Thus the point $\mathsf{L}_E(E(3))$ is on the line $\{\chi(E, -) = 0\}$.

By the Riemann–Roch formula, we have

$$\chi(E, E(3)) = 1 + 9\mathrm{ch}_0^2(E), \quad \chi(E(3), E) = 1, \quad \chi(E(3), E(3)) = 1.$$

Since $[\mathsf{L}_E(E(3))] = [E(3)] - \chi(E, E(3))[E]$ in $K_{\mathrm{num}}(\mathbb{P}^2)$, we have

$$\chi(\mathsf{L}_E(E(3)), \mathsf{L}_E(E(3))) = 1 - \chi(E(3), E)\chi(E, E(3)) = -9\mathrm{ch}_0^2(E) < 0.$$

Then by (4.1), the point E^l is in the region $\Delta_{>\frac{1}{2}}$. In particular, E^l is in the line segment $\overline{E^+ e^l}$. Similarly, E^r stands for the reduced character of $\mathsf{R}_E(E(-3))$. It is in the region $\Delta_{>\frac{1}{2}}$ and in the line segment $\overline{E^+ e^r}$. One can check that both of points E^l and E^r are in the parabola $\frac{1}{2}s^2 - q = \frac{1}{2} + \frac{1}{18\mathrm{ch}_0^4(E)}$.

Definition 4.5. ([Li17, Definition 1.4]) The Le Potier curve C_{LP} is a fractal curve defined in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane as

$$\mathrm{C}_{\mathrm{LP}} := \bigsqcup_{\{E = E(\frac{p}{2^m}) \,|\, p \in \mathbb{Z}, \, m \in \mathbb{Z}_{\geq 0}\}} \left(\overline{E^+ e^l} \cup \overline{E^+ e^r} \right) \bigsqcup \{\text{Cantor pieces of } \Delta_{\frac{1}{2}}\}.$$

The region $\mathrm{Geo}_{\mathrm{LP}}$ is defined as $\mathrm{Geo}_{\mathrm{LP}} := \Big\{ (1, s, q) \in \{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}\text{-plane} \,|\, (1, s, q) \text{ is above}$

the curve C_{LP} and is not on line segment $\overline{EE^+}$ for any exceptional bundle $E \Big\}$.

4.3. Special regions associated to an exceptional triple.

Definition 4.6. For an exceptional triple $\mathcal{E} = \{E_1, E_2, E_3\}$, the region $\mathrm{MZ}_{\mathcal{E}}^c$ is defined as the open region in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane bounded by the line segments $\overline{E_1 E_1^r}$, $\overline{E_1^r E_2}$, $\overline{E_2 E_3^l}$, $\overline{E_3^l E_3}$ and $\overline{E_3 E_1}$ (see Figure 2). The region $\mathrm{MZ}_{\mathcal{E}}$ is defined as the open region in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane bounded by line segments $\overline{E_1 E_1^+}$, $\overline{E_1^+ E_2}$, $\overline{E_2 E_3^+}$, $\overline{E_3^+ E_3}$ and $\overline{E_3 E_1}$ (see Figure 3).
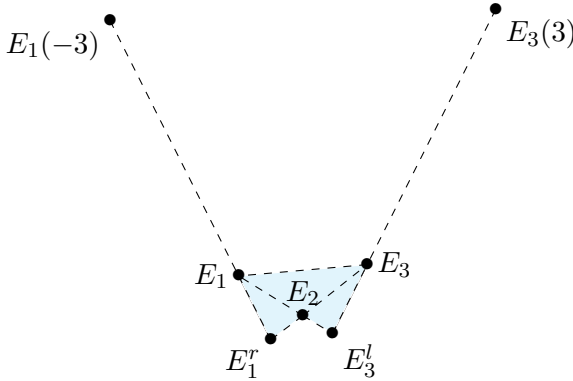
FIGURE 2. The region of $\mathrm{MZ}_{\mathcal{E}}^c$ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane.

We have $\mathrm{MZ}_{\mathcal{E}} \subset \mathrm{Geo}_{\mathrm{LP}}$ and ([Li17, Proposition 2.5])

$$\Theta_{\mathcal{E}}^{\mathrm{Geo}} = \widetilde{\mathrm{GL}^+(2, \mathbb{R})} \cdot \{\sigma_{s,q} \in \mathrm{Stab}^{\mathrm{Geo}}(\mathbb{P}^2) \,|\, (1, s, q) \in \mathrm{MZ}_{\mathcal{E}}\}.$$

Remark 4.7. Let $\mathcal{E} = \{E_1, E_2, E_3\}$ be an exceptional triple. Note that the region $\mathrm{MZ}_{\mathcal{E}}^c$ is a subregion of the $\mathrm{MZ}_{\mathcal{E}}$.

1°. Since E_2 is in the region $\Delta_{[0,\frac{1}{2})}$ by (4.2) and E_1^+, E_3^+ are in the region $\Delta_{>\frac{1}{2}}$ by (4.6), we have

$$e_1^r = \Delta_{\frac{1}{2}} \cap \overline{E_1^+ E_2}, \qquad e_3^l = \Delta_{\frac{1}{2}} \cap \overline{E_2 E_3^+}.$$

2°. The line $L_{E_3(E_3(3))}$ is given as $\{\chi(-, E_3) = \frac{\text{ch}_0(-)}{\text{ch}_0(E_3)}\}$. For every stable vector bundle A with slope between the slopes of E_3 and $E_3(3)$, we have $\chi(A, E_3) \leq 0$. The line segment $\overline{E_3 E_3(3)}$ is contained in Geo_{LP}.

3°. The point E_3^l is on the line segment $\overline{e_3^l E_3^+}$. In particular, The reduced character of any exceptional bundles with slope smaller than that of E_3 is to the left of e_3^l.

4°. By [Li17, Corollary 1.19], the exceptional object $E_3(3)$ is stable with respect to $\sigma_{s,q}$ for any $(1, s, q)$ in $\text{MZ}_\mathcal{E}^c$, and is destabilized by E_3 on the line segment $\overline{E_3 E_3^l}$. In particular, the region $\text{MZ}_\mathcal{E}^c$ is a subregion of $\text{MZ}_\mathcal{E}$ by removing the region that either $E_3(3)$ or $E_1(-3)$ is not stable. In particular, we can identify the region $\text{MZ}_\mathcal{E}^c$ as the following algebraic stability conditions.

Lemma 4.8. *Let $\mathcal{E} = \{E_1, E_2, E_3\}$ be an exceptional triple, then*

$$\Theta_\mathcal{E} \setminus (\Theta_{\mathcal{E},E_1}^{\text{right}} \cup \Theta_{\mathcal{E},E_3}^{\text{left}} \cup \Theta_\mathcal{E}^{\text{Pure}}) = \widetilde{\text{GL}^+(2, \mathbb{R})} \cdot \left\{ \sigma_{s,q} \in \text{Stab}^{\text{Geo}}(\mathbb{P}^2) \mid (1, s, q) \in \text{MZ}_\mathcal{E}^c \right\}.$$

Proof. By the previous Remark 4.7.4, the proof is the same as that for [LZ19a, Lemma 1.29]. $\square$

Definition 4.9. Given an exceptional triple $\mathcal{E} = \{E_1, E_2, E_3\}$ on $\mathcal{D}^b(\mathbb{P}^2)$, we define $\text{MZ}_{E_3}^l$ and $\text{MZ}_{E_1}^r$ as subregions of $\text{MZ}_\mathcal{E}$ as follows:

$$\text{MZ}_{E_3}^l := \left\{ (1, s, q) \in \text{MZ}_\mathcal{E} \mid s < s(E_3), (1, s, q) \text{ is not above the line segment } \overline{E_3 E_3^l} \right\},$$

$$\text{MZ}_{E_1}^r := \left\{ (1, s, q) \in \text{MZ}_\mathcal{E} \mid s > s(E_1), (1, s, q) \text{ is not above the line segment } \overline{E_1 E_1^r} \right\}.$$

Lemma 4.10 (Definition of Θ_E^{left} and Θ_E^{right}). *Given exceptional triples $\mathcal{E}$ and $\mathcal{E}'$ on $\mathcal{D}^b(\mathbb{P}^2)$ with the same $E_3 = E_3' = E$, then $\Theta_{\mathcal{E},E_3}^{\text{left}} = \Theta_{\mathcal{E}',E_3'}^{\text{left}}$. We denote this subspace by Θ_E^{left}. In a similar way, we define the subspace $\Theta_{E_1}^{\text{right}} := \Theta_{\mathcal{E},E_1}^{\text{right}}$. Moreover, we have*

$$\Theta_{E_3}^{\text{left}} = \Theta_{E_3}^- \bigsqcup \widetilde{\text{GL}^+(2, \mathbb{R})} \cdot \left\{ \sigma_{s,q} \in \text{Stab}^{\text{Geo}}(\mathbb{P}^2) \mid (1, s, q) \in \text{MZ}_{E_3}^l \right\}, \qquad (4.7)$$

$$\Theta_{E_1}^{\text{right}} = \Theta_{E_1}^+ \bigsqcup \widetilde{\text{GL}^+(2, \mathbb{R})} \cdot \left\{ \sigma_{s,q} \in \text{Stab}^{\text{Geo}}(\mathbb{P}^2) \mid (1, s, q) \in \text{MZ}_{E_1}^r \right\}. \qquad (4.8)$$

Proof. By Remark 4.7.4, Lemma 4.8 and [Li17, Proposition and Definition 3.1], we have the equation (4.7), where $\Theta_{E_3}^- = \Theta_{\mathcal{E},E_3}^-$ is independent of the choice of E_1 and E_2. Note that by Remark 4.7.3, the boundary segment $\overline{E_3 E_3^l}$ of $\text{MZ}_{E_3}^l$ is also independent of the choice of E_1 and E_2 in the exceptional triple. The subspace Θ_E^{left} is well-defined. Similarly, we have the equation (4.8). $\square$

Remark 4.11. We illustrate the regions in Figure 3. Then we could state Remark 4.7.4 in a precise way, namely for an exceptional triple $\mathcal{E} = \{E_1, E_2, E_3\}$,

$$\text{MZ}_\mathcal{E} = \text{MZ}_{E_1}^r \bigsqcup \text{MZ}_\mathcal{E}^c \bigsqcup \text{MZ}_{E_3}^l. \qquad (4.9)$$
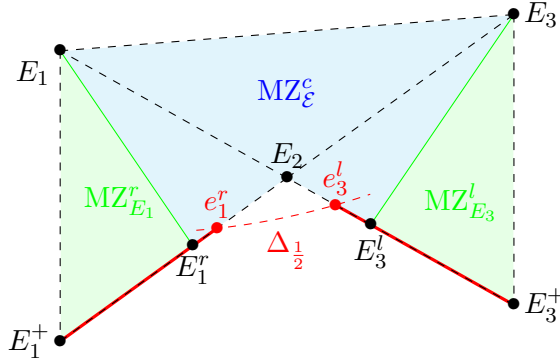
FIGURE 3. The regions of $\mathrm{MZ}_{\mathcal{E}}$, $\mathrm{MZ}_{\mathcal{E}}^c$, $\mathrm{MZ}_{E_3}^l$ and $\mathrm{MZ}_{E_1}^r$ with relation (4.9) in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. The line segments $\overline{E_1^+ e_1^r}$ and $\overline{E_3^+ e_3^l}$ give parts of the Le Potier curve $\mathrm{C_{LP}}$, and $\mathrm{MZ}_{\mathcal{E}} \subset \mathrm{Geo_{LP}}$.

5. CALCULATION OF GLOBAL DIMENSION FUNCTIONS

The main result of this section is to compute the global dimension function on the algebraic stability conditions.

Proposition 5.1. *Let $\mathcal{E} = \{E_1, E_2, E_3\}$ be an exceptional triple on $\mathcal{D}^b(\mathbb{P}^2)$ and $\Theta_{\mathcal{E}}$ be the algebraic stability conditions with respect to $\mathcal{E}$. The value of the global dimension function is*

$$
\mathrm{gldim}(\sigma) = \begin{cases} 2, & when\ \sigma \in \Theta_{\mathcal{E}} \setminus \left(\Theta_{E_1}^{\mathrm{right}} \cup \Theta_{E_3}^{\mathrm{left}} \cup \Theta_{\mathcal{E}}^{\mathrm{Pure}}\right); \\ \phi(\mathsf{R}_{E_1}(\mathbb{S}E_1)) - \phi_1, & when\ \sigma \in \Theta_{E_1}^{\mathrm{right}}; \\ \phi_3 - \phi(\mathsf{L}_{E_3}(\mathbb{S}^{-1}E_3)), & when\ \sigma \in \Theta_{E_3}^{\mathrm{left}}; \\ \phi_3 - \phi_1, & when\ \sigma \in \Theta_{\mathcal{E}}^{\mathrm{Pure}}. \end{cases}
$$

Recall that R and L are the right and left mutations in Section 2.1. The rest of the section is devoted to the proof of the proposition above.

5.1. **The locus with minimum global dimension.** The other three cases are much more subtle, we first discuss the case when $\sigma \in \Theta_{\mathcal{E}} \setminus \left(\Theta_{E_1}^{\mathrm{right}} \cup \Theta_{E_3}^{\mathrm{left}} \cup \Theta_{\mathcal{E}}^{\mathrm{Pure}}\right)$.

Proposition 5.2. *Let σ be a stability condition in $\Theta_{\mathcal{E}} \setminus \left(\Theta_{E_1}^{\mathrm{right}} \cup \Theta_{E_3}^{\mathrm{left}} \cup \Theta_{\mathcal{E}}^{\mathrm{Pure}}\right)$, then* $\mathrm{gldim}(\sigma) = 2$.

The non-trivial part is the '$\leq$' part. As for a brief idea of the proof, we will view σ both as a stability condition in the region $\mathrm{MZ}_{\mathcal{E}}^c$ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane, and as a quiver stability condition. We will show that we only need to concern about $\mathrm{Hom}(F, G[2]) \neq 0$ for two σ stable objects F and G in a same heart with $\phi_{\sigma}(F) < \phi_{\sigma}(G)$. The line segments in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane where F and G are σ stable (i.e. $W_{F\sigma}$ and $W_{G\sigma}$ below) are 'long' enough so that the line segment where $F(-3)$ is stable with respect to

$\sigma(-3)$(i.e. $W_{F\sigma}(-3)$ below) intersects with previous two line segments $W_{F\sigma}$ and $W_{G\sigma}$. Then by the argument as that for stability conditions $\sigma_{s,q}$ above the parabola we show that $\text{Hom}(F, G[2]) = 0$ and get a contradiction. Details of the proof is given as follows.

Proof for Proposition 5.2. By Lemma 4.8, sky-scraper sheaves are all stable with respect to σ. For any closed point $x \in \mathbb{P}^2$, since $\text{Hom}(\mathcal{O}_x, \mathcal{O}_x[2]) = \mathbb{C}$, we have $\text{gldim}(\sigma) \geq 2$.

Step 1: We show that we only need to concern $\text{Hom}(F, G[2]) \neq 0$ for two σ stable objects F and G in a same heart.

By Lemma 4.8, we can view σ as a stability $\sigma_{s,q}$ condition in the region $\text{MZ}_{\mathcal{E}}^c$ in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane. On the other hand, up to a suitable $\mathbb{C}$-action on σ, we may let the heart contain $E_1[2]$, $E_2[1]$ and E_3. Denote this stability condition and its heart by $\tilde{\sigma}$ and $\tilde{\mathcal{A}}$ respectively. As $\{E_1[2], E_2[1], E_3\}$ is an Ext-exceptional collection ([Mac07, Definition 3.10]), an object in the heart is always of the form

$$E_1^{\oplus a_1} \to E_2^{\oplus a_2} \to E_3^{\oplus a_3}$$

for some non-negative integers a_i's.

For any generators $E_i[3-i]$ in $\tilde{\mathcal{A}}$, we always have

$$\text{Hom}(E_i[3-i], E_j[3-j][m]) = 0$$

for every $m \geq 3$. Therefore, for any objects F and G in $\tilde{\mathcal{A}}$, we have

$$\text{Hom}(F, G[m]) = 0$$

for every $m \geq 3$. To prove the '$\leq$' part, we only need to show that for any σ-stable F and G with $\phi(F) < \phi(G)$ in the heart $\tilde{\mathcal{A}}$, we have $\text{Hom}(F, G[2]) = 0$.

Step 2: We show that the phases of F and G are all in $[\phi(E_3(3)), \phi(E_1(-3)[2])]$.

Suppose there are σ-stable F and G with $\phi(F) < \phi(G)$ in the heart $\tilde{\mathcal{A}}$, such that $\text{Hom}(F, G[2]) \neq 0$. Note that $\text{Hom}(E_i[3-i], E_j[3-j][2]) \neq 0$ if and only if $i = 1$ and $j = 3$, we must have

$$\text{Hom}(F, E_3[2]) \neq 0 \text{ and } \text{Hom}(E_1[2], G[2]) \neq 0.$$

By the Serre duality, we have

$$\text{Hom}(E_3(3), F) \neq 0 \text{ and } \text{Hom}(G, E_1(-3)[2]) \neq 0.$$

By [Li17, Corollary 1.19], both objects $E_3(3)$ and $E_1(-3)[2]$ are $\sigma_{s,q}$-stable (hence $\tilde{\sigma}$-stable). Both objects are in the heart $\tilde{\mathcal{A}}$. Therefore, their phases satisfies the inequality:

$$\phi(E_3(3)) \leq \phi(F) < \phi(G) \leq \phi(E_1(-3)[2]). \tag{5.1}$$

Step 3: We show that the walls $W_{F\sigma}$ and $W_{G\sigma}$ are 'long' enough so that the wall $W_{F\sigma}(-3)$ intersects the walls $W_{F\sigma}$ and $W_{G\sigma}$. We compare their slopes and get the contradiction.

Here the *wall* $W_{F\sigma} := \{(1, s, q) \in \{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane| the line segment along the line $L_{F\sigma}$ that is above the Le Potier curve $C_{\text{LP}}\}$ and the wall

$$W_{F\sigma}(-3) := \{(1, s-3, q-3s+\frac{9}{2}) \mid (1, s, q) \in W_{F\sigma}\}.$$

By Bertram's nested wall theorem, [LZ19a, Corollary 1.24], the object F is stable along the wall $W_{F\sigma}$. Let $F_a = (1, s(F_a), q(F_a))$ and $F_b = (1, s(F_b), q(F_b))$ be the two edges of the wall $W_{F\sigma}$ as that in the Figure 4. We denote similar notations for G as that for F.

FIGURE 4. Compare the slopes of the wall $W_{G\sigma}$ and the wall $W_{F\sigma}(-3)$.

By the relation of phases as that in (5.1), counter-clockwisely, one has the line segment $\overline{\sigma_{s,q}(E_3(3))}$, $\overline{\sigma_{s,q}F_b}$, $\overline{\sigma_{s,q}G_b}$ and $\overline{\sigma_{s,q}(E_1(-3))}$. In particular, either the wall $W_{F\sigma}$ is a vertical wall (parallel to the $\frac{\mathrm{ch}_2}{\mathrm{ch}_0}$-axis) or $|s(F_b) - s| > 3$. Same statements holds for $W_{G\sigma}$. In every case, the segment

$$\overline{\sigma_{s,q}(-3)F_b(-3)} = \overline{(1, s-3, q-3s+\frac{9}{2})(1, s(F_b) - 3, q(F_b) - 3s(F_b) + \frac{9}{2})}$$

intersects both segments $\overline{\sigma_{s,q}F_b}$ and $\overline{\sigma_{s,q}G_b}$ at P and Q respectively. The object $F(-3)$ is stable at both P and Q. By comparing the slopes, we have

$$\phi_Q(F(-3)) = \phi_P(F(-3)) < \phi_P(F) = \phi_{s,q}(F) < \phi_{s,q}(G) = \phi_Q(G).$$

By the Serre duality,

$$\mathrm{Hom}(F, G[2]) \cong (\mathrm{Hom}(G, F(-3)))^* = 0.$$

We get the contradiction. $\qquad\square$

5.2. **The global dimension on the leg locus.** We discuss the case that $\sigma \in \Theta_{E_3}^{\mathrm{left}}$. We first recall the following basic properties for an exceptional triple $\mathcal{E} = \{E_1, E_2, E_3\}$. Denote by $\mathrm{rk}(E) = \mathrm{ch}_0(E)$ and $\hom(E, F) = \dim \mathrm{Hom}(E, F)$.

Lemma 5.3. *For an exceptional triple $\mathcal{E} = \{E_1, E_2, E_3\}$, the ranks and homs of these exceptional objects satisfy the following equations.*

$$(\mathrm{rk}E_1)^2 + (\mathrm{rk}E_2)^2 + (\mathrm{rk}E_3)^2 = 3\mathrm{rk}E_1\mathrm{rk}E_2\mathrm{rk}E_3 \quad \text{(Markov equation)},$$
$$\hom(E_1, E_2) = 3\mathrm{rk}E_3, \quad \hom(E_2, E_3) = 3\mathrm{rk}E_1,$$
$$\hom(E_1, E_3) = 9\mathrm{rk}E_1\mathrm{rk}E_3 - 3\mathrm{rk}E_2.$$

The object $\mathsf{L}_{E_3}(E_3(3))[-1]$ *admits a resolution:*

$$0 \to E_1^{\oplus \hom(E_1, E_3)} \to E_2^{\oplus r} \to \mathsf{L}_{E_3}(E_3(3))[-1] \to 0, \tag{5.2}$$

where $r = \hom(E_1, E_3)\hom(E_1, E_2) - \hom(E_2, E_3)$.

Proof. The equations of rank and hom are well-known in [GR87]. As for the last statement, we consider the resolution of $E_3(3)$. Note that $\mathcal{D}_\infty(\mathbb{P}^2)$ has the semiorthorgonal decomposition $\langle E_1, E_2, E_3 \rangle$, so an object A admits a unique filtration

$$0 = F_0 \subset F_1 \subset F_2 \subset F_3 = A$$

such that $\mathrm{Cone}(F_i \to F_{i+1}) \in \langle E_{3-i} \rangle$ for $i = 0, 1, 2$. The term $\mathrm{Cone}(F_0 \to F_1)$ is given by $\bigoplus_i E_3[i] \otimes \mathrm{Hom}(E_3[i], A)$, while the term $\mathrm{Cone}(F_2 \to F_3)$ is given by $\bigoplus_i E_1[i] \otimes \mathrm{Hom}(A, E_1[i])^*$.

When $A = E_3(3)$, we have $\mathrm{Cone}(F_0 \to F_1) = E_3 \otimes \mathrm{Hom}(E_3, E_3(3)) = E_3^{\oplus 9(\mathrm{rk}E_3)^2 + 1}$ and $\mathrm{Cone}(F_2 \to F_3) = E_1^{\oplus \hom(E_1, E_3)}[2]$. The factor $\mathrm{Cone}(F_1 \to F_2)$ can only be $E_2^{\oplus r}[1]$. By the equations of rank and hom in the lemma, the rank

$$r = \frac{9(\mathrm{rk}E_3)^3 + 9\mathrm{rk}E_3(\mathrm{rk}E_1)^2}{\mathrm{rk}E_2} - 3\mathrm{rk}E_1 = \hom(E_1, E_3)\hom(E_1, E_2) - \hom(E_2, E_3).$$

Note that $\mathsf{L}_{E_3}(E_3(3))[-1]$ is the kernel of the map $E_3 \otimes \mathrm{Hom}(E_3, E_3(3)) \xrightarrow{\mathrm{ev}} E_3(3)$, the resolution sequence is clear. □

Lemma 5.4. *Let* σ *be a stability condition in* $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$, *suppose an object* $F = \mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})$ *is stable with respect to* σ, *then* F *is stable everywhere in* $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$.

Proof. For any stability condition in $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$, by a suitable $\mathbb{C}$-action, we may assume that the heart contains $E_1[2]$, $E_2[1]$ and $E_3[n]$ for some $n \leq 0$. As $\{E_1[2], E_2[1], E_3[n]\}$ is an Ext-exceptional collection, an object in the heart is always of the form

$$E_1^{\oplus a_1} \to E_2^{\oplus a_2} \to E_3^{\oplus a_3}.$$

The object $F[1]$ can only be destabilized by some subobjects $F' = \mathrm{Cone}(E_1^{\oplus a'} \to E_2^{\oplus b'})[1]$ in the heart generated by $\{E_1[2], E_2[1], E_3[n]\}$ with larger phase, which means $\frac{a'}{b'} > \frac{a}{b}$. Note that this is independent with choice of σ in $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$, the object F is stable everywhere in $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$. □

Now we are ready to compute the example achieving the value of the global dimension function in the region of $\Theta_{E_3}^{\mathrm{left}}$.

Lemma 5.5. *Let* σ *be a stability condition in* $\Theta_{E_3}^{\mathrm{left}}$, *then* $\mathsf{L}_{E_3}E_3(3)$ *is* σ-*stable and it has a non-zero morphism to* $E_3[2]$. *In particular, we have* $\mathrm{gldim}(\sigma) \geq \phi_3 - \phi(\mathsf{L}_{E_3}E_3(3)) + 2$.

Proof. By [LZ19a, Corollary 3.2], the object $\mathsf{L}_{E_3}E_3(3) = \mathrm{Cone}(E_3 \otimes \mathrm{Hom}(E_3, E_3(3)) \xrightarrow{\mathrm{ev}} E_3(3))$ is $\sigma_{s,q}$-stable for $(1, s, q)$ which is slightly above the line segment $\overline{E_3 E_3(3)}$ in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. The object $\mathsf{L}_{E_3}E_3(3)$ is stable along the line segment $\overline{(1, s, q)E_3^l}$. As this segment intersects $\mathrm{MZ}_{\mathcal{E}}^c$, by Lemma 4.8, the object $\mathsf{L}_{E_3}E_3(3)$ is stable with respect

to some stability condition in $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$. By Lemma 5.3 and 5.4, the object $\mathsf{L}_{E_3}E_3(3)$ is σ-stable for every $\sigma \in \Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$.

By applying $\mathrm{Hom}(-, E_3[2])$ on the distinguished triangle $E_3 \otimes \mathrm{Hom}(E_3, E_3(3)) \xrightarrow{\mathrm{ev}} E_3(3) \to \mathsf{L}_{E_3}E_3(3) \xrightarrow{+}$, we have $\mathrm{Hom}(\mathsf{L}_{E_3}E_3(3), E_3[2]) \cong \mathrm{Hom}(E_3(3), E_3[2]) = \mathbb{C}$. $\qquad\square$

As for the '$\leq$' direction, we first treat with the easier case that the stable objects can be classified.

Proposition 5.6. *Let σ be a stability condition in $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1 < \phi_3 - 1)$, then up to a homological shift, a σ-stable object is either*

- E_3 *or*
- $\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})$

induced by a stable quiver representation $\mathbb{C}^{\oplus a} \xRightarrow{\mathrm{hom}(E_1, E_2)\ \mathit{arrows}} \mathbb{C}^{\oplus b}$. Moreover, we have $\mathrm{gldim}(\sigma) = \phi_3 - \phi(\mathsf{L}_{E_3}E_3(3)) + 2$.

Proof. By a suitable $\mathbb{C}$-action, we may assume that the heart contains $E_1[2]$, $E_2[1]$ and $E_3[n]$ for some $n \leq -1$. As $\{E_1[2], E_2[1], E_3[n]\}$ is an Ext-exceptional collection, an object in the heart is $\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})[1] \bigoplus E_3^{\oplus c}[n]$. An object $\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})[1]$ is σ-stable if and only if for any non-zero proper subobject $\mathrm{Cone}(E_1^{\oplus a_1} \to E_2^{\oplus b_1})[1]$ we have $\frac{a_1}{b_1} < \frac{a}{b}$. The first part of the statement is clear.

As for the second part of the statement, by Lemma 5.5, we only need to show the '$\leq$' side, note that for any two stable objects F and F' in the form of $\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})$, we always have $\mathrm{Hom}(F, F'[m]) = 0$ for $m \geq 2$. By the classification of stable objects, we only need to consider potential non-zero morphisms from $\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})$ to $E_3[m]$ for $m \geq 1$. When $\phi(\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})) < \phi(\mathsf{L}_{E_3}E_3(3)[-1])$, by Lemma 5.3, we have

$$\frac{b}{a} > \mathrm{hom}(E_1, E_2) - \frac{\mathrm{hom}(E_2, E_3)}{\mathrm{hom}(E_1, E_3)}. \tag{5.3}$$

Let ϕ_3' be $\phi(\mathsf{L}_{E_3}E_3(3))$, which is greater than $\phi_1 + 1$. We consider the stability condition σ' in $\Theta_{\mathcal{E}}$ given by $(m_1, m_2, m_3, \phi_1, \phi_2, \phi_3')$. By Lemma 5.4, $\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})$ is σ'-stable and $E(3)$ is σ'-semistable with phase $\phi_3' = \phi'(E_3) = \phi'(\mathsf{L}_{E_3}E_3(3))$. By (5.3), we have $\phi'(\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b})) < \phi'(\mathsf{L}_{E_3}E_3(3)) - 1 = \phi'(E(3)) - 1$. Therefore, for any $m \geq 1$, by the Serre duality, we have

$$\mathrm{Hom}(\mathrm{Cone}(E_1^{\oplus a} \to E_2^{\oplus b}), E_3[m]) \cong (\mathrm{Hom}(\mathrm{Cone}(E_3(3)[m-2], E_1^{\oplus a} \to E_2^{\oplus b})))^* = 0.$$

As a summary, the global dimension at σ is $\phi_3 - \phi(\mathsf{L}_{E_3}E_3(3)) + 2$, and is achieved via the morphism between $\mathsf{L}_{E_3}E_3(3)$ and $E_3[2]$. $\qquad\square$

We finally treat with region $\Theta_{\mathcal{E}}(\phi_2 < \phi_3 - 1 < \phi_1 + 1) \cap \Theta_{E_3}^{\mathrm{left}}$, where the stable objects are more complicated. In this case, the potential stable characters are away from the kernel of central charge of every σ in $\Theta_{\mathcal{E}}(\phi_2 < \phi_3 - 1 < \phi_1 + 1)$. We will think both the stable characters and (kernels of central charges of) stability conditions in the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane. This will allow us to show the vanishing of certain morphisms by comparing slopes.

We first prove a nested wall result for the algebraic stability stability conditions. Denote $\Theta_{\mathcal{E}}^+(\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2) \coloneqq \{\sigma \in \Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2)|$ the kernel of central charge of σ is spanned by $(1, s, q)$ for some $s > s(E_1)\}$.

Lemma 5.7. *Let σ be a stability condition in $\Theta_{\mathcal{E}}^+(\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2)$ and G be a σ-stable object. Then G is σ'-stable for every σ' in $\Theta_{\mathcal{E}}^+(\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2)$ with kernel of central charge on the line through G and σ.*

Proof. In the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane, the kernel of the central charge of σ is in the region bounded by rays through $E_1 E_3$, $E_1 E_2$ as shown in the Figure 5 (Area I).
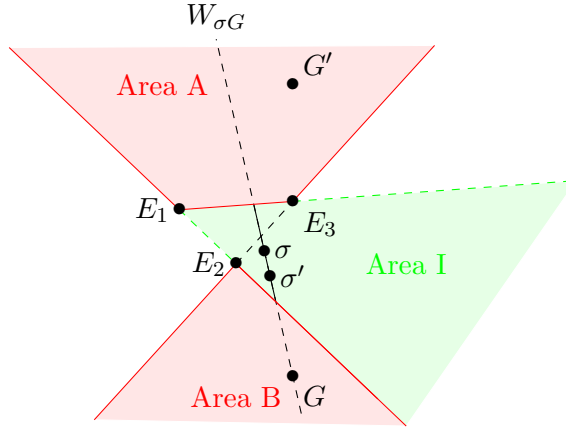

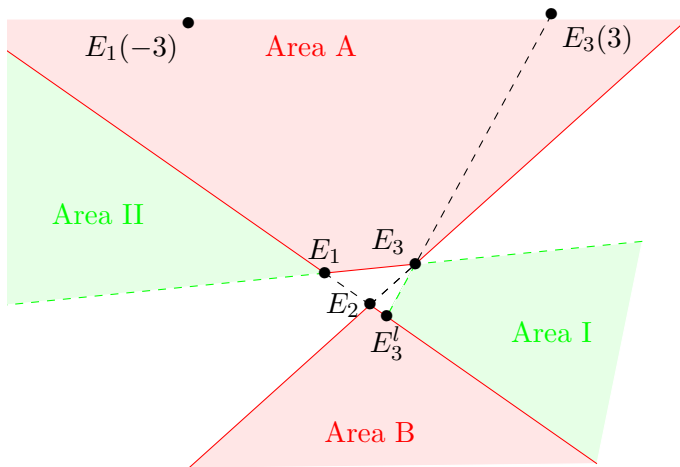FIGURE 5. Stability conditions through $W_{\sigma G}$.

By a suitable $\mathbb{C}$-action, we may assume that the heart contains $\{E_1[2], E_2[1], E_3\}$. Denote this heart by $\tilde{\mathcal{A}}$, then an object in $\tilde{\mathcal{A}}$ is of the form $E_1^{\oplus a_1} \to E_2^{\oplus a_2} \to E_3^{\oplus a_3}$. In particular, the reduced character of a stable object is in the closed region (Area A $\cup$ Area B in Figure 5) bounded by the rays through $E_1 E_2$, $E_2 E_3$ and line segment through $E_1 E_3$.

The phase of G is determined by the slope of line through σ and $\mathrm{v}(G)$. As for another object G', its phase $\phi(G') < \phi(G)$ if and only if the line through σ and $\mathrm{v}(G')$ rotates counter-clockwisely to the line through σ and $\mathrm{v}(G)$ without passing though the line through σ and $E_1[2]$.

For every non-zero proper subobject G' of G in $\tilde{\mathcal{A}}$, since G is stable, G' has smaller phase than that of G. In the $\{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane, that is equivalent to the following description for $\mathrm{v}(G')$:

The reduced character of G' is either to the right of the line through G and σ when it is in Area A, or it is to the left of the line through G and σ when it is in Area B.

Note that for every stability condition σ' in $\Theta_{\mathcal{E}}^+(\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2)$ with kernel of central charge on the line through G and σ, the line through G' and σ' rotates counter-clockwisely to the line through σ, σ' and G. The object G is σ'-stable. $\square$

Proposition 5.8. *Let σ be a stability condition in $\Theta_{\mathcal{E}}(\phi_2 < \phi_3 - 1 < \phi_1 + 1) \cap \Theta_{E_3}^{\text{left}}$.
Then* $\text{gldim}(\sigma) = \phi(E_3) - \phi(\mathsf{L}_{E_3}E_3(3)) + 2$.

Proof. By Lemma 5.5, we only need to show the '$\leq$' part.

Step 1: We show that we only need to concern $\text{Hom}(F, G[2]) \neq 0$ for F and G in a same heart.

By a suitable $\mathbb{C}$-action, we may assume that the heart contains $\{E_1[2], E_2[1], E_3\}$. Denote this heart by $\tilde{\mathcal{A}}$, we have the same description for objects in $\tilde{\mathcal{A}}$ as that in Lemma 5.7.

In the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane, the kernel of the central charge of σ is in the region bounded by rays through $E_1 E_3$, $E_1 E_2$ and line segment $\overline{E_3 E_3^l}$ as shown in the Figure 6 (Area I $\cup$ Area II). Recall that the point E_3^l and $\mathsf{L}_{E_3}(E_3(3))$ are the same point in the $\{1, \frac{\text{ch}_1}{\text{ch}_0}, \frac{\text{ch}_2}{\text{ch}_0}\}$-plane.

FIGURE 6. Stable characters are in Area A $\cup$ Area B. The kernels of the central charges are in Area I $\cup$ Area II.

As for any generators in $\{E_1[2], E_2[1], E_3\}$, we have $\text{Hom}(-, -[m]) = 0$ for any $m \geq 3$. For any objects F and G in the heart, we have $\text{Hom}(F, G[m]) = 0$ for any $m \geq 3$. To prove the '$\leq$' part of the statement, we only need consider $\text{Hom}(F, G[2]) \neq 0$ for stable objects F, G in the heart with $\phi(F) < \phi(G)$.

Suppose there are σ-stable objects F and G with

$$\phi(G) - \phi(F) > \phi(E_3) - \phi(\mathsf{L}_{E_3}E_3(3)) \tag{5.4}$$

in the heart $\tilde{\mathcal{A}}$, such that $\text{Hom}(F, G[2]) \neq 0$. By the same argument as that in Proposition 5.2, we must have

$$\text{Hom}(E_3(3), F) \neq 0 \text{ and } \text{Hom}(G, E_1(-3)[2]) \neq 0. \tag{5.5}$$

Step 2: We show that $\phi(G) > \phi(E_3)$.

Suppose $\phi(F) < \phi(\mathsf{L}_{E_3}E_3(3))$, then $\phi(F) < \phi(E_3)$. Therefore, the object F is of the form $\text{Cone}(E_1^{\oplus a_F} \to E_2^{\oplus b_F})[1]$. By Lemma 5.4, F is stable everywhere in $\Theta(\phi_1 + 1 > \phi_2)$.

In particular, it is stable with every stability condition σ' on the line segment $\overline{E_3 E_3^l}$, where $E_3(3)$ is σ'-semistable. Since $\mathrm{Hom}(E_3(3), F) \neq 0$, we have $\phi'(F) \geq \phi'(E_3(3)) = \phi'(\mathsf{L}_{E_3} E_3(3))$. Therefore, we have

$$\frac{b_F}{a_F} \leq \mathrm{hom}(E_1, E_2) - \frac{\mathrm{hom}(E_2, E_3)}{\mathrm{hom}(E_1, E_3)},$$

and $\phi'(F) \geq \phi'(\mathsf{L}_{E_3} E_3(3))$, which contradicts the assumption that $\phi(F) < \phi(\mathsf{L}_{E_3} E_3(3))$.

By (5.4), we must have

$$\phi(G) > \phi(E_3). \tag{5.6}$$

Step 3: We show that the kernel of the central charge of σ is in Area I and is below the line through $E_1(-3)[2]$ and E_3, i.e. the open region bounded by line segments $\overline{R E_3}$, $\overline{E_3 E_3^l}$ and $\overline{E_3^l R}$, with $R := L_{E_1(-3)E_3} \cap L_{E_1 E_2}$ as in Figure 7.

Let the central charge of $E_i[3 - i]$ be z_i for $i = 1, 2, 3$. Let the object G be of the form $E_1^{\oplus n_1} \to E_2^{\oplus n_2} \to E_3^{\oplus n_3}$.

By Lemma 5.4 and a same argument as that in Lemma 5.5, the object $\mathsf{L}_{E_3}(E_1(-3)[2])$ is of the form $\mathrm{Cone}(E_1^{\oplus r_1} \to E_2^{\oplus r_2})$ in the heart $\tilde{\mathcal{A}}$, and is stable with respect to every stability condition in $\Theta_{\mathcal{E}}(\phi_2 < \phi_1 + 1)$. By (5.6), we have $\mathrm{Hom}(G, E_3) = 0$. By (5.5), we have $\mathrm{Hom}(G, \mathsf{L}_{E_3}(E_1(-3)[2])) \neq 0$. Therefore, we have

$$\phi(\mathsf{L}_{E_3}(E_1(-3)[2])) > \phi(G) > \phi(E_3). \tag{5.7}$$

Therefore, the kernel of the central charge of σ is in Area I and is below the line $L_{E_1(-3)E_3}$ as in Figure 7.
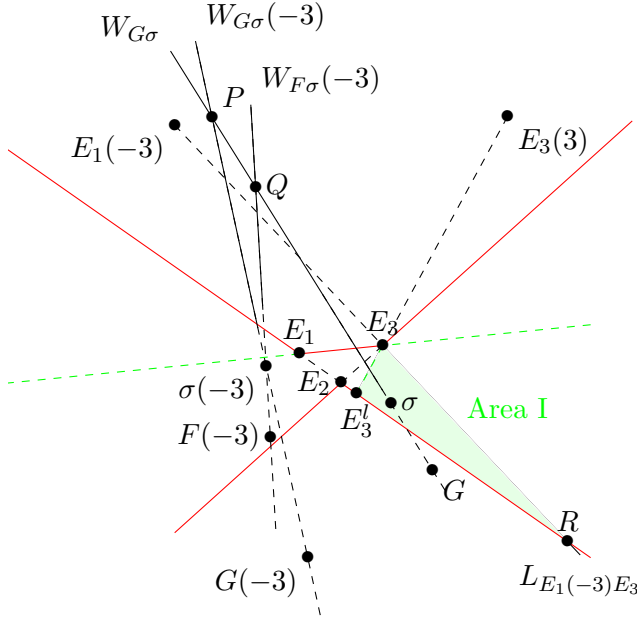

FIGURE 7. Comparing the phases of G and $F(-3)$.

Step 4: We show that the wall $W_{G\sigma}$ intersects the wall $W_{G\sigma}(-3)$. We denote the intersection point by $P := W_{G\sigma} \cap W_{G\sigma}(-3)$.

Consider the line $L_{G\sigma}$ through the reduced character of G and the kernel of the central charge of σ, which is in Area I. In particular, the stability condition

$$\sigma \in \Theta_{\mathcal{E}}^+ (\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2).$$

By (5.7), the line $L_{G\sigma}$ intersects the line segment $\overline{E_1(-3)E_3}$. Therefore, the line $L_{G\sigma}$ intersects the region MZ$_{\mathcal{E}}^c$.

By Lemma 5.7, the object G is stable with respect to every stability condition in σ in $\Theta_{\mathcal{E}}^+ (\phi_2 < \phi_1 + 1, \phi_3 < \phi_1 + 2)$ with kernel on $L_{G\sigma}$. Note that there exists a point $(1, s_0, q_0)$ in MZ$_{\mathcal{E}}^c \cap L_{G\sigma}$, by Lemma 4.8, the object G is σ_{s_0,q_0}-stable.

Recall the wall $W_{G\sigma} := \{(1, s, q) \in \{1, \frac{\mathrm{ch}_1}{\mathrm{ch}_0}, \frac{\mathrm{ch}_2}{\mathrm{ch}_0}\}$-plane$|$ the line segment along the line $L_{G\sigma}$ that is above the Le Potier curve $C_{\mathrm{LP}}\}$. By the Bertram's nest wall theorem [LZ19a, Corollary 1.24], the object G is $\sigma_{s,q}$-stable for every $(1, s, q)$ on $W_{G\sigma}$. Note that $W_{G\sigma}$ intersects the line segment $\overline{E_1 E_3}$, but does not intersects the line segment $\overline{E_1(-3)E_1}$ or $\overline{E_3 E_3(3)}$, both of which are above the Le Potier curve C_{LP}. Therefore the horizontal length of $W_{G\sigma}$ is greater than 3 when $W_{G\sigma}$ is not the vertical wall. Let

$$W_{G\sigma}(-3) := \{(1, s - 3, q - 3s + \frac{9}{2}) \mid (1, s, q) \in W_{G\sigma}\},$$

then $G(-3)$ is $\sigma_{s,q}$-stable for every $(1, s, q)$ on $W_{G\sigma}(-3)$. The wall $W_{G\sigma}(-3)$ intersects the wall $W_{G\sigma}$ at some point P and

$$\phi_P(G(-3)) < \phi_P(G). \tag{5.8}$$

As for the only exceptional case that $W_{G\sigma}$ is the vertical wall, we can view that the point P is at $(0, 0, 1)$. This will not affect the statement in the next step.

Step 5: When $s(F) > s(E_3)$, we show that the wall $W_{F\sigma}(-3)$ intersects the wall $W_{G\sigma}$. We denote the intersection point by $Q := W_{G\sigma} \cap W_{F\sigma}(-3)$.

By (5.7), we have the same bounds for F

$$\phi(\mathsf{L}_{E_3}(E_1(-3)[2])) > \phi(G) > \phi(F) > \phi(E_3). \tag{5.9}$$

The horizontal length of $W_{F\sigma}(-3)$ is greater than 3 when it is not vertical. Note that the slope of $W_{F\sigma}(-3)$ is less than that of $W_{G\sigma}(-3)$, the segment $W_{F\sigma}(-3)$ intersects $W_{G\sigma}$ at Q on the line segment $\overline{P\sigma}$. The fact that $\phi_\sigma(G) > \phi_\sigma(F)$ implies $\phi_{\sigma(-3)}(G(-3)) > \phi_{\sigma(-3)}(F(-3))$. Both $F(-3)$ and G are σ_Q stable. We then compare their phases at Q by using (5.8) as follows:

$$\phi_Q(G) = \phi_P(G) > \phi_P(G(-3)) = \phi_{\sigma(-3)}(G(-3)) > \phi_{\sigma(-3)}(F(-3)) = \phi_Q(F(-3)).$$

So $\mathrm{Hom}(G, F(-3)) = 0$. By the Serre duality, we have $\mathrm{Hom}(F, G[2]) = 0$.

Step 6: When $s(F) \leq s(E_3)$, we reduce this case to Proposition 5.2 .

Note that F is of the form $E_1^{\oplus a_1} \to E_2^{\oplus a_2} \to E_3^{\oplus a_3}$, we have $\mathrm{Hom}(E_3, F) \neq 0$ when $a_3 \neq 0$. The object F is either of the form $\mathrm{Cone}(E_1^{\oplus a_1} \to E_2^{\oplus a_2})[1]$ or E_3. Let $(1, s_0, q_0)$ be a point in MZ$_{\mathcal{E}}^c \cap L_{G\sigma}$. By Lemma 5.4, in any case, F is σ_{s_0,q_0}-stable. By Lemma 5.7 and Lemma 4.8, the object G is also σ_{s_0,q_0}-stable and has phase

$$\phi_{s_0,q_0}(G) > \phi_{s_0,q_0}(F).$$

By Proposition 5.2, we have $\mathrm{Hom}(F, G[2]) = 0$.

As a summary, we have shown that $\mathrm{Hom}(F, G[2]) = 0$ when $\phi(F) < \phi(\mathsf{L}_{E_3} E_3(3))$ or $\phi(G) > \phi(E_3)$. In particular, we have $\mathrm{gldim}(\sigma) = \phi(E_3) - \phi(\mathsf{L}_{E_3} E_3(3)) + 2$. $\qquad\square$

Proof for Proposition 5.1. When $\sigma \in \Theta_{\mathcal{E}} \setminus \left(\Theta_{E_1}^{\mathrm{right}} \cup \Theta_{E_3}^{\mathrm{left}} \cup \Theta_{\mathcal{E}}^{\mathrm{Pure}} \right)$, the global dimension is computed in Proposition 5.2.

When $\sigma \in \Theta_{E_3}^{\mathrm{left}}$, the global dimension is computed in Propositions 5.6 and 5.8.

When $\sigma \in \Theta_{E_1}^{\mathrm{right}}$, we take the derived dual stability condition $\sigma^{\vee} \in \Theta_{\mathcal{E}^{\vee}, E_1^{\vee}}^{\mathrm{left}}$, where $\mathcal{E}^{\vee}$ is the dual exceptional triple $\{E_3^{\vee}, E_2^{\vee}, E_1^{\vee}\}$. We reduce to the previous case and have

$$\mathrm{gldim}(\sigma) = \mathrm{gldim}(\sigma^{\vee}) = \phi^{\vee}(E_1^{\vee}) - \phi^{\vee}(\mathsf{L}_{E_1^{\vee}}(\mathbb{S}^{-1} E_1^{\vee}))$$
$$= -\phi(E_1) - \phi^{\vee}\left((\mathsf{R}_{E_1}(\mathbb{S} E_1))^{\vee} \right) = \phi(\mathsf{R}_{E_1}(\mathbb{S} E_1)) - \phi_1.$$

When $\sigma \in \Theta_{\mathcal{E}}^{\mathrm{Pure}}$, by Lemma 4.4, the only stable objects are $E_i[m]$ for $E_i \in \mathcal{E}$ and $m \in \mathbb{Z}$. As $\mathcal{E}$ is a strong exceptional collection, we have $\mathrm{Hom}(E_i, E_j[m]) \neq 0$ if and only if $j \geq i$ and $m = 0$. So the result is clear. $\qquad\square$

Remark 5.9. Following the notations in Remark 4.11, for any exceptional bundle E, we associate two regions MZ_E^l and MZ_E^r, which consist with geometric stability conditions. Moreover, if $\sigma \in \mathrm{MZ}_E^l \setminus \overline{E^l E}$ or $\sigma \in \mathrm{MZ}_E^r \setminus \overline{E E^r}$, we have $2 < \mathrm{gldim}(\sigma) < 3$.

Corollary 5.10. *The global dimension function*

$$\mathrm{gldim} \colon \mathrm{Stab}^{\dagger}(\mathbb{P}^2) \to \mathbb{R}_{\geq 0}$$

has minimum value 2 and $\mathrm{gldim}\, \mathrm{Stab}^{\dagger}\, \mathbb{P}^2 = [2, \infty)$. *Moreover, the subspace* $\mathrm{gldim}^{-1}(2)$ *is contained in* $\overline{\mathrm{Stab}^{\mathrm{Geo}}(\mathbb{P}^2)}$, *and is contractible.*

Proof. The image of gldim follows from Proposition 3.4, Proposition 5.1 and the description of $\mathrm{Stab}^{\dagger}\, \mathbb{P}^2$ (2.6). The contractibility of $\mathrm{gldim}^{-1}(2)$ is clear. $\qquad\square$

6. Contractibility via global dimension

We denote $\mathrm{gldim}^{-1}(I)$ by the space of all stability conditions in the component $\mathrm{Stab}^{\dagger}(\mathbb{P}^2)$ with global dimension in I for an interval $I \subset \mathbb{R}$. Base on Proposition 5.1 and the cell-decomposition description for $\mathrm{Stab}^{\dagger}(\mathbb{P}^2)$, our main result shows that the connected component $\mathrm{Stab}^{\dagger}(\mathbb{P}^2)$ is contractible via the global dimension:

Theorem 6.1. *For any $x > 2$, the space $\mathrm{gldim}^{-1}([2, x))$ contracts to $\mathrm{gldim}^{-1}(2)$.*

Proof. By Proposition 5.1, Remark 4.7 and [Li17, Corollary 3.5, Theorem 3.9], the space $\mathrm{gldim}^{-1}([2, x))$ has a cell decomposition as

$$\mathrm{gldim}^{-1}(2) \bigcup \left(\bigsqcup_{E \text{ excep. bundles}} \left(\Theta_E^{\mathrm{left}}(x) \bigsqcup \Theta_E^{\mathrm{right}}(x) \right) \bigsqcup \left(\bigsqcup_{\mathcal{E} \text{ excep. triples}} \Theta_{\mathcal{E}}^{\mathrm{Pure}}(x) \right) \right). \tag{6.1}$$

Here each $\Theta_*^{\dagger}(x)$ stands for $\Theta_*^{\dagger} \cap \mathrm{gldim}^{-1}([2, x))$. By Proposition 5.1, we have $\Theta_{\mathcal{E}}^{\mathrm{Pure}}(x) = \Theta_{\mathcal{E}}^{\mathrm{Pure}}(\phi_3 - \phi_1 < x)$. Each $\Theta_{\mathcal{E}}^{\mathrm{Pure}}(x)$ has an open neighborhood, say, $\Theta_{\mathcal{E}}(\phi_3 - \phi_2 > \frac{1}{2}, \phi_2 - \phi_1 > \frac{1}{2}, \phi_3 - \phi_1 < x)$, in $\Theta_{\mathcal{E}}(x)$ which does not intersect any other $\Theta_{\mathcal{E}'}^{\mathrm{Pure}}(x)$. As $\mathrm{Stab}^{\dagger}(\mathbb{P}^2)$ admits a metric, we may then choose open neighborhoods

of $\Theta_{\mathcal{E}}^{\mathrm{Pure}}(x)$'s which do not intersect with each other. By the cell decomposition, the space gldim$^{-1}([2, x))$ contracts to its subspace

$$A(x) \coloneqq \mathrm{gldim}^{-1}(2) \bigcup \left(\bigsqcup_{E \text{ excep. bundles}} \left(\Theta_E^{\mathrm{left}}(x) \bigsqcup \Theta_E^{\mathrm{right}}(x) \right) \right).$$

For each exceptional object E, let $\mathcal{E} = \{E_1, E_2, E_3\}$ be an exceptional collection such that $E_3 = E$. By Proposition 5.1 and Lemma 5.3, we have

$$\Theta_E^{\mathrm{left}}(x) \cong \{(m_1, m_2, m_3, \phi_1, \phi_2, \phi_3) \in (\mathbb{R}_{>0})^3 \times \mathbb{R}^3 | \phi_1 < \phi_2 < \phi_1 + 1,$$

$$\phi_3 > \phi_2 + 1 + \frac{1}{\pi} \arctan \left(\frac{\sin((\phi_1 + 1 - \phi_2)\pi)}{\cos((\phi_1 + 1 - \phi_2)\pi) + \frac{m_2}{m_1} h} \right) > \phi_3 - x + 2\},$$

where $h = \hom(E_1, E_2) - \frac{\hom(E_2, E_3)}{\hom(E_1, E_3)}$. Therefore, the space $\Theta_E^{\mathrm{left}}(x)$ contracts to $\Theta_E^{\mathrm{left}}(x) \cap$ gldim$^{-1}(2)$.

By Remark 4.7 and [Li17, Lemma 3.7], each $\Theta_E^{\mathrm{left}}(x)$ has an open neighborhood in $A(x)$, which does not intersects any other $\Theta_{E'}^{\mathrm{left}}(x)$ or $\Theta_{E'}^{\mathrm{right}}(x)$. Same argument works for all $\Theta_E^{\mathrm{right}}(x)$, we may therefore contract all $\Theta_E^{\mathrm{left}}(x)$ and $\Theta_E^{\mathrm{right}}(x)$ in $A(x)$ simultaneously to gldim$^{-1}(2)$, which is a contractible space. $\qquad\square$

7. Inducing stability conditions from projective plane to the local projective plane

Let Y be the total space of the canonical bundle of $\mathbb{P}^2$, and $i : \mathbb{P}^2 \hookrightarrow Y$ be the inclusion of the zero-section. We write $\mathcal{D}_{\mathbb{P}^2}^b(Y)$ for the subcategory of $\mathcal{D}^b(Y)$ of complexes with bounded cohomology, such that all of its cohomology sheaves are supported on the zero-section. The space of Bridgeland stability conditions on $\mathcal{D}_{\mathbb{P}^2}^b(Y)$ has been studied by Bayer and Macrì [BM11]. In this section, we prove that the stability conditions in gldim$^{-1}(2) \subset \overline{\mathrm{Stab}}^{\mathrm{Geo}}(\mathbb{P}^2)$ can be used to induce stability conditions on $\mathcal{D}_{\mathbb{P}^2}^b(Y)$ by Ikeda-Qiu's inducing theorem, via q-stability conditions on Calabi–Yau-$\mathbb{X}$ categories.

Following the notion in [IQ18a], we have the Calabi–Yau-$\mathbb{X}$ version of $\mathcal{D}_\infty(\mathbb{P}^2)$

$$\mathcal{D}_{\mathbb{X}}(\mathbb{P}^2) \coloneqq \mathcal{D}_{c,\mathbb{C}^*}^b(Y). \tag{7.1}$$

By [IQ18a, Proposition 3.14], we have $\mathcal{D}_{\mathbb{X}}(\mathbb{P}^2) \cong \mathcal{D}_{\mathrm{fd}}(\Gamma_{\mathbb{X}}(\widetilde{Q}_{\mathrm{gr}}, W_{\mathrm{gr}}))$ with $\mathbb{Z} \oplus \mathbb{Z}[\mathbb{X}]$ graded quiver $\widetilde{Q}_{\mathrm{gr}}$ as follows and potential $W_{\mathrm{gr}} = \sum_{i=1}^3 (x_i y_i z_i - x_i z_i y_i)$,
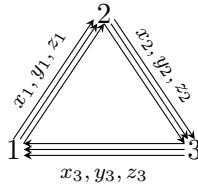
where $\deg x_3, y_3, z_3 = 3 - \mathbb{X}$ and gradings of other arrows are zero. Here $\Gamma_{\mathbb{X}}(\widetilde{Q}_{\mathrm{gr}}, W_{\mathrm{gr}})$ is the Calabi–Yau-$\mathbb{X}$ Ginzburg dg algebra [IQ18a, IQ18b]. Note that there is a canonical fully faithful embedding

$$\mathcal{D}_\infty(\mathbb{P}^2) \to \mathcal{D}_{\mathbb{X}}(\mathbb{P}^2)$$

whose image is an $\mathbb{X}$-baric heart of $\mathcal{D}_{\mathbb{X}}(\mathbb{P}^2)$ in the sense of [IQ18a, Definition 2.17].

Finally, we have the 3-reduction of $\mathcal{D}_{\mathbb{X}}(\mathbb{P}^2)$ (see [IQ18a, Example 3.16])

$$\mathcal{D}_3(\mathbb{P}^2) := \mathcal{D}_{\mathbb{X}}(\mathbb{P}^2) \,/\!/\, [\mathbb{X} - 3] \cong \mathcal{D}_{\mathbb{P}^2}^b(Y) \tag{7.2}$$

which is equivalent to the derived category of coherent sheaves on the local $\mathbb{P}^2$.

Lemma 7.1. *Consider the composition of functors in the inducing process*

$$\Phi : \mathcal{D}_{\infty}(\mathbb{P}^2) \xrightarrow{\sim} \mathcal{D}^b(\mathbf{k}Q/R) \to \mathcal{D}_{\mathrm{fd}}(\Gamma_{\mathbb{X}}(\widetilde{Q}_{\mathrm{gr}}, W_{\mathrm{gr}})) \to$$
$$\mathcal{D}_{\mathrm{fd}}(\Gamma_{\mathbb{X}}(\widetilde{Q}_{\mathrm{gr}}, W_{\mathrm{gr}})) \,/\!/\, [\mathbb{X} - 3] \xrightarrow{\sim} \mathcal{D}^b(\mathrm{mod}-J(\widetilde{Q}, W)) \xrightarrow{\sim} \mathcal{D}_{\mathbb{P}^2}^b(Y).$$

Then $\Phi = i_* : \mathcal{D}_{\infty}(\mathbb{P}^2) \to \mathcal{D}_{\mathbb{P}^2}^b(Y)$.

Proof. Let $E_i = \mathcal{O}_{\mathbb{P}^2}(i)$. Then the first equivalence $\mathcal{D}_{\infty}(\mathbb{P}^2) \cong \mathcal{D}^b(\mathbf{k}Q/R)$ in Φ is given by

$$\mathrm{Hom}^{\bullet}(\bigoplus_{i=0}^{2} E_i, -) : \mathcal{D}_{\infty}(\mathbb{P}^2) \to \mathcal{D}^b(\mathbf{k}Q/R),$$

and the last equivalence $\mathcal{D}_{\mathbb{P}^2}^b(Y) \cong \mathcal{D}^b(\mathrm{mod}-J(\widetilde{Q}, W))$ in Φ is given by

$$\mathrm{Hom}^{\bullet}(\bigoplus_{i=0}^{2} \pi^* E_i, -) : \mathcal{D}_{\mathbb{P}^2}^b(Y) \to \mathcal{D}^b(\mathrm{mod}-J(\widetilde{Q}, W)),$$

where $\pi : Y \to \mathbb{P}^2$ is the projection [Bri05]. The lemma then follows from

$$\mathrm{Hom}^{\bullet}(\pi^*\mathcal{E}, i_*\mathcal{F}) = \mathrm{Hom}^{\bullet}(\mathcal{E}, \pi_* i_*\mathcal{F}) = \mathrm{Hom}^{\bullet}(\mathcal{E}, \mathcal{F}).$$

$\square$

Now we recall the inducing construction of stability conditions from the projective plane to the local projective plane, through the 'q-stability conditions' introduced by Ikeda and Qiu [IQ18a].

Construction 7.2. Let $\sigma_{\infty} = (Z_{\infty}, \mathcal{P}_{\infty})$ be a stability condition in $\mathrm{gldim}^{-1}(2) \subset \overline{\mathrm{Stab}^{\mathrm{Geo}}}(\mathbb{P}^2)$.

- By [IQ18a, Theorem. 2.25], there is an induced q-stability conditions (σ, s) in $\mathrm{QStab}\,\mathcal{D}_{\mathbb{X}}(\mathbb{P}^2)$ with parameter $s = 3$, as constructed in [IQ18a, Cons. 2.18].
- By [IQ18a, Theorem. 2.16], (σ, s) projects to a stability condition σ_3 in the principal (connected) component $\mathrm{Stab}^{\dagger}\,\mathcal{D}_3(\mathbb{P}^2)$.

Denote by

$$\iota_3 : \mathrm{gldim}^{-1}(2) \to \mathrm{Stab}^{\dagger}\,\mathcal{D}_3(\mathbb{P}^2)$$

the map of the above inducing process.

Proposition 7.3. *The inducing map ι_3 is injective. Moreover, it factors through the isomorphism between the spaces of geometric stability conditions on $\mathbb{P}^2$ and local $\mathbb{P}^2$:*

$$\iota_3 : \mathrm{gldim}^{-1}(2) \hookrightarrow \overline{\mathrm{Stab}^{\mathrm{Geo}}}(\mathbb{P}^2) \xrightarrow{\sim} \overline{\mathrm{Stab}^{\mathrm{Geo}}}(\mathcal{D}_{\mathbb{P}^2}^b(Y)) \hookrightarrow \mathrm{Stab}^{\dagger}\,\mathcal{D}_3(\mathbb{P}^2).$$

Proof. Recall that a stability condition σ on $\mathcal{D}_{\mathbb{P}^2}^b(Y)$ is called *geometric* if all skyscraper sheaves $i_*\mathcal{O}_x$ of closed points $x \in \mathbb{P}^2$ are σ-stable of the same phase. By Lemma 7.1, the inducing map ι_3 maps geometric stability conditions on $\mathbb{P}^2$ with global dimension 2 to geometric stability conditions on local $\mathbb{P}^2$.

Let $\sigma = (Z, P) \in \text{gldim}^{-1}(2)$ and $\iota_3(\sigma) = (\widetilde{Z}, \widetilde{P}) \in \text{Stab}^{\text{Geo}}(\mathcal{D}_{\mathbb{P}^2}^b(Y))$. By Lemma 7.1, we have $Z = \widetilde{Z} \circ [i_*]$, where $[i_*] : K_0(\mathcal{D}_\infty(\mathbb{P}^2)) \xrightarrow{\sim} K_0(\mathcal{D}_{\mathbb{P}^2}^b(Y))$. By [BM11, Theorem 2.5] and [Li17, Proposition 1.12], any geometric stability condition on $\mathbb{P}^2$ or local $\mathbb{P}^2$ is uniquely determined by its central charge. Moreover, the open set $U \subset \text{Hom}(K_0(\mathcal{D}_\infty(\mathbb{P}^2)), \mathbb{C})$ consists of central charges of geometric stability conditions on $\mathbb{P}^2$ and the open set $\widetilde{U} \subset \text{Hom}(K_0(\mathcal{D}_{\mathbb{P}^2}^b(Y)), \mathbb{C})$ of central charges of geometric stability conditions on local $\mathbb{P}^2$ coincide via the isomorphism $[i_*]$. This proves the proposition. $\square$

Finally, we remark that most of the stability conditions in $\text{Stab}^\dagger \mathcal{D}_3(\mathbb{P}^2)$ can be obtained by inducing from stability conditions on $\mathcal{D}_\infty(\mathbb{P}^2)$ and autoequivalences, since the translates of $\overline{\text{Stab}^{\text{Geo}}(\mathcal{D}_{\mathbb{P}^2}^b(Y))}$ under the group of autoequivalences cover the whole connected component $\text{Stab}^\dagger \mathcal{D}_3(\mathbb{P}^2)$ [BM11, Theorem 1].

References

[AW19] J. August and M. Wemyss, Stability on contraction algebras implies $K(\pi, 1)$, preprint. arXiv:1907.12756

[BB17] A. Bayer and T. Bridgeland, Derived automorphism groups of K3 surfaces of Picard rank 1, *Duke Math. J.* **166(1)** (2017) 75–124. arXiv:1310.8266

[Bei83] A. A. Beĭlinson, The derived category of coherent sheaves on $\mathbf{P}^n$, *Selecta Math. Soviet.* **3(3)** (1983/84) 233–237.

[BM11] A. Bayer and E. Macrì, The space of stability conditions on the local projective plane, *Duke Math. J.* **160(2)** (2011) 263–322. arXiv:0912.0043

[BMS16] A. Bayer and E. Macrì and P. Stellari, The Space of Stability Conditions on abelian threefolds, and on some Calabi-Yau threefolds, *Invent. Math.* **206(3)** (2016) 869–933. arXiv:1410.1585

[BK90] A. I. Bondal and M. M. Kapranov, Representable functors, Serre functors, and mutations, *Mathematics of the USSR-Izvestiya* **35(3)** (1990) 519–541.

[Bri07] T. Bridgeland, Stability conditions on triangulated categories, *Ann. Math.* **166** (2007) 317–345. arXiv:math/0212237

[Bri08] T. Bridgeland, Stability conditions on K3 surfaces, *Duke. Math. J.* **141** (2008) 241–291. arXiv:math/0307164

[Bri05] T. Bridgeland, T-structures on some local Calabi–Yau varieties, *J. Algebra* **289** (2005) 453–483. arXiv:math/0502050

[DK16] G. Dimitrov and L. Katzarkov, Bridgeland stability conditions on the acyclic triangular quiver, *Adv. Math.* **288** (2016) 825–886. arXiv:1410.0904

[DK19] G. Dimitrov and L. Katzarkov, Bridgeland stability conditions on wild Kronecker quivers, *Adv. Math.* **352** (2019) 27–55. arXiv:1602.09117v1

[GR87] A. L. Gorodentsev and A. N. Rudakov, Exceptional vector bundles on projective spaces, *Duke Math. J.* **54(1)** (1987) 115–130.

[Huy06] D. Huybrechts, *Fourier-Mukai transforms in algebraic geometry* Clarendon Press, 2006.

[IQ18a] A. Ikeda and Y. Qiu, q-Stability conditions on Calabi–Yau-$\mathbb{X}$ categories and twisted periods, preprint. arXiv:1807.00469

[IQ18b] A. Ikeda and Y. Qiu, q-Stability conditions via q-quadratic differntials for Calabi–Yau-$\mathbb{X}$ categories, preprint. arXiv:1812.00010

[IUU10] A. Ishii and K. Ueda and H. Uehara, Stability conditions on A_n-singularities, *J. Differential Geom.* **84(1)** (2010) 87–126. arXiv:math/0609551

[Li17]   C. Li, The space of stability conditions on the projective plane, *Selecta Math. New Ser.* **23** (2017) 2927–2945. arXiv:1611.02087

[LZ19a]   C. Li and X. Zhao, Birational models of moduli spaces of coherent sheaves on the projective plane, *Geom. Topol.* **23** (2019) 347–426. arXiv:1603.05035

[LZ19b]   C. Li and X. Zhao, Smoothness and Poisson structures of Bridgeland moduli spaces on Poisson surfaces, *Math. Z.* **291** (2019) 437–447. arXiv:1612.02729

[Liu18]   W. Liu, Bayer-Macrì decomposition on Bridgeland moduli spaces over surfaces, *Kyoto J. Math.* **58** (2018) 595–621. arXiv:1501.06397

[LLM19]   W. Liu and J. Lo and C. Martinez, Fourier-Mukai transforms and stable sheaves on Weierstrass elliptic surfaces, preprint. arXiv:1910.02477

[Mac07]   E. Macrì, Stability conditions on curves, *Math. Res. Lett.* **14(4)** (2007) 657–672.

[MS17]   E. Macrì and B. Schmidt, Bridgeland stability conditions on the acyclic triangular quiver, *Lect. Notes Unione Mat. Ital.* **21** (2017) 139–211. arXiv:1607.01262

[Oka06]   S. Okada, Stability manifold of $\mathbb{P}^1$, *J. Algebraic Geom.* **15(3)** (2006) 487–505.

[Qiu16]   Y. Qiu, Decorated marked surfaces: Spherical twists versus braid twists, *Math. Ann.* **365** (2016) 595–633. arXiv:1407.0806

[Qiu18]   Y. Qiu, Global dimension function on stability conditions and Gepner equations, preprint. arXiv:1807.00010

[Qiu20]   Y. Qiu, Contractible flow of stability conditions via global dimension function, preprint. arXiv:2008.00282

[QW18]   Y. Qiu and J. Woolf, Contractible stability spaces and faithful braid group actions, *Geom. Topol.* **22** (2018) 3701–3760. arXiv:1407.5986

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