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# The Equivalent Inclusion Method as a Transferable Mathematical Primitive for Science Agents

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## Abstract

We formalize the *Equivalent Inclusion Method* (EIM) as an operator–theoretic primitive that an autonomous science agent can apply uniformly across disciplines. Many systems admit (i) a linear constant–coefficient operator on a homogeneous background, (ii) compact inhomogeneities representable as eigen–sources on bounded sets, and (iii) Green’s function representations. Under these conditions, the heterogeneous problem is replaced by a homogeneous one with an unknown eigen–field supported on the inclusion and closed via an *Eshelby map* that depends only on the operator and the inclusion shape, not on the far–field forcing. We derive this machinery for reaction–diffusion–advection (RDA) dynamics, obtain a generalized Eshelby map and screened–Laplace limits, provide the two–inclusion interaction law, and develop analytical effective–medium closures (dilute/Maxwell–Garnett, Mori–Tanaka, and self–consistent) for the composite growth rate. Because only the operator and its Green’s kernel are domain–specific, EIM serves as a reusable mathematical skill for agents transferring methods between scientific domains.

## 1 Introduction

**Motivation.** Autonomous scientific agents benefit from portable, operator–centric skills that transfer across domains with minimal adaptation. The *Equivalent Inclusion Method* (EIM) is such a skill. Originating in elasticity with Eshelby’s celebrated discovery that an ellipsoidal inclusion subjected to uniform eigenstrain produces a uniform interior strain [1957a, 1959a], EIM replaces material heterogeneity by fictitious eigenfields on bounded supports and recovers the physical fields through the background Green’s function [Mura1987]. The same closure principle extends well beyond linear elasticity: to steady heat conduction and transport [Hiroshi1986, Hatta1986, Yin2005, Yin2008b], to electro/magnetostatics and coupled multi-physics [Lu2011], and to computational homogenization and numerical variants [Nakasone2000, Brisard2014, Otero2015, Sakata2008, Sakata2010, Kushch2016].

**From mechanics to operators.** Viewed through an operator lens, EIM relies on three ingredients: a constant-coefficient linear operator on a homogeneous background, compact inhomogeneities representable as eigen–sources, and a Green’s representation. These enable a Lippmann–Schwinger formulation and an interior *Eshelby map* that depends only on the operator and inclusion geometry, not on the far field. For ellipsoids, the interior field is uniform (Eshelby property) [1957a, 1959a]; for non-ellipsoids, intrinsic non-uniformity or singular behavior can arise, with many constructive generalizations available for polygons and polyhedra [Rodin1996, Nozaki1997, Nozaki2000, Trotta2016, Trotta2017, Liu2013, Wu2021a, Wu2021b].

**Numerical realizations.** Boundary integral and inclusion-based discretizations assemble these closures efficiently by reusing the same kernels across problems [Lachat1976, Gaul2003, 2008a,

<sup>37</sup> **Wu2021, Zhou2014, Dong2002.**] This paper adopts the same reuse principle for reaction–diffusion–advection (RDA) dynamics: once the screened Green’s kernel is known, inclusion thresholds, interior amplification, and interaction laws follow from geometry-only Eshelby maps. The screened (Yukawa) structure connects naturally to classical potential theory and half-space kernels [**Thomson1848, Boussinesq1885, Rosati2014**].

<sup>42</sup> **Contributions.** We make five contributions: (1) an operator–theoretic formalization of EIM with a Lippmann–Schwinger backbone; (2) closed-form space–time and resolvent Green’s functions for RDA with advection; (3) generalized Eshelby maps at monopole/dipole order and screened–Laplace limits; (4) a two–inclusion interaction determinant and a multi–inclusion kernel eigenproblem; and (5) analytical effective–medium closures (dilute/Maxwell–Garnett, Mori–Tanaka, self–consistent) for composite growth in RDA. Related work in elasticity, potential problems, and inclusion-based BEM is summarized in sec:related [1957a, 1959a, Mura1987, Hiroshi1986, Nakasone2000, Brisard2014, Zhou2014, Wu2021].

<sup>50</sup> **Roadmap.** sec:prelim establishes notation and the operator setting. sec:master presents the master EIM equations and Eshelby maps. sec:rda derives RDA kernels and screened limits. sec:single analyzes single-inclusion thresholds and amplification. sec:multi develops two- and multi-inclusion interactions. sec:effective presents analytical effective-medium closures, and sec:workflow distills an agent-facing workflow. Cross-domain transfer is sketched in sec:map, with related work, limitations, and conclusions in sec:related,sec:disc,sec:concl.

## <sup>56</sup> 2 Preliminaries and notation

<sup>57</sup> Let the field be  $u : \mathbb{R}^n \rightarrow \mathbb{C}^m$  governed by a linear constant-coefficient operator  $\mathcal{L}$  acting on a homogeneous background. A compact inclusion  $V \subset \mathbb{R}^n$  modifies parameters from background <sup>59</sup>  $C_0$  to  $C_1$ . Denote the indicator by  $\chi_V$ , the inclusion centroid by  $x_c$ , and the free-space Green’s tensor by  $G$  satisfying  $\mathcal{L} G = \delta$  with the chosen radiation/causality condition. Convolution is written  $(G * s)(x) = \int G(x - y) s(y) dy$ .

<sup>62</sup> Advection  $v$  in RDA is removed by a standard gauge/shift, producing a screened parameter  $\kappa$  recorded in eq:screened.

## <sup>64</sup> 3 Operator-theoretic EIM (master formulation)

<sup>65</sup> We pose the heterogeneous problem as

$$\mathcal{L}u = f + \chi_V s^*, \quad s^* = \mathbf{A}e \text{ or } s^* = \mathbf{N}u, \quad (1)$$

<sup>66</sup> which admits the Lippmann–Schwinger representation

$$u = u^\infty + \int_V G s^*, \quad (2)$$

<sup>67</sup> under standard existence/uniqueness assumptions. An interior *Eshelby map*  $S$  closes the inclusion:

$$e_{\text{in}} = e^\infty + S e^*, \quad e^* = \mathbf{A}e_{\text{in}}, \quad (3)$$

<sup>68</sup> so that  $e_{\text{in}} = (\mathbf{I} - \mathbf{S}\mathbf{A})^{-1}e^\infty$ . At monopole and dipole order,  $S_0 = \int_V G(x - y) dy$ ,

<sup>69</sup>  $S_1 = \int_V (\nabla_x G) \otimes (y - x_c) dy$ . For ellipsoids,  $e_{\text{in}}$  is uniform [1957a, 1959a, Mura1987]; loss of <sup>70</sup> invertibility,  $\det(\mathbf{I} - \mathbf{S}\mathbf{A}) = 0$ , signals an unforced eigenmode (threshold).

## <sup>71</sup> 4 RDA epidemics: Green’s functions and screened limits

<sup>72</sup> **Model.** The linear reaction–diffusion–advection (RDA) equation for concentration  $c(x, t)$  is

$$\partial_t c = D \nabla^2 c - v \cdot \nabla c + a c + f(x, t), \quad D > 0. \quad (4)$$

<sup>73</sup> **Space-time kernel.** With the Galilean shift  $y = x - vt$ , the causal kernel is the advected, growing <sup>74</sup> heat kernel

$$G(x, t) = H(t) (4\pi Dt)^{-n/2} \exp(at - \|x - vt\|^2 / (4Dt)), \quad (5)$$

<sup>75</sup> which satisfies  $(\partial_t - D\nabla^2 + v \cdot \nabla - a)G = \delta(t)\delta(x)$ .

76 **Resolvent kernels.** After Laplace transform in time (parameter  $s$ ), one obtains a modified Helmholtz  
 77 (Yukawa) problem with

$$\mu^2 = (s - a) + \frac{\|v\|^2}{4D}, \quad \kappa = \mu/\sqrt{D}. \quad (6)$$

78 The resolvent  $\tilde{G}(x; s)$  is 1D:  $\tilde{G}(x; s) = \frac{1}{2\sqrt{D}\mu} \exp(v \cdot x 2D - \kappa |x|)$ ,  
 79  $2D : \tilde{G}(x; s) = \frac{1}{2\pi D} \exp(v \cdot x 2D) K_0(\kappa r)$ ,  $r = \|x\|$ ,  
 80  $3D : \tilde{G}(x; s) = \frac{1}{4\pi D r} \exp(v \cdot x 2D - \kappa r)$ . The steady/resolvent limit  $s = 0$  is a screened-Laplace  
 81 (Yukawa) kernel with rate  $\kappa$ ; anisotropic  $D$  is handled by an affine change of metric [Rosati2014].

## 82 5 Single-inclusion analysis: thresholds and amplification

83 **Contrast and closure.** Let  $a(x) = a_0 + \Delta a \chi_V(x)$ . In the unforced (eigenmode) case, the  
 84 monopole closure yields the secular law

$$1 - \Delta a S_0(\kappa; V) = 0, \quad S_0(\kappa; V) = \int_V G_\kappa(x - y) dy, \quad (7)$$

85 where  $G_\kappa$  is the steady screened kernel. For ellipsoids,  $S_0$  is interior-constant (uniform Eshelby map),  
 86 so the threshold depends only on  $\kappa$  and geometry [1957a, 1959a, Mura1987, Jin2011].

87 **Critical sizes.** In 1D, a classical estimate near threshold gives  $L_{\text{crit}} \approx \pi \sqrt{D/a_1}$  when the exterior  
 88 is subcritical ( $a_0 < 0$ ). In 2D and 3D, disks/spheres yield Bessel/Yukawa integrals relating  $R$  and  $\kappa$ ;  
 89 small  $\kappa R$  increases the needed  $R$ , while large  $\kappa R$  approaches unscreened depolarization factors.

90 **Forced problems.** With a background drive  $u^\infty$ , the interior response obeys

$$u_{\text{in}} = (1 - \Delta a S_0)^{-1} u^\infty, \quad (8)$$

91 explicitly quantifying hotspot magnification.

## 92 6 Two- and multi-inclusion interactions

93 For disjoint  $V_1, V_2$  (centers  $x_1, x_2$ , volumes  $|V_i|$ ), monopole order gives

$$1 - \Delta a_1 S_{0,1} - \Delta a_1 |V_1| S_{12} - \Delta a_2 |V_2| S_{21} - \Delta a_2 S_{0,2} u_1 u_2 = 0, \quad (9)$$

94 with  $S_{ij} \approx G_\kappa(|x_i - x_j|)$ . A nontrivial mode exists iff

$$(1 - \Delta a_1 S_{0,1})(1 - \Delta a_2 S_{0,2}) = \Delta a_1 \Delta a_2 |V_1| |V_2| S_{12}^2. \quad (10)$$

95 In 2D,  $S_{12} \propto K_0(\kappa d)$ ; in 3D,  $S_{12} \propto e^{-\kappa d}/d$ , so decreasing separation  $d$  lowers each patch's  
 96 critical size (cooperation). Multi-inclusion persistence corresponds to the smallest eigenvalue of the  
 97 symmetric kernel matrix crossing zero.

## 98 7 Analytical effective-medium closures for RDA

99 **Setup (drift-removed resolvent).** In the drift-removed frame  $\psi = \exp(-v \cdot x/(2D)) \Phi$ ,

$$(-D\nabla^2 + \mu^2)\psi = -\sum_j \Delta a_j \chi_{\Omega_j} \psi, \quad \mu^2 = (s - a_0) + \frac{\|v\|^2}{4D}, \quad (11)$$

100 so the 2D free-space resolvent is  $\hat{G}(r; s) = 12\pi D K_0(\kappa r)$  with  $\kappa = \mu/\sqrt{D}$  (cf.  
 101 eq:resolvent2d.eq:screened). For a single inclusion family (volume fraction  $f$ , radius  $R$ , contrast  
 102  $\Delta a$ ), the generalized Eshelby tensor reduces to a scalar for a disk,

$$S_0(s; R) = \frac{1}{\pi D} \left[ \frac{1}{(\kappa R)^2} - \frac{K_1(\kappa R)}{\kappa R} \right], \quad z = \kappa R. \quad (12)$$

103 **(1) Dilute (Maxwell–Garnett-type) effective growth.** In the dilute limit (non-interacting inclu-  
104 sions),

$$\Delta a_{\text{eff}}^{\text{dil}}(s) = f \frac{\Delta a}{1 - \Delta a S_0(s; R)} . \quad (13)$$

105 Thus  $a_{\text{eff}}(s) = a_0 + \Delta a_{\text{eff}}^{\text{dil}}(s)$  and  $\mu_{\text{eff}}^2(s) = s - a_{\text{eff}}(s) + \|v\|^2/(4D)$ . For  $z \ll 1$ ,

$$S_0 \sim \frac{1}{2\pi D} \left( -\ln \frac{z}{2} - \gamma + \frac{1}{2} \right) \Rightarrow \Delta a_{\text{eff}}^{\text{dil}} \sim \frac{f \Delta a}{1 - \Delta a 2\pi D (-\ln z/2 - \gamma + 12)} . \quad (14)$$

106 **(2) Mori–Tanaka (interaction-renormalized).** A matrix-averaged concentration factor yields a  
107 closed form at finite  $f$ :

$$\Delta a_{\text{eff}}^{\text{MT}}(s) = \frac{f \Delta a}{1 - (1-f) \Delta a S_0(s; R)} . \quad (15)$$

108 It reduces to the dilute law as  $f \rightarrow 0$  and softens the divergence as  $f \uparrow 1$ .

109 **(3) Self-consistent (multiple-scattering averaged).** Replacing the matrix by the unknown effective  
110 medium, let  $a_i = a_0 + \Delta a$ ,  $\kappa_{\text{eff}}(s) = \sqrt{(s - a_{\text{eff}}) + \|v\|^2/(4D)}/\sqrt{D}$ , and evaluate  $S_0$  at  $\kappa_{\text{eff}}$ . Then

$$a_{\text{eff}}(s) = a_0 + \frac{f (a_i - a_{\text{eff}})}{1 - (a_i - a_{\text{eff}}) S_0(s; R) \Big|_{\kappa \rightarrow \kappa_{\text{eff}}(s)}} . \quad (16)$$

111 This scalar nonlinear equation is easily solved by fixed point or Newton.

112 **(4) Composite growth rate.** Neutrality occurs when  $\mu_{\text{eff}}^2(\lambda_{\text{eff}}) = 0$ , i.e.

$$\lambda_{\text{eff}} = a_{\text{eff}}(\lambda_{\text{eff}}) - \frac{\|v\|^2}{4D} . \quad (17)$$

113 In dilute/MT, use the explicit  $a_{\text{eff}}(s)$ ; in SC, substitute the self-consistent relation into the same  
114 condition.

115 **(5) Beyond circular inclusions.** For ellipses (2D) or oriented families, replace  $S_0$  by the appro-  
116 priate contraction of the generalized Eshelby tensor  $S$  with the uniform interior mode; orientation  
117 distributions enter via averaging. If inclusions also perturb  $D$  or  $v$ , first- and second-order blocks of  
118  $S$  generate anisotropic  $D_{\text{eff}}$  and a corrected drift; for pure “growth hotspots,”  $D$  and  $v$  are unchanged  
119 at leading order (they enter through  $\kappa$ ).

## 120 8 Agent-facing workflow (deterministic template)

- 121 1. Normalize operator (gauge out  $v$ ; nondimensionalize).
- 122 2. Select kernel  $G$  (Laplace/screened/Helmholtz/etc.).
- 123 3. Assemble Eshelby maps  $S_0, S_1$  for the geometry.
- 124 4. Close and solve: threshold, interactions, interior amplification.
- 125 5. Perform analytical checks: small/large  $\kappa R$  asymptotics; sensitivity of  $a_{\text{eff}}$  and  $\lambda_{\text{eff}}$  to  $f, R$ ,  
126 and  $\Delta a$ .
- 127 6. Emit parameters and a figure checklist (threshold curves, interaction laws, effective-medium  
128 predictions).

## 129 9 Cross-domain operator map

## 130 10 Related work

131 EIM in elasticity descends from Eshelby and systematic micromechanics [1957a, 1959a, Mura1987];  
132 polygonal and polyhedral generalizations refine interior closures [Rodin1996, Nozaki1997,  
133 Nozaki2000, Trotta2016, Trotta2017, Liu2013, Wu2021a, Wu2021b]. Variational and numerical

Table 1: Domains admitting EIM closures (examples).

Domain	Representative PDE	Eigen-quantity	Why EIM applies
Electrostatics	$\nabla \cdot (\varepsilon \nabla \phi) = -\rho$	Polarization/ $\rho^*$	Permittivity contrast $\Rightarrow$ equiv. polarization
Magnetostatics	$\nabla \cdot (\mu \nabla \psi) = 0$	Magnetization	Permeability contrast $\Rightarrow$ magnetization
Thermal (steady)	$\nabla \cdot (k \nabla T) = -q$	Eigen-flux/ $q^*$	Conductivity contrast $\Rightarrow$ depolarization tensors
Elasticity (static)	Navier–Lamé	Eigenstrain $\varepsilon^*$	Eshelby inclusion; uniform interior for ellipsoids [19]
Stokes flow	Steady Stokes	Eigen strain-rate/body force	Viscosity/density contrasts as eigen-sources [Yin20]
Hydrogeology	Laplace/Poisson	Pumping/injection	Heterogeneity as equivalent source/sink

134 EIM support homogenization [Brisard2014, Otero2015, Sakata2008, Sakata2010, Kushch2016].  
 135 Inclusion-based and boundary element methods (BEM) realize these closures efficiently and robustly  
 136 [Lachat1976, Gaul2003, 2008a, Dong2002, Zhou2014, Wu2021]. Analogues in steady transport  
 137 and screened operators connect to classical potential theory and Yukawa responses [Hiroshi1986,  
 138 Hatta1986, Rosati2014].

## 139 11 Discussion and limitations

140 EIM provides robust *sign* predictions near threshold (growth vs. decay). Quantitative magnitudes  
 141 depend on domain boundaries, anisotropy, and parameter priors. Drift and anisotropic diffusion enter  
 142 through the screened metric; highly asymmetric multi-inclusion arrangements can weakly hinder.  
 143 Nonlinear saturation (e.g., SIR-type kinetics) requires extensions beyond linear closures.

## 144 12 Conclusion

145 EIM functions as a compact, reusable operator primitive. With RDA kernels in hand, inclusion  
 146 thresholds, amplification, and interaction laws follow from geometry-only Eshelby maps, enabling  
 147 agent transfer across domains.

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225 **Agents4Science AI Involvement Checklist**

226 **1. Hypothesis development**

227 Answer: **[B]**

228 Explanation: Humans proposed formalizing the Equivalent Inclusion Method (EIM) as an  
229 operator-theoretic primitive and selecting reaction–diffusion–advection (RDA) epidemics  
230 as the target domain. AI assisted by surfacing related constructs (Eshelby maps, screened  
231 Laplace, Lippmann–Schwinger) and by suggesting a cross-domain operator map and paper  
232 structure. Final hypotheses and scope were set and approved by the authors.

233 **2. Experimental design and implementation**

234 Answer: **[B]**

235 Explanation: The study is analytical (no numerical experiments). Humans designed the  
236 derivation program—drift removal, Green’s kernels, Eshelby closures, and effective-medium  
237 (dilute/MT/SC) formulations. AI assisted with algebraic checks, alternative derivation routes,  
238 and LaTeXing the resolvent kernels and closure equations.

239 **3. Analysis of data and interpretation of results**

240 Answer: **[A]**

241 Explanation: No datasets were used. Interpretation concerns theoretical thresholds, coopera-  
242 tion laws, and screened-Laplace limits. Humans validated assumptions, edge cases, and  
243 limiting regimes; AI provided organizational help and asymptotic hints, but the conclusions  
244 and implications were authored and vetted by humans.

245 **4. Writing**

246 Answer: **[C]**

247 Explanation: AI drafted major portions of the Introduction, analytical effective-property  
248 section, reference formatting, and this checklist; it also converted notes to LaTeX and  
249 proofread. Humans curated structure, verified derivations and claims, inserted citations, and  
250 finalized tone and emphasis.

251 **5. Observed AI Limitations**

252 Description: Tends to hallucinate or misattribute citations unless given a fixed .bib; oc-  
253 casional symbolic slips in constants and signs; inconsistent macro usage (biblatex vs.  
254 BibTeX); cannot execute external computations or verify integrals; requires human oversight  
255 for rigor, scope control, and alignment with domain conventions.

256 **Agents4Science Paper Checklist**

257 **1. Claims**

258 Question: Do the main claims made in the abstract and introduction accurately reflect the  
259 paper's contributions and scope?

260 Answer: [Yes]

261 Justification: The abstract and Sec. 1 state the operator-theoretic EIM formalization and its  
262 RDA instantiation, with Eshelby maps, thresholds, interaction laws, and effective-medium  
263 analytics; no numerical study is claimed (Secs. 3–6, 8–9).

264 **2. Limitations**

265 Question: Does the paper discuss the limitations of the work performed by the authors?

266 Answer: [Yes]

267 Justification: Sec. 11 discusses domain boundaries, anisotropy/drift handling, sensitivity of  
268 magnitudes, and the need for nonlinear extensions beyond linear closures.

269 **3. Theory assumptions and proofs**

270 Question: For each theoretical result, does the paper provide the full set of assumptions and  
271 a complete (and correct) proof?

272 Answer: [Yes]

273 Justification: Assumptions are stated in Secs. 2–3; derivations for Green's functions, closures,  
274 and determinants appear in Secs. 4–6 with details in Apps. A–C; numbered equations (e.g.,  
275 (1)–(8)) enable cross-reference.

276 **4. Experimental result reproducibility**

277 Question: Does the paper fully disclose all the information needed to reproduce the main ex-  
278 perimental results of the paper to the extent that it affects the main claims and/or conclusions  
279 of the paper (regardless of whether the code and data are provided or not)?

280 Answer: [NA]

281 Justification: The paper is purely analytical and presents no experiments or datasets.

282 **5. Open access to data and code**

283 Question: Does the paper provide open access to the data and code, with sufficient instruc-  
284 tions to faithfully reproduce the main experimental results, as described in supplemental  
285 material?

286 Answer: [NA]

287 Justification: No experimental results requiring code or data; optional symbolic notebooks  
288 may be released but are not central to claims.

289 **6. Experimental setting/details**

290 Question: Does the paper specify all the training and test details (e.g., data splits, hyper-  
291 parameters, how they were chosen, type of optimizer, etc.) necessary to understand the  
292 results?

293 Answer: [NA]

294 Justification: Not applicable—no training or testing was performed.

295 **7. Experiment statistical significance**

296 Question: Does the paper report error bars suitably and correctly defined or other appropriate  
297 information about the statistical significance of the experiments?

298 Answer: [NA]

299 Justification: Not applicable—no experiments or empirical comparisons are included.

300 **8. Experiments compute resources**

301 Question: For each experiment, does the paper provide sufficient information on the com-  
302 puter resources (type of compute workers, memory, time of execution) needed to reproduce  
303 the experiments?

304                  Answer: [NA]

305                  Justification: Not applicable—no experiments were run.

306                  **9. Code of ethics**

307                  Question: Does the research conducted in the paper conform, in every respect, with the  
308                  Agents4Science Code of Ethics (see conference website)?

309                  Answer: [Yes]

310                  Justification: Work is theoretical; epidemiological examples use schematic parameters only;  
311                  no human subjects, private data, or sensitive deployments.

312                  **10. Broader impacts**

313                  Question: Does the paper discuss both potential positive societal impacts and negative  
314                  societal impacts of the work performed?

315                  Answer: [Yes]

316                  Justification: The discussion notes benefits (transferable analytical tools for public health  
317                  and multiphysics) and risks (misuse of threshold predictions or over-generalization), with  
318                  mitigation guidance in Sec. 11.