

Implications of the Higgs Boson Discovery

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Abstract. In this paper the significance of the discovery of a new boson with the ATLAS detector at the LHC proton-proton at a mass of

$$m_H = 124.3^{+0.6}_{-0.5}(\text{stat.})^{+0.5}_{-0.3}(\text{syst.}) \text{ GeV}$$

is explored. A number of computations that should be rather simple are suggested and imply that this mass may well be a fine-tuned value that is required for galaxy formation. It is argued that it points to the scale where new physics must exist and implies a range of possible new physics symmetries over a variety of scales. This leads to fundamental questions regarding quantities such as action, entropy as well as fundamental concepts of causality and an uncertainty principle regarding causality is suggested. A checklist of computable and verifiable computations is provided.

1 Introduction

The discovery of the Higgs is significant for a number of reasons, not the least of which is the completion of the observation of the last of the fundamental particles in the Standard Model (SM). In this paper the significance of this value is considered and leads to a number of suggestions for how to cope with fine tuning as well as how to choose the fine tuning needed based on cosmological bounds. The problem is illustrated by considering an experiment where collisions occur at a centre-of-mass energy just above, below or equal to the planck mass. Consideration is then given to how one would realize it and what the broad characteristics of the phenomenology would be. The role of black holes is then investigated and it is proposed that there are different classes of black holes but that all black holes contain collision scales that are compressed to that which is of order the Planck Mass, thereby providing an abundant number of accelerators in which to observe the proposed experiment. A new class of black hole is considered to be the instantiation of the Fermi sea of anti matter which contains energy and hence is dark matter. The paper is organized as follows. The significance of the Higgs mass is considered in Section 2 and the fine tuning problem is outlined. The proposed experiment is described in Section 3 where its phenomenology is also described and a connection to black holes is suggested and the melding relativistic quantum field theory with a theory of general relativity and quantum mechanics is considered. In Section consideration is given to how one might observe the phenomenology of Planck Mass scale interactions. In light of cosmological arguments fine tuning is addressed in Section 5. Consideration of the physics within the black hole is given in Section 6 and this leads to a discussion of causality and a possible uncertainty principle that sets the scale in which causality may be violated.

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31 2 Significance of the Higgs Mass

32 There is a deeper significance to the discovery of the Higgs Boson: the value
 33 of the Higgs Boson mass itself. The importance of this lies in the fact that
 34 the SM is valid to scales of order the Planck mass but as one approaches such
 35 scales, the vacuum becomes metastable [?]; however, for the SM to be a valid
 36 description of phenomena at this scale it is necessary to fine tune the Higgs
 37 mass to a level of one part in 10^{52} [?]. This feature of the SM is referred to as
 38 the naturalness problem [?] and has motivated the formation of theories beyond
 39 the SM to try to eliminate this fine tuning. One of the most popular theories
 40 to reduce this problem is Supersymmetry [?]. The theory is usually cast in
 41 terms of the Minimal Supersymmetric Model (MSSM) [?] in order to allow for
 42 computations to be performed with a limited set of parameters as inspired by
 43 assumptions of a Grand Unification Theory that includes gravity and breaks
 44 down to SUSY at the electroweak scale, of order 1 TeV. SUSY is however a
 45 broken symmetry and as such there remains a fine tuning of the Higgs fields in
 46 that theory, but to a much smaller degree than in the SM, being of order 10^{-2}
 47 to 10^{-3} [?]. Furthermore, the boson that was discovered can be the CP-even
 48 scalar in SUSY since it is SM-like in all properties except for 20% differences in
 49 cross section. Currently these differences can be easily accommodated within the
 50 present experimental statistics [?]. However it is true that the loop corrections
 51 in SUSY that correspond to the tree level limit that the lightest Higgs must
 52 have a mass less than that of the Z become quite unstable when accommodating
 53 a Higgs mass of order 125 GeV.

54 3 Proposed Experiment

55 In either the case of the SM or the MSSM it is interesting to consider an exper-
 56 iment that explores the SM processes that one typically measures at a collider
 57 such as cross sections and differential measurements of jets, boson production,
 58 Higgs production, heavy quark production, but to consider these at scales ap-
 59 proaching the Planck Mass. Specifically it is interesting to consider what is
 60 expected from an e^+e^- collider with $\sqrt{s} = 10^{16}$ GeV? Furthermore one could
 61 sweep a range of energies from $\sqrt{s}_{SM} \equiv 10^{15}$ to $\sqrt{s}_{GRQ} \equiv 10^{17}$ GeV. The en-
 62 ergy scan includes values that should be described by the SM to values where
 63 the Planck Mass is exceeded and the theory must include a relativistic quantum
 64 field theory that incorporates General Relativity (GR). This is indicated here
 65 by the subscript GRQ. Candidate models currently include string theory [?].

66 While considering this experiment, one has to examine how the vacuum
 67 predicted by the SM is changing with the scale. For values below \sqrt{s}_{SM} and
 68 around \sqrt{s}_{SM} the vacuum has a minimum and states computed as excitations
 69 above this minimum in the normal Relativistic Quantum Field Theory (RQFT)
 70 will lead to the effects that are normally understood within the context of

71 computations in perturbation theory although accuracy may demand rather
 72 higher order computations in order to match observation. As one approaches
 73 the point where the vacuum has no minimum, then there is no broken symmetry
 74 and all the bosons will be massless, like photons. The unbroken SM would then
 75 be a basis for these phenomenological computations. Once there is no minimum
 76 for the vacuum for values of $\sqrt{s} = \sqrt{s}_{GRQ}$ the framework of RQFT would no
 77 longer be valid and a complete theory based on the computation of QM with
 78 GR would be required. As one does this energy sweep, it would also be true that
 79 the tunneling probability from the region where the vacuum has a minimum
 80 to one outside that minimum would increase and the state may well be one in
 81 which the universe tunnels and the states are free with no minimum.

82 **4 Cosmic accelerator**

83 It is entirely possible to consider the computations suggested for the collider
 84 experiment proposed in Section /refsec:collider and to confront the possibility
 85 of tunneling and how that might manifest itself. If the energy input to a collision
 86 were to be dissipated by tunneling, then that would appear as black hole
 87 formation. While it is impractical to perform the experiment it is possible that
 88 collisions of these energies could be observed in particle astrophysics. There
 89 have been many examples where one could find phenomena that seem to be
 90 odd and rare. Gravitational lensing was considered an interesting thought experiment,
 91 unlikely to be found because it required a very specific arrangement
 92 of galaxies and observers; however, this has been observed and indeed has been
 93 used in measurements of MACHOS. It is commonplace. Also, supernovae were
 94 thought to be hard to find until systematic scanning of the sky using large data
 95 collection methods and CCD cameras has led to the ability to find supernovae
 96 on demand and indeed to measure the accelerating expansion of the universe.
 97 Therefore in order to observe the collisions that are of sufficient energy that
 98 lead to tunneling one needs to look for this.

99 Black holes appear to come in two mass ranges. One is that set by the death
 100 of stars having masses between eight and fifty solar masses leading to black holes
 101 in this mass range. The other is much larger, of order billions of solar masses
 102 and these are found at the centre of galaxies. There is a black hole observed in
 103 the globular cluster Omega Centauri that has a mass of 40,000 solar masses and
 104 could be of particular interest because of its unique intermediate mass value.
 105 There are also various generations of stars in Omega Centauri. It is thought
 106 that supermassive black holes that are at the centres come from the merging of
 107 smaller black holes and that these intermediate sizes are important to show the
 108 process of merging in place. However given the new information on the mass
 109 of the Higgs, and considering the possibility of tunneling when collisions at the
 110 scale of the Planck mass occur, it may be that the larger black holes are in fact

111 due to these tunneling phenomena rather than the fusion of black holes. If one
 112 were to have fusion of black holes the spectrum of masses should be continuous
 113 rather than discrete. If collisions have occurred that cause a tunneling, then it
 114 may be that the process results in some characteristically larger size for black
 115 hole formation and indeed that as the tunneling takes place, it may be possible
 116 to observe the tunneling process building up.

117 This would require that there are simply cases of acceleration in the universe
 118 that can reach centre-of-mass energies large enough for tunneling to occur.
 119 Alternatively the supermassive black holes may have been formed when the
 120 universe was young enough to have a high probability for collisions at high
 121 energy to occur. This would have had to have happened just at the end of the
 122 inflationary period and implies another form of fine tuning. It would have been
 123 necessary that tunneling probability not be too large such that every collision
 124 would lead to black holes so that the universe would not get started; however,
 125 it would have to be large enough to allow for sufficient black holes to be formed
 126 such that they produce the correct numbers of galaxies. In addition the total
 127 mass of the universe should be large enough that these black holes do not
 128 condense and prevent formation of the universe at all.

129 5 Fine Tuning Embraced

130 This implies a set of parameters for the formation of the universe: the proba-
 131 bility of tunneling per high energy interaction, P_{tunnel} , the spectrum of particle
 132 center of mass collisions as a function of the size and density of the universe
 133 $S(t, \rho)$ and the size of the tunneling $D(\sqrt{s})$. The tuning of these parameters is
 134 related to the Higgs potential and hence to the fine tuning of the Higgs mass.
 135 As was stated in Section 2, the amount of fine tuning that is required for the
 136 Higgs may well indicate if the particle needs to be SM or MSSM.

137 If one considers the experiment in Section 3 then the theory to describe the
 138 full scan of masses is that which accomodates Quantum Mechanics and General
 139 Relativity, a theory of GRQ. The energy scan is the same as a scan in size, and
 140 just as in atomic physics there must be a correspondence priciple where the size
 141 or energy regime in which both theories are valid give a consisent answer. In
 142 this case the incomplete theory is not Newtonian mechanics but RQFT based
 143 on Special Relativity and the more complete theory is not quantum mechances
 144 but GRQ. One would imagine that the dynamics of the black hole formation
 145 would be described by RQFT and set the parameters for the rate at which the
 146 tunneling occurs and the dynamics that limit the tunneling such a to have some
 147 turn-off mechanism. In addition it is possible that a structure of a RQFT GUT
 148 would be set by the GRQ, with symmetries such as SO(10) or SU(5) predicted
 149 by the theory. From the point of view of RQFT, it would appear as an ad-hoc
 150 fact that some symmetry and indeed the breaking of the symmetry is described

151 by some set of parameters that have to be set by hand. In the breaking of the
 152 symmetry and any further symmetries a set of Higgs bosons would be generated
 153 and of course a set of new particles. Therefore it remains well worth pursuit of
 154 the search for these high mass particles throughout the scales up to the Planck
 155 mass to try to get a clue on the form of the manifestation of the theory from
 156 which they follow.

157 In addition to this it is possible to formulate various ensembles of broken
 158 symmetries and hence fine tunings that allow a range of fine tunings to be
 159 generated. Thus one is not limited to only a fine tuning of the Higgs mass to
 160 the large degree required by the Standard Model nor to the small degree of
 161 the MSSM, but could choose tunings in a spectrum of values. It is therefore
 162 proposed that a computation of the fine tunings for various scenarios of sym-
 163 metries being broken be examined to see what possibilities exist and how that
 164 then can be compared to the probabilities for the formation of black holes for
 165 galaxies.

166 In consider the formation of the black holes, it is also interesting to under-
 167 stand if these galaxies and their black holes can eventually serve as portals to
 168 the next Big Bang or if indeed each of the provides a transport of matter and
 169 energy to the formation of another Big Bang in smaller universes. The expansion
 170 of the universe implies that these galaxies will not have an opportunity to
 171 coalesce and be the source of the next Big Bang. This would have happened
 172 if the universe were contracting and would provide a means by which matter
 173 and energy could be transformed in the Big Crunch.

174 **6 Faster than Light Neutrinos and Physics Near the Planck Mass**

175 While the faster than light (FTL) neutrino observation turned out to be the
 176 result of an instrumental error, it was useful to think about what would one
 177 do to accomodate such an observation. A fundamental assumption of RQFT
 178 is causality and if one considers FTL neutrinos one should ask if there is a
 179 consistent field theory where causality is formally abandoned. It is possible to
 180 formulate such a theory [?] but the result is considered uninteresting since the
 181 particles of that theory cannot interact with slower than light particles. This
 182 however is considered in the context of RQFT rather than GRQ. Therefore it
 183 could be possible that interactions could arise from gravity. Since causality is
 184 a statement regarding time, it is therefore necessary to abandon the need to
 185 consider time in FTL theory. Given that time slows down as one approaches a
 186 black hole and stops at a black hole, it is possible to consider that the dynamics
 187 within a black hole allow for time to be undefined and indeed correspond to
 188 scales smaller than that suggested by the Planck Mass. It is possible that
 189 the black holes from tunneling are then different from those created by stellar
 190 collapse or that the scale of the black hole is not characterized by the event

191 horizon but necessarily by the Planck Scale. This implies then that the physics
192 at the Planck scale is much more common than one would otherwise suspect
193 and that indeed collisions at that scale can be created in the stellar collapse.
194 This in turn implies that the formation of super massive black holes comes
195 from stellar collapse seeds that then provide the opportunity for a tunnelling
196 black hole.

197 One does not have to abandon causality entirely but it may be that causality
198 can be formulated in an uncertainty principle that requires abandonment for
199 small distances only and that a new constant analagous to Planck's constant
200 can characterize the abandonment of causality.

201 It has been suggested that as one looks at scales near the Planck Mass [?]
202 that space-time becomes foamy. It is natural to consider that at those scales
203 causality can be abandoned. If one were to try to formulate a path integral
204 where the path is foamy, then it is expected there would be issues with causality.
205 Therefore it is sensible to revisit the principle of least action and its definition
206 once again, just as was first done by LaGrange, later by Feynman who accomo-
207 dated the Heisenburg uncertainty principle. The issue with causality becomes
208 an issue of getting lost as one tries to integrate over the path.

209 7 Causality, Vacuum, Dark Matter and Dark Energy

210 Having established that action needs to be redefined has a number of impli-
211 cations. First, since the definition of action is the difference in kinetic and
212 potential energies, it is reasonable that once one considers action in a GRQ
213 theory, there could be terms that correspond in GR to the energy contained in
214 the quantum foam which in turn could be the definition of the vacuum energy
215 above which excitations are computed in RQFT. The definition of the vacuum
216 in special relativity has been a difficulty. It is possible that dark matter and
217 dark energy are manifestations of the energy contained in this foam and are not
218 related in any way to the MSSM. The interaction of particles of dark matter
219 would indeed be entirely through gravitation and not require a RQFT formu-
220 lation at all. By providing an energy density the infinities usually sidestepped
221 in RQFT would become finite since the size and density of the foam is set by
222 the Planck Scale - one need only fill up the universe with this background of
223 foam. The nature of the foam would be that of a third class of black holes that
224 in fact are the source of the negative energy sea of RQFT.

225 8 Summary

226 In this paper it has been proposed to consider the investigations of particle
227 astrophysics in light of the understanding of the implications of the Higgs mass.
228 An investigation of the phenomenology associated with collisions just below and

229 above the centre of mass energy at the Planck Mass scale is suggested. It is also
230 suggested to compute the possible fine tunings required for various scenarios of
231 symmetry breaking. With these numbers in hand, then it is worth considering
232 how these phenomena would be observed in cosmic accelerators and that the
233 likelihood of such observations may be quite large. The physics associated with
234 this scan may also require a new interpretation of causality and the vacuum
235 to be studied may be one that describes the dark energy and dark matter of
236 the universe as well as being the vacuum upon which excitations of particles
237 from RQFT is currently computed. Indeed the vacuum is considered to be
238 a third class of black holes that can provide a source of the Fermi sea of anti
239 matter through Hawking radiation and that matter arises from the evaporation
240 of these black holes.

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