

# Phys 772: Week 5 Thursday

\*  $SU(2)$  weak isospin

\*  $SU(2)_L \times U(1)_Y : \langle \varphi \rangle = v = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$

$\downarrow$   
 $L^i = \frac{\sigma^i}{2}$

$\rightarrow L = Y$

$L^i v = \frac{\sigma^i}{2} \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \neq 0$

$\rightarrow L v = Y \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \neq 0$

but  $(\underbrace{L^3 + Y})v = 0$  symmetry is not broken

$Q$

$\rightarrow QED \quad U(1)_Q$

$\hookrightarrow SU(3)_C \times SU(2)_L \times U(1)_Y$  in  $\{m_3, m_2, m_1\}$  rep

$q_{mL}$  in  $\{3, 2, \frac{1}{6}\} \rightarrow Q = \pm \frac{1}{2} + \frac{1}{6} = \begin{pmatrix} 2/3 \\ -1/3 \end{pmatrix}$

$l_{mL}$  in  $\{1, 2, -\frac{1}{2}\} \rightarrow Q = \pm \frac{1}{2} - \frac{1}{2} = \begin{pmatrix} 0 \\ -1 \end{pmatrix}$

$u_{mR}$  in  $\{3, 1, \frac{2}{3}\} \rightarrow Q = \frac{2}{3}$

$d_{mR}$  in  $\{3, 1, -\frac{1}{3}\} \rightarrow Q = -\frac{1}{3}$

$e_{mR}$  in  $\{1, 1, -1\} \rightarrow Q = -1$

$\nu_{mR}$  in  $\{1, 1, 0\} \rightarrow Q = 0$

$m = 1, \dots, 3$  families

$Q$  is equal for L and R by design

\*  $SU(2)$  Higgs doublet and SSB

$$(D^\mu \psi)^\dagger (D_\mu \psi) \text{ becomes}$$

$$= \frac{1}{2} \left( \frac{1}{2} v g \right)^2 (W_\mu'^2 + W_\mu^2) + \frac{1}{8} v^2 (g' B_\mu - g W_\mu^3)^2$$

Define:  $w_{\mu}^{\pm} = \frac{1}{\sqrt{2}} (w_{\mu}^1 \mp i w_{\mu}^2)$

Define  $\begin{cases} Z_\mu = \frac{-g' B_\mu + g W_\mu^3}{\sqrt{g'^2 + g^2}} = -\sin \theta_W \cdot B_\mu \\ A_\mu = \cos \theta_W \cdot B_\mu + \sin \theta_W \cdot W_\mu^3 \end{cases}$

$$\rightarrow (D^\mu \varphi)^\dagger (D_\mu \varphi) = M_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} M_Z^2 Z_\mu Z^\mu$$

$$(+ O. A_\mu A^\mu)$$

$$M_W = g \frac{v}{2}, \quad M_Z = \sqrt{g^2 + g'^2} \frac{v}{2}, \quad M_A = 0 \text{ (photon)}$$

$$\sin \theta_w = \frac{g'}{\sqrt{g^2 + g'^2}}, \quad \cos \theta_w = \frac{g}{\sqrt{g^2 + g'^2}}$$

$$\rightarrow \tan \theta_w = \frac{g'}{g}$$

$$\sin^2 \theta_w = 1 - \frac{M_W^2}{M_Z^2} \approx 0,2397$$

→  $\theta_w$  = electronak mixing angle.

Since we know that  $A_\mu$  interacts with strength  $e$  given by electron charge

$$\rightarrow A_\mu = \cos \theta_w \underset{\substack{\uparrow \\ g'}}{B_\mu} + \sin \theta_w \underset{\substack{\uparrow \\ g}}{W_\mu^3}$$

$$\rightarrow e = g \sin \theta_w = g' \cos \theta_w$$

$$\begin{aligned} \mathcal{L}_{\text{Higgs}} &= (D^\mu \varphi)^\dagger (D_\mu \varphi) - \mu^2 \varphi^\dagger \varphi - \lambda (\varphi^\dagger \varphi)^2 \\ &= M_W^2 W_\mu^+ W^\mu - \left(1 + \frac{H}{v}\right)^2 + \frac{1}{2} M_Z^2 Z_\mu Z^\mu \left(1 + \frac{H}{v}\right)^2 \\ &\quad + \frac{1}{2} (\partial^\mu H)^2 + \mu^2 H^2 - \lambda v H^3 - \frac{\lambda}{4} H^4 + \frac{\mu^4}{4\lambda} \end{aligned}$$

## \* Standard Model Lagrangian

$$\mathcal{L}_{SM} = \mathcal{L}_{\text{fermions}} + \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}}$$

$$\begin{aligned} \mathcal{L}_{\text{fermions}} = \sum_{m=1}^3 & \left( \bar{q}_{mL} i \not{D} q_{mL} + \bar{l}_{mL} i \not{D} l_{mL} \right. \\ & + \bar{u}_{mR} i \not{D} u_{mR} + \bar{d}_{mR} i \not{D} d_{mR} \\ & \left. + \bar{l}_{mR} i \not{D} l_{mR} + \bar{\nu}_{mR} i \not{D} \nu_{mR} \right) \end{aligned}$$

$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} W_{\mu\nu}^i W_{\mu\nu}^i - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$

$i = 1, 2, 3$   
 $a = 1, \dots, 8$

$$\text{with } \begin{cases} B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu \\ W_{\mu\nu}^i = \partial_\mu W_\nu^i - \partial_\nu W_\mu^i - g \varepsilon_{ijk} W_\mu^j W_\nu^k \\ G_{\mu\nu}^a = \partial_\mu G_\nu^a - \partial_\nu G_\mu^a - g_s f_{abc} G_\mu^b G_\nu^c \end{cases}$$

$$\text{with } D_\mu = \partial_\mu + i g' \frac{1}{2} \gamma B_\mu + i g \frac{\vec{\sigma}}{2} \cdot \vec{W}_\mu + i g_s \frac{\vec{T}}{2} \cdot \vec{G}_\mu$$

$$\rightarrow \bar{q}_{mL} i \not{D} q_{mL} =$$

$$i \left( \bar{u}_{mL}^{\alpha} \quad \bar{d}_{mL}^{\alpha} \right) \left( \partial_{\mu} + i g_1 \gamma_{\mu} B + i g \frac{\sigma_i}{2} W_{\mu}^i + i g_s \frac{\lambda_{\alpha\beta}}{2} G_{\mu}^a \right) \begin{pmatrix} u_{mL}^{\beta} \\ d_{mL}^{\beta} \end{pmatrix}$$

$\alpha, \beta = 1, 2, 3$        $a = 1, \dots, 8$        $i = 1, \dots, 3$

$SU(3)_c$  commutes with  $SU(2) \times U(1)$  subgroups

$$\mathcal{L}_{\text{Higgs}} = (D_{\mu} \varphi)^{\dagger} (D_{\mu} \varphi) - \mu^2 \varphi^{\dagger} \varphi - \lambda (\varphi^{\dagger} \varphi)^2$$

with  $\varphi = \begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix}$  as  $\{1, 2, \frac{1}{2}\} \rightarrow Q = \begin{pmatrix} +1 \\ 0 \end{pmatrix}$

with  $\mu^2 < 0, \lambda > 0, v = -\frac{\mu^2}{\lambda}, \langle \varphi \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix}$

$$\mathcal{L}_{\text{Yukawa}} = - \sum_{m,n} \left( g_{um} \bar{q}_{mL} \varphi u_{mR} + g_{dm} \bar{q}_{mL} \varphi d_{mR} \right.$$

$$\left. + g_{\ell m} \bar{l}_{mL} \varphi e_{mR} + g_{\nu m} \bar{l}_{mL} \varphi \nu_{mR} \right) + \text{h.c.}$$

with  $g_{\ell m} = \sqrt{2} \frac{m_{\ell m}}{v}$  coupling of  $\ell_m$  to Higgs

$$\bar{\ell}_{mL} \varphi = \bar{\nu}_L \varphi^+ + \bar{e}_L \varphi^0 \text{ is } SU(2) \text{ invariant}$$

$$\varphi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + \eta \end{pmatrix} \rightarrow \text{also } \bar{\ell}_{mL} \eta \ell_{mR} \text{ vertex}$$

The last term  $g_{YM} \bar{L}_m \gamma^\mu \nu_{mR}$  is absent if we do not include  $\nu_{mR}$  in the Standard Model  $\rightarrow$  then  $m_\nu$  must be zero

If  $\varphi$  is a weak isospin doublet under  $SU(2)$  then

$$\tilde{\varphi} = i\sigma_2 \varphi^\dagger = \begin{pmatrix} \varphi^{0+} \\ -\varphi^- \end{pmatrix}$$

is also a  $SU(2)$  doublet with  $\{1, 1^*, -\frac{1}{2}\}$

$$\text{with } \tilde{\varphi} = \frac{1}{\sqrt{2}} \begin{pmatrix} \nu + \eta \\ 0 \end{pmatrix}$$

$\rightarrow \tilde{\varphi}$  operates on upper weak isospin components,  $u_{mR}$   
 $\downarrow$   
 slight variation to  $\mathcal{L}_{\text{Yukawa}}$   
 $\nu_{mR}$

Finally:  $g_{\ell m}, g_{u m}, g_{d m}, g_{\nu m}$  could be a general

$3 \times 3$  matrix  $\Gamma_\ell, \Gamma_u, \Gamma_d, \Gamma_\nu$  if the weak eigenstates and mass eigenstates are different

$$\Rightarrow \mathcal{L}_{\text{Yukawa}} = - \sum_{m,n}^3 \left( \bar{q}_m \Gamma_{mn}^u \tilde{\varphi} u_{nR} + \text{etc} \right) + \text{h.c.}$$