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Chapter 2

Introduction to quantum mechanics

Notes

Notation

For distinct vectors in an orthonormal set, we can write $\langle i|j\rangle = \delta_{ij}$, where δ_{ij} is the Kronecker product and is 1 if i = j and 0 if $i \neq j$.

Matrix - Linear Operator Congruence

For a matrix to a be a linear operator,

$$A\left(\sum_{i} a_{i} |v_{I}\rangle\right) = \sum_{i} a_{i} A |v_{i}\rangle$$

must be true. Note the LHS is the sum of vectors to which A is applied which is certainly equal to the RHS.

Now suppose $A: V \to W$ is a linear operator and that V has basis $|v_i\rangle, \cdots, |v_m\rangle$ and W has basis $|w_i\rangle, \cdots, |w_n\rangle$. Since we know the kth column of a A will be its transformation of $|v_k\rangle$,

$$A|v_j\rangle = \sum_i A_{ij} |w_i\rangle$$

Note this is just saying $A|v_j\rangle$ is equal to the jth column of A, and we can think of $|w_i\rangle$ as the coordinates of the transformed vector. Thus, we can find the matrix representation of any linear operator by finding a matrix A with entries specified by the above equation.

What's so special about the Pauli matrices?

Things to look into:

- they pop up in the Pauli equation
- they form a basis for the vector space of 2 × 2 Hermitian matrices. Hermitian matrices represent observables? Look at Wikipedia page

What does the Completeness Relation say about matrices?

I don't understand why

$$\sum_{ij} \langle w_j | A | v_i \rangle | w_j \rangle \langle v_i |$$

implies A has matrix element $\langle w_j | A | v_i \rangle$ in the ith column and jth row, with respect to input basis $|v_i\rangle$ and output basis $|w_j\rangle$. page 68.

Solutions

Exercise 2.1

$$\begin{bmatrix} 1 \\ -1 \end{bmatrix} + \begin{bmatrix} 1 \\ 2 \end{bmatrix} - \begin{bmatrix} 2 \\ 1 \end{bmatrix} = 0$$

Exercise 2.2

$$A = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

because $A|0\rangle$ has coordinate 0 in $|0\rangle$ and coordinate 1 in $|1\rangle$.

If we keep our input bases the same but reorder our output bases as $|1\rangle$ and $|0\rangle$,

$$A = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

Exercise 2.3

We know

$$A|v_i\rangle = \sum_j A_{ji}|w_j\rangle$$

and

$$B|w_j\rangle = \sum_k B_{kj}|x_k\rangle$$

Now we can write

$$\begin{split} BA|v_i\rangle &= B(A|v_i\rangle) = B\sum_j A_{ji}|w_j\rangle = \sum_j A_{ji}(B|w_j\rangle) \\ &= \sum_j A_{ji}\sum_k B_{kj}|x_k\rangle \\ &= \sum_k \sum_j B_{kj}A_{ji}|x_k\rangle \\ &= \sum_k (BA)_{ki}|x_k\rangle \end{split}$$

We know $\sum_k (BA)_{ki}$ is the matrix representation of operator BA, which the preceding step says is equal to $\sum_k \sum_j B_{kj} A_{ji}$, which is the matrix multiplication BA.

Exercise 2.4

For the same input and output basis, we want some I such that

$$I\left|v_{j}\right\rangle = \sum_{i} I_{ij}\left|v_{i}\right\rangle = \left|v_{j}\right\rangle$$

which means $I_{ij} = 0$ for all $i \neq j$ and 1 otherwise.

For $|y\rangle$, $|z_i\rangle \in \mathbb{C}^n$ and $\lambda_i \in C$,

$$(|y\rangle, \sum_{i} \lambda_{i} |z_{i}\rangle) = |y\rangle^{*} \sum_{i} \lambda_{i} |z_{i}\rangle$$
$$= \sum_{i} \lambda_{i} |y\rangle^{*} |z_{i}\rangle$$
$$= \left(\sum_{i} \lambda_{i}^{*} |z_{i}\rangle^{*} |y\rangle\right)^{*}$$

The second and third equalities demonstrate linearity in the second argument and $(|y\rangle, |z\rangle) = (|z\rangle, |w\rangle)^*$. Finally, if $|w\rangle = (w_1, \dots, w_n)$ where $w_i \in \mathbb{C}^n$, then

$$(|w\rangle, |w\rangle) = \sum_{i} w_{i}^{*} w_{i} = \sum_{i} |w_{i}|^{2}$$

which proves the non-degeneracy and non-negativity condition.

Exercise 2.6

$$\left(\sum_{i} \lambda_{i} |w_{i}\rangle, |v\rangle\right) = \left(|v\rangle^{*}, \sum_{i} \lambda_{i}^{*} |w_{i}\rangle^{*}\right)^{*}$$

$$= \sum_{i} \lambda_{i}^{*} \left(|v\rangle^{*}, |w_{i}\rangle^{*}\right)^{*}$$

$$= \sum_{i} \lambda_{i}^{*} \left(|w_{i}\rangle, |v\rangle\right)$$

Exercise 2.7

$$(|w\rangle, |v\rangle) = \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ -1 \end{bmatrix} = 1 - 1 = 0$$

To normalize, divide each vector by $\sqrt{2}$.

Exercise 2.8

Since at each step, we divide by the norm of the vector being added, the set is orthonormal. The set is a basis because at step i, we add the basis vector $|w_i\rangle$ but subtract out the portion that was already in span($|v_1\rangle, \dots, |v_{i-1}\rangle$), so we still end up spanning the full vector space.

Exercise 2.9

$$\sigma_{x} = |1\rangle \langle 0| + |0\rangle \langle 1|$$

$$\sigma_{y} = i |1\rangle \langle 0| - i |0\rangle \langle 1|$$

$$\sigma_{z} = |0\rangle \langle 0| - |1\rangle \langle 1|$$

Exercise 2.10

$$\begin{split} |v_{j}\rangle \, \langle v_{k}| &= I \, |v_{j}\rangle \, \langle v_{k}| \, I \\ &= \sum_{a} |v_{a}\rangle \, \langle v_{a}|v_{j}\rangle \sum_{b} \, \langle v_{k}|v_{b}\rangle \, \langle v_{b}| \\ &= \sum_{a,b} \delta_{aj} \delta_{kb} \, |v_{a}\rangle \, \langle v_{b}| \end{split}$$

so the element $(|v_j\rangle \langle v_k|)_{ab} = \delta_{aj}\delta_{kb}$.

Each of the Pauli matrices has eigenvalues ± 1 .

For σ_{χ} ,

$$\sigma_{x+} = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$
 and $\sigma_{x-} = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$

For σ_{v} ,

$$\sigma_{y+} = \begin{bmatrix} 1 \\ i \end{bmatrix}$$
 and $\sigma_{y-} = \begin{bmatrix} i \\ 1 \end{bmatrix}$

For σ_z ,

$$\sigma_{z+} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
 and $\sigma_{z-} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$

The diagonalization easily follows.

Exercise 2.12

The characteristic equation is $(1 - \lambda)^2$, so we have eigenvalue 1. Solving $(A - 1I)|v\rangle = 0$ gives us $|v\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$. Our diagonal form should be

$$A = |v\rangle\langle v| = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \neq \begin{bmatrix} 1 & 0 \\ 1 & 1 \end{bmatrix}$$

Exercise 2.13

$$\left(\left|w\right\rangle\left\langle v\right|\right)^{\dagger}=\left\langle v\right|^{\dagger}\left|w\right\rangle^{dagger}=\left|v\right\rangle\left\langle w\right|$$

Exercise 2.14

Since we know $(a + b)^{\dagger} = a^{\dagger} + b^{\dagger}$, so

$$\left(\sum_{i} a_{i} A_{i}\right)^{\dagger} = \sum_{i} \left(a_{i} A_{i}\right)^{\dagger} = \sum_{i} \left(a_{i}^{*} A_{i}^{\dagger}\right)$$

Exercise 2.15

$$(|v\rangle, A|w\rangle) = (A^{\dagger}|v\rangle, |w\rangle) = (|v\rangle, (A^{\dagger})^{\dagger}|w\rangle)$$

since this holds for all $|v\rangle$, $|w\rangle$, $A = (A^{\dagger})^{\dagger}$.

Exercise 2.16

$$P^2 = \sum_i |i\rangle\,\langle i|\sum_j |j\rangle\,\langle j| = \sum_{ij} |i\rangle\,\langle i|j\rangle\,\langle j| = \sum_{ij} \delta_{ij}\,|i\rangle\,\langle j|$$

Intuitively, projecting some $|v\rangle \in P$ wouldn't change $|v\rangle$ at all.

Exercise 2.17

Since A is normal, we can diagonalize it. If the eigenvalues are all positive, then the adjoint of this diagonal matrix of eigenvalues equals itself. If A is Hermitian, then so is its diagonal form. The adjoint of a diagonal matrix consists of the conjugates of its diagonal entries, so if $A = A^{\dagger}$, then the diagonal entries (eigenvalues) must all be positive.

For an eigenvector $|v\rangle$, we have $A|v\rangle = \lambda |v\rangle \rightarrow \langle v|A^{\dagger} = \lambda^* \langle v|$. Multiplying these two gives us $\langle v|A^{\dagger}A|v\rangle = \lambda^* \lambda \langle v|v\rangle$, and because $A^{\dagger}A = I$.

$$||v||^2 = |\lambda|^2 ||v||^2 \to |\lambda| = 1$$

Exercise 2.19

Omitted because it's just mechanical.

Exercise 2.20

$$\begin{split} A_{ij}^{'} &= \langle v_i | A | v_j \rangle \\ &= \langle v_i | U^{\dagger} U A U U^{\dagger} | v_j \rangle \\ &= \sum_{a} \sum_{b} \sum_{c} \sum_{d} \langle v_i | w_a \rangle \langle v_a | v_b \rangle \langle w_b | A | w_c \rangle \langle v_c | v_d \rangle \langle w_d | v_j \rangle \\ &= \sum_{a} \sum_{b} \sum_{c} \sum_{d} \delta_{ab} \delta_{cd} \langle v_i | w_a \rangle A_{bc}^{''} \langle w_d | v_j \rangle \end{split}$$

This tells us a = b and c = d, so

$$A_{ij}^{'} = \sum_{a} \sum_{c} \langle v_i | w_a \rangle A_{ac}^{''} \langle w_c | v_j \rangle$$

Exercise 2.21

We will prove any Hermitian operator M is diagonal with respect to some orthonormal basis V.

We proceed by induction on dimension d on V.

Base case: d = 1

Trivially, M is diagonal.

Inductive hypothesis: Assume d = n - 1

Inductive step: Prove d = n

Let λ be an eigenvalue of M, P be a projection onto the λ eigenspace, and Q be P's orthogonal complement.

We know M=(P+Q)M(P+Q). First, note that $QMP=\lambda QP=0$. Now for some $|v\rangle\in P$, $M|v\rangle=M^{\dagger}|v\rangle=\lambda|v\rangle$ because M is Hermitian, which means $|v\rangle$ is in the eigenspace λ of M^{\dagger} . Now we have $QM^{\dagger}P|v\rangle=QM^{\dagger}|v\rangle=\lambda Q|v\rangle=0$. Taking the adjoint of this gives us PMQ=0. Now we have M=PMP+QMQ.

Since $PMP = \lambda P$, QMQ must be nonzero also for M = PMP + QMQ to hold, so they both have dimension less than n. Finally, $PMP = (PMP)^{\dagger} = PM^{\dagger}P$ and similarly for QMQ. Since they're both Hermitian, our inductive hypothesis proves the theorem.

Exercise 2.22

For a Hermitian operator A, suppose $A|v\rangle = \lambda |v\rangle$ and $A|w\rangle = \mu |w\rangle$.

Since $\langle v|A = \lambda \langle v|$, we can write

$$\langle v|A^2|w\rangle = \mu^2 \langle v|w\rangle = \lambda \mu \langle v|w\rangle$$

where the first equality follows from $A^2|w\rangle = \mu^2|w\rangle$. Since $\lambda \neq \mu$, $\langle v|w\rangle = 0$.

Exercise 2.23

Suppose $|v\rangle$ is an eigenvector of P with eigenvalue λ , $P|v\rangle = \lambda |v\rangle$. Then

$$P|v\rangle = P^2|v\rangle = \lambda^2|v\rangle$$

where the first equality follows from the property $P^2 = P$. Since $\lambda^2 = \lambda$, P's eigenvalues must be 1 or 0.

Let $B = \frac{A+A^{\dagger}}{2}$ and $C = \frac{A-A^{\dagger}}{2i}$. This means

$$\langle v|A|v\rangle = \langle v|B|v\rangle + i\langle v|C|v\rangle$$

Since we cannot have an imaginary term in a positive operator, C = 0, so $A = A^{\dagger}$.

Exercise 2.25

We can write $\langle Av|Av\rangle = \langle A^{\dagger}Av|v\rangle$. But since $\langle Av|Av\rangle$ can also be written as $||Av||^2 \ge 0$, $A^{\dagger}A$ must be positive.

Exercise 2.26

As a tensor product:

$$|\psi\rangle^{\otimes 2} = \frac{1}{2}\Big(|0\rangle + |1\rangle\Big) \otimes \Big(|0\rangle + |1\rangle\Big) = \frac{1}{2}\Big(|00\rangle + |01\rangle + |10\rangle + |11\rangle\Big)$$

As a Kronecker product: Since $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) = \begin{bmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{bmatrix}$,

$$\begin{bmatrix} \frac{1}{\sqrt{2}} |\psi\rangle \\ \frac{1}{\sqrt{2}} |\psi\rangle \end{bmatrix} = \begin{bmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{bmatrix}$$

As a tensor product:

$$|\psi\rangle^{\otimes 3} = |\psi\rangle\otimes|\psi\rangle^{\otimes 2} = \frac{1}{\sqrt{2}}\Big(|0\rangle + |1\rangle\Big) \otimes \frac{1}{2}\Big(|00\rangle + |01\rangle + |10\rangle + |11\rangle\Big) = \frac{1}{2\sqrt{2}}\Big(|000\rangle + |001\rangle + |010\rangle + |011\rangle + |100\rangle + |101\rangle + |111\rangle\Big)$$

As a Kronecker product:

$$\begin{bmatrix} \frac{1}{\sqrt{2}} |\psi\rangle^{\otimes 2} \\ \frac{1}{\sqrt{2}} |\psi\rangle^{\otimes 2} \end{bmatrix} = \begin{bmatrix} \frac{1}{2\sqrt{2}} & \frac{1}{2\sqrt{2}} \end{bmatrix}^T$$

Exercise 2.27

$$X \otimes Z = \begin{bmatrix} 0 \cdot Z & 1 \cdot Z \\ 1 \cdot Z & 0 \cdot Z \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{bmatrix}$$

$$I \otimes X = \begin{bmatrix} 1 \cdot X & 0 \cdot X \\ 0 \cdot X & 1 \cdot X \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

$$X \otimes I = \begin{bmatrix} 0 \cdot I & 1 \cdot I \\ 1 \cdot I & 0 \cdot I \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}$$

Since $I \otimes X \neq X \otimes I$, the tensor product is not commutative.

Writing the Kronecker product,

$$(A \otimes B)^{\dagger} = \begin{bmatrix} A_{11}B & \cdots & A_{1n}B \\ \vdots & \ddots & \vdots \\ A_{n1}B & \cdots & A_{nn}B \end{bmatrix}^{\dagger} = \begin{bmatrix} A_{11}^{*}B^{\dagger} & \cdots & A_{n1}^{*}B^{\dagger} \\ \vdots & \ddots & \vdots \\ A_{1n}^{*}B^{\dagger} & \cdots & A_{nn}^{*}B^{\dagger} \end{bmatrix} = A^{\dagger} \otimes B^{\dagger}$$

Proving the transpose is similar and proving the complex conjugate requires only using B^* instead of B^{\dagger} .

Exercise 2.29

Let U_1 and U_2 be unitary.

$$(U_1 \otimes U_2)^{\dagger}(U_1 \otimes U_2) = (U_1^{\dagger} \otimes U_2^{\dagger})(U_1 \otimes U_2) = U_1^{\dagger}U_1 \otimes U_2^{\dagger}U_2 = I \otimes I = I$$

Exercise 2.30

Let $A = A^{\dagger}$ and $B = B^{\dagger}$.

$$(A \otimes B)^{\dagger} = A^{\dagger} \otimes B^{\dagger} = A \otimes B$$

Exercise 2.31

Let *A* and *B* be positive operators.

$$\left((A \otimes B)(|v\rangle \otimes |w\rangle), (|v\rangle \otimes |w\rangle\right) = \left(A|v\rangle \otimes B|w\rangle, |v\rangle \otimes |w\rangle\right) = \langle v|A^{\dagger}|v\rangle \langle w|B^{\dagger}|w\rangle = \langle v|A|v\rangle \langle w|B|w\rangle$$

since we know positive operators are Hermitian. Since $\langle v|A|v\rangle$ and $\langle w|B|w\rangle$ are both non-negative, their product is also non-negative.

Exercise 2.32

Let P and Q be projectors. Recall that if $P^2 = P$, P is a projector.

$$(P \otimes Q)(P \otimes Q) = P^2 \otimes Q^2 = P \otimes Q$$

Exercise 2.33

We proceed by induction on n.

Base case: n = 1

e case.

We know

$$H^{\otimes 1} = H = \frac{1}{\sqrt{2}} \left[(|0\rangle + |1\rangle) \langle 0| + (|0\rangle - |1\rangle) \langle 1| \right] = \frac{1}{\sqrt{2}} \sum_{x,y \in \{0,1\}^1} (-1)^{x \cdot y} |x\rangle \langle y|$$

Inductive hypothesis: Assume n = k - 1

$$H^{\otimes k-1} = \frac{1}{2^{k-1/2}} \sum_{x,y \in \{0,1\}^{k-1}} (-1)^{x \cdot y} |x\rangle \langle y|$$

Inductive step: Prove n = k

$$H^{\otimes k} = H \otimes H^{\otimes k-1} = \frac{1}{2^{k/2}} \sum_{x_1, y_1 \in \{0,1\}} (-1)^{x_1 \cdot y_1} \left| x_1 \right\rangle \left\langle y_1 \right| \\ \hspace{0.5cm} \otimes \sum_{x_2, y_2 \in \{0,1\}^{k-1}} (-1)^{x_2 \cdot y_2} \left| x_2 \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_1, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2 \right| \\ \hspace{0.5cm} \otimes \left| x_2, y_2 \in \{0,1\}^{k-1} \right\rangle \left\langle y_2$$

Since this tensor product only flips the sign of the $H^{\otimes k-1}$ if $x_1 = y_1 = 1$, it is easy to see that concatenating x_1, x_2 and y_1, y_2 , would yield the dot product required to flip the signs when we want.