

MEASUREMENT OF ELECTROWEAK PRODUCTION OF SAME-SIGN W BOSON PAIRS WITH ATLAS

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ABSTRACT

62 MEASUREMENT OF ELECTROWEAK PRODUCTION OF SAME-SIGN W BOSON PAIRS
 63 WITH ATLAS

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66 This thesis presents two studies of electroweak same-sign $W^\pm W^\pm jj$ scattering with the ATLAS
 67 experiment. The first is a measurement of the fiducial cross section at $\sqrt{s} = 13$ TeV using 36.1 fb^{-1}
 68 of data recorded in 2015 and 2016. The electroweak production is observed with a signal significance
 69 of 6.9σ , and the fiducial cross section is measured to be $\sigma_{\text{meas}}^{\text{fid}} = 2.91^{+0.51}_{-0.47}(\text{stat})^{+0.28}_{-0.29}(\text{sys}) \text{ fb}$. The
 70 second is a study on the future prospects for the $W^\pm W^\pm jj$ process at the planned High-Luminosity
 71 LHC, with a projected $\sqrt{s} = 14$ TeV and 3000 fb^{-1} of data. The expected electroweak production
 72 cross section is determined with a total uncertainty of 6%, and the purely longitudinal scattering
 73 component is extracted with an expected significance of 1.8σ . Additionally, some time is taken to
 74 detail the alignment of the ATLAS Inner Detector subsystems, as good alignment performance is
 75 essential for making high-quality physics measurements.

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Preface

570 This thesis presents the major highlights of my work with the ATLAS experiment as a graduate
571 student at the University of Pennsylvania from Fall of 2013 until early Spring of 2019.

572 The first step of working on the experiment is to complete a *qualification task* in order to be
573 included on the author list of ATLAS publications. These tasks are an opportunity to contribute
574 to the experiment as a whole, such as maintaining detector hardware or monitoring physics perfor-
575 mance. For my qualification task, I worked with the Inner Detector Alignment group which works
576 to make sure we have accurate knowledge of the locations of each and every sensor in the tracking
577 detector. My qualification task involved investigating a possible momentum bias in the Monte Carlo
578 (MC) simulated data. The MC is supposed to be reconstructed with a perfect detector geometry
579 which should in principle be free of any momentum biases. Ultimately I determined that the size of
580 the biases were small enough to be negligible compared to what is seen in the real data, and that
581 they could be corrected for if necessary.

582 My work with the alignment group would continue for the duration of my time here at Penn. In
583 early 2015, at the start of the LHC’s second data-taking run (Run 2), I assisted with the validation
584 of the first set of alignment constants using $\sqrt{s} = 13$ TeV proton-proton collision data. At this point
585 I took over the responsibility of alignment of the TRT subdetector. The TRT was aligned to high
586 accuracy in Run 1, and over the course of my time working on alignment, the TRT never required
587 a straw-by-straw alignment; however it did require a module-level alignment at the end of 2015.
588 My final responsibility in the alignment group was monitoring momentum biases using the energy-
589 momentum ratio (E/p) of electrons. For the large data reprocessing, the E/p method served as a
590 cross check to a similar method using Z boson events for monitoring and aligning out momentum
591 biases in the detector. The results from both methods were also used in the uncertainties for the
592 tracking measurements.

593 On the analysis side, I had previous experience in Standard Model (SM) electroweak physics
594 from my time as an undergraduate at Duke University, and it remained a point of interest for
595 me in graduate school. As such, I was happy to work with fellow Penn students on the cross
596 section measurement of SM WZ diboson production with the early $\sqrt{s} = 13$ TeV ATLAS data.
597 My contribution to the analysis was primarily on the software side, as I maintained and updated
598 the analysis framework. While the WZ measurement is not covered by this thesis, it provided me
599 with invaluable analysis experience in electroweak physics, as well as a detailed understanding of a
600 major background to many diboson processes. The results for this analysis can be found published
601 in Physics Letters B in 2016 [1].

602 The final two analyses I worked on involved the scattering of same-sign W bosons, and they make
603 up the majority of this thesis. The first analysis is a measurement of the $W^\pm W^\pm jj$ cross section
604 at $\sqrt{s} = 13$ TeV. This measurement along with that of the CMS collaboration represent the first
605 observation of the $W^\pm W^\pm jj$ scattering process. My primary contribution to the analysis is in the
606 estimation of the fake lepton background, where we implemented a brand new version of the fake-
607 factor method using particle isolation variables. I also did a preliminary study of the interference
608 between electroweak and strong production of $W^\pm W^\pm jj$ events, assisted in the production of data
609 samples for use with the analysis framework, and used my familiarity with the WZ process to
610 optimize the rejection of this background. Ultimately the results of this optimization were not
611 included in the final result; however, it is still covered in the thesis in the hopes that it will be useful
612 for similar analyses in the future. The formal publication for this measurement will likely be coming
613 out within the next few months.

614 The second $W^\pm W^\pm jj$ analysis is a study on the prospects for a measurement of the process at the
615 upgraded High-Luminosity LHC, scheduled to begin operation in 2026. Here my main contribution
616 was an optimization of the event selection using a Random Grid Search algorithm. Through the
617 optimization we expect to take advantage of the higher center of mass energy and greater volume of
618 data and tighten certain selection cuts to increase the strength of the $W^\pm W^\pm jj$ signal. In addition, I
619 once again maintained and updated the analysis framework and produced the group's data samples,
620 and I also developed a truth-based particle isolation criteria in order to reduce contributions from
621 backgrounds involving the top quark. The results of this prospects study will be published as a part
622 of the annual Yellow Report for the High-Luminosity LHC.

623

Will K. DiClemente
Philadelphia, February 2019

CHAPTER 1

Introduction

626 The Large Hadron Collider (LHC) at CERN is the most powerful collider experiment in the world.
627 At the time of its construction, the largest unanswered question in the Standard Model (SM) was
628 the mechanism behind electroweak symmetry breaking (EWSB). As a result, one of the primary
629 goals of the experiment is to learn as much as possible about this mechanism. Thus far, the LHC
630 has succeeded in discovering a particle consistent with the long-awaited Higgs boson. In addition,
631 measurements of many SM processes have been performed for the first time or at better precision
632 than before thanks to the high collision energy and large volume of data collected by the LHC.

633 Processes involving the scattering of two massive electroweak (EWK) gauge bosons are of par-
634 ticular interest at the LHC for two main reasons. Firstly, they allow for tests of the self-interactions
635 predicted by the EWK gauge theory through triple and quartic gauge couplings. While the triple
636 couplings have been studied by previous experiments as well as at the LHC, the quartic couplings
637 of the massive gauge bosons have not been accessible previously. Thus, processes involving these
638 couplings can be measured and compared to the SM predictions for the first time. Secondly, the scat-
639 tering of two massive gauge bosons is sensitive to the underlying EWSB mechanism. The W^\pm and
640 Z bosons are given non-zero masses—and consequently a longitudinal polarization mode—through
641 the Higgs mechanism, and thus their interactions serve as a direct probe of the symmetry breaking
642 sector.

643 This thesis presents two separate analyses dealing with the scattering of two same-sign W^\pm
644 bosons with the LHC’s ATLAS experiment. The $W^\pm W^\pm jj$ process is a particularly interesting one:
645 it has access to the $WWWW$ quartic gauge coupling, production modes that involve the exchange
646 of a Higgs boson, and relatively low backgrounds. Evidence of EWK $W^\pm W^\pm jj$ production was first
647 seen by the ATLAS and CMS experiments at $\sqrt{s} = 8$ TeV, however the data set was too small to

648 claim observation of the process. The first analysis covered here is the follow up to the above ATLAS
649 measurement, measuring the EWK fiducial cross section at $\sqrt{s} = 13$ TeV with a larger data sample.
650 The second analysis explores the prospects for future measurements of the $W^\pm W^\pm jj$ process at the
651 planned High-Luminosity LHC (HL-LHC). A measurement of the production cross section as well
652 as sensitivity to the purely longitudinal component of the $W^\pm W^\pm$ scattering is presented.

653 In addition to the SM measurements, a part of this thesis is devoted to alignment of the detector
654 components making up ATLAS's Inner Detector (ID). Precise knowledge of the locations of detector
655 elements is essential for accurate particle track reconstruction, which in turn results in improved
656 resolutions for physics measurements. The ATLAS alignment algorithm determines the positions
657 of each ID sensor through minimizing the distance between reconstructed particle tracks and the
658 sensor hit position. Special emphasis is given to the monitoring of momentum biases that may exist
659 in the ID even after alignment.

660 The first few chapters of this thesis are intended to provide context for the main topics. Chapter 2
661 gives a brief introduction to the Standard Model with a focus on the EWSB mechanism and vector
662 boson scattering. The experimental apparatus—the LHC and the ATLAS detector—are detailed in
663 Chapter 3. The next three chapters present the main body of work. Chapter 4 covers the alignment
664 of the ATLAS Inner Detector. Chapters 5 and 6 detail the ATLAS $\sqrt{s} = 13$ TeV $W^\pm W^\pm jj$ cross
665 section measurement and the $\sqrt{s} = 14$ TeV HL-LHC $W^\pm W^\pm jj$ prospects study, respectively.

666

CHAPTER 2

667

Theoretical Framework

668 This chapter outlines the theoretical groundwork for the rest of the thesis. An overview of the Stan-
669 dard Model of particle physics (SM) is given in Section 2.1, followed by the electroweak symmetry
670 breaking mechanism involving the Higgs boson in Section 2.2. Finally, Section 2.3 will go into detail
671 on the interests of electroweak vector boson scattering (VBS).

672 **2.1 Introduction to the Standard Model**

673 The Standard Model of particle physics serves as a mathematical description of the fundamental
674 particles of the universe and their interactions. It has been developed over the course of the past
675 century, incorporating both predictions from theory and results from experiments. All in all, the
676 SM has proven to be very accurate in describing the particle interactions seen in experiments, as can
677 be seen in the summary plot of ATLAS standard model measurements through 2018 in Figure 2.1.

678 The SM is a quantum field theory (QFT) [3, 4] in which the fundamental particles are represented
679 as excited states of their corresponding fields. The spin- $\frac{1}{2}$ fermionic fields give rise to the quarks
680 and leptons comprising ordinary matter, the spin-1 fields correspond to the electroweak bosons and
681 the gluon which mediate the electroweak and strong forces, respectively, and finally the scalar Higgs
682 field is responsible for electroweak symmetry breaking. The excitations and interactions of the fields
683 are governed by the SM Lagrangian, which is invariant under local transformations of the group
684 $SU(3) \times SU(2) \times U(1)$.

685 The first quantum field theory to be developed was quantum electrodynamics (QED) [5], which
686 describes the electromagnetic interaction. The theory predicts the existence of a $U(1)$ gauge field
687 that interacts with the electrically charged fermions. This field corresponds to the photon. A key

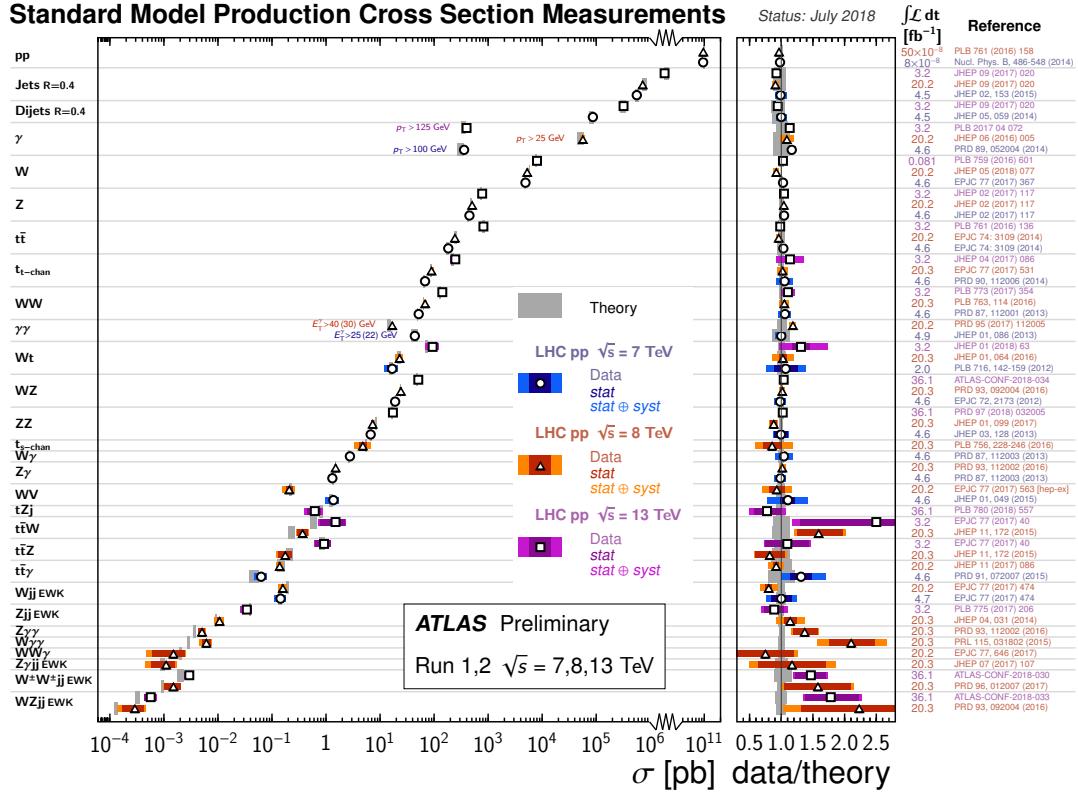


Figure 2.1: Summary of several Standard Model total and fiducial production cross section measurements compared to the corresponding theoretical predictions calculated at NLO or higher. The dark colored error bar represents the statistical uncertainty, and the light colored error bar represents the full uncertainty (including systematic and luminosity uncertainties). The data/theory ratio, luminosity used and reference for each measurement are also shown [2].

aspect of QED is that it is perturbative. The coupling constant $\alpha = e^2/4\pi$ is small, where e is electrical charge of the field, allowing for the use of perturbation theory in calculations. In this case, calculations can be written as a power series in α , where successive higher order terms contribute less to the final result. The accuracy of perturbative calculations is an essential tool for being able to make predictions from the SM.

The strong interaction—the theory of quarks and gluons—has also been described using QFT as quantum chromodynamics (QCD). The symmetry group for QCD is $SU(3)$, and its eight generators correspond to the eight differently charged, massless gluons [6]. Unlike in QED, which has positive and negative charges, the strong force has three “colors”. Color charge combined with the non-Abelian nature of $SU(3)$, which allows the gluons to interact with each other, result in the

most well-known property of QCD: color confinement. In order to increase the separation between two color-charged quarks, the amount of energy required increases until it becomes energetically favorable to pair-produce a new quark-antiquark pair, which then bind to the original quarks. The end result of this is that only color-neutral objects exist in isolation. What this means for the strong coupling constant α_s is that its value at the low energies where confinement occurs is large, on the order of $\alpha_s \sim 1$. The consequence of this is that perturbation theory cannot be used to accurately approximate interactions. While this appears at first to be a critical problem for predictions, fortunately it turns out that α_s “runs”, or decreases in magnitude at higher energy [7, 8]. This so-called “asymptotic freedom” allows QCD to be calculated perturbatively [9] at energies accessible by collider experiments including the LHC.

The last gauge field corresponds to the weak interaction. Ultimately, the weak SU(2) and the electromagnetic U(1) mix to form the $SU(2) \times U(1)$ *electroweak* (EWK) interaction [10, 11]. A more detailed description of the mixing will be discussed in conjunction with electroweak symmetry breaking (EWSB) in Section 2.2; however, a summary of the resulting EWK interaction is presented here, at the risk of some repeated information to follow. There are three weak isospin bosons arising from the SU(2) group (W_μ^1 , W_μ^2 , and W_μ^3) and one weak hypercharge boson from the U(1) group (B_μ). The W_3 and B bosons mix according to the weak mixing angle θ_W to form the Z boson and the photon according to:

$$\begin{pmatrix} \gamma \\ Z \end{pmatrix} = \begin{pmatrix} \cos \theta_W & \sin \theta_W \\ -\sin \theta_W & \cos \theta_W \end{pmatrix} \begin{pmatrix} B_\mu \\ W_\mu^3 \end{pmatrix} \quad (2.1)$$

The value of θ_W is not predicted by the SM; it is one example of an experimental input to the theory, measured to be $\sin^2 \theta_W = 0.23153 \pm 0.00016$ [12]. The charged W^\pm bosons are a mixture of the remaining W_μ^1 and W_μ^2 bosons:

$$W^\pm = \frac{1}{\sqrt{2}}(W_\mu^1 \mp iW_\mu^2) \quad (2.2)$$

Unlike the photon (and the gluon of QCD), the W^\pm and Z bosons are massive. This means that even though SU(2) is non-Abelian, the range of interaction is short and confinement does not occur. Lastly, the EWK interaction is chiral, only coupling to the left-handed component of the fermion fields.

One final field remains within the SM: the scalar Higgs field. It was originally proposed in the 1960’s to explain the masses of the W^\pm and Z bosons [13, 14, 15] and is the mechanism for the EWSB process. The particle associated with the field is a massive scalar boson, the Higgs boson, which was recently discovered by ATLAS and CMS in 2012 [16, 17] with a mass of 125 GeV.

⁷²⁷ **2.2 Electroweak symmetry breaking and the Higgs boson**

⁷²⁸ The results of electroweak mixing and the implications of the Higgs field have been introduced
⁷²⁹ in the previous section. If the EWK theory were an unbroken symmetry, the associated W^\pm and
⁷³⁰ Z bosons would be massless; however, when observed experimentally, they were found to be quite
⁷³¹ heavy [18, 19], at around 80 GeV and 91 GeV, respectively [20]. Here, a more detailed explanation
⁷³² of the Higgs mechanism and how it “spontaneously breaks” the EWK symmetry, resulting in the
⁷³³ three massive bosons (W^\pm and Z) and one massless boson (photon), is presented.

⁷³⁴ To see how the Higgs mechanism results in the massive vector bosons and a massless photon,
⁷³⁵ consider the following. Beginning with a complex scalar doublet ϕ defined as:

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = \sqrt{\frac{1}{2}} \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_3 + i\phi_4 \end{pmatrix} \quad (2.3)$$

⁷³⁶ a Lagrangian \mathcal{L} can be written:

$$\mathcal{L} = (\mathcal{D}_\mu \phi)^\dagger (\mathcal{D}^\mu \phi) - \mu^2 \phi^\dagger \phi - \lambda(\phi^\dagger \phi)^2 \quad (2.4)$$

⁷³⁷ where $\lambda > 0$ and \mathcal{D}_μ is the covariant derivative defined such that \mathcal{L} is invariant under a local
⁷³⁸ $SU(2) \times U(1)$ gauge transformation:

$$\mathcal{D}_\mu \phi = \left(\partial_\mu + \frac{ig}{2} \tau_a W_\mu^a + \frac{ig'}{2} B_\mu \right) \phi \quad (2.5)$$

⁷³⁹ Here W_μ^a ($a = 1, 2, 3$) are the $SU(2)$ fields with generators τ_a and coupling constant g , and B_μ is the
⁷⁴⁰ $U(1)$ field with coupling constant g' .

⁷⁴¹ Isolating the potential term of the Lagrangian:

$$V(\phi) = \mu^2 \phi^\dagger \phi + \lambda(\phi^\dagger \phi)^2 \quad (2.6)$$

⁷⁴² a choice must be made on the sign of μ^2 , and the case of interest is for $\mu^2 < 0$. This results in
⁷⁴³ the famous “Mexican hat potential” shown in Figure 2.2, which is minimized along the collection of
⁷⁴⁴ points:

$$\phi^\dagger \phi = \frac{1}{2} (\phi_1^2 + \phi_2^2 + \phi_3^2 + \phi_4^2) = -\frac{\mu^2}{2\lambda} \quad (2.7)$$

⁷⁴⁵ This means that the minimum of the potential is not at $\phi = 0$ (as it would be in the case where
⁷⁴⁶ $\mu^2 > 0$), but rather at a value:

$$v \equiv \sqrt{-\frac{\mu^2}{\lambda}} \quad (2.8)$$

747 With no loss of generality due to the SU(2) symmetry, $\phi_1 = \phi_2 = \phi_4 = 0$ can be imposed on
 748 Equation 2.7 leaving $\phi_3^2 = v^2$. Finally, the *vacuum expectation value* (VEV) of the field can be
 749 written as:

$$\langle\phi\rangle = \sqrt{\frac{1}{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \quad (2.9)$$

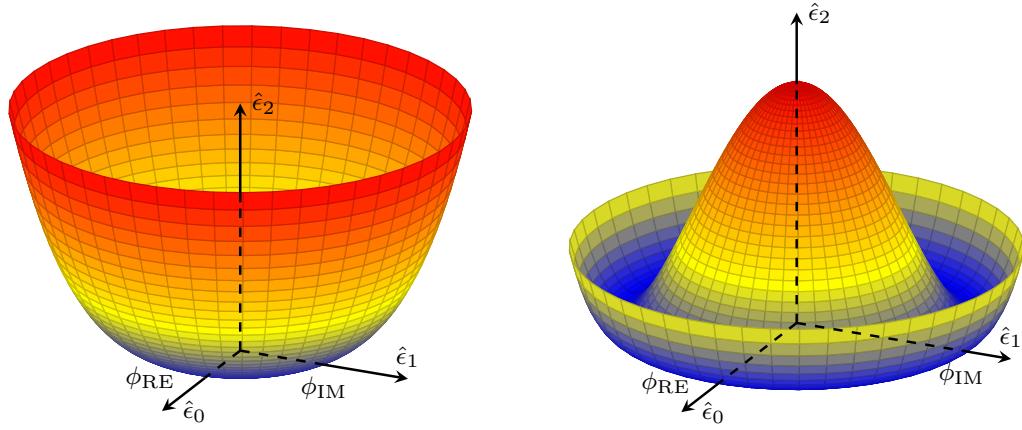


Figure 2.2: An illustration of the potential term $V(\phi)$ in Equation 2.6 for the cases where $\mu^2 > 0$ (left) and $\mu^2 < 0$ (right). The right-hand plot shows the Higgs potential, or ‘Mexican hat potential’, with the minimum at $|\phi| = \sqrt{-\frac{\mu^2}{\lambda}}$ rather than at $|\phi| = 0$ as in the left-hand plot.

750 The VEV can be substituted back into the original Lagrangian in Equation 2.4, and, following
 751 quite a bit of math [21], a collection of mass terms can be identified:

$$\mathcal{L} \subset \mathcal{L}_M \equiv \frac{1}{8}v^2g^2 \left[(W_\mu^1)^2 + (W_\mu^2)^2 \right] + \frac{1}{8}v^2 \left[g^2(W_\mu^3)^2 - 2gg'W_\mu^3B^\mu + g'^2(B_\mu)^2 \right] \quad (2.10)$$

752 Focusing on the first term for the moment, if Equation 2.2 for the physical W^\pm bosons is substituted
 753 in, the mass term can be seen clearly:

$$M_W^2 W^+ W^- = \left(\frac{1}{2}vg \right)^2 W^+ W^- \quad (2.11)$$

754

$$M_W = \frac{1}{2}vg \quad (2.12)$$

755 With a bit of clever forward-thinking, the second term of Equation 2.10 can be rewritten as:

$$\frac{1}{8}v^2 \left[gW_\mu^3 - g'B_\mu \right]^2 + 0 \left[g'W_\mu^3 - gB_\mu \right]^2 = \frac{1}{2}M_Z^2 Z_\mu^2 + \frac{1}{2}M_A^2 A_\mu^2 \quad (2.13)$$

756 where Z_μ^2 and A_μ^2 represent the physical Z boson and photon, respectively, and are defined as:

$$Z_\mu = \frac{gW_\mu^3 - g'B_\mu}{\sqrt{g^2 + g'^2}} \quad (2.14)$$

757

$$A_\mu = \frac{g'W_\mu^3 - gB_\mu}{\sqrt{g^2 + g'^2}} \quad (2.15)$$

758 From this, it can be seen that the photon is massless ($M_A = 0$ in Equation 2.13), and the mass of
759 the Z boson is identified as:

$$M_Z = \frac{1}{2}v\sqrt{g^2 + g'^2} \quad (2.16)$$

760 Lastly, the Higgs field can couple directly to the fermions. Taking the electron as an example,
761 an additional Lagrangian term can be written:

$$\mathcal{L}_e = -G_e[\bar{e}_L\phi e_R + \bar{e}_R\phi^\dagger e_L] \quad (2.17)$$

762 where e_L and e_R are the left-handed doublet and right-handed singlet, respectively, and ϕ is as in
763 Equation 2.3. The symmetry can be spontaneously broken by a perturbation about the VEV:

$$\phi = \sqrt{\frac{1}{2}} \begin{pmatrix} 0 \\ v + h(x) \end{pmatrix} \quad (2.18)$$

764 which, when substituted into \mathcal{L}_e gives:

$$\begin{aligned} \mathcal{L}_e &= -\frac{G_e}{\sqrt{2}}v(\bar{e}_L e_R + \bar{e}_R e_L) - \frac{G_e}{\sqrt{2}}(\bar{e}_L e_R + \bar{e}_R e_L)h \\ &= -m_e \bar{e}e - \frac{m_e}{v} \bar{e}eh \end{aligned} \quad (2.19)$$

765 for electron mass $m_e = \frac{G_e v}{\sqrt{2}}$. From the second term, it can be seen that the strength of the Higgs
766 coupling to the electron is proportional to the mass of the electron. The rest of the fermion couplings
767 follow from this example.

768 What is accomplished here is quite remarkable. The weak and electromagnetic interactions have
769 been unified into a single $SU(2) \times U(1)$ interaction, and the physical bosons observed in nature
770 arise as mixtures of the four gauge fields. Three of the four degrees of freedom in the scalar field
771 ϕ of Equation 2.3, are absorbed (or “eaten”) by the W^\pm and Z bosons, and the fourth generates
772 the Higgs boson. Additionally, it is shown that the Higgs couples to fermions in proportion to
773 their mass. From experimental measurements, the value of the VEV has been determined to be
774 $v \approx 246$ GeV [20]. However, it should be noted that the theory does not predict the mass of the
775 Higgs boson or of the fermions; these must all be determined from experiment.

776 **2.3 Electroweak vector boson scattering**

777 Due to the non-Abelian nature of the EWK interaction, the associated gauge bosons are allowed
778 to self-interact. This results in triple and quartic couplings of gauge bosons (TGCs and QGCs,
779 respectively). The SM allowed TGCs are the $WW\gamma$ and WWZ vertices, which can be measured
780 experimentally via diboson production or through vector boson fusion (VBF). QGCs predicted by the
781 model include $WWZ\gamma$, $WW\gamma\gamma$, $WWZZ$, and $WWWW$ vertices accessible in triboson production
782 or via vector boson scattering (VBS)¹ [22]. VBS processes are defined by a $VV \rightarrow VV$ signature,
783 where V represents one of the EWK gauge bosons (W^\pm , Z , or γ), as shown in Figure 2.3. The actual
784 interaction between the incoming and outgoing vector bosons can be mediated by the exchange of
785 a virtual V , directly via a QGC (as in Figure 2.4), or by the exchange of a Higgs boson (as in
786 Figure 2.5).

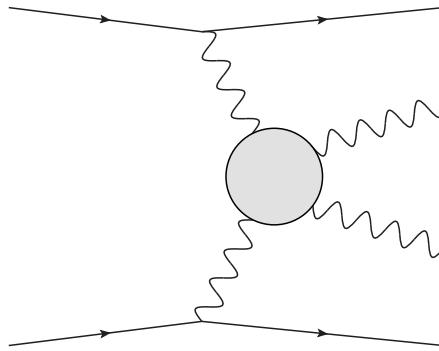


Figure 2.3: Feynman diagram of a generic VBS process. The gray circle represents any interaction with two incoming and two outgoing vector bosons, including any of the diagrams shown in Figures 2.4 and 2.5.

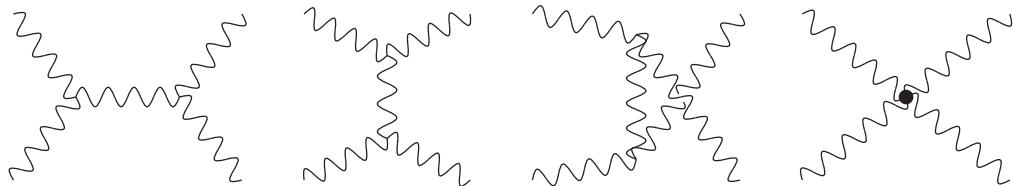


Figure 2.4: Leading order $VV \rightarrow VV$ Feynman diagrams involving EWK bosons. From left to right: s -channel, t -channel, u -channel, and the quartic gauge coupling.

¹Vector boson fusion and scattering typically refer to the s -channel and t -channel exchanges of a vector boson, respectively. Since both deal with a similar $VV \rightarrow VV$ process, for the remainder of this thesis, *vector boson scattering* will refer to both VBF and VBS.

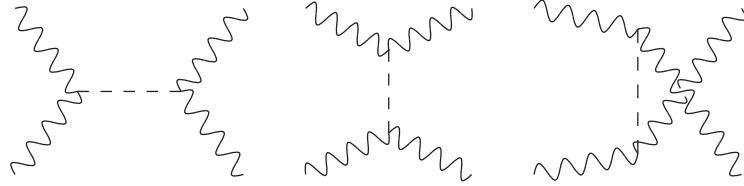


Figure 2.5: Leading order $VV \rightarrow VV$ Feynman diagrams involving the exchange of a Higgs boson. From left to right: s -channel, t -channel, and u -channel.

As detailed in the previous section, the Higgs mechanism produces three Goldstone bosons and a Higgs boson. The Goldstone bosons are then “eaten” by the physical gauge bosons, giving them mass and consequently a longitudinal polarization². In fact, according to the Electroweak Equivalence Theorem, the high-energy interactions of longitudinal gauge bosons can be accurately described by the Goldstone bosons of the EWSB mechanism [23]. Thus, the scattering of the massive gauge bosons are inextricably linked to EWSB.

It turns out that without a light SM Higgs boson, the scattering amplitude of longitudinally polarized vector bosons grows with center-of-mass energy and ultimately violates unitarity above $\sqrt{s} \approx 1.2$ TeV [24, 25]. Writing down the equations for the transverse and longitudinal polarization vectors for a gauge boson of mass M_V [26]:

$$\epsilon_{\pm}^{\mu} = \frac{1}{\sqrt{2}}(0, 0, \pm i, 0) \quad (2.20)$$

$$\begin{aligned} \epsilon_L^{\mu} &= \frac{1}{M_V}(|\vec{p}|, 0, 0, E) \\ &= \frac{p^{\mu}}{M_V} + v^{\mu} \end{aligned} \quad (2.21)$$

where v^{μ} is of the order M_V/E and becomes small in the high energy limit, it can be seen that ϵ_L^{μ} grows with the momentum of the boson p^{μ} . Therefore, the dominant contribution to the VBS process at high energy comes from the longitudinally polarized gauge bosons [27].

The high-energy behavior of longitudinally polarized vector boson scattering can be explored in the case of opposite-sign $W^+W^- \rightarrow W^+W^-$ scattering. In the high-energy limit ($s \gg M_W^2, M_H^2$), the amplitude of W^+W^- scattering without considering the Higgs contributions (the relevant diagrams in Figure 2.4) can be written as [25]:

$$\mathcal{M}_{\text{gauge}} = -\frac{g^2}{4M_W^2}u + \mathcal{O}\left(\left[\frac{E}{M_W}\right]^0\right) \quad (2.22)$$

²A massless spin-1 boson can have one of two transverse polarization states, while a massive spin-1 boson can also be longitudinally polarized. As a result, only the massive W^{\pm} and Z bosons, and not the massless photon, are sensitive to EWSB.

where g is the EWK coupling and u is one of the Mandelstam variables (the others being s and t). The $\mathcal{O}(E^4)$ terms cancel out between the TGC and QGC diagrams [27]. What is left is an amplitude proportional to E^2 that diverges as $E/M_W \rightarrow \infty$. However, the amplitude from the diagrams involving the Higgs boson (the relevant diagrams in Figure 2.5) is:

$$\mathcal{M}_{\text{Higgs}} = -\frac{g^2}{4M_W^2} \left[\frac{(s - M_W^2)^2}{s - m_H^2} + \frac{(t - M_W^2)^2}{t - M_H^2} \right] \quad (2.23)$$

which, in the high-energy limit, reduces to:

$$\mathcal{M}_{\text{Higgs}} = \frac{g^2}{4M_W^2} u + \mathcal{O}\left(\left[\frac{E}{M_W}\right]^0\right) \quad (2.24)$$

Adding the two equations together cancels out the E^2 term and leaves only terms constant in energy. Therefore, with a SM Higgs, the scattering amplitude for longitudinally polarized W bosons no longer diverges. Plots of the cross section of several $VV \rightarrow VV$ scattering processes are shown in Figure 2.6 with and without a SM Higgs boson.

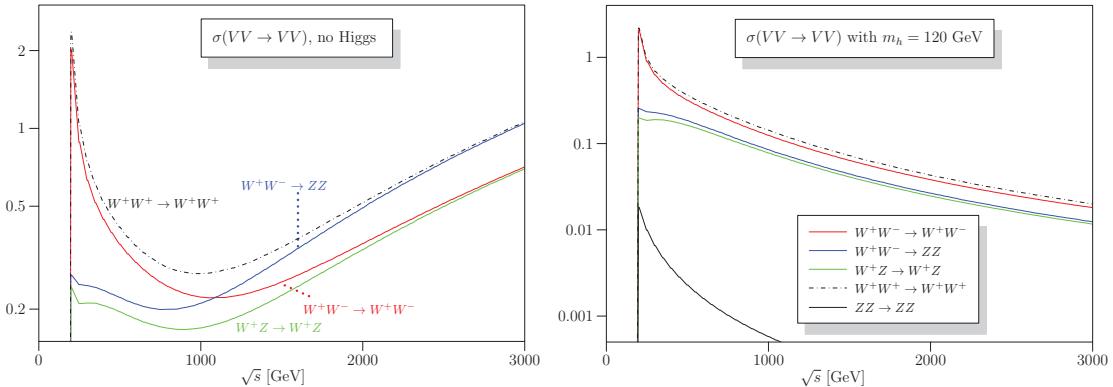


Figure 2.6: Cross sections in nanobarns for five different longitudinally polarized VBS processes as a function of center of mass energy \sqrt{s} . Without a Higgs boson (left), the cross sections grow unbounded with \sqrt{s} . With a 120 GeV Higgs boson (right), the cross sections no longer diverge. Plots taken from [28].

814

CHAPTER 3

815

LHC and the ATLAS Detector

816 This chapter covers the experimental apparatus relevant to this thesis: the Large Hadron Col-
817 lider (LHC) and the ATLAS detector in Sections 3.1 and 3.2, respectively. Some time is taken in
818 Section 3.2.4 to overview the methods used to identify and measure various particle types within
819 ATLAS.

820 3.1 The Large Hadron Collider

821 The Large Hadron Collider (LHC) [29] is the most powerful particle accelerator in the world, colliding
822 two beams of protons at a center of mass energy of $\sqrt{s} = 13$ TeV. It is operated by the European
823 Organization for Nuclear Research (CERN) and consists of a 27 km ring located beneath the France–
824 Switzerland border. A chain of smaller accelerators incrementally boost the protons³ up to higher
825 and higher energies before they reach the final collision energy within the main LHC ring. Collisions
826 occur at each of four detector experiments situated around the ring: ATLAS [30], ALICE [31],
827 CMS [32], and LHCb [33].

828 Protons are obtained from hydrogen atoms stripped of their electrons by an electric field. A beam
829 of protons is first accelerated up to 50 MeV in the Linac 2 accelerator, then to 1.4 GeV in the Proton
830 Synchrotron Booster (PSB), 25 GeV in the Proton Synchrotron (PS), and finally to 450 GeV in the
831 Super Proton Synchrotron (SPS). They are then injected into the LHC ring in two beams running in
832 opposite directions where they accelerate up to the collision energy of 6.5 TeV. The beams consist of
833 bunches containing on the order of 10^{11} protons separated by 25 ns [34]. A schematic of the CERN
834 accelerator complex, including the chain of accelerators mentioned above, is shown in Figure 3.1.

³The LHC can also collide beams of heavy ions; however, this thesis focuses exclusively on the proton-proton collisions.

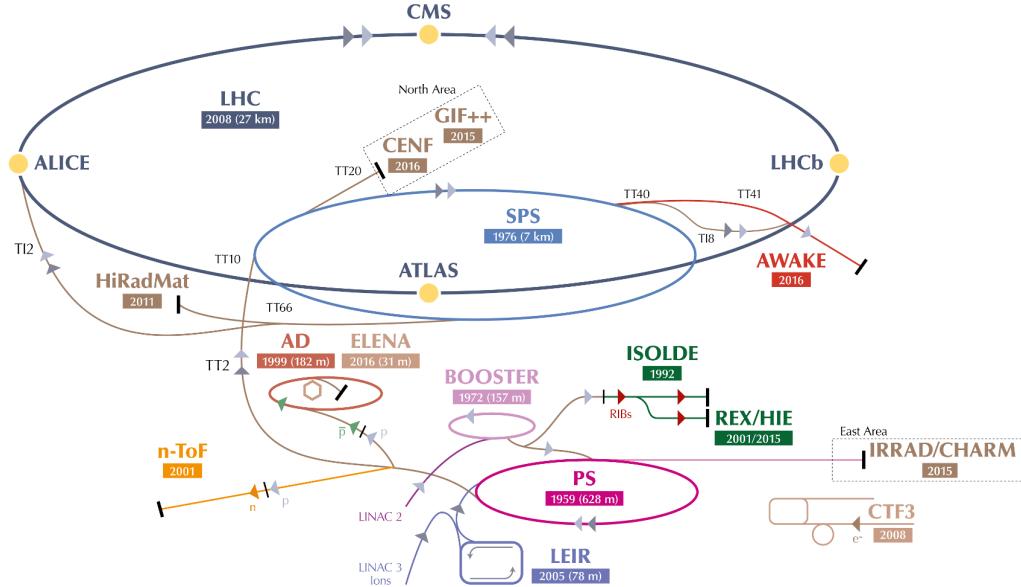


Figure 3.1: The CERN accelerator complex. For LHC collisions, protons are accelerated by the Linac 2 (purple), the PSB (light purple), the PS (magenta), and the SPS (light blue) before entering the LHC ring (dark blue) [35].

835 In addition to a high center of mass energy, the LHC must also deliver enough data to measure
 836 rare processes. The amount of data collected is measured in terms of *luminosity*. The instantaneous
 837 luminosity \mathcal{L} is defined in terms of the number of events per second $\frac{dR}{dt}$ and the production cross
 838 section σ_p :

$$\mathcal{L} = \frac{1}{\sigma_p} \frac{dR}{dt} \quad (3.1)$$

839 The calculation itself can be quite tricky, as it depends on a number of factors including the number
840 of particles per bunch, the spread of the beam, and the crossing angle of the beams [36].

The LHC was originally designed to operate at an instantaneous luminosity of $1.0 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$; however, this number was exceeded by the end of the 2016 data taking period, with a peak luminosity of $1.38 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$. By the end of Run 2 in December 2018, the LHC was running at more than twice the design luminosity [37]. The instantaneous luminosity of proton-proton collisions as a function of time in 2015 and 2016 are shown in Figure 3.2.

The total amount of data collected is reported in terms of *integrated* luminosity, which is simply the time integral of the instantaneous luminosity. By the end of Run 2 (2015-2018), approximately 140 fb^{-1} of 13 TeV data collected by the ATLAS detector is available for physics, as shown in Figure 3.3. The 36.1 fb^{-1} collected during the first two years (2015 and 2016) is used for the

analysis later in this thesis.

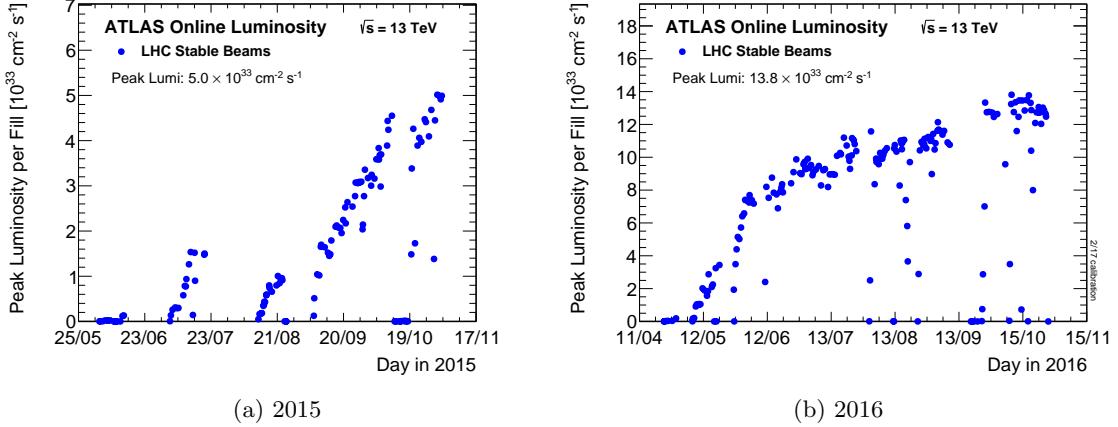


Figure 3.2: Peak instantaneous luminosity delivered to ATLAS during 13 TeV pp data taking as a function of time [37].

Due to the high instantaneous luminosity, more than one pp interaction occurs in a single bunch crossing, referred to as *pileup*. During the 2016 data taking campaign, the average number of interactions per bunch crossing $\langle\mu\rangle$ was approximately 24 but has increased to upwards of 37 in 2017 and 2018 [37]. Figure 3.4 contains the average μ for the 2015-2016 data set used for analysis in this thesis. The high pileup is a challenge for accurately reconstructing an individual collision.

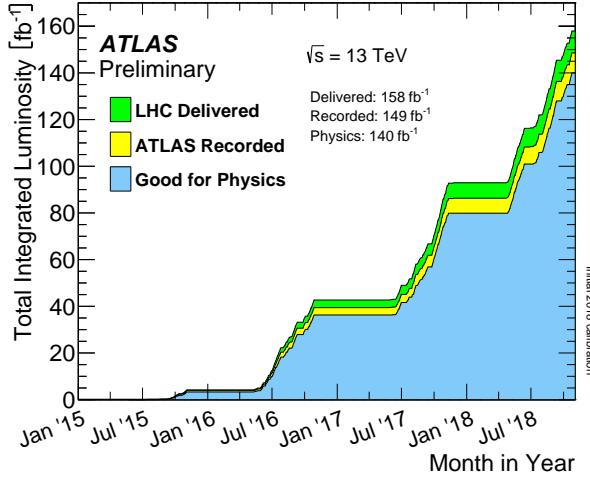


Figure 3.3: Integrated luminosity collected by ATLAS as a function of time at 13 TeV from 2015-2018 [37].

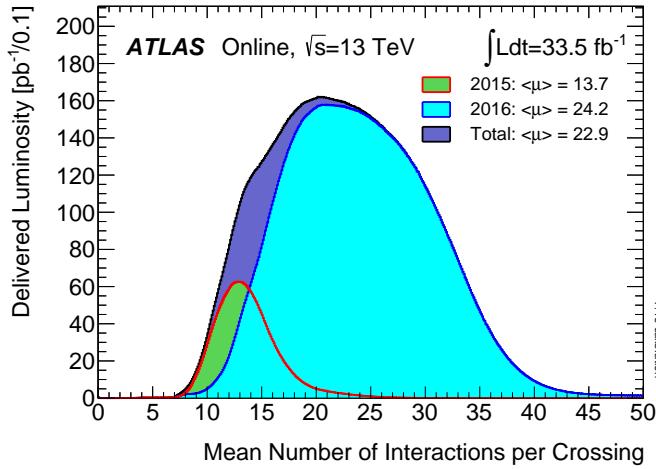


Figure 3.4: Distribution of the mean number of interactions per bunch crossing for the 2015 and 2016 pp collision data at 13 TeV [37].

856 3.2 The ATLAS Detector

857 ATLAS (A Toroidal LHC ApparatuS) is a general-purpose particle detector. It is made up of
 858 several different subdetector systems designed to measure different types of particles. Starting from
 859 the beam line and working outwards, the Pixel Detector (PIX), Semiconductor Tracker (SCT),
 860 and Transition Radiation Tracker (TRT) make up the Inner Detector (ID) and are responsible for
 861 measuring the trajectories and momenta of charged particles. Next are two calorimeters, the Liquid
 862 Argon Calorimeter (LAr) and the Tile Calorimeter (TileCal), which stop electromagnetic and hadronic
 863 objects and measure their energies. Finally, the outermost Muon Spectrometer (MS) measures muon
 864 tracks as they leave the detector, as they are too heavy to be stopped by the calorimeters. The
 865 ATLAS detector and its subsystems are shown in Figure 3.5.

866 ATLAS uses a global, right-handed coordinate system with the origin at the center of the detector
 867 (the nominal interaction point). The x -axis points from the origin inwards to the center of the LHC
 868 ring, the y -axis points upwards, and the z -axis points along the beam line. Due to the azimuthal
 869 symmetry of the detector, it is useful to use cylindrical coordinates (r, ϕ) in the plane transverse to
 870 the z -axis, where ϕ is the azimuthal angle. Instead of using the polar angle θ to describe particle
 871 trajectories, pseudorapidity η is used instead, defined in terms of θ :

$$\eta = -\ln(\tan(\theta/2)) \quad (3.2)$$

872 Pseudorapidity has the useful property that differences in η are invariant under Lorentz boosts along

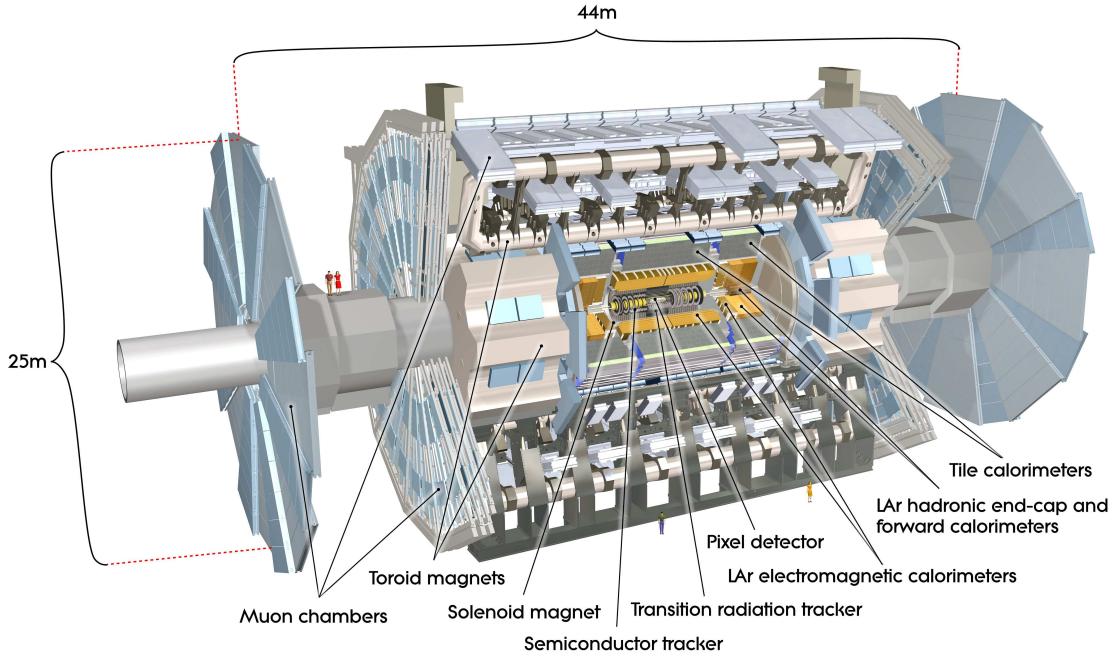


Figure 3.5: Cut-away view of the ATLAS detector [30].

873 the z -axis. The separation between two particles p_i and p_j is often expressed in terms of the quantity
 874 ΔR , defined as:

$$\Delta R(p_i, p_j) = \sqrt{\Delta\eta_{i,j}^2 + \Delta\phi_{i,j}^2} \quad (3.3)$$

875 where $\Delta\phi_{i,j} \in [-\pi, \pi]$ since ϕ is periodic in 2π and “wraps around” the detector in the azimuthal
 876 direction.

877 3.2.1 The Inner Detector

878 The ID [38, 39] is a tracking system that reconstructs the trajectories of charged particles. It spans
 879 just over a meter in radius, with the innermost layer of sensors at a radius of 33.25 cm from the
 880 beam line. Charged particles traveling through the ID leave *hits* in each sensor they pass through,
 881 and a track is fit to the hits to reconstruct the path of the particle according to the techniques that
 882 will be outlined in Section 3.2.4.1. The ID’s pseudorapidity coverage extends out to $|\eta| < 2.5$. A
 883 solenoid magnet outside the ID produces a 2 T magnetic field that bends the particles, allowing for
 884 their momenta in the direction transverse to the field to be measured according to:

$$p_T = q \cdot B \cdot r \quad (3.4)$$

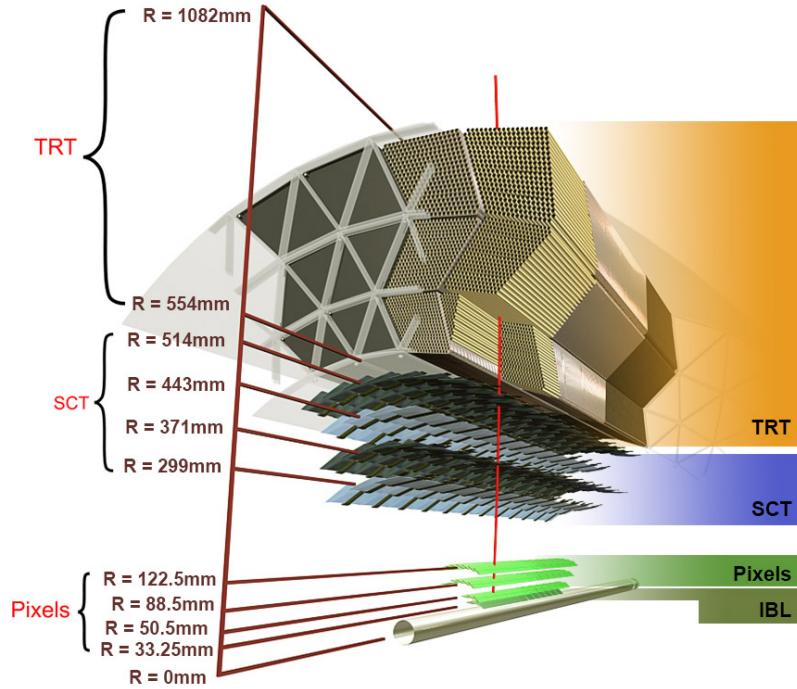


Figure 3.6: The barrel layers of the Pixel, SCT, and TRT detectors making up the Inner Detector.

885 where q is the charge of the particle (± 1), B is the strength of the magnetic field, and r is the radius
 886 of the track's curvature. A cut-away view of the barrel region of the ID is shown in Figure 3.6.

887 3.2.1.1 Pixel Detector

888 The Pixel Detector consists of four cylindrical barrel layers⁴ and three endcap disks on either side.
 889 It is the innermost subdetector of the ID with coverage up to $|\eta| < 2.5$. The individual sensors
 890 measure $50 \mu\text{m} \times 400 \mu\text{m}$ and are installed on silicon wafers that make up the layers. All in all, there
 891 are 1744 wafers with 80 million readout channels. The sensors themselves are silicon semiconducting
 892 diodes that provide a signal when a charged particle passes through. The Pixel Detector has the
 893 finest resolution of all the ID subdetectors, at $10 \mu\text{m}$ in the $r\text{-}\phi$ plane and $40 \mu\text{m}$ in the z direction.

894 During the upgrade period between Run 1 and Run 2, a new innermost layer was added to the
 895 Pixel detector barrel: the Insertable B-Layer (IBL) [40]. The IBL lies closest to the interaction
 896 point, at a radius of 33.25 cm from the beam line, and it is relied upon to provide high-precision

⁴For now, the outer three barrel layers will be covered in conjunction with the endcaps; the innermost layer will be described separately.

897 measurements close to the interaction point. Its addition allows better precision in detecting dis-
898 placed vertices from b -jets, for example. It consists of 280 silicon pixel modules arranged on 14
899 staves that run parallel to the beam line. Each stave consists of 12 two-chip planar modules in the
900 middle ($|\eta| < 2.7$) with four 3D sensors [41] on either side ($2.7 < |\eta| < 3.0$). The IBL's pixel sensors
901 are $50 \mu\text{m} \times 250 \mu\text{m}$ in size and have a resolution of $10 \mu\text{m}$ in $r\text{-}\phi$ and $75 \mu\text{m}$ in z [42].

902 **3.2.1.2 Semiconductor Tracker**

903 The next subdetector of the ID is the SCT, which has four barrel layers and nine endcap disks per
904 side and provides coverage within $|\eta| < 2.5$. The SCT operates on the same principle as the Pixel
905 Detector, but the sensitive elements are larger silicon “strips” placed on the wafers. This shape
906 change assissts in covering the larger surface area required by the increasing detector radius. Each
907 detector layer is actually made up of two layers of wafers, placed back-to-back with an angle of
908 40 mrad between them. The resolution in the $r\text{-}\phi$ plane is very fine at $17 \mu\text{m}$, but, due to the strip
909 shape, the resolution along z is rather poor at $580 \mu\text{m}$.

910 **3.2.1.3 Transition Radiation Tracker**

911 The outermost component of the ID is the TRT [43, 44, 45], which uses a completely different
912 technology from the Pixel and SCT to identify particle hits. The TRT is unique in that it combines
913 a drift tube tracker with transition radiation detection to assist with electron identification. The
914 TRT's sensitive elements are drift tubes (referred to as “straws”) that are 4 mm in diameter and
915 consist of a cylindrical cathode with an anode wire running through the center. Each straw is filled
916 with a gas mixture including xenon or argon which provides ionizing radiation when high energy
917 particles pass through them. The resulting electrons drift to the anode and register a voltage,
918 indicating a hit in the detector element.

919 Between the straws are polyethelene fibers in the barrel and polypropylene foil in the endcaps
920 in order to encourage particles to emit transition radiation photons. These photons also ionize the
921 gas within the straws, leading to a higher signal. The TRT takes advantage of the fact that lighter
922 particles are more likely to emit transition radiation by using a ternary output: zero, low-threshold,
923 and high-threshold. High-threshold hits are generally caused by electrons due to their low mass,
924 and this can help in identifying electron tracks from backgrounds.

925 There are over 100,000 straws in the barrel of the TRT, and nearly 250,000 in the endcaps. The
926 TRT provides pseudorapidity coverage up to $|\eta| < 2.0$ with a resolution in the $r\text{-}\phi$ plane of $130 \mu\text{m}$.

927 Since the drift tubes are insensitive along the direction of the wire, the TRT does not provide a
 928 measurement along the z direction.

929 3.2.2 The Calorimeters

930 ATLAS utilizes two different calorimeters, the Liquid Argon and Tile Calorimeters [46, 47], in order
 931 to measure electromagnetic and hadronic objects. The general principle behind both calorimeters is
 932 the same: an incoming particle showers as it passes through and eventually stops, and the resulting
 933 energy deposits are read out. Both are sampling calorimeters, which consist of alternating layers of a
 934 dense material to induce the showering (called the *absorber*) and a second material which measures
 935 the energy (called the *active material*). An advantage to this type of calorimeter is that a very
 936 dense absorber can be used in order to produce a shower in a limited space, even if it is unsuitable
 937 for measuring the energy from the shower. However, as a result, some of the energy is deposited in
 938 the absorbers, and the total shower energy must be estimated. ATLAS's calorimeter systems are
 939 shown in Figure 3.7.

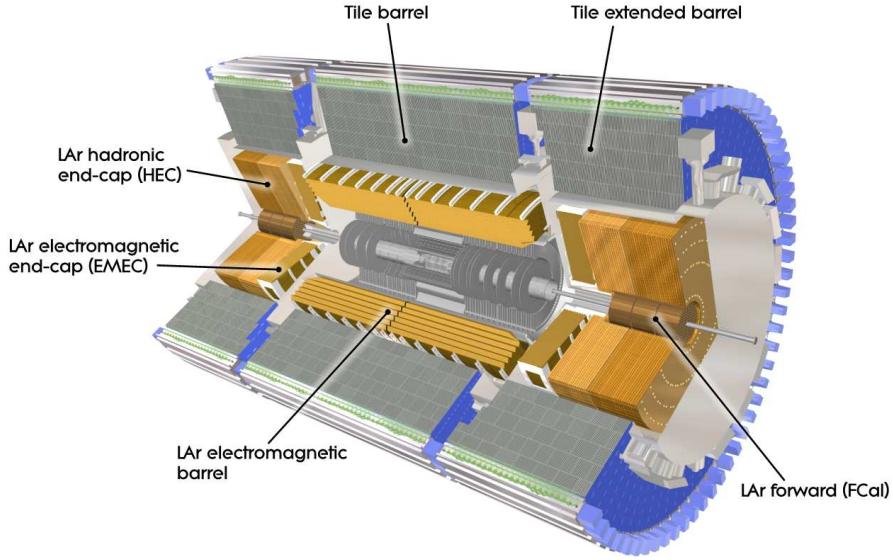


Figure 3.7: Cut-away view of the ATLAS calorimeter systems [48].

940 Electromagnetic objects, such as electrons and photons, shower via cascades of bremsstrahlung
 941 photons and e^+e^- pairs. The radiation length X_0 is defined as the mean distance over which an
 942 electron's energy is reduced to $1/e$ of its original value, or $E(x) = E_0 e^{-x/X_0}$. The majority of the

943 shower energy is deposited in the first few radiation lengths. The longitudinal shower depth scales
 944 logarithmically with particle energy, and the transverse shower width is described by the Molière
 945 radius⁵ of the material.

946 Hadronic showers (referred to as *jets*) are the result of quarks or gluons which hadronize and
 947 shower primarily via the strong interaction. Hadronic showers are generally wider than the electro-
 948 magnetic showers described above. The longitudinal depth of the hadronic shower scales with the
 949 nuclear interaction length of the material λ , defined as the mean distance for the number of particles
 950 in a hadronic jet to be reduced to $1/e$ of the initial number. In addition, about 1/3 of the shower
 951 products are neutral pions π^0 which decay electromagnetically via the process described above.

952 3.2.2.1 Liquid Argon Calorimeter

953 The LAr Calorimeter contains four individual calorimeters: the electromagnetic barrel (EMB) and
 954 endcaps (EMEC), and the hadronic endcap (HEC) and forward calorimeter (FCal). The calorimeter
 955 is surrounded by a cryostat held at a temperature around 90 K.

956 Focusing on the electromagnetic components first, the EMB covers $|\eta| < 1.475$ and the two
 957 EMECs cover $1.375 < |\eta| < 3.2$. They consist of alternating layers of lead absorber and liquid
 958 argon. The exact thickness of the lead depends on the location within the detector, ranging from
 959 1.1-2.2 mm. The absorbers are folded into an accordion shape, where the folding angles are varied
 960 in order to keep the thickness of the liquid argon gap constant across the barrel (about 2.1 mm).
 961 The electromagnetic calorimeter is thick enough that the minimum number of radiation lengths a
 962 particle travels through is $24 X_0$, including the material from other subdetectors.

963 There are four layers within the EMB and EMEC including an innermost pre-sampler that helps
 964 correct for energy lost before the shower reaches the calorimeter. The next three layers consist of
 965 differently shaped cells successively reducing in granularity. The first layer consists of narrow strips
 966 for fine-grained η resolution, while the majority of the shower energy is deposited in the second
 967 layer. The accordion shape as well as the sizes of the cells in the EMB are shown in Figure 3.8.

968 The HEC is located directly behind the EMEC and covers $1.5 < |\eta| < 3.2$. It uses thick copper
 969 plates as the absorber (25 mm in the front wheels and 50 mm in the rear wheels) separated by 8.5 mm
 970 gaps filled with liquid argon. Rather than the accordion shape, the HEC cells are rectangular.

971 The FCal provides coverage for hadronic jets over the range $3.2 < |\eta| < 4.9$. Each FCal endcap
 972 consists of three layers. The first is an electromagnetic calorimeter with a copper absorber, while the

⁵A cone with a radius equal to the Molière radius (M_R) will contain approximately 90% of the shower energy. At a radius of $2M_R$, 95% of the energy will be contained.

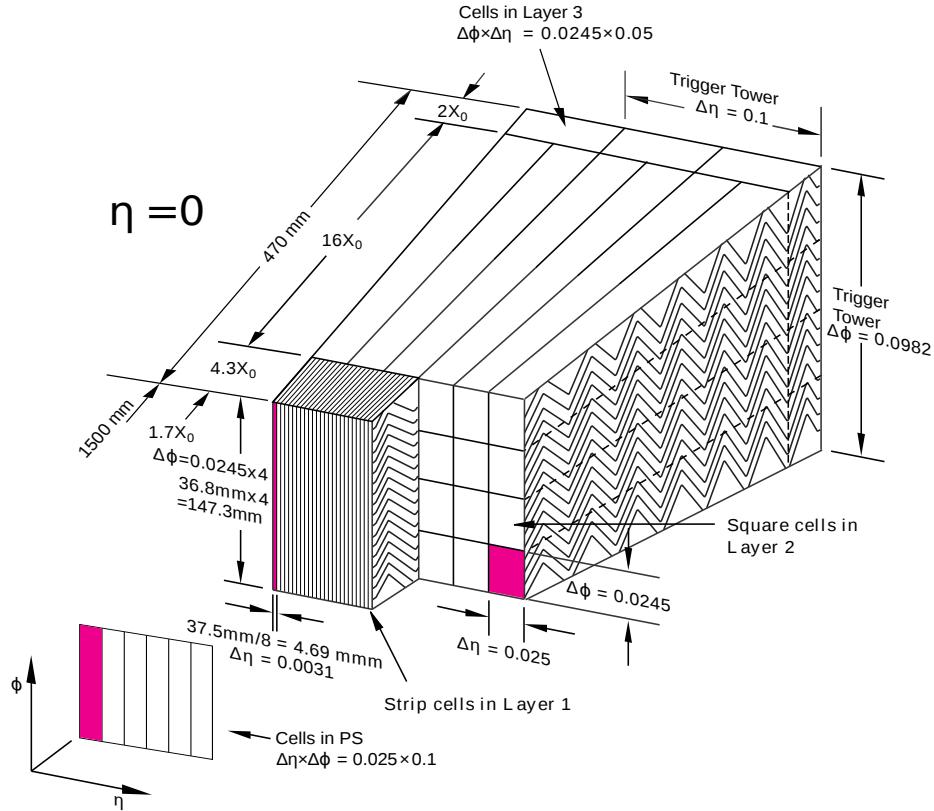


Figure 3.8: Diagram of the cells within the LAr barrel. The accordion structure can be seen in the cut-away view [46].

973 other two hadronic layers use a tungsten absorber. Due to the high particle flux entering the FCal,
 974 the liquid argon gaps are very narrow, and electrodes are embedded into the absorbers parallel to
 975 the beam line.

976 3.2.2.2 Tile Calorimeter

977 The TileCal consists of a barrel and two “extended barrel” sections which cover the range $|\eta| < 1.7$.
 978 It consists of alternating layers of steel plates and polystyrene scintillator tiles as the absorbers and
 979 active material, respectively. The total thickness of the TileCal is approximately 9λ . As the shower
 980 passes through the scintillators, photons are emitted that are picked up by wavelength shifting fibers
 981 and passed to photomultiplier tubes.

982 3.2.3 The Muon Spectrometer

983 The outermost subdetector in ATLAS is the Muon Spectrometer [49]. Due to the high mass of
984 muons compared to electrons, they pass through the calorimeters, necessitating their own detector.
985 The MS is a high-resolution spectrometer which provides tracking for muon reconstruction within
986 $|\eta| < 2.7$. A set of toroid magnets generate an azimuthal magnetic field that bends the muons for
987 momentum measurements, much like in the ID. Four different technologies are used in the MS:

- 988 • Monitored Drift Tubes (MDT) are used across the entire η range for precision measurements
989 of the tracks with a per-hit resolution in the range of 60-80 μm . These consist of an aluminum
990 tube filled with a gas mixture containing argon and an anode wire running through the middle
991 of the tube. When a muon passes through, the gas is ionized, and the electrons are collected
992 on the wire.
- 993 • Cathode Strip Chambers (CSC) are used for the forward regions of the endcaps ($|\eta| >$
994 2.0). They operate on a similar principle to the MDTs, with strips containing a mesh of anode
995 wires running in parallel instead of tubes with a single wire each.
- 996 • Resistive Plate Chambers (RPC) in the barrel are primarily used to provide input for the
997 muon trigger system. They consist of pairs of plastic resistive plates with a 2 mm gap between
998 them filled with a gas mixture. Electrodes are attached to the plates to create a potential
999 between them, and muons passing through ionize the gas and lead to electric discharges which
1000 in turn reduce the potential.
- 1001 • Thin Gap Chambers (TGC) are used for triggering in the endcaps. The TGCs are arranged
1002 on circular disks consisting of two rings, and are similar in function to the CSCs but with a
1003 different gas mixture.

1004 3.2.4 Particle reconstruction

1005 In order to convert the raw detector readouts to information about a given particle, various sets of
1006 reconstruction algorithms are run. This includes building particle trajectories in the ID, as well as
1007 identifying and measuring electrons and muons. A brief overview of these reconstruction methods
1008 follow in order to provide context for when these objects are used later in the thesis.

1009 **3.2.4.1 Track reconstruction**

1010 *Track reconstruction* is the process by which a particle’s trajectory is reconstructed from the raw
1011 measurements recorded in the ID. The ATLAS track reconstruction algorithm [50] follows three
1012 main steps: clusterization, track finding, and ambiguity solving.

1013 The first step, clusterization, uses hits from the Pixel and SCT detectors. Neighboring pixels or
1014 silicon strips that registered a hit in a sensor are grouped together into clusters using a connected
1015 component analysis. Each cluster represents a *space-point*, or a three-dimensional measurement
1016 corresponding to the point where the particle intersected the sensor. Since the SCT sensors consist
1017 of two strip layers on top of each other (as described in Section 3.2.1.2), clusters from each layer
1018 combine to form a single space-point. It is possible to have overlapping clusters from multiple
1019 particles in a single sensor, and care is taken that these *merged clusters* are identified and handled
1020 accordingly.

1021 Next, track seeds are formed using sets of three space-points. This number allows for a first
1022 momentum estimate to be made while still allowing for as large a number of track combinations as
1023 possible. The impact parameters of a track seed—the distance of closest approach to the collision—
1024 are estimated by assuming a perfect helical trajectory in a uniform magnetic field. In order to ensure
1025 the quality of a track seed, criteria on the momentum and impact parameters are imposed as well as
1026 a requirement that at least one additional space-point lies along the preliminary trajectory. Finally,
1027 a combinatorial Kalman filter [51] builds track candidates from the track seeds by incorporating
1028 additional space-points lying along the preliminary trajectory. The filter allows for multiple track
1029 candidates to be fit to the same track seed if more than one set of space-points is compatible.

1030 The final step of the reconstruction process is the ambiguity solving. Each collected track
1031 candidate is processed individually before being sorted by its *track score*, a metric for quantifying
1032 the likelihood that a given track candidate correctly represents the particle’s trajectory. The track
1033 score is determined from a number of factors. Each cluster along the track candidate increases
1034 the score by a weighted amount that incorporates the intrinsic resolution of the relevant detector’s
1035 sensors in addition to other factors. Conversely, if a track candidate passes through a sensor but
1036 there is no associated cluster (called a *hole*), the score is reduced. The χ^2 of the track fit contributes
1037 as well in order to promote tracks with high quality fits. Finally, the logarithm of the momentum
1038 adds to the score to suppress tracks with incorrectly assigned clusters, as these tracks typically have
1039 low momenta. A track candidate is rejected by the ambiguity solver if its score is too low, or if it
1040 fails to meet a basic set of quality criteria. If a cluster would be shared by more than one track, at

1041 most two tracks are allowed to pass through it. In this case, preference is given to tracks already
 1042 passing through the ambiguity solver, which by construction results in the two highest scoring tracks
 1043 using the shared cluster being kept.

1044 Following this procedure, TRT hits can be incorporated into the track fit through *TRT track*
 1045 *extension* [52]. Compatible sets of TRT measurements are found for tracks found in the silicon
 1046 detectors surviving the ambiguity solving. The algorithm requires that the original silicon-only
 1047 track not be modified by the inclusion of the TRT hits; it is simply an extension of the existing
 1048 track.

1049 What is described above is the *inside-out* reconstruction algorithm; there is also an *outside-in*
 1050 reconstruction that begins in the TRT. This algorithm is not covered in detail here, as much of the
 1051 process is similar to the above. The general workflow begins with finding track segments in the
 1052 TRT, constructing the track candidates including the silicon hits, and finally ambiguity solving.

1053 3.2.4.2 Electron reconstruction

1054 Electron reconstruction [53] uses information from both the ID and the electromagnetic calorimeters.
 1055 The characteristic signature of an electron in ATLAS is a charged particle track in the ID that is
 1056 matched in $\eta\text{-}\phi$ to localized clusters of energy deposited in the calorimeter.

1057 Calorimeter cluster candidates are seeded from localized energy deposits according to a sliding-
 1058 window algorithm [54]. The clusters are a 3×5 rectangle of calorimeter towers in $\eta \times \phi$ with a
 1059 total transverse energy greater than 2.5 GeV. In the event that two clusters overlap, if the E_T of
 1060 the clusters vary by more than 10%, the highest E_T cluster is kept, otherwise the candidate with
 1061 the highest E_T in the central tower is kept. The ID tracks are those generated using the algorithm
 1062 detailed above in Section 3.2.4.1, and tracks with a nearby calorimeter cluster are re-fit using a
 1063 Gaussian-sum filter designed to take into account bremsstrahlung effects [55].

1064 To reconstruct the final electron candidate, the refit track and the cluster are subject to the final
 1065 matching criteria:

$$|\eta_{\text{cluster}} - \eta_{\text{track}}| < 0.05 \quad \text{and} \quad (3.5)$$

$$-0.10 < q \times \Delta\phi_{\text{cluster, track}} < 0.05 \quad \text{or} \quad -0.10 < q \times \Delta\phi_{\text{res}} < 0.05$$

1066 where q is the charge of the track and $\Delta\phi_{\text{res}}$ is the azimuthal separation between the cluster position
 1067 and the track after rescaling its momentum to the energy of the cluster. If multiple tracks satisfy the
 1068 above criteria, the primary track is selected by an algorithm that takes into account the center of

1069 each cluster relative to the parameters of the candidate track. Finally, the clusters are reconstructed
 1070 about the seed cluster using a larger window size, 3×7 in the barrel and 5×5 in the endcaps, and
 1071 the energy is calibrated to the original electron energy using techniques described in [56, 57].

1072 **Electron identification** To determine whether an electron candidate is a signal electron or back-
 1073 ground object, an identification criteria (ID) is implemented, covered in more detail in [53, 58]. Elec-
 1074 tron ID is performed using a multivariate likelihood technique (LH) that simultaneously evaluates
 1075 a list of measurements of an electron candidate, including both tracking and calorimeter clustering
 1076 information. To cover the different requirements of various physics and performance studies, four
 1077 different likelihood working points are constructed corresponding to increasing thresholds for the
 1078 LH discriminant: `VeryLooseLH`, `LooseLH`, `MediumLH`, and `TightLH`. Each successive working point is
 1079 a subset of its predecessors. The efficiencies of the `LooseLH`, `MediumLH`, and `TightLH` working points
 1080 as a function of electron E_T are shown in Figure 3.9.

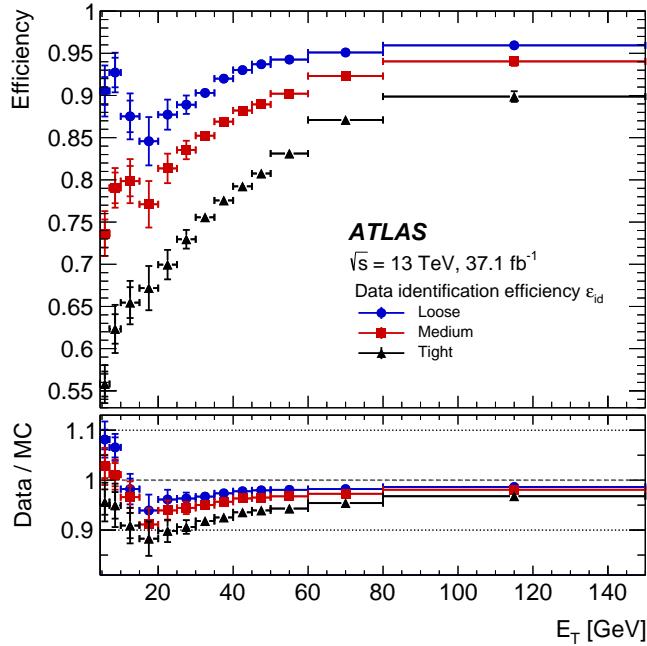


Figure 3.9: Measured LH electron ID efficiencies in $Z \rightarrow ee$ events for the `LooseLH` (blue), `MediumLH` (red), and `TightLH` (black) working points as a function of electron E_T . The bottom plot shows data-to-simulation ratios. Plot taken from [53].

1081 **Electron isolation** Signal electrons, such as those from the decay of a W boson, tend to have
 1082 little detector activity nearby in both the ID and the calorimeters. Background electrons, such as
 1083 those from photon conversions or jets, are often produced in association with other particles. To
 1084 take advantage of this, variables are constructed to quantify the amount of activity within a cone of
 1085 a specified radius in ΔR about an electron in both the tracking systems and the calorimeters.

1086 The track-based isolation consists of the sum of the transverse momentum of tracks within a cone
 1087 of a specified radius about an electron (not including the electron itself). The tracks are required to
 1088 have $p_T > 1$ GeV, satisfy basic quality requirements, and be associated with the vertex from which
 1089 the electron originated. Additionally, particles from bremsstrahlung radiation are considered part of
 1090 the original electron and are subtracted from the isolation cone. A variable cone radius dependent
 1091 on the p_T of the electron is used in order to compensate for busy detector environments:

$$\Delta R = \min\left(\frac{10 \text{ GeV}}{p_T}, R_{\max}\right) \quad (3.6)$$

1092 where R_{\max} is the largest allowed cone size, typically set at 0.2 for electrons.

1093 Calorimeter isolation is a bit more difficult due to the size of the energy deposits relative to the
 1094 cone size, as parts of an energy cluster can lie outside of the cone. As such, topological clusters [59]
 1095 (topo clusters) are seeded by calorimeter cells with deposited energy greater than four times the
 1096 expected noise-level of that cell. The cluster is then expanded to incorporate electromagnetic and
 1097 hadronic cells recording an energy greater than two times their expected noise-levels until no adjacent
 1098 clusters remain satisfying the requirement. The isolation cone is then the sum of the E_T of all
 1099 positive-energy topo clusters whose barycenters fall within a cone of radius $\Delta R < 0.2$. The electron's
 1100 energy is subtracted by removing the cells within a rectangle around the electron.

1101 When applying the isolation selection, *relative isolation* is used, defined as the ratio of the track-
 1102 or calorimeter-based isolation variable divided by the electron p_T . Four isolation working points are
 1103 defined targeting specific values of efficiency. **Loose** and **LooseTrackOnly** target a fixed efficiency
 1104 value across the p_T and η spectrum of the electrons, with the latter not applying a cut on calorimeter
 1105 isolation. **Gradient** and **GradientLoose** target a p_T -dependent fixed efficiency that is uniform in
 1106 η . The efficiencies for these working points as a function of electron E_T are shown in Figure 3.10.
 1107 Additional working points are also provided that instead use fixed values for the relative track and
 1108 calorimeter isolation cuts.

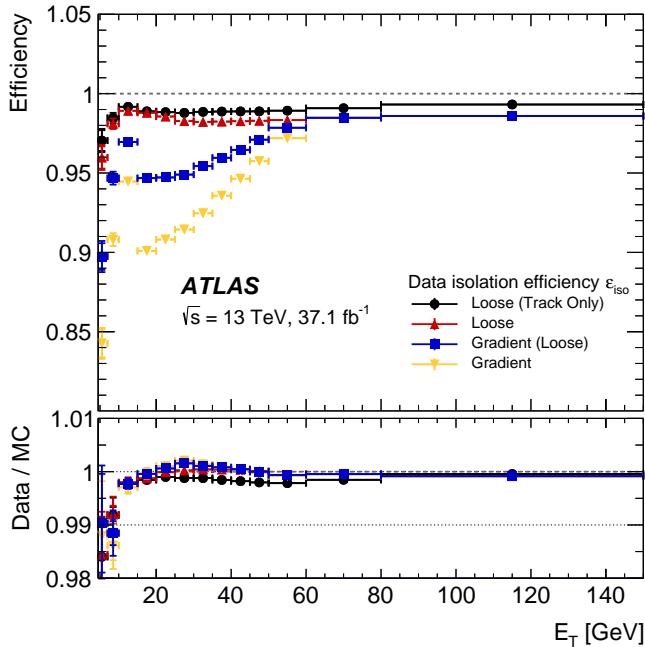


Figure 3.10: Measured isolation efficiencies for the `LooseTrackOnly` (black), `Loose` (red), `GradientLoose` (blue), and `Gradient` (yellow) working points as a function of electron E_T . The bottom panel shows data-to-simulation ratios. Plot taken from [53].

3.2.4.3 Muon reconstruction

Muon reconstruction [60, 61] first occurs independently in the ID and the MS, and then the information from both is combined to form the final muon tracks used in analysis. Muons in the ID are reconstructed as a charged particle track following the standard procedure detailed in Section 3.2.4.1.

In the MS, hits within each muon chamber are collected into segments, with separate algorithms for each of the four different detector components described in Section 3.2.3. Muon track candidates are then built by fitting together segments using a combinatorial search seeded by segments in the middle layers first, and then expanding to the outer and inner layers. A track must contain at least two matching segments, except in the transition region between the barrel and endcap where a single segment can be used. In the event that a segment is shared by multiple tracks, an overlap removal algorithm decides which track should keep the segment, or if it should remain shared. Finally, the hits within each track candidate are fit using a global χ^2 , and the track is accepted if the χ^2 satisfies a set of selection criteria.

The ID and MS tracks are then combined according to several different reconstruction criteria

1123 determined by the available information from the ID, MS, and calorimeters. This results in four
 1124 different muon “types”:

- 1125 • Combined (CB): Independent tracks in the ID and MS are combined with a global track refit
 1126 using both sets of hits. MS hits may be added or removed if it improves the quality of the
 1127 track fit. Muons are typically reconstructed outside-in, matching a MS track to one in the ID;
 1128 however, an inside-out matching is also used as a complementary approach.
- 1129 • Segment-tagged (ST): An ID track is classified as a muon if it can be extrapolated to at least
 1130 one track segment in the MDT or CSC chambers. These are used when a muon only crosses
 1131 one layer of the MS, either due to low p_T or falling in a region of reduced acceptance.
- 1132 • Calorimeter-tagged (CT): An ID track is classified as a muon if it can be matched to a calorime-
 1133 ter energy deposit compatible with a minimum ionizing particle. These muons are generally
 1134 lower in purity than the other types, but they are useful for recovering acceptance in the region
 1135 $|\eta| < 0.1$ where the MS is only partially instrumented.
- 1136 • Extrapolated (ME): The muon is reconstructed using a MS track that is loosely compatible
 1137 with having originated from the interaction point. The MS track must transverse at least
 1138 two layers of the MS in the central region and three in the forward region. These muons are
 1139 generally used to extend the acceptance of muon reconstruction in the forward regions not
 1140 covered by the ID ($2.5 < |\eta| < 2.7$).

1141 **Muon identification** Muon identification serves to select signal muons with a high quality mo-
 1142 mentum measurement from backgrounds (mainly from pions and kaons). For CB muons, three
 1143 variables are used:

- 1144 • q/p significance: The absolute value of the difference in q/p of the muons (where q is the
 1145 muon’s charge) as measured by the ID and the MS divided by the corresponding uncertainties
 1146 added in quadrature.
- 1147 • ρ' : The absolute value of the difference between the p_T measurements in the ID and the MS
 1148 divided by the p_T of the combined track.
- 1149 • The normalized χ^2 of the combined track fit.

1150 Additional requirements are imposed on the number of hits within the ID to ensure the tracks’
 1151 momenta are well measured.

1152 There are three primary muon identification working points⁶ of increasing background rejection:
 1153 **Loose**, **Medium**, and **Tight**. **Loose** muons include all four types listed above, but CT and ST muons
 1154 are restricted to the region $|\eta| < 0.1$. The default recommendation for analysis are **Medium** muons,
 1155 which only include CB and ME tracks. Finally, **Tight** muons are made up of CB muons with hits in
 1156 at least two components of the MS and that pass the **Medium** requirements. Each successive working
 1157 point is a subset of the previous one, and the cuts on the three variables listed above are tightened
 1158 in each step. The muon reconstruction efficiency for each of these three working points in $Z \rightarrow \mu\mu$
 1159 events is shown in Figure 3.11.

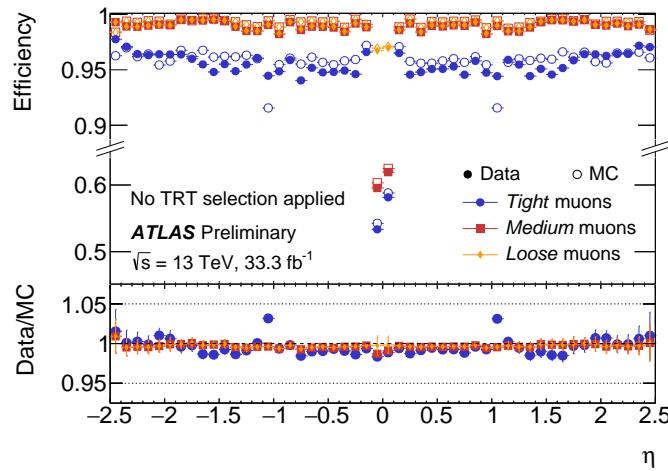


Figure 3.11: Muon reconstruction efficiencies in $Z \rightarrow \mu\mu$ events for the **Loose** (yellow), **Medium** (red), and **Tight** (blue) ID working points as a function of muon η . The drop in efficiency for **Medium** and **Tight** muons in the range $|\eta| < 0.1$ is due to the exclusion of CT and ME muons, which are used to recover acceptance in this region. Collision data (solid points) is compared to simulation (open points) and the ratio is in the lower panel. Plot taken from [62].

1160 **Muon isolation** Isolation for muons is handled in much the same way as for electrons (see Sec-
 1161 tion 3.2.4.2). A track-based variable is computed by summing the transverse momenta of tracks with
 1162 $p_T > 1$ GeV within a cone of variable radius as in Equation 3.6 with $R_{\max} = 0.3$. The calorimeter-
 1163 based isolation again uses topo clusters of radius $\Delta R < 0.2$ with the energy deposit corresponding
 1164 to the muon removed. Similar isolation working points are constructed using the same criteria as
 1165 for electrons, and the efficiency of the **Loose** and **GradientLoose** working points as a function of
 1166 muon p_T are shown in Figure 3.12.

⁶A fourth working point, **Highpt**, is optimized for high mass searches, such as W' and Z' resonances, and is not covered here.

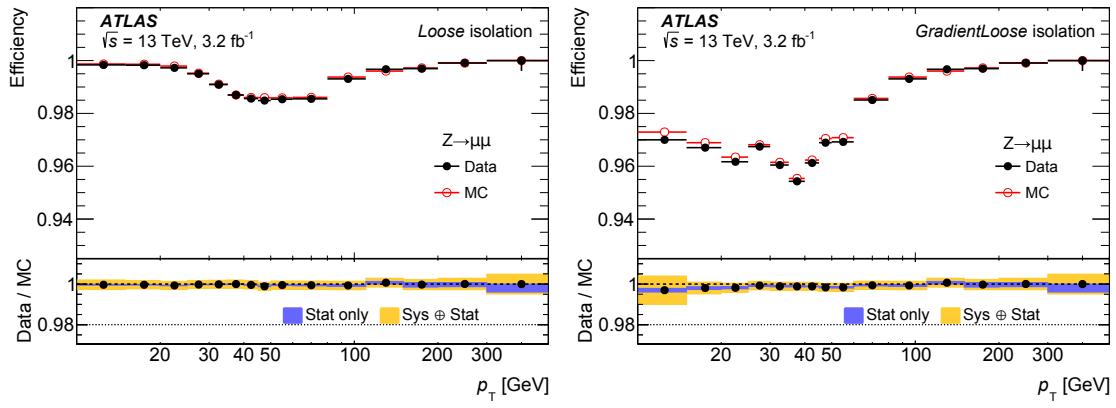


Figure 3.12: Muon isolation efficiencies in $Z \rightarrow \mu\mu$ events for the Loose (left) and GradientLoose (right) working points as a function of muon p_T . Collision data (black) is compared to simulation (red) and the ratio is in the lower panel. Plots taken from [61].

CHAPTER 4

Alignment of the ATLAS Inner Detector

1169 When a charged particle passes through the ATLAS ID, it leaves hits in the sensors along its
1170 path. In order to reconstruct an accurate particle track from these hits, it is necessary to know
1171 where in space these hits occurred as precisely as possible, which in turn requires knowledge of the
1172 physical location of the sensor that registered each hit. If one of the sensors is displaced relative
1173 to its expected position in the known detector geometry, or *misaligned*, the assumed location of
1174 the corresponding hit will not match its actual location, resulting in an incorrect track fit. These
1175 misalignments can occur for any number of reasons, including but not limited to elements shifting
1176 during maintenance periods or from cycles in ATLAS's magnetic field, or small movements during
1177 normal detector operations. A visualization of how a misaligned detector element can affect the
1178 track reconstruction is shown in Figure 4.1.

1179 In order to correct the misalignments, the ID alignment procedure is applied to accurately
1180 determine the physical position and orientation of each detector element. The baseline accuracy of
1181 the alignment is required to be such that the track parameter resolutions are not degraded by more
1182 than 20% with respect to those derived from a perfect detector geometry⁷. This corresponds to a
1183 precision of better than $10 \mu\text{m}$ in the positioning of the elements of the silicon detectors [63].

1184 This chapter outlines the ID alignment procedure, the alignment of the detector during the 2015
1185 data taking period, and the steps taken to measure momentum biases in the alignment.

⁷The so-called *perfect geometry* refers to the description of the ATLAS detector in which every sensor precisely matches its design specifications. The perfect geometry contains no misalignments, and the position of each sensor is known exactly.

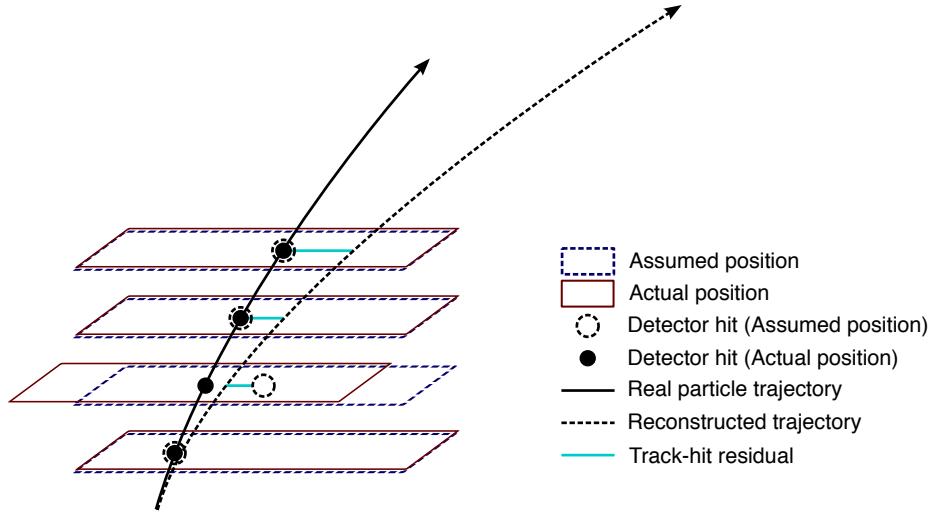


Figure 4.1: Graphical representation of the effect of a misaligned detector element. The reconstructed particle track (dashed arrow) differs from the actual trajectory of the particle (solid arrow) due to a shift in the third detector element. The cyan lines represent the track-to-hit residuals.

1186 4.1 The alignment method

1187 The alignment procedure uses a track-based algorithm that updates the locations of detector ele-
 1188 ments in order to minimize the set of track-hit *residuals*. These residuals are defined as the distance
 1189 between the where fitted track intersects a given detector element and the position of the actual hit
 1190 recorded by the same element, shown by the cyan lines in Figure 4.1.

1191 Tracks in ATLAS are parameterized as five-dimensional vectors [64]:

$$\vec{\tau} = (d_0, z_0, \phi_0, \theta, q/p) \quad (4.1)$$

1192 where d_0 and z_0 are the transverse and longitudinal impact parameters with respect to the origin,
 1193 ϕ_0 is the azimuthal angle of the track at the point of closest approach to the origin, θ is the polar
 1194 angle, and q/p is the charge of the track divided by its momentum. The residual for the i^{th} hit of a
 1195 given track can then be written in terms of the track parameters $\vec{\tau}$ and a set of alignment parameters
 1196 \vec{a} that describe the hit location [65]:

$$r_i(\vec{\tau}, \vec{a}) = (\vec{m}_i - \vec{e}_i(\vec{\tau}, \vec{a})) \cdot \hat{k} \quad (4.2)$$

1197 where \vec{e}_i is the intersection point of the extrapolated track with the sensor, \vec{m}_i is the position of the
 1198 associated hit within the sensor, and \hat{k} is the unit vector defining the direction of the measurement
 1199 within the sensor. \vec{r} then is the vector of residuals for the track.

1200 A χ^2 function can be built from the residuals of all collected tracks:

$$\chi^2 = \sum_{\text{tracks}} \vec{r}^T V^{-1} \vec{r} \quad (4.3)$$

1201 where V is the covariance matrix of the hit measurements. The χ^2 function is then minimized with
 1202 respect to the alignment parameters \vec{a} , which contain all degrees of freedom being aligned. The
 1203 minimization condition with respect to \vec{a} is:

$$\frac{d\chi^2}{d\vec{a}} = 0 \rightarrow 2 \sum_{\text{tracks}} \left(\frac{d\vec{r}}{d\vec{a}} \right)^T V^{-1} \vec{r} = 0 \quad (4.4)$$

1204 This equation can be difficult to solve exactly, so the residual is rewritten as a first order Taylor
 1205 expansion:

$$\vec{r} = \vec{r}_0 + \frac{d\vec{r}}{d\vec{a}} \delta\vec{a} \quad (4.5)$$

1206 where \vec{r}_0 is dependent on an initial set of track and alignment parameters \vec{r}_0 and \vec{a}_0 , respectively;
 1207 the track parameter dependence has also been folded into the total derivative $\frac{d\vec{r}}{d\vec{a}}$. Equation 4.5 can
 1208 then be inserted into the minimization condition from Equation 4.4 to give:

$$\left[\sum_{\text{tracks}} \left(\frac{d\vec{r}}{d\vec{a}} \right)^T V^{-1} \left(\frac{d\vec{r}}{d\vec{a}} \right) \right] \delta\vec{a} + \sum_{\text{tracks}} \left(\frac{d\vec{r}}{d\vec{a}} \right)^T V^{-1} \vec{r}_0 = 0 \quad (4.6)$$

1209 From this equation, the alignment matrix \mathcal{M}_a and alignment vector $\vec{\nu}_a$ can be defined:

$$\mathcal{M}_a = \sum_{\text{tracks}} \left(\frac{d\vec{r}}{d\vec{a}} \right)^T V^{-1} \left(\frac{d\vec{r}}{d\vec{a}} \right) \quad (4.7)$$

1210

$$\vec{\nu}_a = \sum_{\text{tracks}} \left(\frac{d\vec{r}}{d\vec{a}} \right)^T V^{-1} \vec{r}_0 \quad (4.8)$$

1211 Finally, the alignment corrections $\delta\vec{a}$ can be solved for by inverting the alignment matrix:

$$\delta\vec{a} = -\mathcal{M}_a^{-1} \vec{\nu}_a \quad (4.9)$$

1212 which is a linear system of equations with a number of equations equal to the number of alignment
 1213 degrees of freedom [66].

1214 Inverting the full matrix and solving the resulting system of equations is referred to as *Global*
 1215 χ^2 alignment [65]. This can be useful, as \mathcal{M}_a contains all the correlations between the alignable
 1216 structures. However, inverting the matrix is difficult when the number of degrees of freedom becomes
 1217 large—as the number of alignable structures increases, so too does the size of the matrix \mathcal{M}_a .
 1218 Eventually, inverting the matrix becomes too computationally intensive to be practical.

1219 This problem is solved by the *Local χ^2* algorithm [67]. In this case, the alignment matrix is
 1220 constructed to be block-diagonal, allowing for it to be easily inverted even for large numbers of
 1221 degrees of freedom. This is achieved by replacing the full derivative in Equation 4.6 with the partial
 1222 derivative $\frac{\partial \vec{r}}{\partial \vec{a}}$. The new alignment matrix \mathcal{M}'_a and alignment vector $\vec{\nu}'_a$ become:

$$\mathcal{M}'_a = \sum_{\text{tracks}} \left(\frac{\partial \vec{r}}{\partial \vec{a}} \right)^T V^{-1} \left(\frac{\partial \vec{r}}{\partial \vec{a}} \right) \quad (4.10)$$

1223

$$\vec{\nu}'_a = \sum_{\text{tracks}} \left(\frac{\partial \vec{r}}{\partial \vec{a}} \right)^T V^{-1} \vec{r}_0 \quad (4.11)$$

1224 Inverting \mathcal{M}'_a is considerably faster and less memory intensive compared to \mathcal{M}_a , even for large
 1225 numbers of degrees of freedom; however, the correlations between the alignable structures is lost.
 1226 Due to the Taylor expansion used in Equation 4.6, several iterations of the alignment algorithm may
 1227 be necessary to converge on a final set of alignment constants. The Local χ^2 alignment typically
 1228 requires more iterations due to the loss of the correlation information [68].

1229 In practice, the ATLAS reconstruction is run over a set of events, and the resulting tracks are
 1230 fed to the alignment algorithm. The residuals are calculated, the alignment matrix is built and
 1231 inverted, and a new set of alignment constants is obtained. The convergence is checked in two ways:

- 1232 1. Measuring change in the χ^2 with respect to the previous iteration. If it is near zero, then the
 1233 χ^2 is approaching its minimum.
- 1234 2. Looking at the residual distributions for different alignable structures. A well aligned detector
 1235 will have a residual distribution that is approximately Gaussian, with a mean of zero with a
 1236 width approximating the intrinsic resolution of the detector.

1237 If the above checks are satisfied, the process is finished and the final alignment constants are read
 1238 out; if not, another iteration is performed. A visual representation of the alignment chain is shown
 1239 in Figure 4.2.

1240 Since a χ^2 minimization is used to align the detector, if there is a systematic misalignment in
 1241 the detector that does not adversely affect the χ^2 , the algorithm will be insensitive to it. These
 1242 misalignments are referred to as *weak modes*, and special care is taken to remove them [69]. One
 1243 potential impact of weak modes is a bias in the track momentum of reconstructed particles. This
 1244 particular effect is the subject of Section 4.4.

1245 The detector is aligned both in “real-time” as data is collected and during dedicated offline
 1246 alignment campaigns. The real-time alignment is run in ATLAS’s *calibration loop*, which comprises

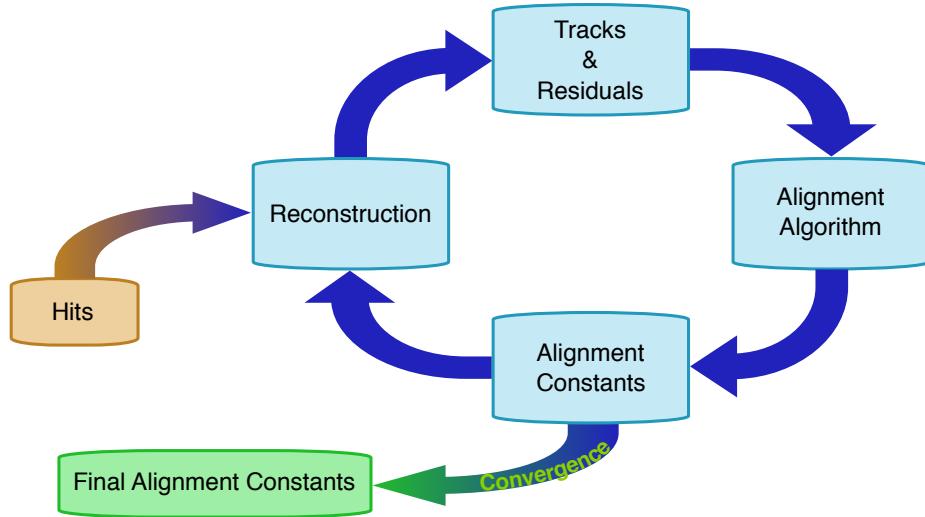


Figure 4.2: Graphical representation of the ID alignment chain.

the first stage in the preparation of data for physics analysis. The calibration loop requires the alignment—as well as various other detector calibrations—to be available within 48 hours for initial data processing. A fast, coarse-grained alignment⁸ is run on a subset of the available data containing full tracking information, and the results are propagated to the reconstruction of that particular run [70]. Due to the time constraints of the calibration loop, a full sensor-by-sensor alignment is not possible.

The more thorough and finely tuned alignments are reserved for dedicated alignment campaigns, typically near the beginning and at the end of data taking campaigns. The former generally occurs once a sufficient amount of data is collected after a detector shutdown, in order to obtain a good baseline alignment for use in the remainder of the data collection period. One such alignment campaign, the initial offline alignment of the ATLAS detector at the beginning of Run 2, is the subject of Section 4.2. The latter occurs once data taking is complete. A new set of alignment constants is derived for the full set of available data, and it is typically divided into several “blocks” to account for potential run-by-run misalignments. The data is then reprocessed using the newly derived detector geometry.

⁸The calibration loop runs up to a Level 2 alignment in the silicon detectors, which involves treating each layer of sensors as a single object. The alignment levels are defined in greater detail in Table 4.1.

1262 **4.1.1 Alignment levels**

1263 The alignment of the detector is performed at several levels of increasing granularity. This adds
 1264 flexibility in being able to align only as finely as needed, and it also allows for global, detector-level
 1265 misalignments to be corrected first before dealing with finer adjustments. The main alignment levels
 1266 are as follows:

- 1267 • Level 1 (L1) alignment involves moving entire subdetector components as a single unit, such
 1268 as the entire Pixel detector, or the SCT barrel. These often have the largest misalignments,
 1269 but they are easily corrected and do not require large volumes of data to do so.
- 1270 • Level 2 (L2) alignment treats individual layers in the silicon detectors (modules in the TRT)
 1271 and end cap disks as individual alignable objects.
- 1272 • Level 2.7 (L27) alignment was introduced with the addition of the IBL to the ID in Run 2. It
 1273 involves the stave-by-stave alignment of the IBL and Pixel barrel⁹.
- 1274 • Level 3 (L3) alignment treats each module in the silicon detectors and each straw in the TRT
 1275 as an individual alignable object. It is the finest grained alignment available but also the most
 1276 computationally intensive due to the large number of degrees of freedom. Due to the large
 1277 number of individual detector elements being aligned, high statistics are required.

1278 The different alignment levels are listed in more detail in Table 4.1, including the number of alignable
 1279 structures and associated degrees of freedom for each detector component.

1280 The implementation of the alignment algorithm in the software is flexible enough to allow each
 1281 subsystem to be aligned individually at a specified level. Each alignable structure has six degrees of
 1282 freedom: 3 translations (T_x, T_y, T_z) and 3 rotations (R_x, R_y, R_z)¹⁰; however individual degrees of
 1283 freedom may be turned on and off as required. In a typical alignment job, L1 and L2 contain few
 1284 enough degrees of freedom that the Global χ^2 algorithm can be used, but L3 alignments (which can
 1285 contain over 36,000 degrees of freedom in the silicon detectors alone) require the Local χ^2 algorithm.

⁹Even though the IBL is considered a part of the Pixel Detector, for the purposes of this Chapter, “Pixel” will refer to the original three layers of the Pixel barrel and its endcaps, and the IBL will be referenced separately.

¹⁰The TRT is an exception, as the subdetector does not have any resolution along the length of the straw. Therefore, for the barrel, T_z is omitted. Similarly for the straws themselves, only two parameters are defined: translation with respect to the radial direction (T_ϕ) and rotation with respect to the radial axis (R_r for the barrel and R_z for the end-caps) [71].

Level	Description of alignable structure	Structures	DoF
1	IBL detector	1	6
	Whole Pixel detector	1	6
	SCT barrel and 2 end-caps	3	18
	TRT barrel and 2 end-caps (T_z fixed)	3	17
Total:		8	47
2	IBL detector	1	6
	Pixel barrel layers	3	18
	Pixel end-cap disks	2×3	36
	SCT barrel layers	4	24
	SCT end-cap disks	2×9	108
	TRT barrel 32 modules (T_z fixed)	3×32	480
	TRT end-cap wheels	2×40	480
Total:		208	792
2.7	IBL staves	14	84
	Pixel barrel staves	$22+38+52$	672
	Pixel end-cap disks	2×3	18
	Total:		132
3	IBL modules	280	1,680
	Pixel modules	1,744	10,464
	SCT modules	4,088	24,528
	TRT barrel wires (T_ϕ, R_r only)	105,088	210,176
	TRT end-cap wires (T_ϕ, R_Z only)	245,760	491,520
	Total silicon sensors:		6,112
	Total TRT wires:		350,848
36,672		701,696	

Table 4.1: The four alignment levels for each of the detector subsystems. The total number of alignable structures and degrees of freedom (DoF) to be aligned are given for each level.

4.1.2 Alignment coordinate systems

The global coordinate system (x, y, z) used by the ID alignment matches that of the ATLAS detector in general. The positions and orientations of individual detector modules of the ID are defined by a right-handed local coordinate system (x', y', z') where the origin is defined as the geometrical center of the module. The x' -axis for each silicon module is defined to point along the most sensitive direction of the module, the y' -axis is oriented along the long side of the module, and the z' -axis is orthogonal to the (x', y') plane. For the TRT straws, the x' -axis is perpendicular to both the wire and the radial direction, defined from the origin of the global frame to the straw center, the y' -axis points along the straw, and once again the z' -axis is orthogonal to the (x', y') plane. A depiction of the global and local coordinate systems for the ID is shown in Figure 4.3.

When considering the alignment degrees of freedom listed earlier in Section 4.1.1, grouped collections of modules, layers, or entire subdetectors use the global coordinate system; individual modules

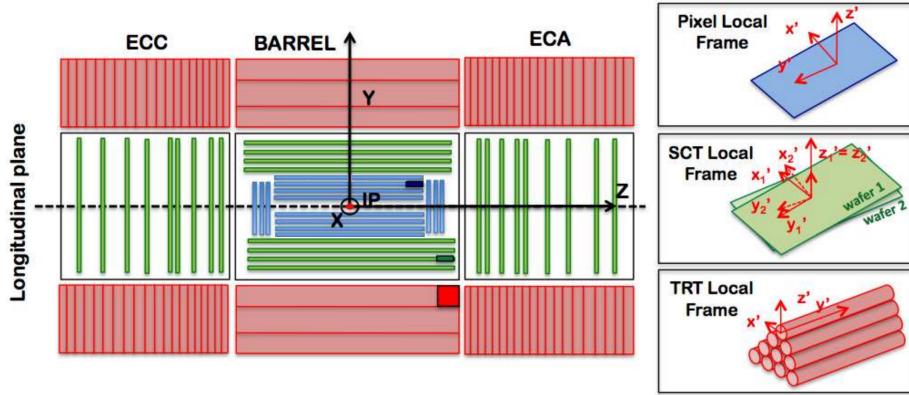


Figure 4.3: A schematic representation of the Inner Detector in the longitudinal plane with the global coordinate system overlaid on top. The Pixel detector and IBL are shown in blue, the SCT in green, and the TRT in red. The local coordinates for each subdetector module are inset on the right. Image taken from [72].

use their respective local coordinate systems. The translations T_i are with respect to the origin of the given reference frame, and the rotations R_i are taken about the Cartesian axes.

4.2 Early 2015 alignment of the ATLAS detector

At the end of Run 1, the LHC was shut down for upgrades and maintenance. During this time, a number of upgrades were performed on the ATLAS detector, including the installation of the Insertable B-Layer (IBL), mounted on a new beam pipe. These changes to the ID required some detector components to be removed temporarily, and many elements shifted relative to each other over the course of the maintenance process. In order to correct for these large detector movements prior to 13 TeV data taking, an alignment was performed using cosmic ray data collected in early 2015 [72]. This alignment was able to correct for the majority of the large detector-wide misalignments as well as determine the global position of the IBL at the micron level.

In June of 2015, shortly after the data taking period began, the first track-based alignment of the refurbished ID was performed using the initial 7.9 pb^{-1} of $\sqrt{s} = 13 \text{ TeV}$ proton-proton collision data [73]. Starting from the geometry determined by the cosmic ray alignment, referred to hereafter as the *March alignment*, an improved set of alignment constants (the *June alignment*) was derived from a data set of approximately 1.4 million selected tracks. A sample of Monte Carlo (MC) events containing approximately 2.7 million tracks from dijet events simulated using a perfect detector geometry is used as a comparison to the data; the MC events are reweighted to match the η and p_T

1316 distributions found in the data. Additional validation of the alignment results uses a set of cosmic
1317 ray data collected by the detector during the LHC collisions.

1318 **4.2.1 June alignment procedure**

1319 The data set used as the input for the alignment is made up of a subset of physics events used for
1320 prompt reconstruction recorded at a rate of 10 Hz. To ensure that only high quality tracks are used
1321 for the alignment, each track is required to have transverse momentum $p_T > 3$ GeV, contain at least
1322 one hit in the Pixel detector, at least seven hits in the combined silicon detectors (IBL, Pixel, and
1323 SCT), and at least 25 hits in the TRT.

1324 A full L3 alignment of the IBL was included in the March alignment; however, a realignment was
1325 still necessary. Since cosmic rays pass through the detector top-down, the staves on the sides of the
1326 IBL recorded fewer hits and thus could not be aligned as precisely as those on the top and bottom.
1327 Additionally, the IBL was operating at a temperature of -20°C during the cosmic data taking, and
1328 it was set to -10°C for collision data taking. This proved to be significant, as it was observed that
1329 the IBL staves experience a temperature-dependent, parabolic bowing in the local x -direction of
1330 approximately $-10\mu\text{m}/\text{K}$ [74]. As a result, a full L3 alignment of the IBL was essential in order to
1331 correct for the bowing. Due to it being a brand new element of the detector as well as its importance
1332 in vertexing and b -jet tagging, aligning the IBL sensors with a high degree of precision was one of
1333 the main goals of the June alignment.

1334 The June alignment was performed in two stages, with the first pass focusing on relative move-
1335 ments of the big structures and correcting for the bowing of the IBL. The March alignment corrected
1336 for these larger movements as well; however, it was observed during Run 1 that these sort of mis-
1337 alignments are introduced by changing conditions in the detector [69], such as adjustments in the
1338 cooling system settings or magnet power cycling, which may have occurred between the early cos-
1339 mic data taking and the first $\sqrt{s} = 13$ TeV collisions. The silicon detectors were aligned at several
1340 different levels and the IBL was aligned at the module level; the TRT detector was kept fixed to
1341 act as a global reference frame. The full alignment chain for the first pass consisted of the following
1342 steps:

- 1343 1. The IBL, Pixel, and SCT detectors were aligned at L1. The SCT barrel was not aligned in
1344 T_z in order to constrain global displacements along the z -axis, as the TRT is not sensitive to
1345 that degree of freedom.

- 1346 2. The IBL, and Pixel barrel and end-caps, and SCT barrel were aligned at L2. The SCT end-caps
 1347 were aligned at L1.
- 1348 3. The IBL and pixel barrel were aligned at L27, using all six degrees of freedom. The Pixel
 1349 end-cap disks were only aligned in the plane (T_x , T_y , and R_z). The SCT was treated the same
 1350 as in the previous step.
- 1351 4. The IBL was aligned at L3 using all six degrees of freedom for each module.

1352 The primary goal for the second pass was to remove a bias in the transverse impact parameter
 1353 d_0 found in the March alignment. In addition to the bias, the resolution of d_0 was also observed
 1354 to be poorer than expected. In order to correct for this, an additional constraint was passed to
 1355 the alignment which added the impact parameter with respect to the beam spot as a pseudo-
 1356 measurement [75]. With this setup, when the alignment algorithm minimizes the χ^2 , it will take
 1357 care of the impact parameter minimization as well. Only the IBL and Pixel detectors were aligned
 1358 in the second pass. The different alignment stages are listed below, with the beam spot constraint
 1359 being used in each:

- 1360 1. The IBL and Pixel detectors were aligned at L2 with the SCT fixed.
- 1361 2. The IBL was aligned at L27
- 1362 3. The IBL and Pixel barrel and end-caps were aligned at L3.
- 1363 The set of alignment constants obtained at the end of the second pass represents the June alignment.
 1364 The highest level of alignment each subdetector received over the course of the two passes is listed
 1365 in Table 4.2.

Detector		Highest level of alignment
IBL		L3
Pixel	Barrel	L3
	End-caps	L3 (T_x , T_y , and R_z only)
SCT	Barrel	L2 (except T_z)
	End-caps	L1
TRT		None

Table 4.2: Summary of the highest level of alignment applied to each ID subsystem when deriving the June alignment.

1366 **4.2.2 Alignment results**

1367 Alignment quality is primarily assessed by looking at the track-hit residual distributions. If the
 1368 detector is well aligned, the residuals will be Gaussian-distributed with a mean of zero and a width
 1369 approximating the detector's resolution. The residual distributions are constructed from the same
 1370 selection of tracks that were used to perform the alignment, and are the focus of Section 4.2.2.1.
 1371 A second check on the alignment involves observables sensitive to the track parameter resolution.
 1372 In this case, cosmic rays are used, making use of a “split track” technique that takes advantage of
 1373 the top-to-bottom cosmic ray trajectory (compared to the center-out trajectory of collision tracks).
 1374 This method and the corresponding tests of the alignment are detailed in Section 4.2.2.2

1375 Additionally, the effect of the beam spot constrained alignment on the impact parameter d_0
 1376 needed to be checked. The d_0 distributions for both the March and June alignments are compared
 1377 to the MC simulation using a perfect geometry in Figure 4.4. In the March alignment, there is
 1378 a clear bias of $18 \mu\text{m}$ in the mean of the distribution and the width is nearly twice that of the
 1379 perfect geometry. After the second pass of the June alignment, the mean has shifted to $1 \mu\text{m}$ and
 1380 the distribution has narrowed considerably. From this, it appears that the constrained alignment
 1381 successfully removed the d_0 bias.

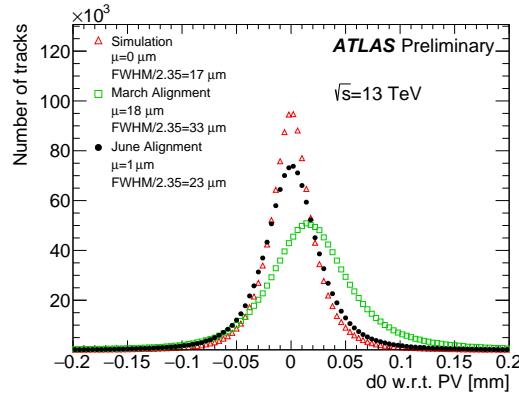


Figure 4.4: The d_0 distributions with respect to the reconstructed primary vertex using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red). The distributions are normalized to the same number of entries.

1382 **4.2.2.1 Residual distributions from collisions**

1383 As mentioned previously, the primary focus of the June alignment campaign was on the IBL and the
1384 Pixel detectors. These subdetectors are the closest to the beam line and have the finest resolutions
1385 within the ID. The residual distributions in local x and y of the IBL planar sensors¹¹ are shown in
1386 Figure 4.5. These and subsequent figures in this section compare the June and March alignments to
1387 the perfectly-aligned MC simulation. Noticeable improvement in the distribution widths can be seen
1388 in both the local x - and y -directions, nearly matching the simulation in local x , the most sensitive
1389 direction.

1390 Due to the temperature-dependent bowing of the IBL, it is also interesting to look at the means of
1391 the residual distributions for each ring of IBL sensors along the beam line, as shown in Figure 4.6. A
1392 deformation is clearly visible in the March alignment in both measurement directions, and the shape
1393 in the local x -direction is consistent with an average stave bowing due to the different operating
1394 temperature of the IBL during the March alignment and the 13 TeV collisions. This feature was
1395 nearly eliminated in both directions through the L3 alignment of the IBL sensors.

1396 The local x and y residual distributions for the Pixel detector barrel and end-caps are shown
1397 in Figures 4.7 and 4.8. Even though the IBL is not included in the plots of the barrel, some of
1398 the noticeable improvement in the more sensitive local x -direction is an effect of the improved IBL
1399 alignment. Similarly, the relatively broad local y residual distribution in the barrel likely indicates
1400 that further refinement of the IBL alignment was needed along that direction. Even so, the June
1401 alignment outperforms the March alignment and rivals the simulation in most of the plots.

1402 Similar distributions for the SCT and TRT barrel and end-caps are shown in Figures 4.9 and
1403 4.10, respectively. Much like with the Pixel residuals, there is a reduction in the width of the TRT
1404 residuals between the March and June alignments due to the alignment of the other subdetectors
1405 improving the quality of the track fit. Even though neither subdetector was aligned at module-
1406 level, the residuals indicate that the previous L3 alignment performed in Run 1 did not degrade
1407 significantly during the upgrade and maintenance period.

¹¹The IBL contains 12 planar sensors in the center of a stave, with four 3D sensors on either end. Only the planar sensors are shown here due to low statistics in the 3D sensors as well as poor MC modeling of these sensors.

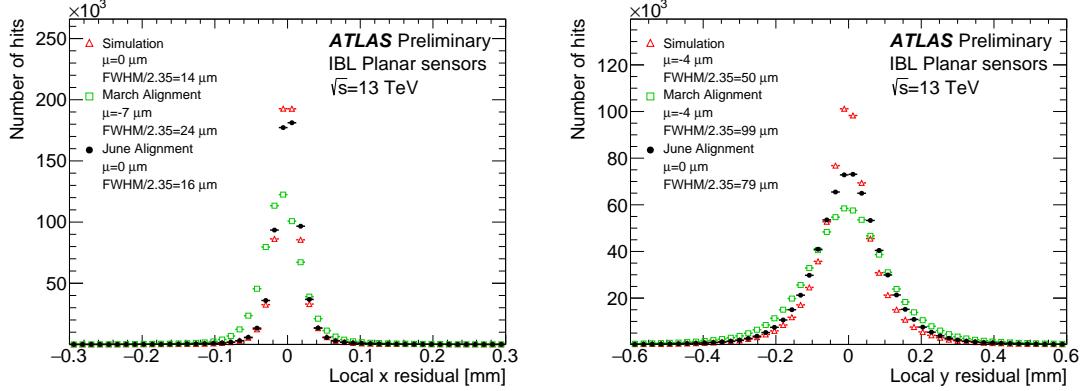


Figure 4.5: Local x (left) and local y (right) residual distributions of the IBL planar sensors using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red). The distributions are normalized to the same number of entries.

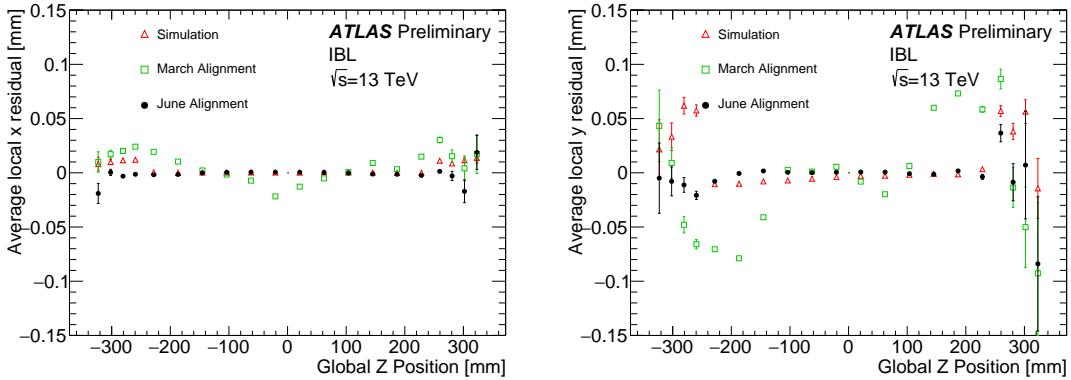


Figure 4.6: The mean of the local x (left) and local y (right) residual distributions as a function of the global z position of each IBL module using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red).

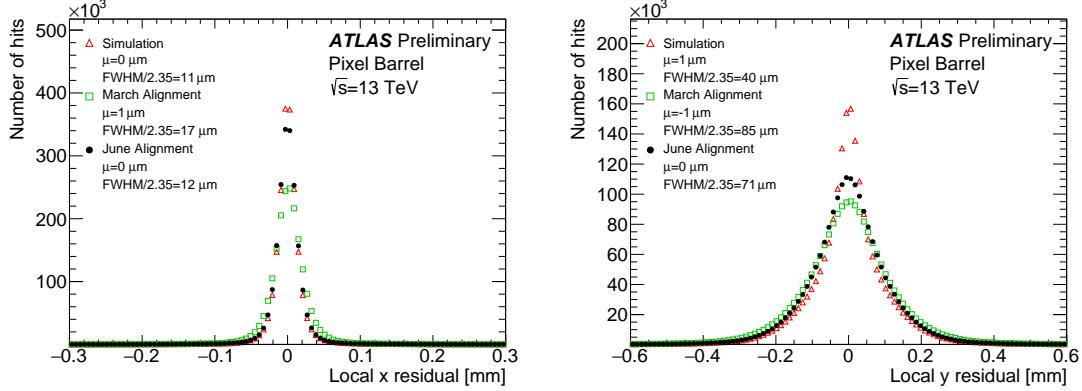


Figure 4.7: Local x (left) and local y (right) residual distributions for the Pixel barrel (excluding the IBL) using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red). The distributions are normalized to the same number of entries.

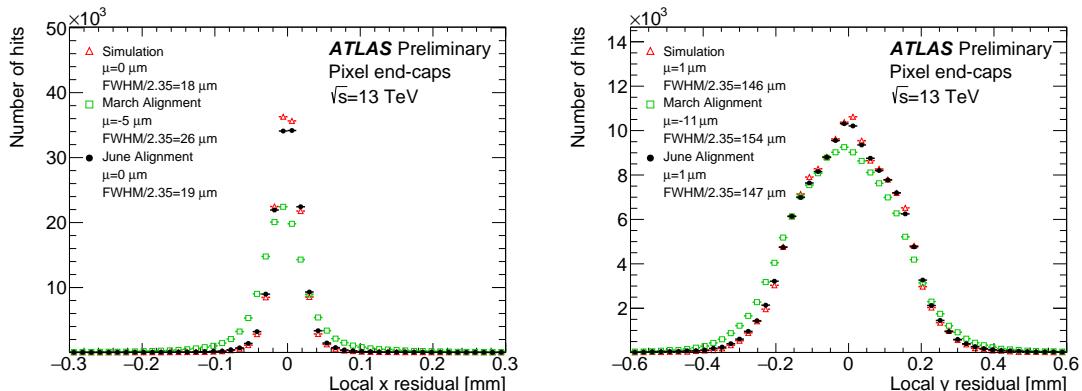


Figure 4.8: Local x (left) and local y (right) residual distributions for the Pixel end-caps using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red). The distributions are normalized to the same number of entries.

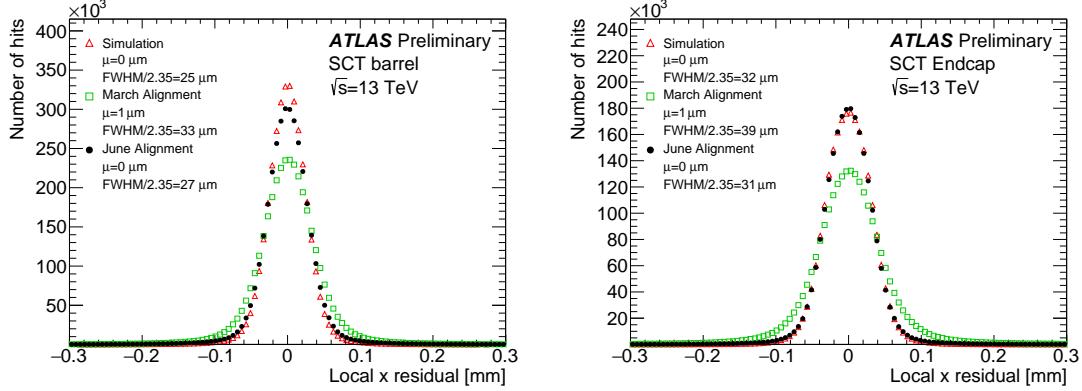


Figure 4.9: Local x residual distributions for the SCT barrel (left) and end-caps (right) using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red). The distributions are normalized to the same number of entries.

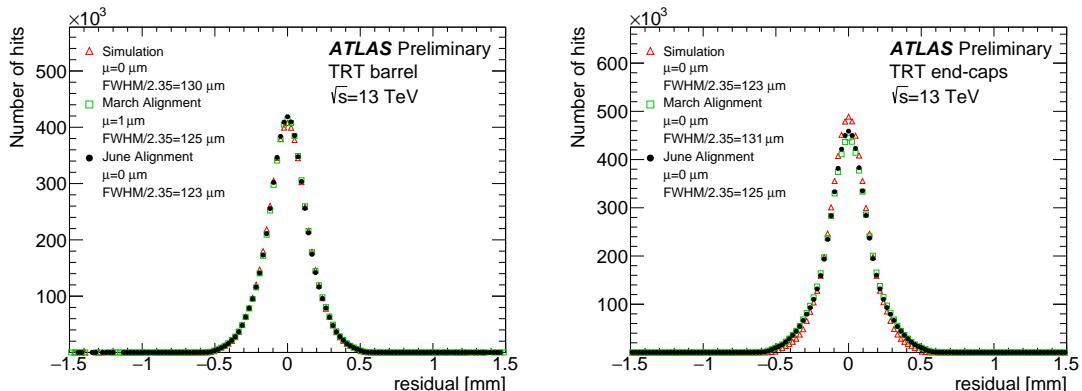


Figure 4.10: Residual distributions for the TRT barrel (left) and end-caps (right) using 13 TeV collision data sample reconstructed using the June (black) and March (green) alignments. The data is compared to a MC simulation using a perfect detector geometry (red). The distributions are normalized to the same number of entries.

1408 **4.2.2.2 Track parameter resolution from cosmic rays**

1409 Cosmic ray data is very useful as an independent check on the alignment in the barrel of the detector.
 1410 While tracks from proton-proton collisions originate within the detector and travel outwards, a
 1411 cosmic ray that passes through the center of the detector leaves a track in both halves of the
 1412 detector. If the cosmic ray is split in half, as in Figure 4.11, then it can be treated as two separate
 1413 tracks each with nearly identical track parameters (some differences arise due to energy loss as the
 1414 particle passes through the detector). The distribution of the difference in a given track parameter
 1415 $\Delta\tau$ is approximately Gaussian with a variance $\sigma^2(\Delta\tau)$. Since both tracks come from the same
 1416 particle, each track individually has a variance equal to $\sigma^2(\Delta\tau)/2$. The resolution of the track
 1417 parameter is then given by the root mean square of the distribution divided by $\sqrt{2}$.

1418 Cosmic rays whose split tracks each have transverse momentum $p_T > 2$ GeV and at least one,
 1419 eight, and 25 hits in the barrels of the Pixel, SCT, and TRT detectors, respectively, were selected to
 1420 measure a collection of track parameters. Figure 4.12 shows the difference in the impact parameter
 1421 Δd_0 and the charge divided by the transverse momentum $\Delta q/p_T$ of the selected split-track cosmic
 1422 rays for both the March and June alignments. Both distributions show a reduction in width in the
 1423 June alignment, corresponding to an improvement in the resolution of each track parameter. The
 1424 Δd_0 plot in particular shows significant improvement with the June alignment, further validating
 1425 the removal of the bias in the impact parameter.

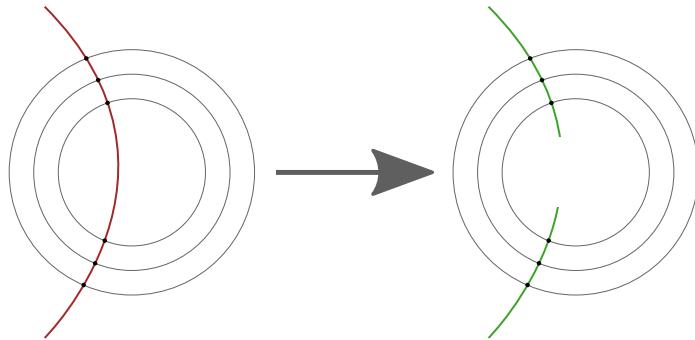


Figure 4.11: Representation of splitting a single cosmic ray track passing through the entire detector (left) into two separate tracks (right).

1426 **4.2.3 Error scaling**

1427 The final step in preparing the new set of June alignment constants deals with the adjustment of
 1428 the hit errors, or *error scaling*. Knowledge of the exact position of a hit measurement on a track

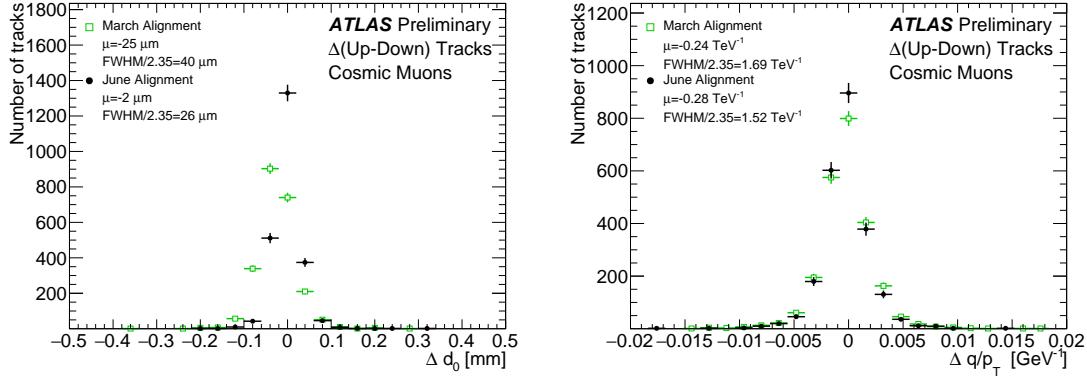


Figure 4.12: Distribution of the difference in the impact parameter Δd_0 (left) and charge over transverse momentum $\Delta q/p_T$ (right) between the two cosmic ray split tracks. The June (black) and March (green) alignments are compared. The distributions are normalized to the same number of entries.

is limited by the accuracy with which the sensors' positions are known. Let σ represent the hit uncertainty used in track fitting, and σ_0 be the detector's intrinsic uncertainty. If $\sigma = \sigma_0$, the pull of the track-hit residual distributions should form a Gaussian distribution centered at zero with unit width [66]. In the case of residual misalignment, the pull distributions' standard deviations will stray from unity. The hit uncertainty can be written as:

$$\sigma = a \cdot \sigma_0 \oplus b \quad (4.12)$$

where a is a scaling factor, and b is a constant term which can be interpreted as a measure of any remaining misalignment of the detector elements. In the June alignment campaign, the value of a was fixed at $a = 1$, and b was evaluated from the residual pull distributions for each subdetector in its sensitive directions.

Once the value of b is determined, pull distributions derived from the new value of σ should have unit width. The error scaling values for each subdetector are listed in Table 4.3, and the pull distributions for the IBL after error scaling are shown in Figure 4.13

4.3 Level 2 alignment of the TRT

During validation of the final end-of-year reprocessing of the 2015 data, a misalignment was found in the barrel of the TRT detector, as several modules (triangular clusters of straws) showed rotations in the local y coordinate. The then-best available constants included a full L3 alignment of the silicon detectors and a separate L2 alignment of the TRT. However, not all degrees of freedom were

Detector	Coordinate	$b(\mu\text{m})$
IBL	x	6.4
	y	43.6
Pixel	Barrel x	5.2
	Pixel y	28.6
	End-caps x	7.5
SCT	End-caps y	0
	Barrel x	10.8
	End-caps x	8.6
TRT	Barrel $r\phi$	0
	End-caps $r\phi$	0

Table 4.3: Estimated value of the error scaling term b for each subdetector component with the June alignment.

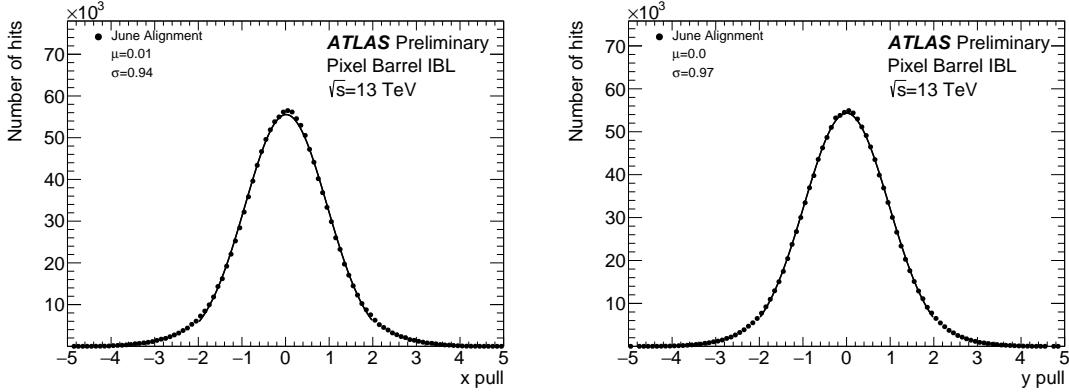


Figure 4.13: Pull distributions in local x (left) and y (right) for the IBL using 13 TeV collision data sample after applying the error scaling.

1446 enabled when the TRT was aligned. To correct for these tilts, an additional four iterations of L2
 1447 alignment was performed on the TRT enabling all available degrees of freedom (T_x , T_y , R_x , R_y , and
 1448 R_z in the barrel, and T_x , T_y , and R_z for the endcaps). Plots of the residual means from TRT barrel
 1449 ϕ -sectors containing modules affected by the tilt misalignment are shown in Figure 4.14 before and
 1450 after the L2 alignment.

1451 Following the L2 alignment, some additional time was taken to determine if a full wire-by-wire
 1452 L3 alignment of the TRT was necessary. The TRT was last aligned at L3 during Run 1, but
 1453 initial alignment campaigns in Run 2 did not show signs of misalignment (see, for example, the
 1454 residual distributions shown earlier in Figure 4.10). In order to assess the alignment more carefully,
 1455 two dimensional residual maps in ϕ and z were constructed for each layer in the TRT barrel and
 1456 endcaps using the current alignment. These maps were compared to a similar set using the L3

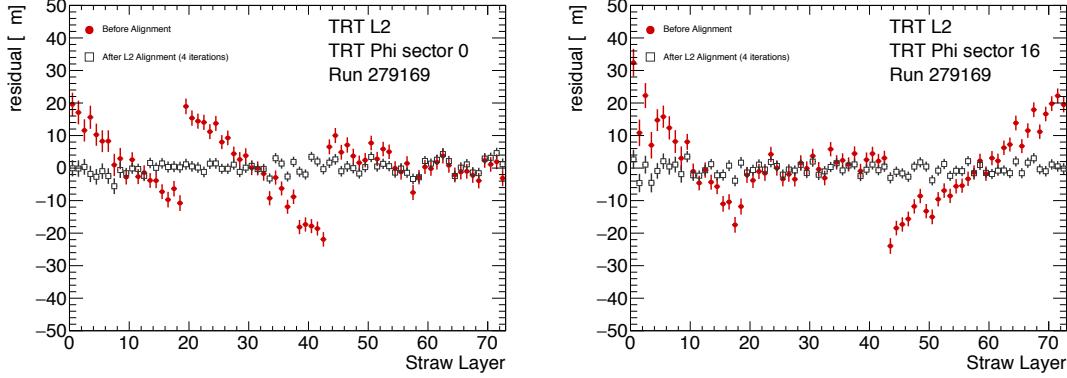


Figure 4.14: Residual means by straw layer in two TRT ϕ -sectors affected by a tilt misalignment. The tilts in each of the three modules are clearly visible in the red points representing the reconstructed data prior to alignment. After four iterations of L2 alignment, the residual means in the gray points are flat.

1457 alignment from 2010, from which it was determined that the straw-level alignment indeed hadn't
 1458 degraded, and a new L3 alignment was not needed. The maps for the first layer of the TRT barrel
 1459 are shown in Figure 4.15 for both sets of alignment constants.

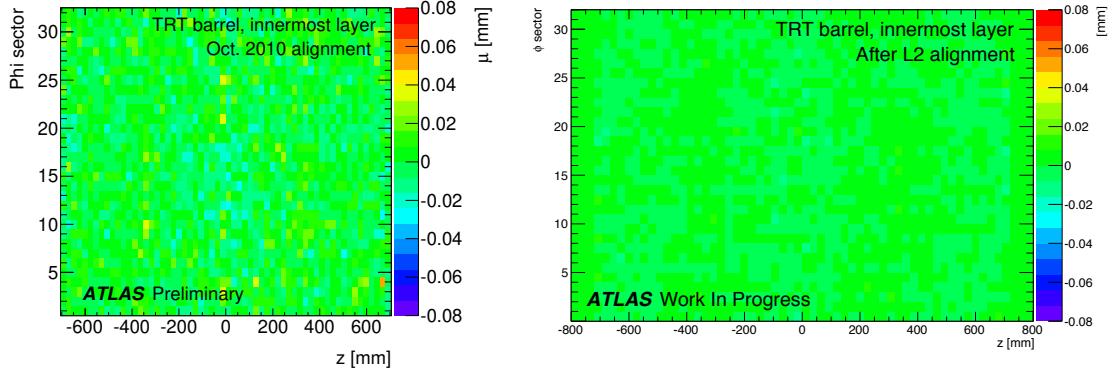


Figure 4.15: Two dimensional map of residuals in the first layer of the TRT barrel vs z and ϕ . Each bin represents the mean of a Gaussian fit to the TRT residuals in that bin. The map on the left is after the L3 (wire-by-wire) alignment of the TRT performed in 2010, and the map on the right is after the L2 alignment at the end of 2015. The z -axis for both plots use the same scale. Left figure taken from [66].

1460 **4.4 Momentum bias from sagitta deformations**

1461 A variety of weak mode deformations can exist in the detector even after alignment. As mentioned
 1462 previously, these weak modes consist of misalignments which don't affect the χ^2 of the residuals and
 1463 thus are not handled by the unconstrained alignment algorithm. In the presence of a weak mode,
 1464 the description of the detector geometry can still provide efficient and high quality track fits, but
 1465 there may also be systematic biases in one or more track parameters. Several weak modes, their
 1466 impacts on the reconstruction, and the steps taken to eliminate them are detailed in [69, 76]. This
 1467 section focuses specifically on *sagitta* deformations that result in a bias in the reconstructed track
 1468 momentum.

1469 These sagitta distortions consist of detector movements orthogonal to the trajectory of the outgoing
 1470 particle. The effect on the reconstructed track curvature is different for positively and negatively
 1471 charged particles, resulting in a charge-antisymmetric bias. An example of this is illustrated by the
 1472 curl deformation in Figure 4.16.

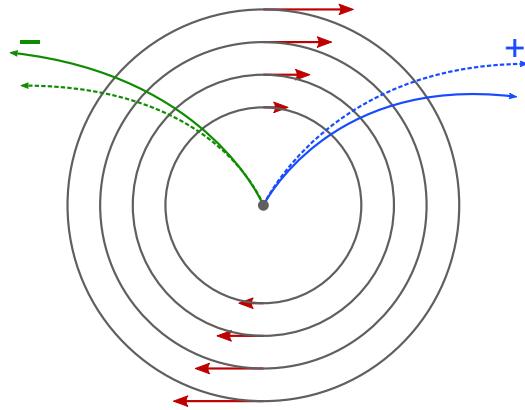


Figure 4.16: Representation of a curl distortion in the detector. The image represents a cutaway view of the transverse plane of the barrel region. The deformation is represented by the red arrows, and the impact on the reconstructed positive (blue) and negative (green) tracks are shown. The dashed lines represent the true particle trajectories, and the solid lines represent the reconstructed trajectories.

1473 In the plane transverse to ATLAS's magnetic field, outgoing particle tracks form circular arcs.
 1474 The sagitta is defined as the distance from the center of this arc to the center of its base, as shown in
 1475 Figure 4.17, and it represents the "amount of bending" in the track. In the case where the sagitta s
 1476 is considerably smaller than the detector radius R_0 , which is a valid assumption when working with

¹⁴⁷⁷ high momentum tracks, the transverse momentum of a particle of charge q can be written as [77]:

$$p_T \propto qB \frac{R_0^2}{8s} \quad (4.13)$$

¹⁴⁷⁸ where B is the strength of the detector's magnetic field. If a sagitta bias is present, the track's
¹⁴⁷⁹ transverse momentum shifts by [76]:

$$q/p_T \rightarrow q/p_T + \delta_s \quad \text{or} \quad p_T \rightarrow p_T \cdot (1 + qp_T\delta_s)^{-1} \quad (4.14)$$

¹⁴⁸⁰ where δ_s is a universal bias parameter that uniquely defines the deformation. Finally, since the
¹⁴⁸¹ reconstructed polar angle does not change under a sagitta deformation, the longitudinal component
¹⁴⁸² of the momentum scales along with the transverse component, and an equivalent equation can be
¹⁴⁸³ written for the total momentum:

$$p \rightarrow p \cdot (1 + qp_T\delta_s)^{-1} \quad (4.15)$$

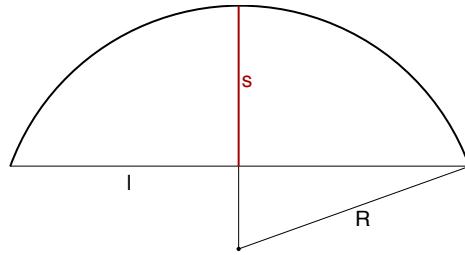


Figure 4.17: Geometric definition of the sagitta s in relation to the length of the chord l and the radius R of a circular arc.

¹⁴⁸⁴ 4.4.1 Sagitta bias monitoring with electron E/p

¹⁴⁸⁵ Since a sagitta bias results in changes in the momenta of particle tracks as measured by the ID, they
¹⁴⁸⁶ can be identified using independent measurements from other systems in the detector. One such
¹⁴⁸⁷ method involves using the energy-momentum ratio of electrons (E/p). Since the electron's energy
¹⁴⁸⁸ is measured in ATLAS's calorimeter systems, it is not sensitive to any sagitta bias that may exist
¹⁴⁸⁹ in the ID and in the corresponding measurement of the track momentum. Under the assumption
¹⁴⁹⁰ that the calorimeter response is independent of the charge of incoming particles, a charge-dependent
¹⁴⁹¹ momentum bias in the ID will manifest as a difference in the E/p ratio for electrons and positrons.

¹⁴⁹² In the presence of a sagitta bias, the momentum will change according to Equation 4.15 and the
¹⁴⁹³ average measured $\langle E/p \rangle$ can be written as:

$$\langle E/p \rangle^\pm \rightarrow \langle E/p \rangle^\pm \pm \langle E_T \rangle \delta_s \quad (4.16)$$

1494 where the approximation $p_T \approx E_T$ is used. Assuming that $\langle E/p \rangle^+ = \langle E/p \rangle^-$ in the absence of a
 1495 bias, the sagitta bias parameter can be written as:

$$\delta_s = \frac{\langle E/p \rangle^+ - \langle E/p \rangle^-}{2\langle E_T \rangle} \quad (4.17)$$

1496 If the kinematic selections for electrons and positrons are identical, the energy scale of the calorimeter
 1497 will not factor into the $\langle E/p \rangle$ difference; however, it will affect $\langle E_T \rangle$, which would scale the measured
 1498 δ_s . This is expected to be a small effect, as the energy scale for electrons has been measured at
 1499 13 TeV with uncertainties on the per-mil level across the entire detector [78].

1500 4.4.1.1 Measuring $\langle E/p \rangle$

1501 The E/p ratio is measured using electrons and positrons from $Z \rightarrow ee$ events in order to obtain a
 1502 high purity sample of candidate particles. They are required to pass a basic selection criteria to
 1503 ensure they are well measured in both the ID and the calorimeters:

- 1504 • $E_T > 25$ GeV
- 1505 • $|\eta| < 2.47$, excluding the calorimeter’s barrel-to-endcap transition region in $1.37 < |\eta| < 1.52$
- 1506 • Pass the MediumLH identification working point detailed in Section 3.2.4.2
- 1507 • Pass a selection of quality cuts, including a requirement that the electron be identified using
 1508 cluster information in the calorimeter
- 1509 • The associated track must have at least one hit in the IBL, three in the Pixel detector, and
 1510 five in the SCT detector.

1511 Events containing exactly two opposite-charge electrons passing this selection with an invariant mass
 1512 within 30 GeV of the Z boson mass are then used for the E/p calculation.

1513 Since the size of the sagitta bias δ_s is not expected to be constant across the entire detector,
 1514 a two-dimensional rectangular grid is constructed binned in detector η and ϕ . From the selected
 1515 events, separate distributions of E/p are made for electrons and positrons within each bin. Each
 1516 distribution is fit with Crystal Ball function¹², and the peak of the distribution is taken as the value
 1517 of $\langle E/p \rangle$. If there is no bias on the track momentum in the bin, the peaks for electrons and positrons
 1518 should agree. Example E/p distributions including the Crystal Ball fits are shown in Figure 4.18.

¹²The Crystal Ball function is a probability density function consisting of a Gaussian core and a power-law tail.

1519 It is important to emphasize that deviations from one in the *ratio* of $\langle E/p \rangle^+$ to $\langle E/p \rangle^-$ points
 1520 to a potential momentum biases. The value of $\langle E/p \rangle$ itself is not expected to equal one exactly,
 1521 as the track momentum on average tends to be slightly lower than the energy measurement in the
 1522 calorimeter. This is due to the fact that if the electron were to radiate a photon, its momentum
 1523 would change slightly, while it is likely that both the electron and the emitted photon would leave
 1524 energy deposits near each other in the calorimeter and be reconstructed into the same object.

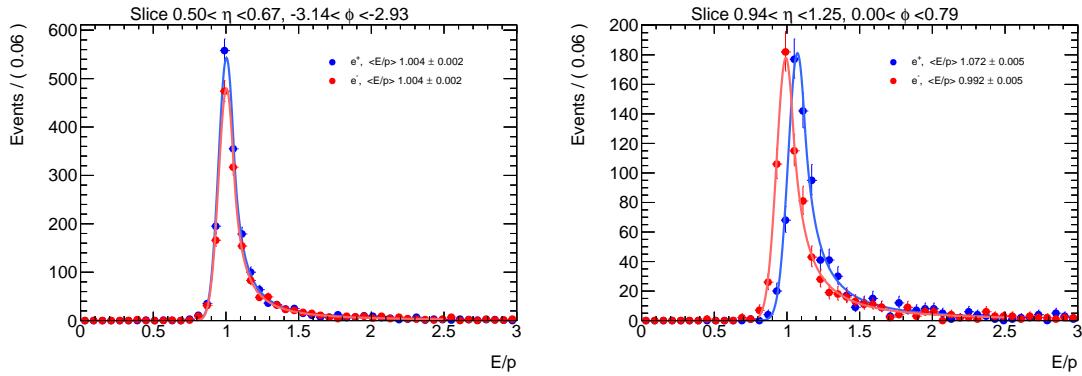


Figure 4.18: E/p distributions of electrons and positrons in two different $\eta\text{-}\phi$ bins of the detector. The left hand plot is taken from a region with no momentum bias where $\langle E/p \rangle^+ = \langle E/p \rangle^-$, while the right hand plot shows an 8% disagreement in $\langle E/p \rangle$ between electrons and positrons.

1525 Once the $\langle E/p \rangle^\pm$ distributions in each $\eta\text{-}\phi$ bin have been extracted from the fits, a two dimensional
 1526 map of δ_s can be constructed using Equation 4.17. The map gives an overview of sagitta
 1527 biases that may be present in the detector, and can be used by the alignment algorithm to reduce
 1528 the bias in the next iteration. In this case, the tracks fed to the alignment have their momenta
 1529 corrected according to [76]:

$$q/p_{\text{corr}} = q/p_{\text{reco}}(1 - qp_T\delta_s) \quad (4.18)$$

1530 where p_{reco} is the reconstructed momentum of the track. The corrected momentum is then con-
 1531 strained in the alignment.

1532 4.4.1.2 Results in 13 TeV data

1533 The E/p method has been used to monitor sagitta biases in the detector several times over the
 1534 course of Run 2. During this time, it has primarily served as an independent cross-check to a second
 1535 method using $Z \rightarrow \mu\mu$ events [76]. The $Z \rightarrow \mu\mu$ method identifies individual track momentum
 1536 biases through shifts in the reconstructed Z mass, which leaves it relatively insensitive to global
 1537 sagitta biases. For this reason, the sagitta bias maps produced using this technique are normalized

1538 to those from the E/p method before being used to constrain the alignment. The results of two
 1539 implementations of the E/p method are presented here.

- 1540 1. The first follows the end-of-year reprocessing of the entire ATLAS 2016 data set. Two sets
 1541 of alignment constants are compared: the *prompt* alignment, which was derived shortly after
 1542 each run was recorded, and the *reprocessed* alignment. The maps of the sagitta bias in each
 1543 alignment calculated using the E/p method are shown in Figure 4.19, and the comparison of
 1544 the η projection of the maps is shown in Figure 4.20.
 1545
 1546 2. The second uses the 2017 data after reprocessing, and compares the effects of multiple iterations
 1547 of the E/p method. In each iteration, the momenta of the electrons and positrons are corrected
 1548 according to Equation 4.15 using the value of δ_s computed in the previous iteration, and a
 1549 new sagitta bias map is calculated. If the method is indeed characterizing the sagitta biases
 1550 correctly, the corrections should converge quickly. The initial sagitta bias map is compared to
 1551 the map after two such iterations in Figure 4.21, and the sagitta bias projected along η for
 1552 each iteration is shown in Figure 4.22. Indeed, after just two iterations, δ_s is consistent with
 zero in nearly all bins.

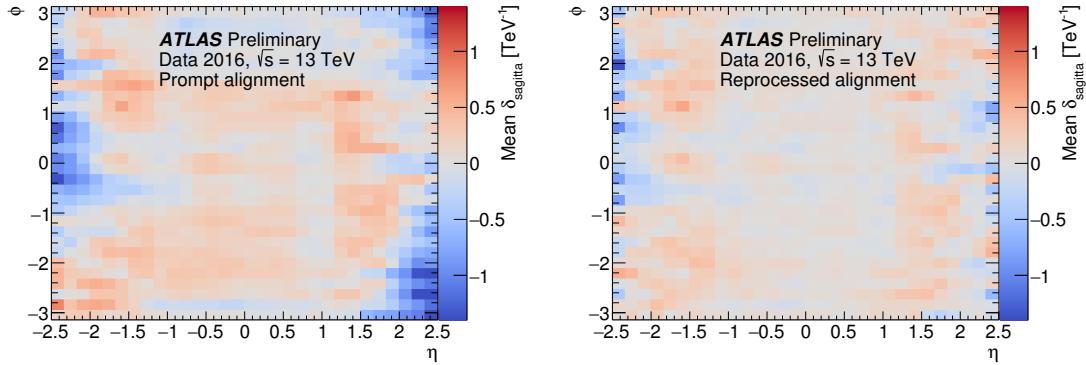


Figure 4.19: Sagitta bias in the $\sqrt{s} = 13 \text{ TeV}$ data collected by ATLAS in 2016 as a function of η and ϕ for the prompt (left) and reprocessed (right) alignments using the E/p method.

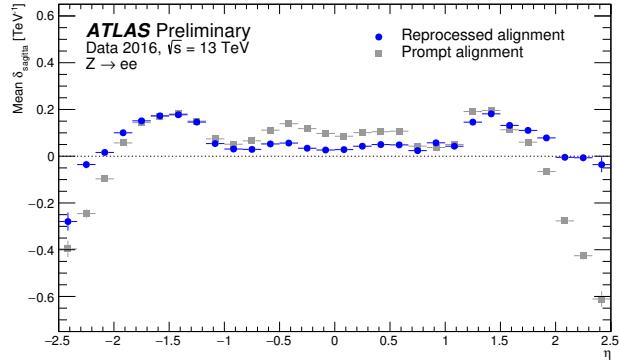


Figure 4.20: Sagitta bias in the $\sqrt{s} = 13 \text{ TeV}$ data collected by ATLAS in 2016 projected along η for the prompt (gray) and reprocessed (blue) alignments using the E/p method.

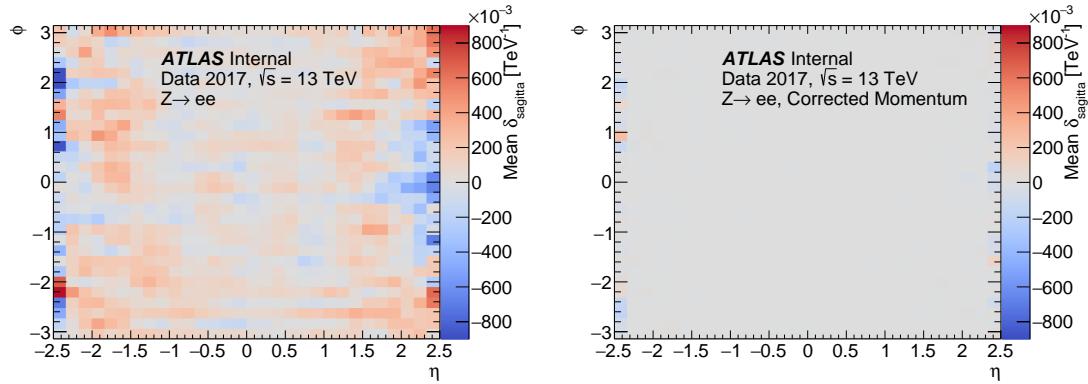


Figure 4.21: Sagitta bias in the $\sqrt{s} = 13 \text{ TeV}$ data collected by ATLAS in 2017 as a function of η and ϕ in reconstructed electrons (left) and after two iterations of momentum corrections (right) from the E/p method.

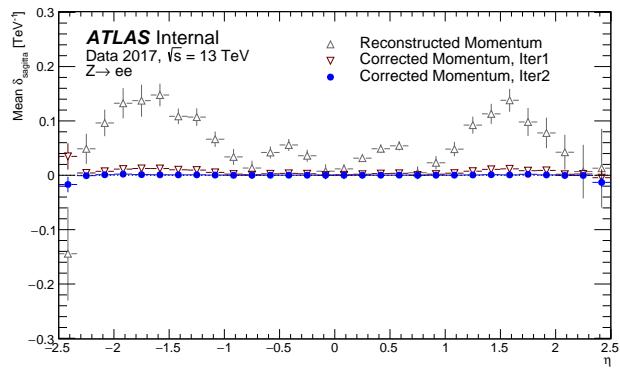


Figure 4.22: Sagitta bias in the $\sqrt{s} = 13$ TeV data collected by ATLAS in 2017 projected along η in reconstructed electrons (gray) and after one (red) and two (blue) iterations of momentum corrections from the E/p method.

1553

CHAPTER 5

1554

1555

Measurement of same-sign WW production at $\sqrt{s} = 13$ TeV with ATLAS

1556 Production of same-sign W boson pairs is a particularly interesting SM process. When produced
1557 via vector boson scattering (VBS), $W^\pm W^\pm jj$ is sensitive to the electroweak symmetry breaking
1558 (EWSB) mechanism as well as potential Beyond the Standard Model (BSM) physics processes.
1559 $W^\pm W^\pm jj$ events can be produced via electroweak-mediated (EWK) diagrams, of which VBS is a
1560 subset, or QCD-mediated diagrams. The biggest advantage of same-sign $W^\pm W^\pm jj$ over other VBS
1561 processes lies in its ratio of electroweak (EWK) to QCD production cross sections. Despite the
1562 opposite-sign $W^\pm W^\mp$ having a larger total cross section, its EWK-mediated diagrams are much
1563 smaller than its QCD-mediated diagrams, while for same-sign $W^\pm W^\pm$ the EWK production is
1564 considerably larger. This makes $W^\pm W^\pm jj$ one of the premier channels for studying VBS at the
1565 LHC.

1566 The first evidence of electroweak (EWK) $W^\pm W^\pm jj$ production was seen by the ATLAS and CMS
1567 experiments at $\sqrt{s} = 8$ TeV with excesses of 3.6σ [79] and 2.0σ [80] over backgrounds, respectively.
1568 More recently, ATLAS and CMS have both observed the EWK process at $\sqrt{s} = 13$ TeV with
1569 significances of 6.9σ [81] and 5.5σ [82], respectively. The ATLAS $\sqrt{s} = 13$ TeV observation and
1570 cross section measurement of EWK-produced $W^\pm W^\pm jj$ is presented in this chapter [81, 83].

1571 5.0.1 Experimental overview of vector boson scattering

1572 VBS processes are very important to understand due to their sensitivity to the EWSB mechanism.
1573 As explained in Section 2.3, in the absence of a light SM Higgs boson, the scattering amplitude of
1574 longitudinally polarized vector bosons grows with center-of-mass energy. However, once the Higgs

1575 is introduced, the divergences cancel and the cross section no longer grows unbounded.

1576 With the discovery of the Higgs boson in 2012 [16, 17], the EWSB mechanism can now be directly
 1577 studied. Due to the potential exchange of a Higgs boson in the VBS diagrams ($W^\pm W^\pm jj$ itself does
 1578 not contain an s -channel Higgs exchange), VBS processes are directly sensitive to properties of the
 1579 Higgs. For example, the high-mass tail in the VV scattering system allows an approximation of the
 1580 effective coupling strength of the Higgs to vector bosons that is independent of any assumptions
 1581 on the Higgs width [84]. Additionally, the center of mass energy dependence of the VV scattering
 1582 can reveal whether the Higgs boson unitarizes the longitudinal scattering amplitude fully or only
 1583 partially [85].

1584 VBS events are characterized by two quarks from the colliding protons each radiating a vector
 1585 boson which then scatter and decay in the detector. The incoming quarks carry a large amount of
 1586 momentum and only deflect a small amount upon emitting the vector boson; as a result, they often
 1587 enter the calorimeters very close to the beam line. Ignoring the decay products of the scattered
 1588 bosons for now, these VBS events result in a final state of two vector bosons (V) and two jets (j)
 1589 at high pseudorapidities (called *forward jets* or *tag jets*) from the outgoing quarks. The shorthand
 1590 $VVjj$ is used to represent this final state.

1591 $VVjj$ events can be produced via two different physical processes. The first involves purely
 1592 electroweak interactions in the tree-level diagrams, of order $\mathcal{O}(\alpha_{\text{EWK}}) = 6$ in the electroweak coupling
 1593 constant, and will be referred to as *EWK production*. This can be further broken down into VBS
 1594 and non-VBS events. In the VBS EWK production, the scattering occurs via triple or quartic gauge
 1595 couplings, as well as the exchange of a Higgs boson. The non-VBS EWK production contains the
 1596 same final state of two vector bosons and two outgoing quarks, but the bosons do not scatter. Due to
 1597 gauge invariance, it is not possible to separate the VBS from the non-VBS productions [86]; therefore,
 1598 both are included in the signal generation and are indistinguishable from one another. The second
 1599 process involves a mix of the EWK and strong interactions, of order $\mathcal{O}(\alpha_s) = 2 \otimes \mathcal{O}(\alpha_{\text{EWK}}) = 4$, and
 1600 will be referred to as *QCD production*. The tree-level Feynman diagrams for VBS EWK, non-VBS
 1601 EWK, and QCD $VVjj$ production are found in Figures 5.1, 5.2, and 5.3, respectively.

1602 5.0.2 Same-sign $W^\pm W^\pm$ scattering

1603 Same-sign $W^\pm W^\pm jj$ scattering is considered to be one of the best channels for studying VBS at
 1604 the LHC due to its favorable ratio of EWK to QCD production [84]. Since the VBS diagrams are
 1605 the primary interest for analysis, the QCD production is considered a background. Therefore a

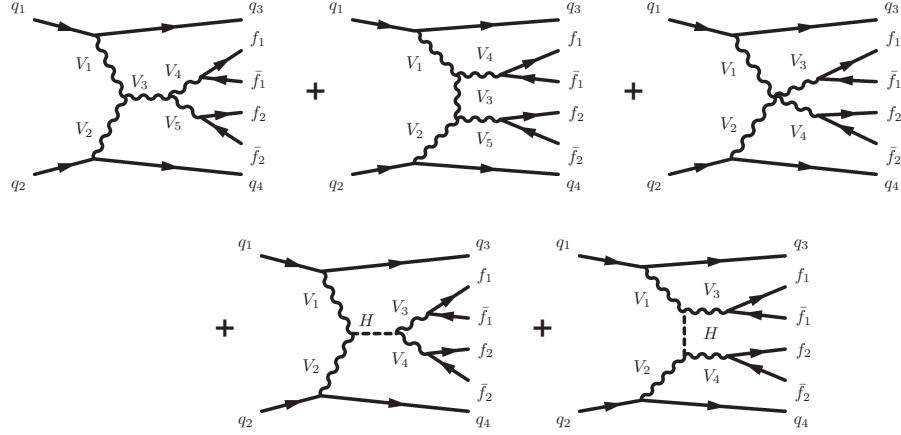


Figure 5.1: Tree-level Feynman diagrams for VBS EWK $VVjj$ production including triple gauge couplings involving W and/or Z bosons (top left and top middle), quartic gauge coupling (top right), or the exchange of a Higgs boson (s -channel bottom left and t -channel bottom right). The labels are quarks (q), fermions (f), and gauge bosons ($V = W, Z$).

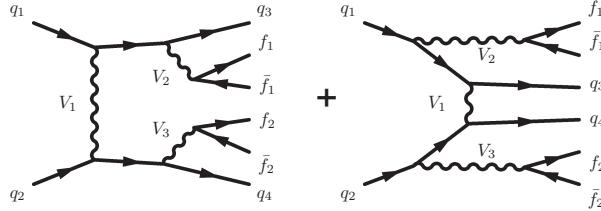


Figure 5.2: Tree-level Feynman diagrams for non-VBS EWK $VVjj$ production. The labels are quarks (q), fermions (f), and gauge bosons ($V = W, Z$).

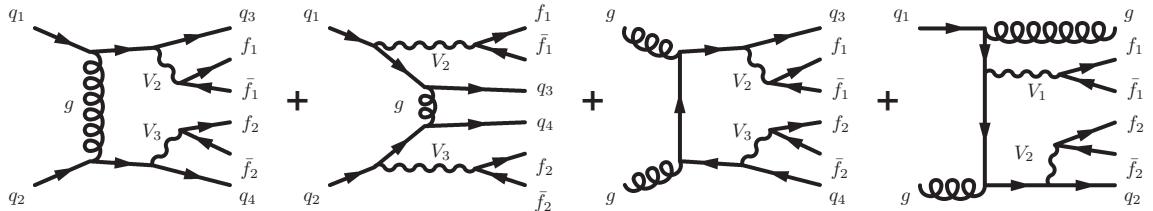


Figure 5.3: Tree-level Feynman diagrams for QCD $VVjj$ production. The labels are quarks (q), fermions (f), and gauge bosons ($V = W, Z$).

higher EWK-to-QCD ratio results in increased sensitivity to VBS. EWK and QCD cross sections at $\sqrt{s} = 13$ TeV for six leptonic $VVjj$ final states were calculated using the **SHERPA** MC generator in a VBS-enriched fiducial phase space in [87]. Despite its relatively low total cross section compared to some other $VVjj$ processes, the EWK-to-QCD ratio for $W^\pm W^\pm jj$ is 10-20 times higher than for other processes after applying VBS-enhancing selection criteria.

Final state	Process	σ_{EWK} [fb]	σ_{QCD} [fb]	$\sigma_{\text{EWK}}/\sigma_{\text{QCD}}$
$l^\pm l^\mp l^\pm l^\mp jj$	ZZ	0.098	0.100	0.98
$l^\pm l^\pm l^\mp \nu\nu jj$	$W^\pm Z$	2.34	4.38	0.53
$l^\pm l^\mp \nu\nu jj$	$W^\pm W^\mp, ZZ$	12.3	21.8	0.56
$l^\pm l^\pm \nu\nu jj$	$W^\pm W^\pm$	3.97	0.346	11.47
$l^\pm \nu\nu jj$	$W^\pm Z$	7.64	15.5	0.49
$\nu\nu\nu\nu jj$	ZZ	1.68	1.38	1.22

Table 5.1: Predicted cross sections for EQK and QCD production of diboson processes relevant to VBS at $\sqrt{s} = 13$ TeV using the **SHERPA** MC generator. The numbers for the $W^\pm W^\pm jj$ process are bolded. Leptons are required to have $p_T \geq 25$ GeV and lie within $|\eta| \leq 2.5$ with $m_{ll} > 20$ GeV, and at least two jets are required with $p_T \geq 30$ GeV and $|\eta| < 4.5$. The VBS contributions are enhanced by requiring the dijet invariant mass $m_{jj} > 500$ GeV and dijet separation $\Delta y_{jj} > 2.4$. Numbers taken from [87].

This analysis studies $W^\pm W^\pm jj$ scattering where both W bosons decay leptonically to $e\nu$ or $\mu\nu$ ¹³. The $W^\pm W^\pm jj$ VBS final state consists of two leptons with the same electric charge, two neutrinos, and two high energy forward jets with a large invariant mass. Tree-level Feynman diagrams of VBS $W^\pm W^\pm jj$ production can be found in Figure 5.4 and a visual representation of the VBS topology can be found in Figure 5.5.

The two tag jets in the characteristic VBS signature also serve as a powerful tool to suppress the QCD production mode. In EWK events, the two jets tend to have much higher separation and a larger combined invariant mass than the two leading jets in a QCD event. The two plots shown in Figure 5.6 highlight the differences in these dijet quantities between the two production modes. Finally, an ATLAS event display of a real $W^\pm W^\pm jj$ candidate event is shown in Figure 5.7.

5.0.3 Overview of backgrounds

In addition to QCD production of $W^\pm W^\pm jj$ events, there are several other processes with a final state of two same-sign leptons, two neutrinos, and two jets. However, due to the ± 2 final state charge,

¹³Throughout the rest of this chapter, unless stated otherwise, l denotes either electrons (e) or muons (μ), and ν denotes a neutrino. Additionally, e , μ , and ν with no charge or anti-particle designation refer interchangeably to either the particle or anti-particle.

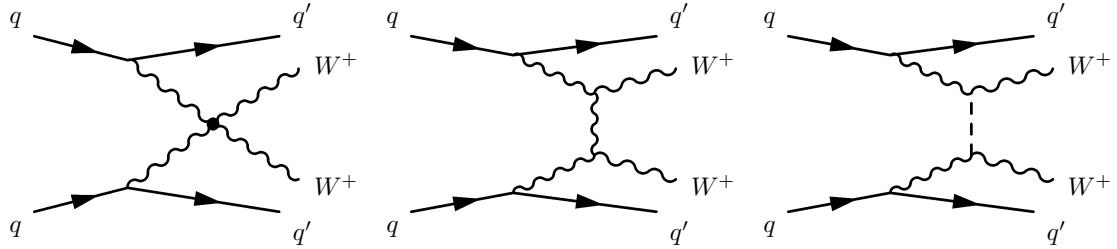


Figure 5.4: Leading order Feynman diagrams for VBS EWK production of $W^\pm W^\pm jj$ events. The leftmost diagram contains a quartic gauge coupling vertex, and the rightmost diagram contains an exchange of a Higgs boson.

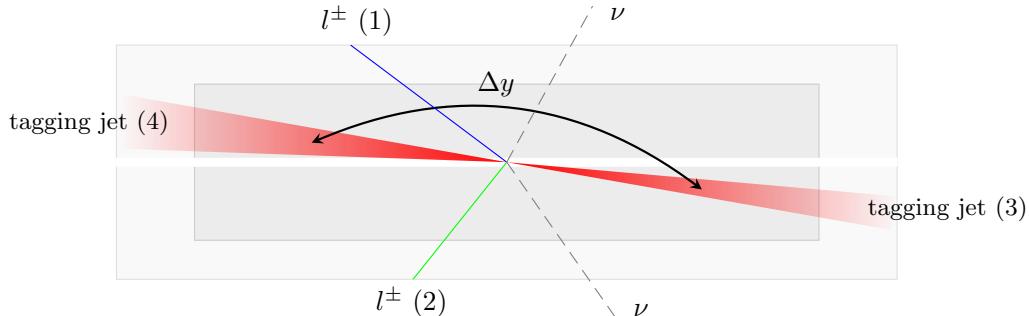


Figure 5.5: $W^\pm W^\pm jj$ VBS event topology containing two leptons (1 and 2) with the same electric charge, two neutrinos, and two forward tagging jets (3 and 4) with large rapidity separation Δy .

1624 there is a considerable reduction in SM backgrounds (such as Z boson events) when compared to
 1625 an analysis like opposite-sign $W^\pm W^\mp jj$.

1626 One of the largest background sources involves processes with prompt leptons¹⁴. These are events
 1627 that contain two leptons with the same electric charge and one or more additional leptons that are
 1628 “lost”, either by failing the selection criteria or falling outside of the detector’s acceptance. The
 1629 number of processes that can contribute is limited by the requirement of same-sign leptons, and as a
 1630 result this background is dominated by multi-boson processes, with the largest contribution coming
 1631 from WZ events and smaller contributions from ZZ and $t\bar{t} + V$ events. Triboson events where one
 1632 boson decays hadronically also contribute to this background; however, the jets are generally softer
 1633 and more central than in a typical VBS event, and the cuts applied on the forward jets suppress
 1634 these contributions.

¹⁴Prompt leptons are those that are produced in the primary collision and are a direct decay product of the process of interest. Non-prompt leptons originate from some secondary process, such as a b -hadron decay. Jets that are mis-reconstructed as a lepton are also considered non-prompt.

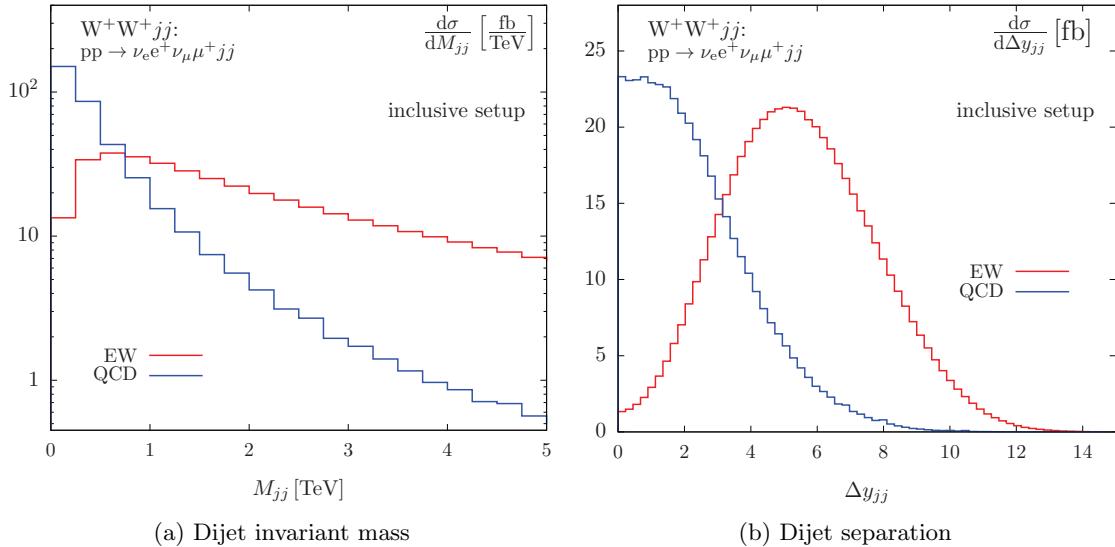


Figure 5.6: Generator level comparisons at $\sqrt{s} = 7$ TeV of dijet invariant mass (M_{jj} , left) and dijet rapidity (Δy_{jj} , right) in EWK (red) and QCD (blue) $W^\pm W^\pm jj$ events with no selection cuts applied. Plots taken from [88].

1635 The other dominant background comes from non-prompt, or “fake”, leptons. Here one or more
 1636 leptons originate from the decay of another particle unrelated to the signal process, such as a heavy-
 1637 flavor decay or photon conversion, or come from a jet that is misidentified as a lepton. This back-
 1638 ground is mostly made up of events from $t\bar{t}$ and $W+jets$ processes, with a much smaller contribution
 1639 from conversions in $V\gamma$ events.

1640 Finally, opposite-sign lepton pairs can enter the signal region if one of the leptons is reconstructed
 1641 with the wrong charge (called *charge misidentification*¹⁵). In practice, this only affects events with
 1642 electrons, as the charge misidentification rate for muons is negligible [89]. This is a major background
 1643 in events with two electrons, dominated by $Z \rightarrow ee$ events; it is a much smaller contribution for
 1644 events with one electron and one muon, where the primary contribution comes from $t\bar{t}$ events.

1645 5.1 Data and Monte Carlo samples

1646 This analysis uses 36.1 fb^{-1} of $\sqrt{s} = 13$ TeV proton-proton collisions recorded by ATLAS during
 1647 the 2015 and 2016 data taking periods. The uncertainty in the combined integrated luminosity is
 1648 2.1%. It is derived following a methodology similar to that detailed in [90] and using the LUCID-2

¹⁵Charge misidentification is also referred to interchangeably as *charge mis-ID* and *charge flip*.

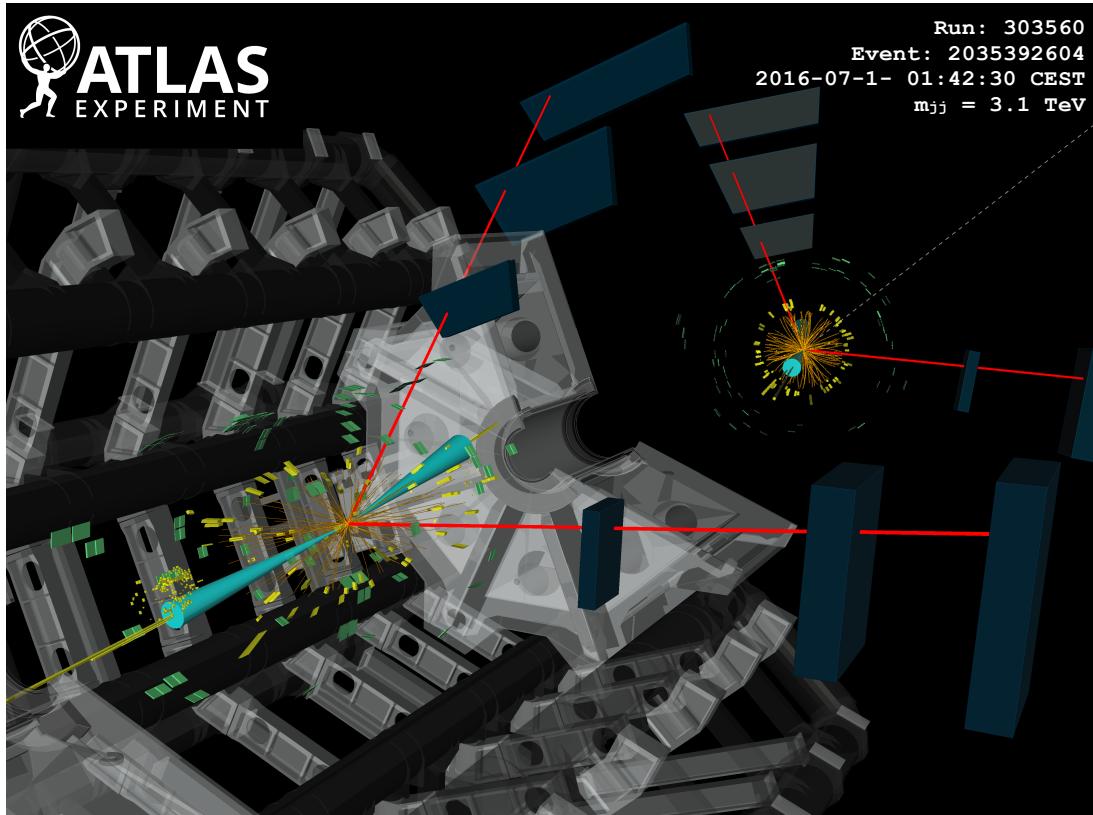


Figure 5.7: ATLAS event display of a $pp \rightarrow W^+W^+ \rightarrow \mu^+\nu_\mu\mu^+\nu_\mu jj$ event. The muons are represented by the red lines travelling from the ID through the MS, and the forward jets are represented by the blue cones with yellow energy deposits in the calorimeters. The direction of the E_T^{miss} in the transverse plane is indicated by the gray dashed line in the inset image. Event display taken from [81].

1649 detector for the baseline luminosity measurements [91] from calibration of the luminosity scale using
 1650 x - y beam-separation scans.

1651 5.1.1 Monte Carlo samples

1652 A number of Monte Carlo (MC) simulations are employed to model signal and background processes.
 1653 In order to model the real collision data as closely as possible, each MC sample has been passed
 1654 through a full simulation of the ATLAS detector in GEANT4 [92, 93], and events have been recon-
 1655 structed using the same algorithms as the data. The simulation reproduces as closely as possible
 1656 the momentum resolutions and calorimeter responses of the detector, and also includes the effects
 1657 of pileup by including soft QCD interactions using PYTHIA v8.1 [94]. The MC samples used in this

Process	Generator	Comments
$W^\pm W^\pm jj$ (EWK)	SHERPA v2.2.2	Signal sample
$W^\pm W^\pm jj$ (EWK)	POWHEG-BOX v2	Systematics sample
$W^\pm W^\pm jj$ (QCD)	SHERPA v2.2.2	
Diboson	SHERPA v2.2.2	Both bosons decay leptonically ($llll$, $ll\nu\nu$, $l\nu\nu\nu$)
	SHERPA v2.2.1	One boson decays leptonically, the other hadronically
Triboson	SHERPA v2.1.1	
$W+jets$	SHERPA v2.2.1	
$Z+jets$	Madgraph5_aMC@NLO	
$V\gamma$	SHERPA v2.1.1	
$V\gamma jj$ (EWK)	SHERPA v2.2.4	
$t\bar{t}V$	Madgraph5_aMC@NLO	
$t\bar{t}$	POWHEG-BOX v2	
Single top	POWHEG-BOX v1	EWK t -, s -, & Wt -channels

Table 5.2: Summary of MC samples used in the analysis.

analysis are detailed in this section and summarized in Table 5.2.

The $W^\pm W^\pm jj$ samples are modeled using SHERPA v2.2.2 [95, 96, 97] with the NNPDF3.0 PDF set [98]. The EWK signal samples were generated by fixing the electroweak coupling constant to $\mathcal{O}(\alpha_W) = 6$, and a QCD background sample was also generated with $\mathcal{O}(\alpha_W) = 4$. SHERPA includes up to one parton at next-to-leading order (NLO) and up to three at leading order (LO) in the strong coupling constant α_s . A second $W^\pm W^\pm jj$ EWK sample was generated using POWHEG-BOX v2 [99] with the NNPDF3.0 PDF set and at NLO accuracy. This sample is only used for systematic studies, as POWHEG-BOX does not include resonant triboson contributions in its matrix element, which are non-negligible at NLO [100].

Diboson processes (VV where $V = W, Z$) are simulated with SHERPA v2.2.2 for mixed hadronic and leptonic decays and SHERPA v2.2.1 for fully leptonic decays of the bosons. Similarly, triboson (VVV) and $V\gamma$ processes are simulated using SHERPA v2.1.1 with up to one parton at NLO and up to three at LO. $W+jets$ processes are simulated with SHERPA v2.2.1 with up to two partons at NLO and four at LO. All the above SHERPA samples use the NNPDF3.0 PDF set and SHERPA's own parton showering. The $Z+jets$ events are generated with Madgraph5_aMC@NLO [101] at LO and interfaced with PYTHIA v8.1 for parton showering.

$t\bar{t}$ events are generated using POWHEG-BOX v2 with the CT10 PDF set [102]. $t\bar{t}V$ samples are generated at NLO with Madgraph5_aMC@NLO and the NNPDF3.0 PDF set interfaced with PYTHIA v8 for parton showering. Finally, single top events are generated with POWHEG-BOX v1 and the CT10f4 PDF set interfaced with PYTHIA v6 [103] for parton showering.

1678 **5.2 Object and event selection**

1679 This section details the selection criteria for objects used in the analysis as well as the selection for
1680 signal events.

1681 **5.2.1 Object selection**

1682 Muons, electrons, and jets all must pass strict selection requirements to ensure that only high quality,
1683 well measured objects are used. For leptons, a baseline selection is defined (called the *preselection*),
1684 which all leptons must pass in order to be considered for the analysis. This preselection is an
1685 intentionally loose set of criteria designed to have high acceptance in order to reject backgrounds
1686 with additional leptons (such as $WZ \rightarrow 3l\nu jj$). Signal leptons are then required to satisfy a much
1687 tighter *signal selection* aimed at suppressing backgrounds from non-prompt or fake leptons. A third
1688 set of lepton selection criteria, the *loose selection*, defines a sample enriched in non-prompt leptons,
1689 and it is used in the fake-factor method for estimating the non-prompt background, discussed in
1690 detail in Section 5.3.4. Jets are only required to pass one set of selection criteria. These selections
1691 are outlined in the following subsections and summarized in Table 5.3 for muons, Table 5.4 for
1692 electrons, and Table 5.5 for jets.

1693 **5.2.1.1 Muon candidate selection**

1694 Cuts on muon momentum serve to reject low momentum leptons from background processes and from
1695 additional collisions occurring in pileup events. Preselected muons must have transverse momentum
1696 $p_T > 6$ GeV, and the signal muons must pass $p_T > 27$ GeV. The p_T requirement for loose muons
1697 is lower than for signal muons, at $p_T > 15$ GeV, for reasons that are discussed in Section 5.3.4.3.
1698 Muons are required to fall within the detector's η acceptance: $|\eta| < 2.7$ for preselected muons, which
1699 is tightened to $|\eta| < 2.5$ for the signal muons.

1700 Cuts on the transverse and longitudinal impact parameters are applied to ensure that the can-
1701 didate muon originated from the primary particle interaction and not some other source. The
1702 preselection and the loose selection both have relaxed requirements on the transverse impact pa-
1703 rameter significance (d_0/σ_{d_0}) than the signal selection; all three have the same requirement on the
1704 transverse impact parameter ($|z_0 \times \sin \theta|$).

1705 Finally, the muon candidates are required to pass a particle identification and an isolation criteria
1706 as defined in [61]. The methods used in constructing the identification and isolation working points
1707 are described in more detail in Section 3.2.4.3. The muon identification serves to select prompt

1708 muons with high efficiency and well measured momenta. This analysis uses two different working
 1709 points: **Loose** for preselected muons and **Medium** for loose and signal muons, where **Medium** muons
 1710 are a tighter subset of those that pass the **Loose** requirement. Muon isolation is a measurement
 1711 of detector activity around the muon candidate, and it is measured with both track-based and
 1712 calorimeter-based variables. The isolation working point used for the signal muons, **Gradient**, is
 1713 defined such that there is 90% or better background rejection efficiency for 25 GeV muons, and 99%
 1714 efficiency at 60 GeV. There is no minimum isolation requirement for preselected or loose muons.
 1715 Loose muons are additionally required to fail one or both of the signal transverse impact parameter
 1716 cut and signal isolation requirement.

Muon preselection	
Momentum cut	$p_T > 6$ GeV
Angular acceptance	$ \eta < 2.7$
Longitudinal impact parameter	$ z_0 \times \sin \theta < 0.5$ mm
Transverse impact parameter	$d_0/\sigma_{d_0} < 10$
Particle identification	Loose

Muon signal selection	
Momentum cut	$p_T > 27$ GeV
Angular acceptance	$ \eta < 2.5$
Longitudinal impact parameter	$ z_0 \times \sin \theta < 0.5$ mm
Transverse impact parameter	$d_0/\sigma_{d_0} < 3$
Particle identification	Medium
Particle isolation	Gradient

Muon loose selection	
Momentum cut	$p_T > 15$ GeV
Angular acceptance	$ \eta < 2.5$
Longitudinal impact parameter	$ z_0 \times \sin \theta < 0.5$ mm
Transverse impact parameter	$d_0/\sigma_{d_0} < 10$
Particle identification	Medium
Fail signal transverse impact parameter and/or isolation cuts	

Table 5.3: Muon selection criteria. All muons are required to pass the preselection (top), and then either the signal (middle) or loose (bottom) criteria is applied to the preselected electrons.

1717 5.2.1.2 Electron candidate selection

1718 The electron candidate selections are very similar to those for muons. The momentum cut starts
 1719 at $p_T > 6$ GeV for the preselection, increases to $p_T > 20$ GeV for loose electrons, and finally to

1720 $p_T > 27$ GeV for signal electrons. The $|\eta|$ cut for electrons requires $|\eta| < 2.47$ for all electrons,
1721 with the region $1.37 \leq |\eta| \leq 1.52$ removed from loose and signal electrons. This is where the
1722 electromagnetic calorimeter transitions from the barrel to the endcaps and is not fully instrumented.
1723 Both the transverse and longitudinal impact parameter cuts are the same for all electron selections.

1724 The electron particle identification uses a multivariate likelihood technique (LH) detailed in
1725 Section 3.2.4.2. Preselected electrons must pass the `LooseLH` working point with an additional
1726 requirement that there be a reconstructed track hit in the first layer of the pixel detector (a so-
1727 called B -layer hit). The LH requirement for the loose and signal electrons increases in tightness
1728 using `MediumLH` and `TightLH` electrons, respectively. As for isolation, the `Gradient` working point is
1729 required for signal electrons only. The loose electrons must fail one or both of the signal identification
1730 and isolation requirements.

Electron preselection	
Momentum cut	$p_T > 6$ GeV
Angular acceptance	$ \eta < 2.47$
Longitudinal impact parameter	$ z_0 \times \sin \theta < 0.5$ mm
Transverse impact parameter	$d_0/\sigma_{d_0} < 5$
Particle identification	<code>LooseLH</code> + B -layer hit

Electron signal selection	
Momentum cut	$p_T > 27$ GeV
Angular acceptance	$ \eta < 2.47$, excluding $1.37 \leq \eta \leq 1.52$
Longitudinal impact parameter	$ z_0 \times \sin \theta < 0.5$ mm
Transverse impact parameter	$d_0/\sigma_{d_0} < 5$
Particle identification	<code>TightLH</code>
Particle isolation	<code>Gradient</code>

Electron loose selection	
Momentum cut	$p_T > 20$ GeV
Angular acceptance	$ \eta < 2.47$, excluding $1.37 \leq \eta \leq 1.52$
Longitudinal impact parameter	$ z_0 \times \sin \theta < 0.5$ mm
Transverse impact parameter	$d_0/\sigma_{d_0} < 5$
Particle identification	<code>MediumLH</code>
Fail signal identification and/or isolation cuts	

Table 5.4: Electron selection criteria. All electrons are required to pass the preselection (top), and then either the signal (middle) or loose (bottom) criteria is applied to the preselected electrons.

1731 **5.2.1.3 Jet candidate selection**

1732 The final objects that need to pass selection are jets. Jets are clustered using the anti- k_t algo-
1733 rithm [104] within a radius of $\Delta R = 0.4$. The jets are then calibrated using E_T - and η -dependent
1734 correction factors that are trained using MC simulations [105]. The calibrated jets are required
1735 to have $p_T > 30$ GeV if they lie in the forward regions of the detector ($2.4 < |\eta| < 4.5$) and
1736 $p_T > 25$ GeV in the central region ($|\eta| \leq 2.4$). In order to suppress pileup jets, the so-called jet-
1737 vertex-tagger (JVT) discriminant associates a jet with the primary interaction vertex [106]; central
1738 jets with $p_T > 60$ GeV are required to pass the **Medium** JVT working point, which corresponds to
1739 an average efficiency of over 92%. Finally, the jets are required to be separated from the selected
1740 leptons by at least $\Delta R(j, l) > 0.3$.

Jet selection	
Momentum cut	$p_T > 30$ GeV for $2.4 < \eta < 4.5$ $p_T > 60$ GeV for $ \eta < 2.4$
JVT cut	Medium
Jet-lepton separation	$\Delta R(j, l) > 0.3$

Table 5.5: Jet selection criteria. All jets are required to pass the above selection in order to be used in the analysis.

1741 **5.2.1.4 Treatment of overlapping objects**

1742 In the event that one or more objects are reconstructed very close to each other, there is the possibility
1743 for double-counting if both originated from the same object. The procedure by which this ambiguity
1744 is resolved is called *overlap removal* (OR). The standard ATLAS recommendation for OR [107, 108]
1745 is implemented in this analysis and is summarized in Table 5.6.

1746 Since electrons leave a shower in the EM calorimeter, every electron has a jet associated with
1747 it. Therefore, any jets close to an electron (within $\Delta R(e, j) < 0.2$) are rejected due to the high
1748 probability that they are the same object. On the other hand, when jets and electrons overlap
1749 within a larger radius of $0.2 < \Delta R(e, j) < 0.4$, it is likely that the electron and jet both are part of
1750 a heavy-flavor decay, and the electron is rejected.

1751 High energy muons can produce photons via bremsstrahlung radiation or collinear final state
1752 radiation which result in nearby energy deposits in the calorimeters. Non-prompt muons from
1753 hadronic decays produce a similar signature; however, in this case the jet has a higher track multi-
1754 plicity in the ID. It is possible to address both cases simultaneously by rejecting the jet when the

1755 ID track multiplicity is less than three, and otherwise rejecting the muon, for jets and muons within
 1756 $\Delta R(\mu, j) < 0.4$.

1757 In addition to the case above where muon bremsstrahlung results in a nearby reconstructed
 1758 jet, the ID track from the muon and the calorimeter energy deposit from the photon can lead to an
 1759 electron being reconstructed. In this case, if both a muon and an electron share a track in the ID, the
 1760 muon is kept and the electron is rejected, unless the muon is calorimeter-tagged (see Section 3.2.4.3),
 1761 in which case the muon is removed in favor of the electron.

Overlap	Check	Result (remove → keep)
Electron & Jet	$\Delta R(e, j) < 0.2$	Jet → Electron
	$0.2 < \Delta R(e, j) < 0.4$	Electron → Jet
Muon & Jet	$\Delta R(\mu, j) < 0.4$ and Jet N_{ID} tracks < 3	Jet → Muon
	$\Delta R(\mu, j) < 0.4$ and Jet N_{ID} tracks ≥ 3	Muon → Jet
Electron & Muon	Shared ID track	Electron → Muon
	Shared ID track & muon is calo-tagged	Muon → Electron

Table 5.6: Summary of the overlap removal procedure used in the analysis. If the criteria in the “check” column is met, in the “result” column, the object on the left of the arrow is removed in favor of the object on the right.

1762 5.2.2 Signal event selection

1763 After the objects have been selected, cuts are applied on a per-event level to select $W^\pm W^\pm jj$ signal
 1764 events. The event selection is summarized in Table 5.8.

1765 The initial event selection chooses events that pass one or more of the trigger requirements listed
 1766 in Table 5.7. At least one signal lepton is “matched” to a passed trigger in order to ensure that
 1767 it was indeed a signal lepton that fired the trigger. A collection of *event cleaning* cuts must also
 1768 be passed in order to remove events collected during periods in which one or more components of
 1769 the detector were not operating optimally. Finally, the events are required to contain at least one
 1770 interaction vertex. An event can have multiple reconstructed vertices from additional proton-proton
 1771 collisions that occurred in the same bunch crossing. In this case, the *primary vertex* is determined
 1772 by choosing the vertex with the largest sum of the p_T^2 of its associated tracks.

1773 Events are then required to contain exactly two signal leptons with the same electric charge.
 1774 The dilepton pair must have a combined invariant mass of $m_{ll} \geq 20$ GeV in order to suppress
 1775 low mass Drell-Yan backgrounds. Two additional selections are applied to events containing two
 1776 electrons: both are required to have $|\eta| < 1.37$ with an invariant mass at least 15 GeV away from
 1777 the Z -boson mass to reduce events where one electron is reconstructed with the wrong charge (this

	2015 data	2016 data
Electrons	$p_T > 24$ GeV and Medium ID	$p_T > 26$ GeV and Tight ID and Loose isolation
	$p_T > 60$ GeV and Medium ID	$p_T > 60$ GeV and Medium ID
	$p_T > 120$ GeV and Loose ID	$p_T > 140$ GeV and Loose ID
Muons	$p_T > 20$ GeV and Loose isolation $p_T > 50$ GeV	$p_T > 26$ GeV and Medium isolation $p_T > 50$ GeV

Table 5.7: Summary of trigger requirements for electrons and muons for $\sqrt{s} = 13$ TeV data collected in 2015 and 2016. At least one of the triggers must be satisfied.

background will be discussed in more detail in Section 5.3.3). To suppress backgrounds from final states with more than two leptons, such as WZ or ZZ , events with more than two leptons passing the preselection are vetoed.

Missing transverse energy (E_T^{miss}) represents any particles that escape the detector without being measured, such as neutrinos, and it is defined as the magnitude of the vector sum of transverse momenta of all reconstructed objects. It can be difficult to calculate accurately, as it involves measurements from all subsystems within the detector, and it is sensitive to any corrections that may be applied to the reconstructed physics objects [109]. These corrections, including the momentum smearing for muons, energy scale and smearing for electrons, and jet calibrations, are propagated to the E_T^{miss} calculation. Events are required to contain $E_T^{\text{miss}} > 30$ GeV in order to account for the two neutrinos from the W boson decays.

At least two jets are required. The leading and subleading jets must have $p_T > 65$ GeV and $p_T > 35$ GeV, respectively, and are referred to as the *tagging jets*. Events are vetoed if they contain one or more jets that have been tagged as a b -jet to suppress backgrounds from heavy flavor decays (especially top quark events). The b -tagging algorithm used by ATLAS is a boosted decision tree (BDT) called MV2c10, and this analysis uses a working point with 85% efficiency [110].

Finally, cuts are applied on the VBS signature outlined in Section 5.0.2. The tagging jets are required to have a dijet invariant mass $m_{jj} > 200$ GeV and be separated in rapidity by $|\Delta y_{jj}| > 2.0$. This preferentially selects the VBS EWK events over the QCD-produced $W^\pm W^\pm jj$ events.

5.3 Background estimations

The major sources of background events are summarized in Section 5.0.3, and the methods used to estimate them are detailed in this section. Prompt backgrounds from ZZ and $t\bar{t} + V$ are estimated directly from MC simulations. The shape of the WZ and $V\gamma$ backgrounds are taken from MC, and the predicted yeilds are normalized to the data predictions in dedicated control regions, as outlined

Event selection	
Event preselection	Pass at least one trigger with a matched lepton Pass event cleaning At least one reconstructed vertex
Lepton selection	Exactly two leptons passing signal selection Both signal leptons with the same electric charge Dilepton mass $m_{ll} > 20$ GeV $ \eta < 1.37$ and $ M_{ee} - M_Z > 15$ GeV (ee-channel only) Veto events with more than two preselected leptons
Missing transverse energy	$E_T^{\text{miss}} \geq 30$ GeV
Jet selection	At least two jets Leading jet $p_T > 65$ GeV Subleading jet $p_T > 35$ GeV $m_{jj} > 200$ GeV $N_{b\text{-jet}} = 0$ $ \Delta y_{jj} > 2.0$

Table 5.8: The signal event selection.

in Sections 5.3.1 and 5.3.2, respectively. Opposite sign events with a charge misidentified electron are estimated by a data-driven background method which is summarized in Section 5.3.3. Finally, a *fake-factor* method is used to estimate the contributions from non-prompt backgrounds and is the subject of Section 5.3.4.

5.3.1 Estimation of the WZ background

The dominant background involving prompt leptons comes from $WZ + \text{jets}$ events. The contribution is estimated from MC simulation and normalized to data in a control region enriched in WZ events. This region is defined by the same event selection as the signal region in Table 5.8, with the following changes applied to increase the purity of the WZ process:

- The third lepton veto is inverted, requiring a third lepton with $p_T > 15$ GeV
- Two of the leptons must make a same-flavor opposite-sign pair. If more than one pair exists, the one with m_{ll} closest to the Z boson mass is chosen.
- The trilepton invariant mass is required to be $m_{lll} > 106$ GeV to reduce contributions from $Z\gamma$ and $Z + \text{jets}$

Once the event yields in the control region are calculated, they are propagated to the final signal region fit, detailed in Section 5.4.1, in a single bin combining all the lepton channels. The systematic

1818 uncertainties of the WZ background are also calculated at this time. The event yields for the WZ
 1819 control region are listed in Table 5.9, and distributions of the leading lepton p_T and η as well as
 1820 trilepton invariant mass m_{lll} are found in Figures 5.9 and 5.8, respectively.

Event yields in the WZ control region	
WZ	197.9 ± 1.4
ZZ	14.1 ± 0.3
Triboson	1.26 ± 0.1
top	10.8 ± 1.1
$Z\gamma$	3.1 ± 1.1
$Z+jets$	2.5 ± 1.4
Total prediction	229.7 ± 2.5
Data	201 ± 14.2

Table 5.9: Event yields in the WZ control region before normalization. All lepton flavor channels are combined.

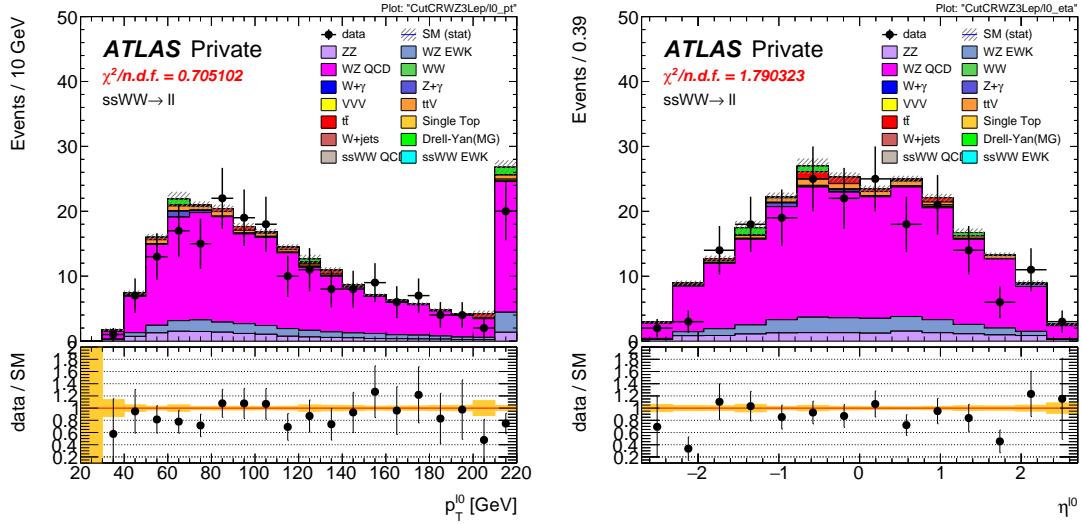


Figure 5.8: Leading lepton p_T (left) and η (right) distributions in the WZ control region before normalization. All lepton channels are combined.

1821 5.3.2 Estimation of the $V\gamma$ background

1822 Events from $V\gamma$ processes can pass selection if the photon converts into an e^+e^- pair and one of the
 1823 electrons passes the selection criteria. The background is estimated from MC simulations which are
 1824 then scaled by a normalization factor calculated from a control region enriched in $Z \rightarrow \mu\mu + \gamma$ events.

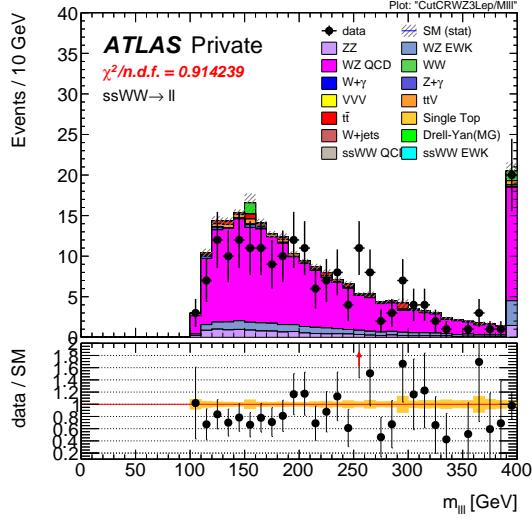


Figure 5.9: Trilepton invariant mass m_{lll} distribution in the WZ control region before normalization. All lepton channels are combined.

1825 This control region selects two opposite-sign muons and an additional electron that is assumed to
 1826 come from the photon conversion. The full event selection is detailed in Table 5.10.

$V\gamma$ control region
Exactly two muons with $p_T > 27$ GeV and $p_T > 20$ GeV
Exactly one additional electron with $p_T > 15$ GeV
Remove overlap between $Z+jets$ and $Z\gamma$
Di-muon + photon invariant mass $75 < M_{\mu\mu\gamma} < 100$ GeV
$E_T^{\text{miss}} < 30$ GeV

Table 5.10: Selection criteria for the $V\gamma$ control region.

1827 The $Z\gamma$ MC samples available do not cover the full range of p_T^γ and $\Delta R(\gamma, l)$; thus, additional
 1828 Drell-Yan samples ($Z+jets$) are used to fill out the phase space. Overlap between the two samples
 1829 are removed to avoid double counting. Events with final state photons are checked at truth level¹⁶
 1830 to ensure that the photon did not originate from a hadronic decay. Cuts on $p_T^\gamma > 10$ GeV and
 1831 $\Delta R(\gamma, l) > 0.1$ are then applied at generator level, and $Z\gamma$ events that fail this additional selection
 1832 and $Z+jets$ events that pass it are removed.

1833 The normalization factor is calculated directly from the event yields in the $V\gamma$ control region
 1834 rather than in the signal fit, as is done for the WZ background. The event yields are listed in

¹⁶Truth particles are the particles produced directly by the MC generator before being passed through the full detector simulation, at which point they are considered *reconstruction-level* (or *reco-level*) particles.

1835 Table 5.11, and the normalization factor is determined to be 1.77. No MC events from $Z\gamma$ processes
 1836 survive the full event selection; thus, the scaling is only applied to the $W\gamma$ background in the signal
 1837 region. A systematic uncertainty of 44% is assigned to the background based off of the uncertainties
 1838 in the calculation of the normalization factor.

Event yields in the $V\gamma$ control region	
$Z\gamma$	24.6 ± 3.3
$Z+jets$	3.0 ± 1.5
diboson + triboson	6.7 ± 0.3
top	1.5 ± 0.5
Total prediction	35.8 ± 3.7
Data	57 ± 7.6

Table 5.11: Event yields in the $V\gamma$ control region. The $V\gamma$ scale factor of 1.77 is calculated by scaling up the $Z\gamma$ and $Z+jets$ backgrounds to account for the difference between the data and predicted total background.

1839 5.3.3 Estimation of backgrounds from charge misidentification

1840 If an electron's charge is mis-reconstructed, it can lead to a real opposite-sign lepton pair passing
 1841 the same-sign requirement in the event selection. There are two primary reasons this can occur:

- 1842 1. An electron emits a photon via bremsstrahlung which then converts into an electron-positron
 1843 pair, and the conversion track with the wrong electric charge is matched to the original electron.
 1844 This is the dominant process leading to charge flip, and it is highly dependent on the electron
 1845 η due to the different amount of detector material the electron passes through.
- 1846 2. The curvature of the electron's track is mis-measured, resulting in the wrong charge being
 1847 assigned. This process is dependent on the momentum of the electron, as its track becomes
 1848 more straight as the momentum of the electron increases.

1849 In order to estimate this background, the rate at which an electron's charge is misidentified is
 1850 calculated from $Z \rightarrow ee$ MC simulation. It is known that the MC does not perfectly model the
 1851 material effects leading to charge flip; as a result, scale factors are applied to the MC in order for it
 1852 to better reflect the real performance. These scale factors are obtained from the ratio of charge
 1853 mis-ID rates in data and uncorrected MC in [83] following the method outlined in [111]. Once the
 1854 scale factors are applied, the charge misidentification rate ε can be extracted by comparing the

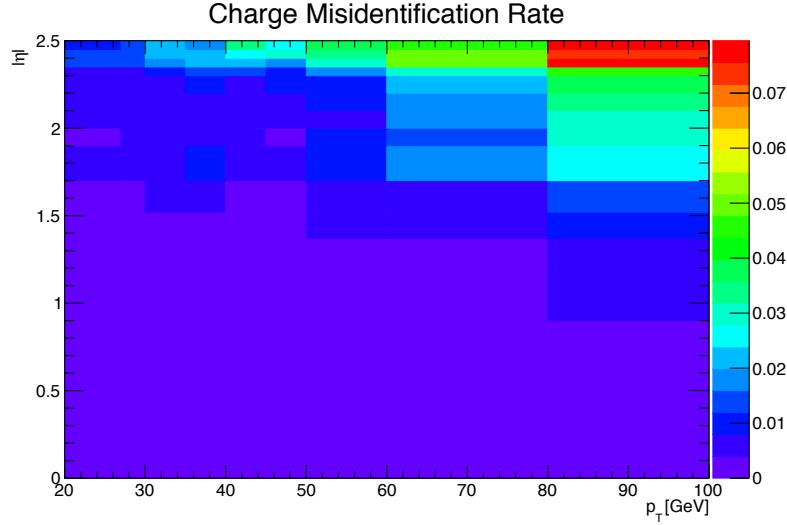


Figure 5.10: Charge misidentification rates for electrons as a function of $|\eta|$ and p_{T} . Rates are calculated from $Z \rightarrow e^+e^-$ MC after applying scale factors to approximate the charge mis-ID rates in data.

1855 electron's reconstructed charge with the charge of its truth particle:

$$\varepsilon(\eta, p_{\text{T}}) = \frac{N_{\text{wrong charge}}}{N_{\text{prompt electrons}}} \quad (5.1)$$

1856 The charge mis-ID rate is calculated in bins of electron $|\eta|$ and p_{T} , and it varies from below 0.1%
 1857 in the central region of the detector up to 8% in the forward regions for high p_{T} (above 80 GeV)
 1858 electrons. A two-dimensional plot of ε can be found in Figure 5.10.

1859 Given the charge flip rate $\varepsilon(\eta, p_{\text{T}})$, the rate at which an electron has its charge correctly recon-
 1860 structed is $(1 - \varepsilon)$. Thus there are three possible combinations of charge identification, assuming a
 1861 two-electron event:

- 1862 1. Both electrons are reconstructed correctly: $(1 - \varepsilon)^2$
- 1863 2. Both electrons are mis-reconstructed: ε^2
- 1864 3. Only one electron is mis-reconstructed: $2\varepsilon(1 - \varepsilon)$

1865 In order to estimate the size of the background from charge misidentification, opposite-sign events
 1866 are selected using the default event selection for a given signal or control region with the same-sign
 1867 requirement inverted. These events are then weighted by the probability for one of the electrons to

1868 be reconstructed with the wrong charge:

$$\omega = \frac{\varepsilon_1(1 - \varepsilon_2) + \varepsilon_2(1 - \varepsilon_1)}{(1 - \varepsilon_1)(1 - \varepsilon_2) + \varepsilon_1\varepsilon_2} \quad (5.2)$$

1869 where the subscripts 1 and 2 refer to the leading and subleading electrons, respectively, and ε_i is a
 1870 function of the η and p_T of the i^{th} electron. In the case of an event with one electron and one muon
 1871 ($\varepsilon_\mu = 0$), Equation 5.2 simplifies:

$$\omega = \frac{\varepsilon}{1 - \varepsilon} \quad (5.3)$$

1872 This method assumes that there is little contamination from fake electrons in the opposite-sign
 1873 sample, and this has been verified with MC simulation.

1874 Additionally, charge-flipped electrons tend to be reconstructed with lower energy when compared
 1875 to electrons with the correct charge. This is due to energy loss from the material interactions that
 1876 can cause the charge to be misidentified in the first place. A correction factor is calculated from
 1877 MC simulations, comparing the p_T of the truth electron to its reconstructed counterpart:

$$\alpha = \frac{\left(\frac{p_T^{\text{reco}}}{p_T^{\text{truth}}} - 1\right)_{\text{correct charge}}}{\left(\frac{p_T^{\text{reco}}}{p_T^{\text{truth}}} - 1\right)_{\text{wrong charge}}} \quad (5.4)$$

1878 The correction is then applied to the p_T of the charge-flipped electron via

$$p_T = p_T^0 / (1 + \alpha) + dE \quad (5.5)$$

1879 where p_T^0 is the uncorrected p_T of the electron and dE is a gaussian smearing factor centered at zero
 1880 with a width related to the energy resolution. Since which electron is mis-reconstructed is never
 1881 determined in this method, in the case of a two-electron event, the energy correction is applied
 1882 randomly to one of the two electrons based on the probabilities for them to be charge-flipped. This
 1883 also determines the overall sign of the event; the charge of the electron that does not receive the
 1884 correction is taken to be the charge for both.

1885 Systematic uncertainties on the charge mis-ID rates are calculated by generating two additional
 1886 sets of rates with the uncertainties on the scale factors varied up and down. The size of the esti-
 1887 mated charge flip background without the energy correction applied is also taken as a systematic
 1888 uncertainty. These systematic uncertainties are estimated to be approximately $\pm 15\%$.

1889 5.3.3.1 Validation of the charge misidentification estimate

1890 The performance of the charge misidentification estimation is tested in the same-sign inclusive
 1891 validation region (VR), defined in Table 5.12. For ee events, the mass of the dilepton pair is required

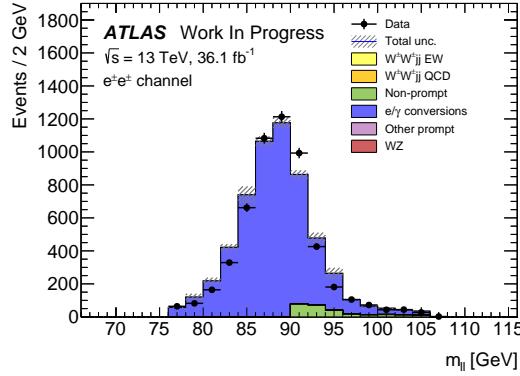


Figure 5.11: Dilepton invariant mass distribution m_{ll} for the ee channel in the same-sign inclusive VR.

1892 to lie within 15 GeV of the Z boson mass to increase the purity of the charge flip background.
 1893 $t\bar{t}$ production, which can contribute to both the charge mis-ID and fake lepton backgrounds, is
 1894 suppressed by the b -jet veto. The di-electron invariant mass is shown in Figure 5.11, and distributions
 1895 of the leading and subleading electron p_T in the ee -channel are shown in Figure 5.12 with the Z
 1896 mass cut inverted. Agreement between data and prediction is seen within the total statistical and
 1897 systematic uncertainties in the VR.

Same-sign inclusive VR
Exactly 2 same-sign signal leptons
$p_T > 27$ GeV for both leptons
$m_{ll} > 20$ GeV
$ m_{ee} - m_Z > 15$ GeV ($e^\pm e^\pm$ -channel only)
$N_{b\text{-jet}} = 0$

Table 5.12: Selection criteria for the same-sign inclusive validation region.

1898 5.3.4 Estimation of non-prompt backgrounds with the fake-factor method

1899 Events with one prompt lepton produced in association with hadronic jets can pass the event selection
 1900 if a jet is misidentified as a charged lepton or if a non-prompt lepton from the decay of a heavy
 1901 flavor particle (such as b - and c -hadrons) passes the signal lepton criteria. These misidentified jets
 1902 and non-prompt leptons are collectively referred to as *fake leptons*, or simply *fakes*. The rate at
 1903 which a fake lepton is misidentified is generally not modelled well enough by the MC to accurately
 1904 estimate their contributions directly from simulation. Therefore, a data-driven technique called the

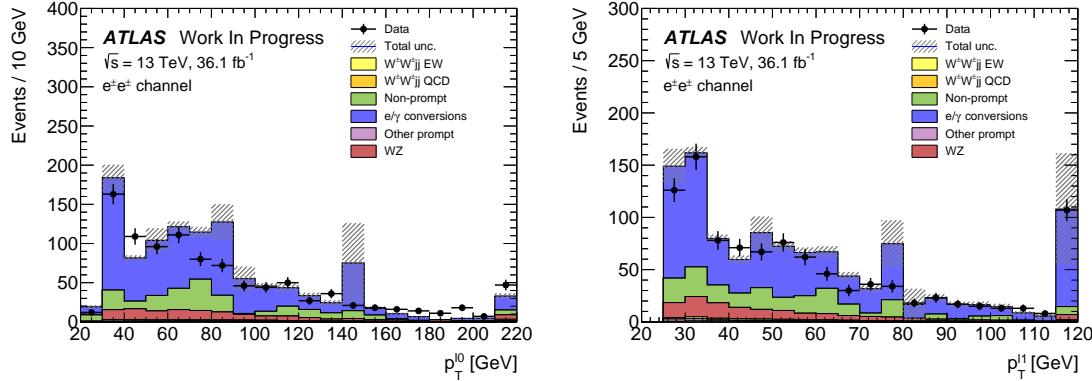


Figure 5.12: p_T distributions for the leading (left) and subleading (right) electron for the ee channel in the same-sign inclusive VR. In these plots, the cut requiring m_{ee} to fall within the Z mass window has been inverted in order to test the modelling away from the Z peak.

1905 *fake-factor* is used to estimate the size and shape of background processes from fake leptons. In this
 1906 analysis, a new modification to the fake-factor is used involving the particle isolation variables; the
 1907 method is outlined first in the context of the *default* fake-factor in Section 5.3.4.1, and the modified
 1908 fake-factor is covered in Section 5.3.4.2.

1909 5.3.4.1 Overview of the default fake-factor method

1910 The goal of the fake-factor method is to measure the fake rate from real collision events in a region
 1911 enriched in fake leptons and use it to estimate the size of the fake lepton background in a chosen
 1912 signal or control region. This is done by creating two samples using different lepton definitions:

- 1913 1. The *nominal* sample is made up of leptons passing the signal selection.
 1914 2. The *loose* sample is made up of leptons that fail the signal selection while still passing a
 1915 loosened set of criteria. This sample is enriched in fake leptons and is orthogonal to the set of
 1916 nominal leptons.

1917 Using the sets of nominal and loose leptons, a fake-factor f can be calculated from a region enriched
 1918 in processes that are prone to producing fake leptons:

$$f = \frac{N_{\text{nominal}}}{N_{\text{loose}}} \quad (5.6)$$

1919 Since the fake rate is not expected to be constant over the entire phase space, the fake-factor can
 1920 be divided into bins:

$$f(b) = \frac{N_{\text{nominal}}(b)}{N_{\text{loose}}(b)} \quad (5.7)$$

1921 where b represents the bin number. In this analysis, the fake-factor is binned in lepton p_T .

1922 In order to estimate the fake background contribution in a given signal or control region, the
 1923 fake-factor is applied to a second control region with a selection identical to the region of interest
 1924 except one of the leptons required to satisfy the loose criteria. The region for which the background
 1925 is estimated contains two nominal leptons and is referred to as *nominal+nominal* (*NN*), and the
 1926 associated control region where the fake-factor is applied contains one nominal and one loose lepton
 1927 and is referred to as *nominal+loose* (*NL*). The fake background in a *NN* region can then be
 1928 calculated as:

$$N_{NN}^{\text{fake bkg.}} = \sum_b f(b) N_{NL}(b) \quad (5.8)$$

1929 Backgrounds containing two prompt leptons can also enter the *NL* region if one of the leptons
 1930 passes the nominal selection and the other passes the loose selection. Since the fake-factor method
 1931 estimates the fake background by scaling the amount of non-prompt events in the *NL* region, if these
 1932 prompt contributions are not removed, they will be included in the scaling, and the background will
 1933 be overpredicted. The final estimate of the fake background becomes:

$$N_{NN}^{\text{fake bkg.}} = \sum_b f(b) (N_{NL}(b) - N_{NL}^{\text{prompt}}(b)) \quad (5.9)$$

1934 A visual representation of the fake background estimation process is shown in Figure 5.13.

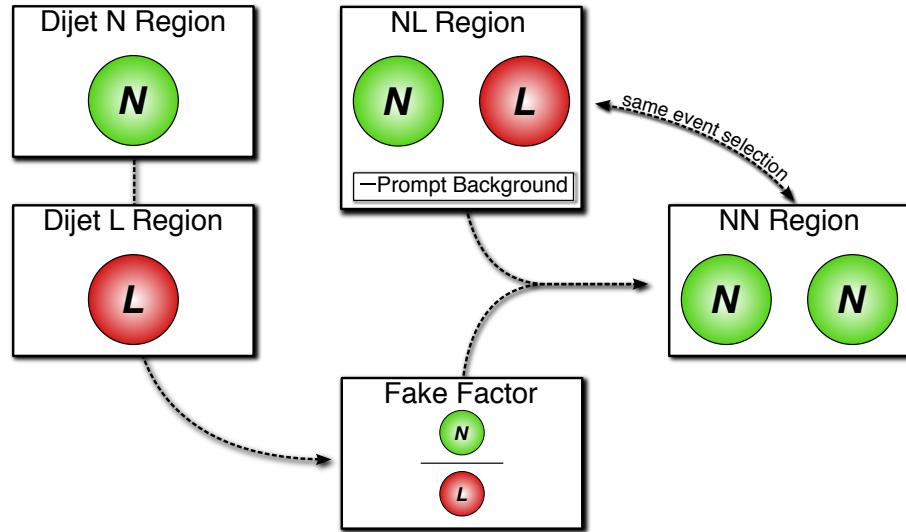


Figure 5.13: Graphical representation of how the fake factor method is used to estimate the fake background in a given *NN* region.

1935 **5.3.4.2 The fake-factor with p_T^{cone}**

1936 When a jet produces a non-prompt lepton, that lepton only carries a fraction of the underlying jet's
 1937 total momentum. Due to the isolation cut applied to the nominal leptons, they typically carry a
 1938 much larger percentage of the underlying jet momentum than the loose leptons. Since the isolation
 1939 essentially sets a limit on the amount of detector activity allowed around the lepton, if other nearby
 1940 particles carried a significant amount of momentum, the lepton would likely fail this cut.

1941 This discrepancy in the underlying jet momentum fraction can cause problems in the calculation
 1942 of the fake-factor f . Consider the case of two separate events with muons of identical momentum,
 1943 but one passes the nominal selection, and the other passes the loose selection. The loose lepton on
 1944 average will originate from a jet with higher p_T than the one corresponding to the nominal lepton
 1945 despite both muons having the same momentum. This can be seen explicitly when comparing the
 1946 p_T of a muon to its associated truth jet:

$$\Delta p_T(\mu, j) = \frac{p_T(j) - p_T(\mu)}{p_T(j) + p_T(\mu)} \quad (5.10)$$

1947 Since muons are not included in the jet reconstruction algorithm, Δp_T approximates the momentum
 1948 of the muon compared to the rest of the jet. For muons that carry more than 50% of the jet's
 1949 momentum, Δp_T will be negative and vice-versa. The Δp_T distributions for nominal and loose
 1950 muons in $t\bar{t}$ MC events is shown Figure 5.14, where a 25 GeV nominal muon on average corresponds
 1951 to a 35 GeV jet, and a 25 GeV loose muon on average corresponds to a 62 GeV jet¹⁷.

1952 Since the default fake-factor defined in Equation 5.7 is binned in lepton p_T , the 25 GeV muons
 1953 in the example above would occupy the same bin despite originating from very different jets. As
 1954 a result, within a given bin, the underlying jet p_T spectrum can differ substantially between the
 1955 numerator and the denominator. Additionally, these differences can vary depending on the process
 1956 producing the non-prompt leptons or on the specific kinematic selections of the signal or control
 1957 regions where the fake-factor is applied.

1958 Fortunately, the majority of the jet momentum not carried by the non-prompt lepton (excluding
 1959 neutrinos) can be recovered using isolation variables. A track-based isolation is chosen, referred
 1960 to as p_T^{cone} , and it contains the sum of the p_T of all particle tracks with $p_T > 1$ GeV originating
 1961 from the primary vertex within a cone of $\Delta R < 0.3$ around the lepton. Thus, the sample of loose
 1962 leptons in the denominator of the fake-factor calculation is binned in $p_T + p_T^{\text{cone}}$ rather than simply

¹⁷To better illustrate the point, here the muon is added back into the jet p_T , and the corresponding muon p_T is obtained via $\Delta p_T(\mu, j) = \frac{(p_T(j) - p_T\mu) - p_T(\mu)}{(p_T(j) - p_T\mu) + p_T(\mu)} = \frac{p_T(j) - 2p_T(\mu)}{p_T(j)}$.

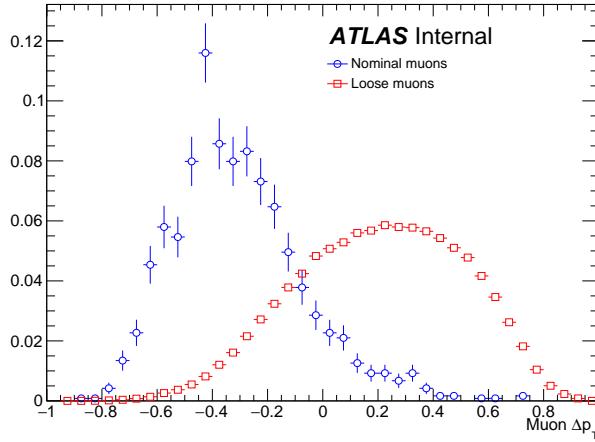


Figure 5.14: Δp_T distributions for nominal (blue) and loose (red) muons in simulated $t\bar{t}$ events. Each muon has been matched to a truth-level jet. Both distributions are normalized to unit area.

lepton p_T . Adding the isolation cone greatly reduces the difference in the fraction of the underlying jet momentum carried by the nominal and loose leptons. To check this, a new Δp_T is calculated between a lepton and its matched truth jet, where the truth jet p_T has been corrected to include all muons within a cone of $\Delta R < 0.4$:

$$p_T(j) = p_T(j_{\text{truth}}) + \sum_{\Delta R < 0.4} p_T(\mu_{\text{truth}}) \quad (5.11)$$

The Δp_T distributions comparing p_T and $p_T + p_T^{\text{cone}}$ for nominal and loose leptons using the corrected jet p_T are found in Figure 5.15, and better agreement is seen between the numerator (nominal) and denominator (loose with $p_T + p_T^{\text{cone}}$) distributions.

The numerator remains binned in lepton p_T , due to the fact that it is meant to mirror the signal region as closely as possible, and the signal lepton selection does not use $p_T + p_T^{\text{cone}}$. The impact of this is expected to be negligible due to the p_T^{cone} isolation being small for signal leptons, as shown for muons in Figure 5.16. Finally, the fake-factor f becomes:

$$f(b) = \frac{N_{\text{nominal}}(b(p_T))}{N_{\text{loose}}(b(p_T + p_T^{\text{cone}}))} \quad (5.12)$$

5.3.4.3 Application of the fake-factor

The fake-factor itself is measured from a sample of collision events passing a dijet selection that requires exactly one lepton (either passing the nominal or loose selections) and at least one jet.

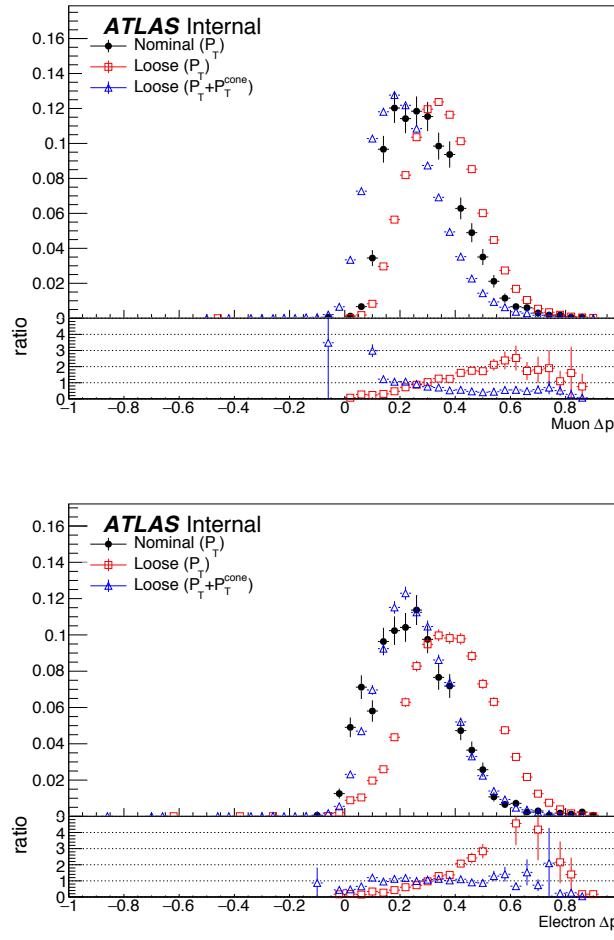


Figure 5.15: Δp_T distributions for muons (top) and electrons (bottom) in simulated $t\bar{t}$ events. Each lepton has been matched to a truth-level jet, and that truth jet has had its p_T corrected to include all truth muons within a cone of $\Delta R < 0.4$. The nominal leptons are in black. Δp_T is calculated for the loose leptons using p_T (red) and $p_T + p_T^{\text{cone}}$ (blue).

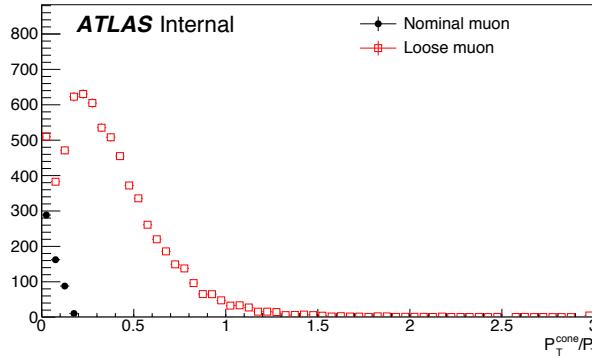


Figure 5.16: Distributions of p_T^{cone}/p_T for nominal (black) and loose (red) muons in simulated $t\bar{t}$ events.

1977 The leading jet must also be b -tagged and approximately back-to-back with the lepton in order to
 1978 enhance non-prompt lepton contributions while reducing contributions from processes involving W
 1979 and Z bosons. W boson events are further suppressed by requiring the sum of the E_T^{miss} and the
 1980 transverse mass of the lepton to be less than 50 GeV. The full event selection for the dijet region is
 1981 summarized in Table 5.13.

Dijet event selection
Event preselection
Exactly one lepton with $p_T > 15$ GeV
$N_{\text{jet}} > 0$
Leading jet is b -tagged
$p_T^{\text{lead. jet}} > 25$ GeV
$p_T^{\text{lead. jet}} > 30$ GeV if $ \eta_j > 2.5$
$ \Delta\phi(l, \text{lead. jet}) > 2.8$
$m_T(l, E_T^{\text{miss}}) + E_T^{\text{miss}} < 50$ GeV

Table 5.13: Event selection for the dijet region used for calculating the fake-factor. The selected lepton can pass either the nominal (signal) or loose selections. In the case of the nominal leptons, the $p_T > 27$ GeV requirement is replaced with $p_T > 15$ GeV.

1982 The numerator sample is constructed from dijet events in which the lepton passes the nominal
 1983 (signal) selection and is binned in the lepton p_T . Similarly, the denominator sample is made up of
 1984 the remaining dijet events where the lepton passes the loose selection and is binned in the lepton
 1985 $p_T + p_T^{\text{cone}}$. The nominal and loose leptons pass the signal selection¹⁸ and loose selection, respectively,
 1986 defined earlier in Table 5.3 for muons and Table 5.4 for electrons. Backgrounds from $W + \text{jets}$, $Z + \text{jets}$,

¹⁸The $p_T > 27$ GeV cut in the signal lepton selection is dropped in favor of the $p_T > 15$ GeV requirement in the dijet selection.

1987 $t\bar{t}$, and single top processes are estimated from MC simulations requiring one lepton to be prompt
 1988 using the truth information; these contributions are subtracted from the dijet data. The fake-factor
 1989 is then calculated using Equation 5.12 for muons and for central and forward electrons separately.
 1990 The muon fake-factor is shown in Figure 5.17, and the two electron fake-factors (one each for central
 1991 and forward electrons) are shown in Figure 5.18. The numerical values of the fake-factors, including
 1992 their systematic uncertainties (which will be discussed in Section 5.3.4.4) are listed in Table 5.14.

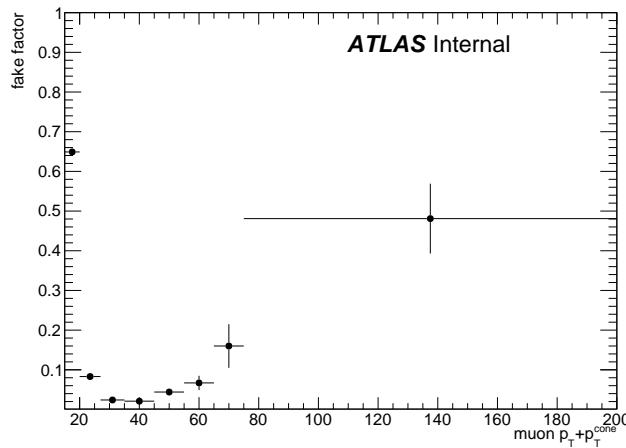


Figure 5.17: The measured fake-factor as a function of muon $p_T + p_T^{\text{cone}}$. The error bars represent the statistical uncertainty only.

1993 In order to properly account for the denominator being binned in $p_T + p_T^{\text{cone}}$, special care needs
 1994 to be taken when estimating the fake background from the *NL* regions. For the purposes of the
 1995 fake-factor calculation, it is perhaps more intuitive to consider a loose *object* with $p_T = p_T + p_T^{\text{cone}}$
 1996 instead of simply a loose lepton, as the lepton and the underlying jet are treated as a whole with this
 1997 method. When the lepton p_T cuts required by a particular signal or control region are applied to
 1998 nominal and loose leptons, the cut is applied to the p_T of the nominal lepton and to the $p_T + p_T^{\text{cone}}$ of
 1999 the loose object. Similarly, when looking up the fake-factor weight for a given *NL* event, the value
 2000 is taken from the bin corresponding to the $p_T + p_T^{\text{cone}}$ of the loose object. Finally, when applying the
 2001 weight to the *NL* event, $p_T + p_T^{\text{cone}}$ is assigned as the p_T of the loose object. This can be visualized
 2002 by referring back to Figure 5.13; every time a loose lepton is used (the red circles in the Figure),
 2003 $p_T + p_T^{\text{cone}}$ is used in place of p_T .

2004 Finally, it should be noted that the addition of p_T^{cone} to the loose object may cause the loose
 2005 leptons in the denominator sample to migrate into higher bins. This results in an overall decrease in

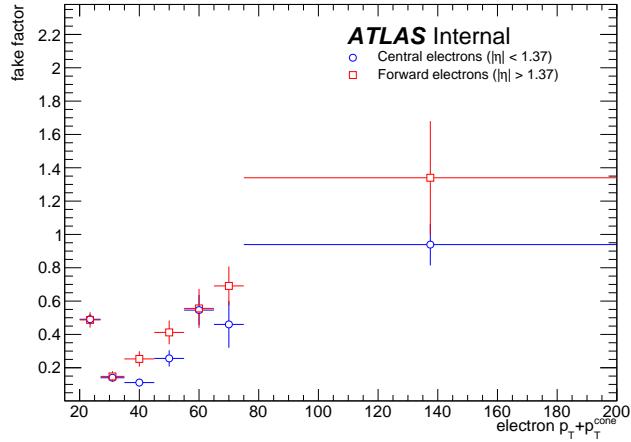


Figure 5.18: The measured fake-factor as a function of electron $p_T + p_T^{\text{cone}}$ in the central ($|\eta| < 1.37$, blue) and forward ($|\eta| > 1.37$, red) regions of the detector. The error bars represent the statistical uncertainty only.

the number of loose objects in the lower $p_T + p_T^{\text{cone}}$ bins due to there not being additional leptons at lower p_T to replace them. Since the fake-factor is a ratio of the number of events in a bin, this effect causes the first few bins of the fake-factor to increase, as can be seen clearly in Figure 5.17. However, the signal and control regions (and their corresponding NL regions) contain a $p_T > 27$ GeV cut that prevents these migrations from negatively impacting the fake estimation.

5.3.4.4 Systematic uncertainties

Four sources of systematic uncertainty are considered: the dijet event selection, prompt background subtraction, jet flavor composition, and residual dependence on the underlying jet p_T spectrum. In order to measure the impact of these systematics, new fake-factors are computed with variations in each of the systematic sources and the differences from the nominal values are taken as the uncertainty.

- In order to estimate uncertainties due to the dijet selection, the cut on $M_T + E_T^{\text{miss}}$ is varied by ± 5 GeV, the jet-lepton separation $\Delta\phi(l, j)$ by ± 0.1 , and the jet p_T cut by ± 5 GeV.
- To estimate the systematic uncertainty on the prompt background subtraction, the MC prediction in a $W+\text{jets}$ control region is compared to data. The discrepancy between data and MC is found to be approximately 10% [83]. Therefore, the prompt background used for the subtraction is scaled up and down by $\pm 10\%$.

- 2023 3. The difference in the jet flavor composition between the dijet events and the events in the
 2024 NL regions can affect the accuracy of the fake background estimation. The dijet sample is
 2025 dominated by light jets, while the NL regions tend to be dominated by heavy flavor from $t\bar{t}$.
 2026 To account for this, the fake-factor is computed with a b -jet veto.
- 2027 4. To measure any residual dependence on the underlying jet p_T spectrum, the leading jet p_T
 2028 distribution is reweighted to match the p_T spectrum of truth jets that produce fake leptons
 2029 in MC simulations. This results in an increase in the number of nominal and loose leptons at
 2030 high momentum [83].

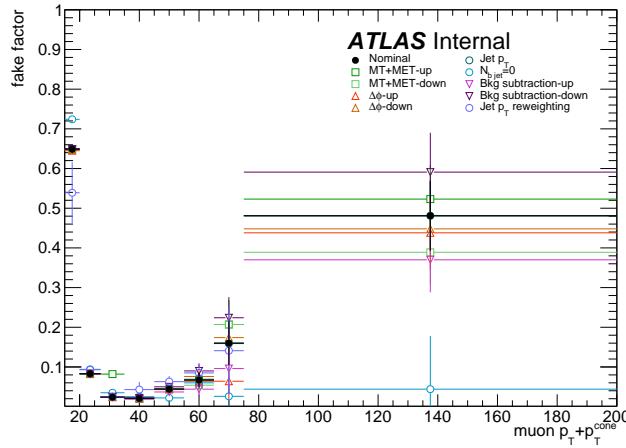


Figure 5.19: Systematic variations in the fake-factor as a function of muon $p_T + p_T^{\text{cone}}$. The individual fake-factors obtained for each systematic variation are displayed with their statistical uncertainties.

2031 5.3.4.5 Results of the fake-factor

- 2032 The fake background contribution in the signal region is estimated by applying the fake-factors
 2033 to the equivalent NL region using Equation 5.9, where the fake-factor used corresponds to the
 2034 flavor of the loose lepton in the event. As usual, the prompt background is subtracted from the
 2035 NL events using MC simulation. Charge misidentification is handled using the same method as
 2036 in Section 5.3.3, with an additional set of charge flip rates calculated for loose leptons. The fake
 2037 background yields in the signal region are listed in Table 5.15. An overall uncertainty of 50% is
 2038 assigned to the fake background estimation in $\mu^\pm\mu^\pm$ events, and between 40% to 90% for $e^\pm e^\pm$ and
 2039 $\mu^\pm e^\pm$ events, including both statistical and systematic effects.

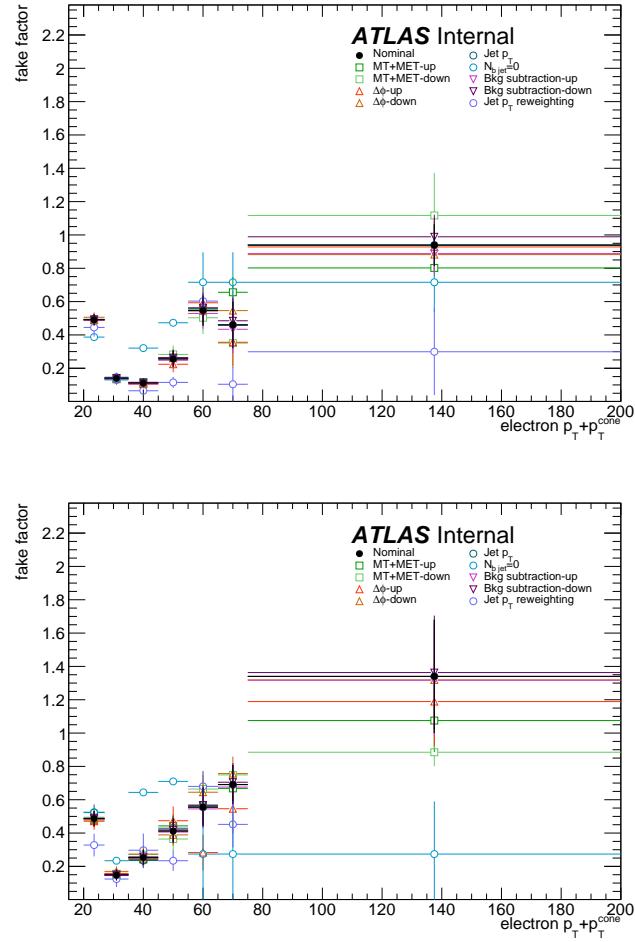


Figure 5.20: Systematic variations in the fake-factor as a function of electron $p_T + p_T^{\text{cone}}$ in the central ($|\eta| < 1.37$, top) and forward ($|\eta| > 1.37$, bottom) regions of the detector. The individual fake-factors obtained for each systematic variation are displayed with their statistical uncertainties.

fake-factor	p_T [15, 20]	p_T [20, 27]	p_T [27, 35]	p_T [35, 45]	p_T [45, 55]	p_T [55, 65]	p_T [65, 75]	p_T [75, 200]
nominal	0.649 ± 0.007	0.083 ± 0.002	0.024 ± 0.002	0.021 ± 0.003	0.044 ± 0.007	0.067 ± 0.018	0.160 ± 0.055	0.481 ± 0.088
MT+MET	0.649 ± 0.007	0.082 ± 0.002	0.082 ± 0.002	0.020 ± 0.003	0.045 ± 0.007	0.068 ± 0.018	0.207 ± 0.062	0.523 ± 0.086
$\Delta\phi(\ell, j)$	0.648 ± 0.007	0.083 ± 0.003	0.024 ± 0.002	0.022 ± 0.004	0.044 ± 0.007	0.054 ± 0.020	0.207 ± 0.060	0.389 ± 0.081
Jet p_T	0.645 ± 0.008	0.083 ± 0.003	0.024 ± 0.002	0.021 ± 0.004	0.045 ± 0.008	0.064 ± 0.021	0.064 ± 0.058	0.438 ± 0.092
$N_{b\text{-jet}} = 0$	0.646 ± 0.006	0.083 ± 0.002	0.024 ± 0.002	0.020 ± 0.003	0.043 ± 0.006	0.076 ± 0.017	0.174 ± 0.050	0.448 ± 0.078
Bkg. subtraction	0.648 ± 0.007	0.083 ± 0.002	0.024 ± 0.002	0.019 ± 0.003	0.037 ± 0.007	0.044 ± 0.019	0.096 ± 0.062	0.370 ± 0.082
Jet p_T Reweighting	0.649 ± 0.007	0.083 ± 0.002	0.025 ± 0.002	0.022 ± 0.003	0.050 ± 0.007	0.090 ± 0.017	0.224 ± 0.052	0.591 ± 0.099

(a) Fake-factor values for muons.

fake-factor	p_T [20, 27]	p_T [27, 35]	p_T [35, 45]	p_T [45, 55]	p_T [55, 65]	p_T [65, 75]	p_T [75, 200]
nominal	0.491 ± 0.031	0.140 ± 0.020	0.111 ± 0.023	0.256 ± 0.049	0.546 ± 0.091	0.460 ± 0.140	0.939 ± 0.125
MT+MET	0.493 ± 0.030	0.138 ± 0.019	0.115 ± 0.022	0.261 ± 0.045	0.559 ± 0.084	0.656 ± 0.091	0.802 ± 0.016
$\Delta\phi(\ell, j)$	0.488 ± 0.032	0.137 ± 0.020	0.110 ± 0.025	0.283 ± 0.053	0.503 ± 0.097	0.351 ± 0.149	1.117 ± 0.255
Jet p_T	0.489 ± 0.035	0.134 ± 0.021	0.105 ± 0.025	0.224 ± 0.048	0.593 ± 0.093	0.356 ± 0.144	0.928 ± 0.177
$N_{b\text{-jet}} = 0$	0.506 ± 0.029	0.140 ± 0.018	0.111 ± 0.022	0.260 ± 0.046	0.545 ± 0.084	0.546 ± 0.120	0.882 ± 0.103
Bkg. subtraction	0.493 ± 0.032	0.146 ± 0.021	0.115 ± 0.024	0.259 ± 0.049	0.550 ± 0.091	0.460 ± 0.140	0.939 ± 0.125
Jet p_T Reweighting	0.387 ± 0.009	0.130 ± 0.008	0.321 ± 0.012	0.473 ± 0.015	0.716 ± 0.180	0.716 ± 0.180	0.716 ± 0.180

(b) Fake-factor values for central electrons ($|\eta| < 1.37$).

fake-factor	p_T [20, 27]	p_T [27, 35]	p_T [35, 45]	p_T [45, 55]	p_T [55, 65]	p_T [65, 75]	p_T [75, 200]
nominal	0.487 ± 0.046	0.148 ± 0.031	0.253 ± 0.046	0.412 ± 0.071	0.556 ± 0.117	0.691 ± 0.117	1.340 ± 0.340
MT+MET	0.483 ± 0.045	0.152 ± 0.031	0.241 ± 0.043	0.443 ± 0.070	0.565 ± 0.106	0.668 ± 0.117	1.075 ± 0.189
$\Delta\phi(\ell, j)$	0.495 ± 0.047	0.156 ± 0.033	0.271 ± 0.052	0.364 ± 0.074	0.664 ± 0.107	0.749 ± 0.056	0.885 ± 0.084
Jet p_T	0.471 ± 0.051	0.158 ± 0.035	0.247 ± 0.051	0.474 ± 0.085	0.283 ± 0.107	0.546 ± 0.149	1.189 ± 0.266
$N_{b\text{-jet}} = 0$	0.478 ± 0.042	0.170 ± 0.031	0.274 ± 0.046	0.389 ± 0.066	0.645 ± 0.104	0.757 ± 0.102	1.319 ± 0.326
Bkg. subtraction	0.523 ± 0.048	0.149 ± 0.033	0.235 ± 0.045	0.429 ± 0.073	0.555 ± 0.117	0.691 ± 0.117	1.340 ± 0.340
Jet p_T Reweighting	0.525 ± 0.011	0.234 ± 0.013	0.644 ± 0.016	0.710 ± 0.014	0.274 ± 0.316	0.274 ± 0.316	0.274 ± 0.316

(c) Fake-factor values for forward electrons ($1.37 < |\eta|$).Table 5.14: Values of the fake-factor in each p_T bin and for each individual systematic source.

	estimated yield	f_e stat. up	f_e stat. dn	f_e syst. up	f_e syst. dn	f_μ stat. up	f_μ stat. dn	f_μ syst. up	f_μ syst. dn
$e^\pm e^\pm$	11.42 ± 3.13	—	—	—	—	—	—	—	—
$\mu^\pm \mu^\pm$	4.82 ± 0.77	—	—	—	—	0.65	-0.65	3.64	-0.61
$\mu^\pm e^\pm$	37.08 ± 5.16	4.90	-4.90	5.59	-14.34	1.39	-1.39	16.10	-1.98

Table 5.15: Estimated yields for the fake lepton background. The estimated yield is shown in the first column together with the statistical uncertainty followed by the systematic uncertainties from variations of the the fake-factors within their statistical (stat.) and systematic (syst.) uncertainties. The labels f_e and f_μ indicate the fake-factors for electrons and muons, respectively.

2040 **5.3.4.6 Validation of the fake-factor**

2041 The accuracy of the fake-factor method is tested in several validation regions, the most sensitive
 2042 of which is the same-sign top fakes VR (SS top VR), defined in Table 5.16. This region inverts
 2043 the signal region's b -jet veto to accept events with exactly one b -jet. Due to this requirement, the
 2044 dominant source of events comes from the $t\bar{t}$ process where a b -jet fakes an isolated lepton. The
 2045 distribution of the subleading lepton p_T in this VR is shown in Figure 5.21 for all lepton flavor
 2046 combinations. There is good agreement between the data and the prediction, even when only taking
 2047 into account the statistical uncertainty and not the large systematic uncertainties assigned to the
 2048 fake estimation.

Same-sign top fakes VR
Exactly 2 same-sign signal leptons
$p_T > 27$ GeV for both leptons
$m_{ll} > 20$ GeV
$ m_{ee} - m_Z > 15$ GeV ($e^\pm e^\pm$ -channel only)
$N_{b\text{-jet}} = 1$
$N_{\text{jet}} \geq 2$
Leading jet $p_T > 65$ GeV
Subleading jet $p_T > 35$ GeV

Table 5.16: Selection criteria for the same-sign top fakes validation region.

2049 **5.3.5 Reduction of WZ background using custom overlap removal**

2050 The dominant source of prompt background in this analysis comes from WZ events where both
 2051 bosons decay leptonically. Traditionally, the background is dealt with by imposing a veto on any
 2052 event with a third lepton passing some loose identification criteria (the so-called *trilepton veto*).
 2053 In the case of this analysis, if one or more leptons in addition to the two signal leptons pass the
 2054 preselection criteria, the event is rejected. However, WZ events can still enter the signal region if
 2055 one of the leptons fails the preselection or falls outside of the detector's acceptance.

2056 In order to understand the sources of WZ events that are not removed by the trilepton veto, a
 2057 study was performed on truth-level leptons in $W^\pm W^\pm jj$ and WZ MC samples. Events with three
 2058 truth leptons were selected, and each was matched to its reconstruction-level partner by finding the
 2059 closest $\Delta R(\text{truth}, \text{reco})$ and $\Delta p_T(\text{truth}, \text{reco})$ match. For events surviving the trilepton veto, the
 2060 two signal leptons were removed, and the remaining leptons represent real leptons that failed to
 2061 be selected for the veto. Between 40-50% of these leptons fall outside of the eta acceptance of the

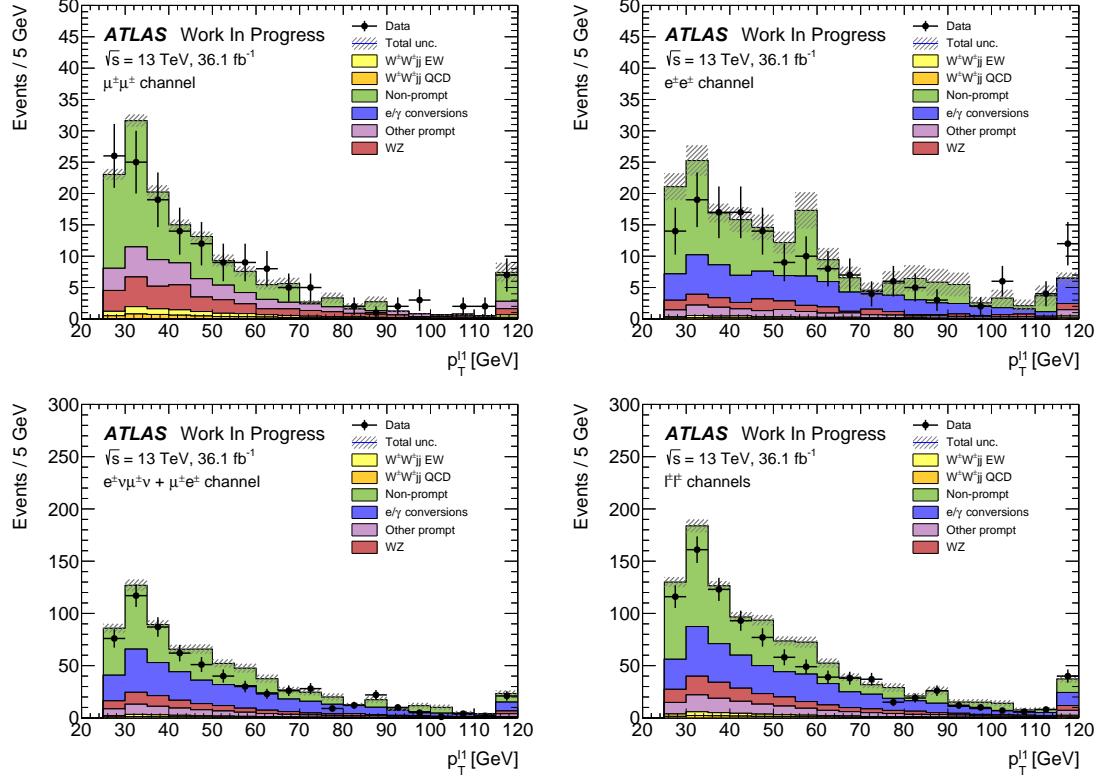


Figure 5.21: Distributions of the subleading lepton p_T in the same-sign top fakes VR for $\mu^\pm\mu^\pm$ events (top right), $e^\pm e^\pm$ events (top left), $\mu^\pm e^\pm$ events (bottom left), and all events combined (bottom right). All errors are statistical only.

analysis (see Figure 5.22) and are unrecoverable. The second largest source of leptons failing the preselection is the OR, defined in Section 5.2.1.4. The standard OR procedure appears to be too aggressive in removing leptons in favor of jets, causing many three lepton events to “lose” their third lepton and pass the trilepton veto. Therefore a *custom OR* is investigated which would replace the standard OR in the preselection and allow for better WZ rejection by removing fewer third leptons.

In order to construct this custom OR, a new quantity is defined between a lepton (l) and a nearby jet (j)

$$p_{T,\text{ratio}}(l, j) = \frac{p_{T,l}}{p_{T,j}} \quad (5.13)$$

which, along with $\Delta R(l, j)$, will make up the custom OR criteria. The idea behind including $p_{T,\text{ratio}}$ is to be able to preferentially remove background leptons originating from jets (those that carry a low percentage of the total jet momentum) instead of removing *any* lepton near a jet. The distributions of $p_{T,\text{ratio}}$ and the associated efficiency curves for muons and electrons can be found in Figures 5.23

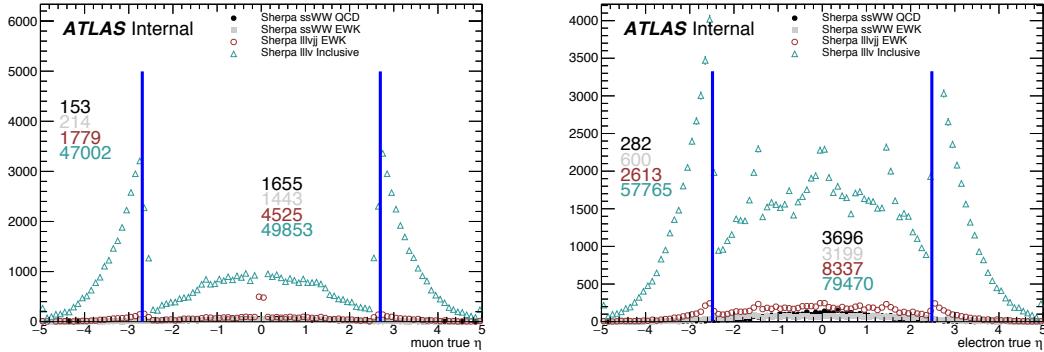


Figure 5.22: Pseudorapidity (η) distributions of truth muons (top) and electrons (bottom) for Sherpa $W^\pm W^\pm jj$ and WZ MC samples. The blue vertical lines represent the allowed η range for each lepton flavor. The numbers correspond to the number of raw MC events that fall within and outside of the allowed η range for each MC sample.

and 5.25, respectively, and the distributions for $\Delta R(\mu, j)$ for muons can be found in Figure 5.24. Since all electrons have an associated jet in the calorimeters, the $\Delta R(e, j)$ variable is not a good quantity to use for this custom OR.

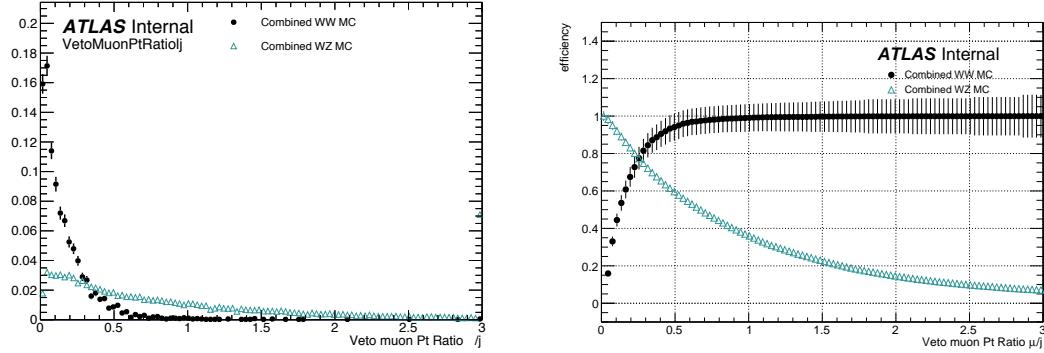


Figure 5.23: Distributions of $p_{T,\text{ratio}}(\mu, j)$ for EWK and QCD $W^\pm W^\pm jj$ signal (black) and WZ background (teal) for truth-matched third muons in events that pass the trilepton veto. Both distributions are normalized to unit area. The associated efficiency curves are on the right where efficiency is defined as the percentage of total events that would pass a cut on $p_{T,\text{ratio}}(\mu, j)$ at a given value on the x -axis.

A working point for the Custom OR was chosen by requiring 90% signal retention for muons and 90% background rejection for electrons. The cut on electrons was allowed to be much tighter because the number of signal events with a third electron is considerably smaller than for muons. It should be emphasized that the signal events present in Figures 5.23-5.25 do not represent the full

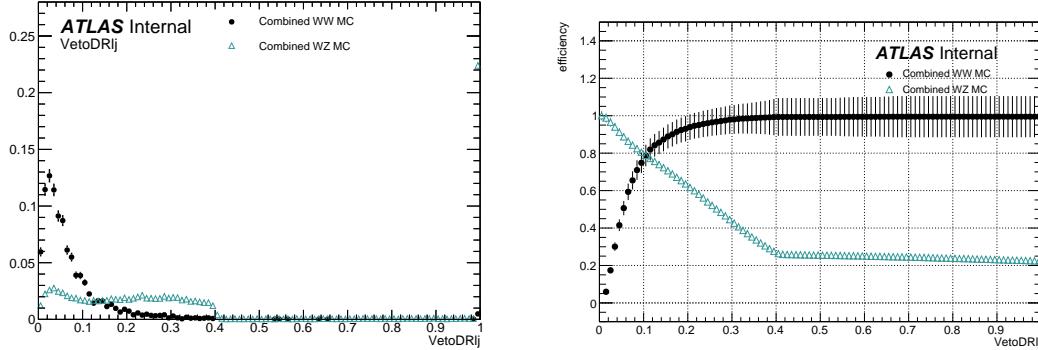


Figure 5.24: Distributions of $\Delta R(\mu, j)$ for EWK and QCD $W^\pm W^\pm jj$ signal (black) and WZ background (teal) for truth-matched third muons in events that pass the trilepton veto. Both distributions are normalized to unit area. The associated efficiency curves are on the right where efficiency is defined as the percentage of total events that would pass a cut on $\Delta R(\mu, j)$ at a given value on the x -axis.

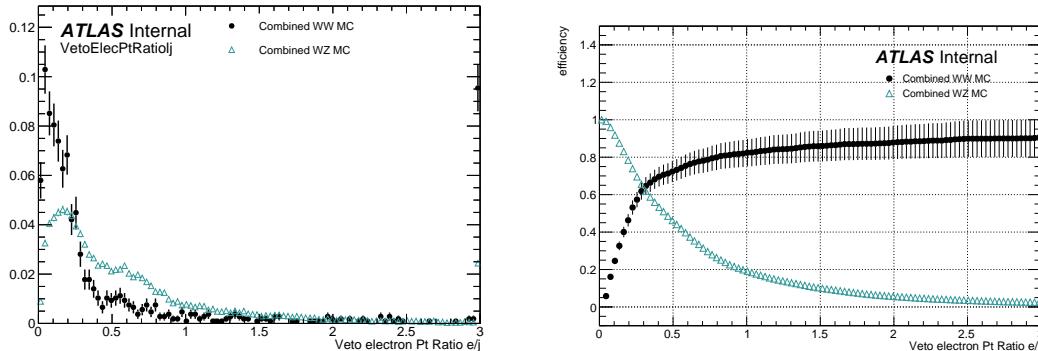


Figure 5.25: Distributions of $p_{T,\text{ratio}}(e, j)$ for EWK and QCD $W^\pm W^\pm jj$ signal (black) and WZ background (teal) for truth-matched third electrons in events that pass the trilepton veto. Both distributions are normalized to unit area. The associated efficiency curves are on the right where efficiency is defined as the percentage of total events that would pass a cut on $p_{T,\text{ratio}}(e, j)$ at a given value on the x -axis.

set of signal events, but only those with a real third lepton (which must come from some source other than the signal $W^\pm W^\pm jj$ process). For muons, a logical ‘or’ of $p_{T,\text{ratio}}(\mu, j)$ and $\Delta R(\mu, j)$ is used to maximize the third lepton acceptance due to correlations between the quantities, as shown in Figure 5.26; for electrons, only a cut on $p_{T,\text{ratio}}(e, j)$ is used. The Custom OR working point is defined in Table 5.17.

Custom OR Definition	
Muons	$p_{T,\text{ratio}}(\mu, j) > 0.40$ or $\Delta R(\mu, j) > 0.15$
Electrons	$p_{T,\text{ratio}}(e, j) > 0.18$

Table 5.17: Custom OR definition. Leptons must pass this selection in order to be counted for the trilepton veto.

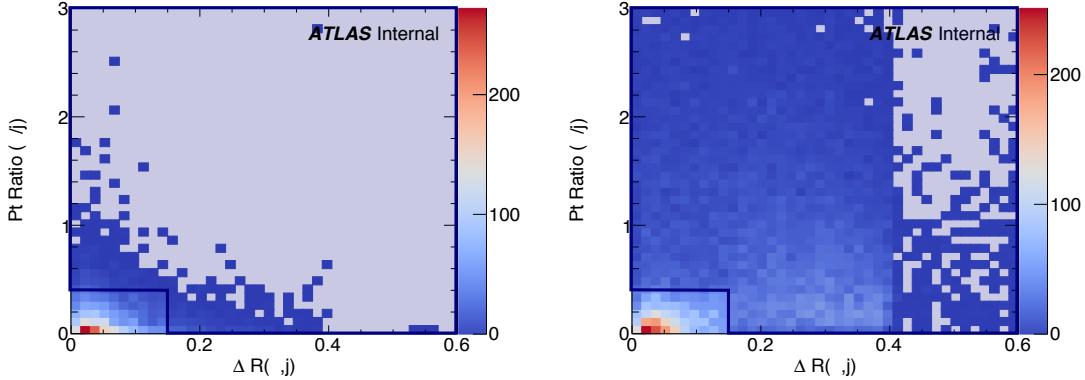


Figure 5.26: Two-dimensional plots of $p_{T,\text{ratio}}(\mu, j)$ vs $\Delta R(\mu, j)$ for truth-matched third muons in events that pass the trilepton veto for EWK and QCD $W^\pm W^\pm jj$ signal (left) and WZ background (right). The blue overlay indicates the area in which the third leptons will pass the custom OR and result in the event failing the trilepton veto.

Initial tests of the performance of the Custom OR yielded promising results, with approximately 20% reduction in WZ background compared to less than 2% signal loss in the signal region. Unfortunately, due to differences between the primary analysis framework and the one used for testing, in practice the gains in WZ rejection were not nearly as substantial, and ultimately the Custom OR was not included in the final analysis. However, it is still a potentially useful tool for improving background rejection based on lepton counting in analyses with overly aggressive OR procedures.

2091 **5.4 Cross section measurement**

2092 The $W^\pm W^\pm jj$ EWK cross section is extracted from the signal region using a maximum-likelihood
 2093 fit applied simultaneously to four m_{jj} bins in the signal region as well as to the low- m_{jj} and WZ
 2094 control regions. For the fit and cross section extraction, the signal region is defined as in Table 5.8
 2095 with the dijet invariant mass requirement raised to $m_{jj} > 500$ GeV. The low- m_{jj} region is defined
 2096 to mirror the signal region exactly with the dijet invariant mass inverted to $200 < m_{jj} < 500$ GeV,
 2097 and the WZ control region is as defined previously in Section 5.3.1.

2098 The signal and low- m_{jj} regions are split into six channels based on the flavor and charge of
 2099 the dilepton pair: $\mu^+\mu^+$, $\mu^-\mu^-$, μ^+e^+ , μ^-e^- , e^+e^+ , and e^-e^- . This split by charge increases the
 2100 sensitivity of the measurement due to the W^+/W^- charge asymmetry favoring the production of
 2101 W^+ bosons [112]. Since the signal events contain two W bosons, the signal strength compared to
 2102 charge-symmetric backgrounds is much greater in the $++$ channels than for both charges combined.
 2103 The WZ control region is included in the fit as a single bin ($l^\pm l^\mp l^\pm$).

2104 The maximum likelihood fit and cross section extractions are outlined in Sections 5.4.1 and
 2105 5.4.3, respectively. The results of the cross section measurement and of the analysis as a whole are
 2106 presented in Section 5.6.

2107 **5.4.1 Maximum likelihood fit**

2108 The number of predicted signal events in each channel c and m_{jj} bin b can be calculated from the SM
 2109 predicted total production cross section $\sigma_{\text{theo}}^{\text{tot}}$ scaled by the total integrated luminosity \mathcal{L} , the signal
 2110 acceptance \mathcal{A} , and the efficiency corrections $\mathcal{C}(\theta)$. Here θ represents the set of nuisance parameters
 2111 that parameterize the effects of each systematic uncertainty on the signal and background expec-
 2112 tations. The acceptance and efficiency corrections will be covered in more detail in Section 5.4.2.
 2113

$$N_{cb}^{\text{sig}}(\theta) = \sigma_{\text{theo}}^{\text{tot}} \mathcal{A}_b \mathcal{C}_b(\theta) \mathcal{L} \quad (5.14)$$

2114 A signal strength parameter μ is defined as the ratio of the measured cross section to the SM
 2115 predicted cross section. The expected number of events in a given channel and bin can then be
 2116 expressed as the sum of the estimated background ($N_{cb}^{\text{bkg}}(\theta)$) and the number of predicted signal
 2117 events scaled by μ :

$$\begin{aligned} N_{cb}^{\text{exp}}(\theta) &= \mu N_{cb}^{\text{sig}}(\theta) + N_{cb}^{\text{bkg}}(\theta) \\ &= \mu \sigma_{\text{theo}}^{\text{tot}} \mathcal{A}_b \mathcal{C}_b(\theta) \mathcal{L} + N_{cb}^{\text{bkg}}(\theta) \end{aligned} \quad (5.15)$$

2118 The nuisance parameters are constrained by Gaussian probability distribution functions, and
 2119 the normalization of the WZ background mentioned in Section 5.3.1 is included in the fit as a free
 2120 parameter. The expected yields for signal and background processes are adjusted by the set of
 2121 nuisance parameters within the constraints of the systematic uncertainties. The yields after the fit
 2122 correspond to the value that best matches the observed data.

2123 The number of events per channel and bin after the fit can be written as a sum of the predicted
 2124 event yields for each sample s :

$$\nu_{cb}(\phi, \theta, \gamma_{cb}) = \gamma_{cb} \sum_s [\eta_{cs}(\theta) \phi_{cs}(\theta) \lambda] h_{cbs}(\theta) \quad (5.16)$$

2125 In this equation, the fitted number of events in a given channel and bin is obtained by weighting
 2126 the histogram of predicted yields h_{cbs} by the product of a given luminosity λ and any normalization
 2127 factors ϕ_{cs} that may be given for each channel and sample. The input histogram and the normaliza-
 2128 tion factors may depend on the nuisance parameters θ taking into account sources of systematic
 2129 uncertainty. Uncertainties on the normalization factors $\eta_{cs}(\theta)$ are also included. Finally, bin-by-bin
 2130 scale factors γ_{cb} are included to parameterize the statistical uncertainties of the MC predictions.

2131 The binned likelihood function is given by a product of Gaussian functions for the luminosity
 2132 and for the background uncertainties and a product of Poisson functions for the number of observed
 2133 events in each bin and channel:

$$L(\mu|\theta) = \mathcal{G}(\mathcal{L}|\theta_{\mathcal{L}}, \sigma_{\mathcal{L}}) \cdot \prod_c \prod_b \mathcal{P}(N_{cb}^{\text{meas.}}|\nu_{cb}(\mu)) \prod_p \mathcal{G}(\theta_p^0|\theta_p) \quad (5.17)$$

2134 where \mathcal{G} and \mathcal{P} are the Gaussian and Poisson functions, respectively. As before, \mathcal{L} represents the
 2135 integrated luminosity with uncertainty $\sigma_{\mathcal{L}}$ and associated nuisance parameter $\theta_{\mathcal{L}}$. The number of
 2136 measured events in a given bin and channel is represented by $N_{cb}^{\text{meas.}}$, and $\nu_{cb}(\mu)$ is the predicted
 2137 number of events defined in Equation 5.16 expressed as a function of the signal strength μ . Finally,
 2138 the set of nuisance parameters θ and any auxiliary measurements used to constrain them (θ^0) are
 2139 multiplied for each parameter p .

2140 The profile likelihood ratio is defined as

$$q_{\mu} = -2 \ln \frac{L(\mu, \hat{\theta}_{\mu})}{L(\hat{\mu}, \hat{\theta})} \quad (5.18)$$

2141 where $\hat{\mu}$ and $\hat{\theta}$ are the unconditional maximum likelihood estimates, and $\hat{\theta}$ is the conditional maxi-
 2142 mum likelihood estimate for a given value of μ . The fitted signal strength $\hat{\mu}$ is obtained by maximiz-
 2143 ing the likelihood function with respect to all parameters. The compatibility of the observed data

with the background-only hypothesis can then be calculated by setting $\mu = 0$. Observation of the $W^\pm W^\pm jj$ EWK process is claimed if the data is found to be inconsistent with the background-only hypothesis by more than 5σ .

5.4.2 Definition of the fiducial volume

Before extracting the cross section, it is necessary to define the fiducial volume, or the phase space of measureable events. It is a subset of the total phase space defined by selection requirements designed to mirror those applied in the analysis as closely as possible. The selection criteria for the fiducial volume are listed in Table 5.18.

Fiducial region selection	
Lepton selection	Two prompt leptons (e, μ) $p_T > 27$ GeV and $ \eta < 2.5$ for both leptons Both leptons with the same electric charge Dilepton invariant mass $m_{ll} > 20$ GeV Dilepton separation $\Delta R(ll) > 0.3$
Missing transverse energy	Two neutrino system with $p_T^{\nu\nu} > 30$ GeV
Jet selection	At least two jets Leading jet $p_T > 65$ GeV Subleading jet $p_T > 35$ GeV Leading and subleading jet $ \eta < 4.5$ Jet-lepton separation $\Delta R(l, j) > 0.3$ Dijet invariant mass $m_{jj} > 500$ GeV Dijet separation $\Delta y_{jj} > 2.0$

Table 5.18: Definition of the fiducial volume.

The full phase space is generated in MC simulations, providing the total theoretical cross section $\sigma_{\text{theo}}^{\text{tot}}$ and the total number of signal events $\mathcal{N}_{\text{sig}}^{\text{tot}}$ ¹⁹. After applying the fiducial selection at truth level, the total number of signal events in the fiducial region $\mathcal{N}_{\text{sig}}^{\text{fid}}$ is obtained. An acceptance factor \mathcal{A} is used to represent the efficiency of events falling inside the fiducial region at truth level:

$$\mathcal{A} = \frac{\mathcal{N}_{\text{sig}}^{\text{fid}}}{\mathcal{N}_{\text{sig}}^{\text{tot}}} \quad (5.19)$$

A correction factor \mathcal{C} is also necessary to translate from the truth level fiducial volume to the reconstruction level signal region and is defined in terms of the number of reconstruction level MC events in the signal region $N_{\text{sig},\text{MC}}^{\text{SR}}$:

$$\mathcal{C} = \frac{N_{\text{sig},\text{MC}}^{\text{SR}}}{\mathcal{N}_{\text{sig}}^{\text{fid}}} \quad (5.20)$$

¹⁹For the purpose of clarity, the number of events at truth level is represented by a script \mathcal{N} , and the number of events at reconstruction level uses a regular N .

2159 Since the fit is binned in m_{jj} , the acceptance and efficiency correction factors must be as well.
 2160 Therefore, \mathcal{A}_i and \mathcal{C}_{ij} are written in terms of truth m_{jj} bins i and reconstruction m_{jj} bins j . A
 2161 graphical representation of these regions and the use of the acceptance and correction factors can
 2162 be seen in Figure 5.27.

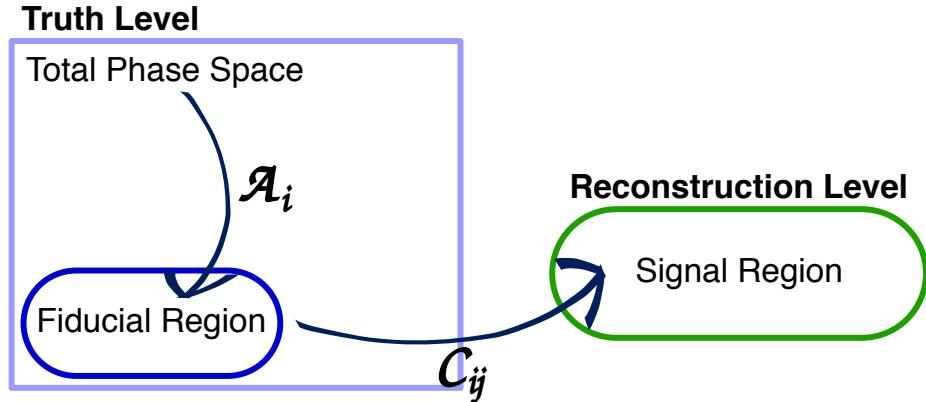


Figure 5.27: Visual representation of the different kinematic regions relevant to the cross section measurement. The acceptance factor \mathcal{A} converts from the truth level total phase space to the truth level fiducial region, and the efficiency correction \mathcal{C} translates the fiducial region into the reconstruction level signal region.

2163 **5.4.3 Cross section extraction**

2164 The $W^\pm W^\pm jj$ EWK fiducial cross section is measured using the signal strength parameter μ that is
 2165 determined by the maximum likelihood fit. This parameter is dependent on the nuisance parameters
 2166 θ and can be written explicitly in terms of the measured and theoretical cross sections as:

$$\mu(\theta) = \frac{\sigma_{\text{meas}}^{\text{SR}}}{\sigma_{\text{theo}}^{\text{SR}}} \quad (5.21)$$

2167 In the simple case with only one bin, the equation for the total number of expected events in the
 2168 signal region first introduced in Equation 5.15 can be written as:

$$N_{\text{exp}}^{\text{SR}}(\theta) = \mu(\theta) \cdot \sigma_{\text{theo}}^{\text{tot}} \cdot \mathcal{L} \cdot \mathcal{A} \cdot \mathcal{C}(\theta) + N_{\text{bkg}}^{\text{SR}}(\theta) \quad (5.22)$$

2169 with the unbinned versions of \mathcal{A} and \mathcal{C} defined in Equations 5.19 and 5.20, respectively.

2170 If the measured fiducial cross section is written as:

$$\sigma_{\text{meas}}^{\text{fid}} = \mu \cdot \mathcal{A} \cdot \sigma_{\text{theo}}^{\text{tot}} \quad (5.23)$$

2171 then Equation 5.22 can be rearranged to read:

$$\sigma_{\text{meas}}^{\text{fid}} = \frac{N_{\text{exp}}^{\text{SR}}(\theta) - N_{\text{bkg}}^{\text{SR}}(\theta)}{\mathcal{L} \cdot \mathcal{C}(\theta)} \quad (5.24)$$

2172 The measured fiducial cross section can finally be rewritten in terms of $\hat{\mu}$, which is the best estimator
2173 of the signal strength as extracted from the fit:

$$\begin{aligned} \sigma_{\text{meas}}^{\text{fid}} &= \hat{\mu}(\theta) \cdot \sigma_{\text{theo}}^{\text{tot}} \cdot \mathcal{A} \\ &= \hat{\mu}(\theta) \cdot \sigma_{\text{theo}}^{\text{fid}} \end{aligned} \quad (5.25)$$

2174 In practice, however, the cross section is not extracted from a single bin, and Equation 5.22
2175 becomes:

$$N_{\text{exp}}^{\text{SR}}(\theta) = \mu(\theta) \cdot \sigma_{\text{theo}}^{\text{tot}} \cdot \mathcal{L} \cdot \sum_i \mathcal{A}_i \sum_j \mathcal{C}_{ij} + \sum_j N_{\text{bkg},j}^{\text{SR}}(\theta) \quad (5.26)$$

2176 for a single channel in truth and reconstruction level m_{jj} bins i and j , respectively, where the binned
2177 versions of \mathcal{A}_i and \mathcal{C}_{ij} are used. This equation can be extended to include all the analysis channels
2178 by increasing the number of bins i and j . Additionally, it can be shown that Equation 5.25 holds
2179 for this more complex case as well [83], provided care is taken to ensure that all the uncertainties
2180 are handled properly.

2181 5.5 Summary of uncertainties

2182 Systematic uncertainties enter the final fit as nuisance parameters which can impact the estimated
2183 signal and background yields and the shapes of the m_{jj} distributions. These uncertainties can arise
2184 from the experimental methods or from the theoretical calculations used in the analysis. This section
2185 summarizes the systematic uncertainties; the experimental uncertainties are detailed in Section 5.5.1,
2186 and the theoretical uncertainties are covered in Section 5.5.2. The impacts of the systematic uncer-
2187 tainties on the final cross section measurement are summarized in Table 5.19.

2188 5.5.1 Experimental uncertainties

2189 Experimental uncertainties include detector effects as well as uncertainties on the background es-
2190 timation methods. Sources of systematic uncertainty on the measurement of physics objects are
2191 listed in Table 5.20, grouped by the relevant object type. For backgrounds estimated from MC
2192 simulations, variations in these sources of uncertainty are propagated through the analysis to obtain
2193 the corresponding uncertainties on the event yields. Additional experimental uncertainties include

Source	Impact [%]
Reconstruction	± 4.0
Electrons	± 0.5
Muons	± 1.2
Jets and E_T^{miss}	± 2.8
b -tagging	± 2.0
Pileup	± 1.5
Background	± 5.0
Misid. leptons	± 3.9
Charge misrec.	± 0.3
WZ	± 1.3
$W^\pm W^\pm jj$ QCD	± 2.8
Other	± 0.8
Signal	± 3.6
Interference	± 1.0
EW Corrections	± 1.3
Shower, Scale, PDF & α_s	± 3.2
Total	± 7.4

Table 5.19: Impact of various systematic effects on the fiducial cross section measurement. The impact of a given source of uncertainty is computed by performing the fit with the corresponding nuisance parameter varied up or down by one standard deviation from its nominal value.

2194 the integrated luminosity, the photon conversion rate from Section 5.3.2, and the data driven charge
 2195 misidentification and fake lepton background estimations from Sections 5.3.3 and 5.3.4.5, respec-
 2196 tively.

2197 The largest sources of experimental uncertainty on the MC estimations come from the jet-related
 2198 uncertainties and the b -tagging efficiency, while the largest uncertainty on the background estimation
 2199 comes from the fake-factor. The effects of the uncertainties on the $W^\pm W^\pm jj$ EWK signal and the
 2200 dominant MC estimated background, WZ , are listed in Tables 5.21 and 5.22, respectively. Since
 2201 the overall contributions from other processes estimated with MC are small, the uncertainties on
 2202 these backgrounds have a lesser impact on the final measurement; these tables can be found in
 2203 Appendix A.1.

2204 5.5.2 Theoretical uncertainties

2205 It is also necessary to consider uncertainties on the theoretical predictions in the fiducial region. They
 2206 include the choice of PDF set, the value of the strong coupling constant α_s , the renormalization
 2207 scale μ_R , the factorization scale μ_F , and the parton showering. The size of these uncertainties are
 2208 measured by generating new samples with variations in a chosen parameters and comparing them

Experimental uncertainties	
Electrons	Energy resolution
	Energy scale
	Identification efficiency
	Isolation efficiency
	Reconstruction efficiency
	Trigger efficiency
Muons	Energy scale
	Identification efficiency
	Inner detector track resolution
	Muon spectrometer resolution
	Trigger efficiency
E_T^{miss}	Resolution
	Scale
Jets	Energy resolution
	Energy scale
	JVT cut efficiency
	b -tagging efficiency
	Jets from pileup

Table 5.20: List of sources of experimental uncertainties on the reconstruction of physics objects.

$W^\pm W^\pm jj$ EWK	$e^\pm e^\pm$ % Yield	$\mu^\pm e^\pm$ % Yield	$\mu^\pm \mu^\pm$ % Yield
Jet-related Uncertainties	2.28	2.22	2.28
b -tagging efficiency	1.81	1.76	1.74
Pile-up	0.48	0.97	2.42
Trigger efficiency	0.02	0.08	0.47
Lepton reconstruction/ID	1.45	1.14	1.83
MET reconstruction	0.26	0.17	0.21

Table 5.21: Impact of experimental uncertainties for the $W^\pm W^\pm jj$ EWK processes in all channels.

WZ	$e^\pm e^\pm$ % Yield	$\mu^\pm e^\pm$ % Yield	$\mu^\pm \mu^\pm$ % Yield
Jet-related Uncertainties	9.58	5.03	8.45
b -tagging efficiency	2.49	2.23	2.40
Pile-up	2.99	3.49	3.33
Trigger efficiency	0.03	0.09	0.43
Lepton reconstruction/ID	1.52	1.24	3.07
MET reconstruction	0.93	0.79	1.63

Table 5.22: Impact of experimental uncertainties for the WZ process in all channels.

2209 to samples using the nominal choice of the parameter.

2210 For the signal sample, internal variations on the PDF sets as well as using a different set entirely
 2211 results in a relative uncertainty of up to 2.25% on the nominal sample. The impact from varying α_s is
 2212 very small, on the order of < 0.01%. The factorization and renormalization scales are independently
 2213 varied between 0.5-2.0 from their nominal values of 1.0. This results in relative uncertainties on the
 2214 prediction of up to 15%. Finally, varying the parameters in the parton showering results in up to
 2215 8% uncertainty.

2216 5.5.2.1 Uncertainties from EWK-QCD interference

2217 As mentioned in Section 5.0.1, $W^\pm W^\pm jj$ production consists of both EWK processes. The two
 2218 production modes cannot be naively separated due to cross terms in the matrix element calculation.
 2219 These cross terms are referred to as *interference* terms. Since the $W^\pm W^\pm jj$ EWK production is
 2220 the focus of the analysis, and the signal region is designed to preferentially select those events, it is
 2221 important to measure the size of the EWK-QCD interference contributions.

2222 The interference effects are estimated using the `MadGraph` MC generator, as it has a feature that
 2223 allows direct modelling of the interference term. This allows four samples to be generated:

2224 1. Inclusive: All available diagrams are used in the matrix element calculation

2225 2. EWK only: Only EWK diagrams ($\mathcal{O}(\alpha_{\text{EWK}}) = 6$) are used

2226 3. QCD only: Only QCD diagrams ($\mathcal{O}(\alpha_s) = 2 \otimes \mathcal{O}(\alpha_{\text{EWK}}) = 4$) are used

2227 4. Interference: Only the interference terms are used

2228 A minimal set of generator level cuts, listed in Table 5.23, is applied in order to avoid biasing the
 2229 sample towards either production mode. The cross sections for each of the four channels can be
 2230 found in Table 5.24. The size of the interference is found to be approximately 6% of the total cross
 2231 section and is taken as a systematic uncertainty.

Generator level cuts
$\Delta\eta_{jj} < 10$
Jet $p_T > 20$ GeV
$M_{jj} > 10$ GeV

Table 5.23: The set of generator level cuts used for generating the interference samples with `MadGraph`.

Sample	σ (fb)
Inclusive	3.646 ± 0.0012
EWK only	2.132 ± 0.0005
QCD only	1.371 ± 0.0008
Interference	0.227 ± 0.0002

Table 5.24: Cross sections for each different $W^\pm W^\pm jj$ production mode (inclusive, EWK only, QCD only, and interference only) generated using `MadGraph`. The cross sections are calculated using a minimal set of generator level cuts from events where the W decays to a muon.

2232 5.6 Results

2233 After running the full analysis chain, the event yields in the signal region, low- m_{jj} control region,
 2234 and WZ control region as well as associated nuisance parameters representing the uncertainties are
 2235 passed to the maximum likelihood fit. From this fit, the normalization factor for the WZ control
 2236 region μ_{WZ} and the signal strength parameter in the signal region μ_{obs} are determined, and the
 2237 predicted yields in each input bin have been shifted according to the process detailed in Section 5.4.1.

2238 The WZ normalization factor is measured to be:

$$\mu_{WZ} = 0.88^{+0.07}_{-0.07}(\text{stat})^{+0.31}_{-0.21}(\text{model sys})^{+0.22}_{-0.11}(\text{experimental sys}) \quad (5.27)$$

2239 and is constrained primarily by the number of data events in the WZ control region. The observed
 2240 signal strength of $W^\pm W^\pm jj$ EWK production, defined in Equation 5.21, is extracted from the fit
 2241 and measured with respect to the prediction of the `SHERPA v2.2.2` MC generator:

$$\mu_{\text{obs}} = 1.45^{+0.25}_{-0.24}(\text{stat})^{+0.27}_{-0.22}(\text{sys}) \quad (5.28)$$

2242 This corresponds to a rejection of the background-only hypothesis with a significance of 6.9σ .

2243 The observed number of data events are compared to the predicted signal and background yields
 2244 in the signal region after applying the fit in Table 5.25. The pre-fit event yields can be found in
 2245 Appendix A.2. 122 candidate events are observed compared to a prediction of 60 signal and 69
 2246 background events. The m_{jj} distributions for data and prediction are shown in Figure 5.28 after
 2247 the fit, and the fitted event yields in the low- m_{jj} and WZ control regions are shown in Figure 5.29.

2248 The last ingredient necessary to measure the $W^\pm W^\pm jj$ EWK cross section is the theory predicted
 2249 cross section in the fiducial region defined in Table 5.18. `SHERPA v2.2.2` is used for the calculation,
 2250 and the cross section in the total generator phase space is 40.81 ± 0.05 fb, and the fiducial cross section
 2251 is 2.01 ± 0.02 fb. This corresponds to an acceptance factor of $\mathcal{A} = 0.0493 \pm 0.0002$. Uncertainties on
 2252 the simulation are estimated using variations of the scale, parton shower, and PDF set. The final

	e^+e^+	e^-e^-	μ^+e^+	μ^-e^-	$\mu^+\mu^+$	$\mu^-\mu^-$	combined
WZ	1.49 ± 0.30	1.10 ± 0.26	11.7 ± 1.7	8.0 ± 1.3	5.0 ± 0.6	3.5 ± 0.6	31 ± 4
Non-prompt	2.2 ± 1.3	1.2 ± 0.7	5.7 ± 2.8	4.5 ± 1.8	0.57 ± 0.06	0.65 ± 0.14	15 ± 6
e/γ conversions	1.6 ± 0.4	1.6 ± 0.5	6.3 ± 1.6	4.3 ± 1.1	—	—	13.8 ± 2.9
Other prompt	0.16 ± 0.04	0.14 ± 0.04	0.90 ± 0.19	0.63 ± 0.13	0.39 ± 0.09	0.22 ± 0.05	2.4 ± 0.5
$W^\pm W^\pm jj$ QCD	0.35 ± 0.13	0.15 ± 0.05	2.9 ± 1.0	1.2 ± 0.4	1.8 ± 0.6	0.76 ± 0.25	7.2 ± 2.4
Expected background	5.8 ± 1.5	4.1 ± 1.1	27 ± 4	18.7 ± 2.6	7.7 ± 0.8	5.1 ± 0.6	69 ± 7
$W^\pm W^\pm jj$ EWK	5.6 ± 1.0	2.2 ± 0.4	24 ± 5	9.4 ± 1.8	13.5 ± 2.5	5.2 ± 1.0	60 ± 11
Data	10	4	44	28	25	11	122

Table 5.25: Table of the data and prediction event yields in the signal region after the fit. Numbers are shown for the six lepton flavor and charge channels and for all channels combined. The background estimations from the fake-factor are included in the “Non-prompt” category, and backgrounds from $V\gamma$ production and electron charge misidentification are combined in the “ e/γ conversions” category. Finally, ZZ , VVV , and $t\bar{t}V$ backgrounds are combined in the “Other prompt” category.

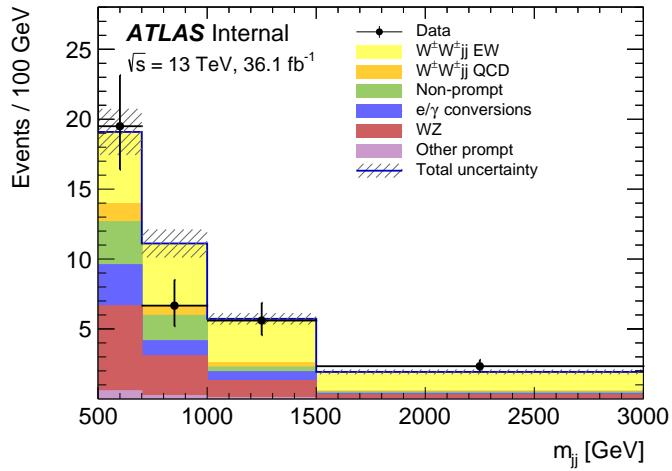


Figure 5.28: The dijet invariant mass m_{jj} distributions for data and predicted signal and background in the signal region after the fit. The shaded band represents the statistical and systematic uncertainties added in quadrature. Note that the bins have been scaled such that they represent the number of events per 100 GeV in m_{jj} . The background estimations from the fake-factor are included in the “Non-prompt” category, and backgrounds from $V\gamma$ production and electron charge misidentification are combined in the “ e/γ conversions” category. Finally, ZZ , VVV , and $t\bar{t} + V$ backgrounds are combined in the “Other prompt” category.

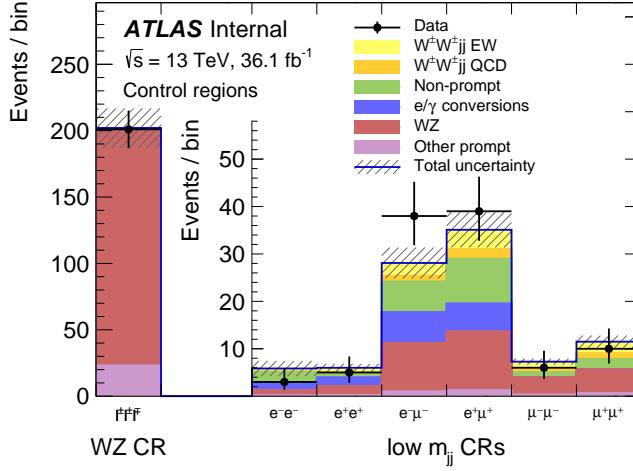


Figure 5.29: The event yields for data and predicted signal and background in the WZ and low- m_{jj} control regions after the fit. The shaded band represents the statistical and systematic uncertainties added in quadrature. The background estimations from the fake-factor are included in the “Non-prompt” category, and backgrounds from $V\gamma$ production and electron charge misidentification are combined in the “ e/γ conversions” category. Finally, ZZ , VVV , and $t\bar{t}V$ backgrounds are combined in the “Other prompt” category.

prediction used in the cross section measurement including uncertainties from Section 5.5.2 is:

$$\sigma_{\text{SHERPA}}^{\text{fid}} = 2.01 \pm 0.02(\text{stat}) \pm 0.29(\text{scale}) \pm 0.16(\text{parton shower}) \pm 0.05(\text{PDF}) \text{ fb} \quad (5.29)$$

Combining this **SHERPA** prediction with the measured signal strength μ_{obs} from Equation 5.28, the measured fiducial cross section $\sigma_{\text{meas}}^{\text{fid}}$ can be calculated using Equation 5.25:

$$\sigma_{\text{meas}}^{\text{fid}} = 2.91^{+0.51}_{-0.47}(\text{stat}) \pm 0.12(\text{model sys}) \pm 0.24(\text{experimental sys}) \pm 0.08(\text{luminosity}) \text{ fb} \quad (5.30)$$

A plot comparing the measured fiducial cross section to two theoretical calculations is shown in Figure 5.30. The measured value is compared to the **SHERPA v2.2.2** prediction used to calculate μ_{obs} as well as to **POWHEG-BOX v2**. As mentioned in Section 5.1.1, this **POWHEG** sample does not include the resonant triboson diagrams and is only used here for a visual comparison.

5.6.1 Beyond the Standard Model extensions of $W^\pm W^\pm jj$

Many so-called *Beyond the Standard Model* (BSM) theories exist that incorporate new physics with what has been experimentally observed. BSM theories often manifest as deviations from the expected SM cross sections, either due to additional decay possibilities affecting branching ratios or modifications of the couplings themselves. One of the most well-known avenues for BSM

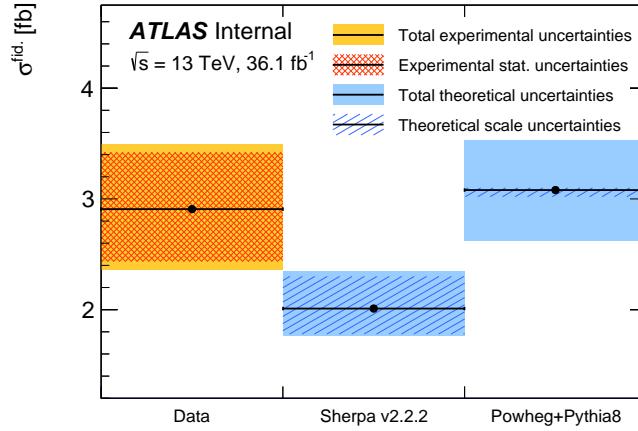


Figure 5.30: Comparison of the measured $W^\pm W^\pm jj$ EWK fiducial cross section with theoretical calculations from **SHERPA** v2.2.2 and **POWHEG-BOX** v2. The light orange band represents the total experimental uncertainty on the measured value, and the dark orange hashed band is the statistical uncertainty. For the simulations, the light blue band represents the total theoretical uncertainty, and the dark blue hashed band are the scale uncertainties. The theory predictions do not include the interference between the EWK and QCD production.

involving new particles is supersymmetry [113]; however, two popular BSM extensions relevant to the $W^\pm W^\pm jj$ process involve a doubly-charged Higgs particle ($H^{\pm\pm}$) and anomalous triple and quartic gauge couplings (aQGC and aTGC, respectively)²⁰. These two BSM theories will be touched on in the context of $W^\pm W^\pm jj$ analyses at the LHC.

5.6.1.1 Doubly charged Higgs bosons

[114, 115]

5.6.1.2 Anomalous quartic couplings

In the event that new physics exists at an energy scale far above what is currently accessible at the LHC, it cannot be directly observed by the experiment; however, its effects can still appear in the interactions between known particles. In this case, the SM is simply the low-energy behavior of a larger *effective field theory* (EFT), which contains additional, higher-dimensional operators that

²⁰The aQGC's are the focus in this section since the $WWWW$ QGC vertex is accessible through $W^\pm W^\pm$ scattering, as well as the fact that aTGC's have been studied in far greater detail due to being accessible through a larger number of processes.

2276 obey the existing SM symmetries:

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_{d>4} \sum_i \frac{\tilde{c}_i}{\Lambda^{d-4}} \mathcal{O}_i \quad (5.31)$$

2277 where \mathcal{O}_i are operators of dimension d with coefficients \tilde{c}_i , and Λ is the energy scale of the new
2278 physics. Here it can be clearly seen that as the energy scale $\Lambda \rightarrow \infty$, the SM behavior dominates. In
2279 the region where $E \ll \Lambda$, operators with high dimensionality contribute less to the total Lagrangian,
2280 and the summation may be truncated above a chosen value of d , at which point \mathcal{L}_{EFT} becomes
2281 predictive and can parametrize any heavy new physics [116].

2282 Only operators with even dimensionality are allowed in order to conserve baryon and lepton
2283 numbers. The largest contributions to \mathcal{L}_{EFT} therefore come from operators with $d = 6$; however,
2284 any of these operators which modify the QGC's also modify the TGC's. As a result, these operators
2285 are better constrained by existing analyses with greater sensitivity to TGC's. Operators with $d = 8$
2286 are the lowest that modify exclusively the QGC's, of which there are 18, and nine of them modify
2287 the $WWWW$ QGC accessible through same-sign $W^\pm W^\pm$ scattering [117, 116]:

$$\begin{aligned} \mathcal{O}_{S,0} &= [(D_\mu \Phi)^\dagger D_\nu \Phi] \times [(D^\mu \Phi)^\dagger D^\nu \Phi] \\ \mathcal{O}_{S,1} &= [(D_\mu \Phi)^\dagger D^\mu \Phi] \times [(D_\nu \Phi)^\dagger D^\nu \Phi] \\ \mathcal{O}_{M,0} &= \text{Tr}[\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times [(D_\beta \Phi)^\dagger D^\beta \Phi] \\ \mathcal{O}_{M,1} &= \text{Tr}[\hat{W}_{\mu\nu} \hat{W}^{\nu\beta}] \times [(D_\beta \Phi)^\dagger D^\mu \Phi] \\ \mathcal{O}_{M,6} &= [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\nu} D^\mu \Phi] \\ \mathcal{O}_{M,7} &= [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\mu} D^\nu \Phi] \\ \mathcal{O}_{T,0} &= \text{Tr}[\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times \text{Tr}[\hat{W}_{\alpha\beta} \hat{W}^{\alpha\beta}] \\ \mathcal{O}_{T,1} &= \text{Tr}[\hat{W}_{\alpha\nu} \hat{W}^{\mu\beta}] \times \text{Tr}[\hat{W}_{\mu\beta} \hat{W}^{\alpha\nu}] \\ \mathcal{O}_{T,2} &= \text{Tr}[\hat{W}_{\alpha\mu} \hat{W}^{\mu\beta}] \times \text{Tr}[\hat{W}_{\beta\nu} \hat{W}^{\nu\alpha}] \end{aligned} \quad (5.32)$$

2288 Each operator is paired with a coupling in the Lagrangian term: $\mathcal{L}_{S,0} = \frac{f_{S,0}}{\Lambda^4} \mathcal{O}_{S,0}$ and so on. The SM
2289 prediction can be compared to simulations generated with chosen values for the anomalous coupling
2290 constants, as shown in Figure 5.31.

2291 Limits on the anomalous couplings generated by the $d = 8$ operators of Equation 5.32 have been
2292 set by CMS in their $W^\pm W^\pm jj$ analyses at $\sqrt{s} = 8$ and 13 TeV [80, 82]. ATLAS has also set limits
2293 at $\sqrt{s} = 8$ TeV [79] using a different parameterization of the anomalous couplings outlined in [118].
2294 The limits set in CMS's 13 TeV analysis are reproduced in Table 5.26. The limits are obtained from

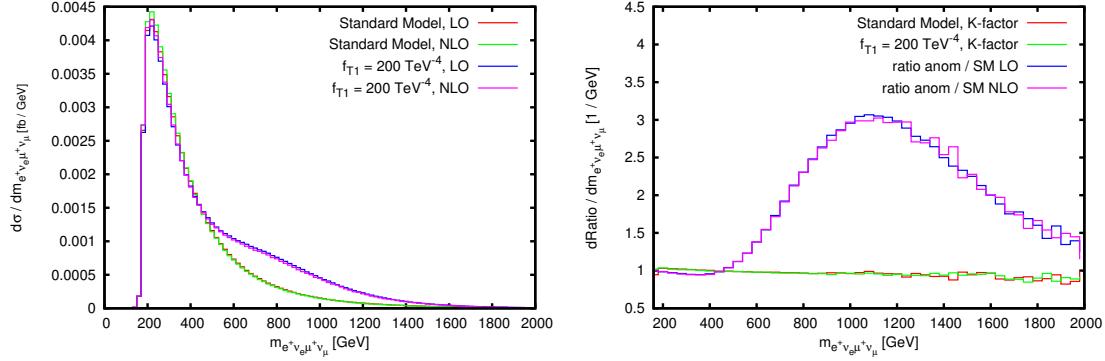


Figure 5.31: Invariant mass distributions of the $2l2\nu$ system in $pp \rightarrow e^+\bar{\nu}_e\mu^+\bar{\nu}_\mu jj$ events at LO and NLO with the VBFNLO MC generator. SM predictions are compared to those with the anomalous coupling $\frac{f_{T,1}}{\Lambda^4} = 200 \text{ TeV}^{-4}$. The left plot shows the differential cross section for each prediction, and the right plot shows the K -factors for the SM and anomalous coupling predictions as well as the cross section ratio between the anomalous coupling and SM predictions at LO and NLO. Plots taken from [116].

2295 fits to the m_{ll} distributions in the signal and WZ control regions, and 95% confidence intervals are
2296 calculated by varying each operator individually.

Coupling	Observed limits [TeV^{-4}]
$f_{S,0}/\Lambda^4$	$[-7.7, 7.7]$
$f_{S,1}/\Lambda^4$	$[-21.6, 21.8]$
$f_{M,0}/\Lambda^4$	$[-6.0, 5.9]$
$f_{M,1}/\Lambda^4$	$[-8.7, 9.1]$
$f_{M,6}/\Lambda^4$	$[-11.9, 11.8]$
$f_{M,7}/\Lambda^4$	$[-13.3, 12.9]$
$f_{T,0}/\Lambda^4$	$[-0.62, 0.65]$
$f_{T,1}/\Lambda^4$	$[-0.28, 0.31]$
$f_{T,2}/\Lambda^4$	$[-0.89, 1.02]$

Table 5.26: Observed 95% confidence limits set by CMS at $\sqrt{s} = 13$ TeV on the nine dimension-eight operators that modify the $WWWW$ QGC listed in Equation 5.32. Table taken from [82].

CHAPTER 6

Prospects for same-sign WW at the High Luminosity LHC

2300 On December 3, 2018, Run 2 of the LHC officially ended, and the collider was shut down to begin
 2301 the first of two scheduled extended maintenance periods [119]. During these two long shutdowns,
 2302 the Phase-I and Phase-II upgrades of the LHC and ATLAS will occur in order to prepare for the
 2303 High-Luminosity LHC (HL-LHC) which is scheduled to begin operation in 2026 [120].

2304 The HL-LHC is planned to run at a center-of-mass energy of $\sqrt{s} = 14$ TeV with an instantaneous
 2305 luminosity of $\mathcal{L} = 5 \times 10^{34}$ cm $^{-2}$ s $^{-1}$ with up to 200 collisions per beam-crossing. Over the course
 2306 of operation, the HL-LHC is expected to collect a total integrated luminosity of $\mathcal{L} = 3000$ fb $^{-1}$ by
 2307 2035 [121].

2308 These run conditions will be much harsher than what ATLAS has experienced so far, and there
 2309 are several upgrades planned to adapt the detector to the high luminosity environment. Most
 2310 notably, the entire ID will be replaced with an all-silicon tracker which will extend the coverage
 2311 from $|\eta| \leq 2.7$ up to $|\eta| \leq 4.0$. This will allow for forward particle tracks to be reconstructed, which
 2312 can in turn be matched to clusters in the calorimeters for use in electron identification or forward
 2313 jet tagging [122].

2314 The upgraded detector, the higher beam energy, and the increased volume of data to be col-
 2315 lected provides the opportunity to measure rare processes to a much higher precision than what
 2316 is possible with the current LCH dataset. Same-sign $W^\pm W^\pm jj$ production is one such process.
 2317 With greater statistics, the accuracy of the cross section measurement can be improved over the
 2318 13 TeV analysis detailed in Chapter 5, and it also will allow for more detailed physics studies, such
 2319 as measuring the polarization state of the scattered W bosons. A measurement of the cross section

of longitudinally polarized $W^\pm W^\pm jj$ scattering is one of the most enticing extensions of the existing $W^\pm W^\pm jj$ measurements due to its sensitivity to the EWSB mechanism [123], and it is expected to be measurable for the first time at the HL-LHC. The analysis detailed in this chapter is based off of the 2018 ATLAS HL-LHC $W^\pm W^\pm jj$ prospects study [124] which extends upon the results of the previous year’s study [125].

6.0.2 Analysis Overview

The experimental signature of interest is identical to the 13 TeV analysis (see Chapter 5): two prompt leptons (either electrons or muons) with the same electric charge, missing transverse energy, and two high energy forward jets. These jets are again required to have a large angular separation and a high combined invariant mass to preferentially select EWK-produced $W^\pm W^\pm jj$ events.

Background processes are not expected to change with respect to the 13 TeV analysis and are summarized here. The dominant source of prompt background comes from $WZ+jets$ events where both bosons decay leptonically. If the lepton from the Z -decay with opposite charge from the W falls outside of the detector acceptance or is not identified, the remaining two leptons will form a same-sign pair, and the event may pass the signal lepton criteria. To a much lesser extent, $ZZ+jets$ events can enter the signal region this way provided two leptons are “lost”. Other prompt sources include $t\bar{t}+V$ and multiple parton interactions, however both contributions are small. Overall, prompt backgrounds are expected to contribute less in HL-LHC analyses than they do currently due to the forward tracking in the upgraded ATLAS detector reducing the probability of leptons falling outside the detector acceptance. Jets mis-reconstructed as leptons or leptons from hadronic decays (such as $t\bar{t}$ and $W+jets$ production) comprise the non-prompt lepton background. Lastly, events with two prompt, opposite-charge electrons can appear as a same-sign event provided one of the electrons is mis-reconstructed and assigned the wrong charge.

In this analysis, the EWK production of $W^\pm W^\pm jj$ is studied in the context of the planned HL-LHC run conditions and upgraded ATLAS detector. An optimized event selection (referred to as the *optimized selection*) is also explored in an effort to gain increased signal significance over the *default selection*. The cross section of the inclusive EWK production is measured for both the default and optimized selections, and the extraction of the expected longitudinal scattering significance is measured with the optimized selection.

2349 **6.1 Theoretical motivation**

2350 The motivation for studying the $W^\pm W^\pm jj$ process as well as VBS in general has been established
 2351 previously in Sections 2.3 and 5.0.1. Since only the longitudinally polarized vector bosons that is
 2352 sensitive to the EWSB mechanism, a direct measurement of this cross section will be very useful for
 2353 understanding how the Higgs unitarizes the scattering amplitude [123].

2354 **6.1.1 Experimental sensitivity to longitudinal polarization**

2355 There are three possible polarization states for a massive vector boson: two transverse (+ or -)
 2356 and one longitudinal (0). Therefore, in a system with two W bosons, the overall polarization can be
 2357 purely longitudinal (00), purely transverse (++, --, and +-), or mixed (+0 and -0). The three
 2358 combinations will be referred to as *LL*, *TT*, and *LT* respectively.

2359 In order extract the longitudinal scattering component, it is necessary to find variables that can
 2360 help distinguish the LL from the TT and LT events. Several were studied, and those with the best
 2361 discriminating power between the different polarization states are the leading and subleading lepton
 2362 transverse momenta as well as the azimuthal separation of the two VBS jets $|\Delta\phi_{jj}|$. Both leptons
 2363 in LL events tend to be softer than the TT and LT events (see Figure 6.1), which motivates keeping
 2364 cuts on the lepton p_T as low as possible in the event selection. In the case of the dijet separation,
 2365 LL events prefer the tag jets to be back-to-back (see Figure 6.2). The $|\Delta\phi_{jj}|$ distribution is chosen
 2366 to be the discriminating variable between the polarizations, and it is ultimately passed to a binned
 2367 likelihood fit to extract the longitudinal scattering significance.

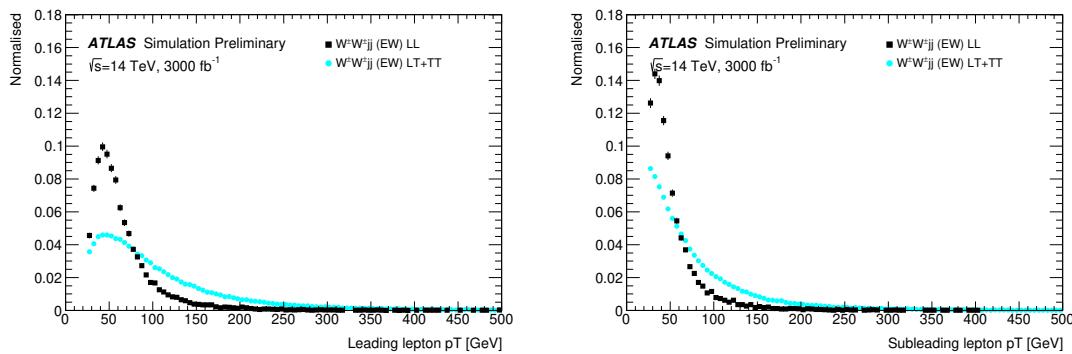


Figure 6.1: Comparison of the leading (left) and subleading (right) lepton p_T distributions for purely longitudinal (LL, black) and mixed polarization (LT+TT, cyan) $W^\pm W^\pm jj$ events.

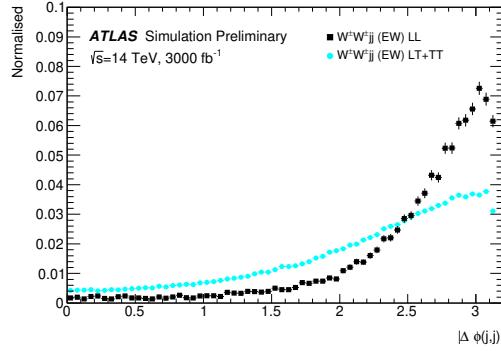


Figure 6.2: Comparison of the azimuthal dijet separation ($|\Delta\phi_{jj}|$) for purely longitudinal (LL, black) and mixed polarization (LT+TT, cyan) $W^\pm W^\pm jj$ events.

2368 6.2 Monte Carlo samples

2369 As this is a prospects study for a future collider, all signal and background processes are modeled
 2370 using MC simulations. The samples are generated at the expected HL-LHC center of mass energy
 2371 $\sqrt{s} = 14$ TeV, and the event yields are scaled to the anticipated integrated luminosity of $\mathcal{L} =$
 2372 3000 fb^{-1} . The MC samples used in the analysis are generated at particle-level and have not been
 2373 run through the full simulation of the ATLAS detector. Instead, smearing functions derived from
 2374 a GEANT4 simulation of the upgraded ATLAS detector are used to estimate detector effects such as
 2375 momentum resolution. In addition, pileup events are fully simulated. The MC samples used in this
 2376 analysis are summarized in Table 6.1.

2377 The signal sample consists of both VBS and non-VBS electroweak (EWK) $W^\pm W^\pm jj$ production,
 2378 and it is simulated with the `Madgraph5_aMC@NLO` generator using the NNPDF3.0 PDF set and in-
 2379 terfaced with PYTHIA v8 [126] for hadronization and parton showering. To study the longitudinal
 2380 polarization more directly, two additional `Madgraph5_aMC@NLO` $W^\pm W^\pm jj$ samples are used: one
 2381 containing only the longitudinal contribution (LL) and a second containing the transverse (TT) and
 2382 mixed (LT) contributions.

2383 There are many other processes that can produce the same final state as the $W^\pm W^\pm jj$ and must
 2384 also be accounted for using MC simulations. WZ events are generated using `SHERPA v2.2.0`, which
 2385 includes up to one parton at NLO in the strong coupling constant and up to three additional partons
 2386 at LO. Both EWK and QCD production are included in these samples. ZZ and triboson VVV
 2387 ($V = W, Z$) events are generated using `SHERPA v2.2.2` with up to two additional partons in the
 2388 final state. For the triboson backgrounds, the bosons can decay leptonically or hadronically. $W+jets$

Process	Generator	Comments
$W^\pm W^\pm jj$ (EWK)	<code>Madgraph5_aMC@NLO</code>	Signal sample
$W^\pm W^\pm jj$ (QCD)	<code>Madgraph5_aMC@NLO</code>	
$W^\pm W^\pm jj$ (LL)	<code>Madgraph5_aMC@NLO</code>	Pure longitudinal polarization sample
$W^\pm W^\pm jj$ (TT+LT)	<code>Madgraph5_aMC@NLO</code>	Mixed and transverse polarization sample
Diboson	<code>SHERPA v2.2.0</code>	WZ events
	<code>SHERPA v2.2.2</code>	ZZ events
Triboson	<code>SHERPA v2.2.2</code>	
$W+jets$	<code>Madgraph5_aMC@NLO</code>	
$Z+jets$	<code>POWHEG-BOX v2</code>	
$t\bar{t}$	<code>POWHEG-BOX</code>	
Single top	<code>POWHEG-BOS</code>	

Table 6.1: Summary of MC samples used in the analysis.

2389 backgrounds are generated for electron, muon, and tau final states at LO with `Madgraph5_aMC@NLO`
 2390 and the `NNPDF3.0` PDF set with showering from `PYTHIA v8`. $Z+jets$ events are produced using
 2391 `POWHEG-BOX v2` and the `CT10` PDF set interfaced with `PYTHIA v8`. Finally, $t\bar{t}$ and single-top events
 2392 are generated using `POWHEG-BOX` with showering from `PYTHIA v6`.

2393 6.3 Background estimations

2394 In this analysis, all background contributions are estimated using MC simulations. Backgrounds
 2395 from electron charge misidentification and fake electrons from jets, which are traditionally estimated
 2396 using data-driven techniques, are instead estimated using a set of parameterization functions applied
 2397 to the MC. These functions calculate the probability that an electron is assigned the wrong charge or
 2398 a jet is mis-reconstructed as an electron parameterized by the p_T and η of the respective electron or
 2399 jet. The probabilities are derived from studies on expected electron performance with the upgraded
 2400 ATLAS detector [127].

2401 Processes involving two W or Z bosons are grouped together as *diboson* backgrounds, with the
 2402 exception of $W^\pm W^\pm jj$ events produced via QCD interactions, which are kept separate. Similarly,
 2403 all backgrounds with three vector bosons are combined and labeled as *triboson*. Any $W+jets$ or top
 2404 events that pass selection and do not contain a fake electron, as well as any $Z+jets$ events without an
 2405 electron identified as having its charge misidentified are combined as *other non-prompt* backgrounds.

2406 **6.3.1 Truth-based isolation**

2407 The canonical isolation variables used in ATLAS analyses require detailed information from several
 2408 detector subsystems including particle tracks and calorimeter responses. Since the MC samples used
 2409 in this analysis have not been run through a full detector simulation, it is not possible to reproduce
 2410 the official isolation variables. For truth-level analysis, this is generally not a serious concern, as
 2411 high- p_T signal leptons tend to be well isolated to begin with. However, isolation is one of the
 2412 most powerful tools for rejecting leptons from non-prompt sources, such as top events, which are
 2413 produced in association with additional nearby particles from b and c hadron decays. It was seen in
 2414 the early stages of this analysis that without any sort of isolation requirement, contributions from
 2415 top backgrounds (including single top, $t\bar{t}$, and $t\bar{t} + V$) were more than an order of magnitude higher
 2416 than expected.

2417 As a result, it was necessary to find one or more quantities that are comparable to the isolation
 2418 information that is available in fully-simulated samples. Analogues to track- and calorimeter-based
 2419 isolation are constructed by summing the momentum and energy, respectively, of stable truth parti-
 2420 cles with $p_T > 1$ GeV within a specified radius of each signal lepton. For the track-based isolation,
 2421 only charged truth particles are used; both charged and neutral particles (excluding neutrinos) are
 2422 included for the calorimeter-based isolation. Ultimately, a set of isolation cuts are chosen that are
 2423 similar to the fixed-cut recommendations for ATLAS Run 2 analyses. The truth-based isolation
 2424 requirements are listed in Table 6.2.

	Electron Isolation	Muon Isolation
Track-based isolation cone size	$\Delta R < 0.2$	$\Delta R < 0.3$
Track-based isolation requirement	$\sum p_T/p_T^e < 0.06$	$\sum p_T/p_T^\mu < 0.04$
Calorimeter-based isolation cone size	$\Delta R < 0.2$	$\Delta R < 0.2$
Calorimeter-based isolation requirement	$\sum E_T/p_T^e < 0.06$	$\sum E_T/p_T^\mu < 0.15$

Table 6.2: Truth-based isolation requirements for electrons and muons.

2425 With no cut on truth-based isolation, 83% of the total background consisted of top events
 2426 (without including top contributions to the fake electron background). The isolation requirement
 2427 reduces the top background by over 99%, and the percentage of the total background from top
 2428 events is reduced to 2%. Additional studies on the truth-based isolation as well as full event yields
 2429 with and without the isolation requirement can be found in Appendix B.1.

2430 **6.4 Object and event selection**

2431 The majority of the default object and event selections were determined in the preceeding $W^\pm W^\pm jj$
 2432 HL-LHC prospects study [125], which focused on the impact of the upgraded detector's forward
 2433 tracking capabilities. Several different combinations of lepton and jet η ranges were tested, and the
 2434 results are used in this study.

2435 **6.4.1 Object selection**

2436 Electrons and muons are preselected to have $p_T > 7$ and 6 GeV, respectively, and lie within $|\eta| \leq 4.0$.
 2437 The likelihood of a given lepton to pass the trigger and identification requirements is estimated by
 2438 calculating an efficiency dependent on the p_T and η of the lepton. The leptons are also required to
 2439 pass the isolation criteria detailed in Table 6.2. Jets that have been tagged as a fake electron by the
 2440 functions described in Section 6.3 are treated as electrons for the purpose of the object selection and
 2441 are subject to the same criteria. In order to be considered a signal lepton, the transverse momentum
 2442 requirement is raised to $p_T > 25$ GeV. The two highest p_T leptons passing this selection are chosen
 2443 to be the leading and subleading signal leptons.

2444 Jets are clustered using the anti- k_t algorithm from final-state particles (excluding muons and
 2445 neutrinos) within a radius of $\Delta R = 0.4$. All jets are required to have $p_T > 30$ GeV and lie within
 2446 $|\eta| < 4.5$; in order to suppress jets from pileup interactions, jets outside of $|\eta| \geq 3.8$ must pass an
 2447 higher momentum cut of $p_T > 70$ GeV. Jets overlapping a preselected electron within $\Delta R(e, j) <$
 2448 0.05 are removed in order to prevent double counting. The two highest p_T jets are defined as the
 2449 leading and subleading tag jets.

2450 **6.4.2 Event selection**

2451 The default event selection is summarized in Table 6.3 and described here. Exactly two signal
 2452 leptons are required with the same electric charge and separated from each other by $\Delta R(ll) > 0.3$.
 2453 In order to suppress contributions from Drell-Yan backgrounds, the two signal leptons must have
 2454 an invariant mass m_{ll} greater than 20 GeV. Additionally, if both signal leptons are electrons, their
 2455 mass must be at least 10 GeV away from the Z -boson mass in order to reduce background from
 2456 Z -boson decays²¹. The event is required to have at least 40 GeV of missing transverse energy (E_T^{miss})

²¹The electron charge misidentification rate in the upgraded ATLAS detector is estimated to be high enough that contributions from $Z \rightarrow ee$ backgrounds are non-negligible.

Selection requirement	Selection value
Lepton kinematics	$p_T > 25 \text{ GeV}$ $ \eta \leq 4.0$
Jet kinematics	$p_T > 30 \text{ GeV}$ for $ \eta \leq 4.5$ $p_T > 70 \text{ GeV}$ for $ \eta > 3.8$
Dilepton charge	Exactly two signal leptons with same charge
Dilepton separation	$\Delta R_{l,l} \geq 0.3$
Dilepton mass	$m_{ll} > 20 \text{ GeV}$
Z boson veto	$ m_{ee} - m_Z > 10 \text{ GeV}$ (ee -channel only)
E_T^{miss}	$E_T^{\text{miss}} > 40 \text{ GeV}$
Jet selection	At least two jets with $\Delta R_{l,j} > 0.3$
b jet veto	$N_{\text{b-jet}} = 0$
Dijet separation	$\Delta \eta_{jj} > 2.5$
Trilepton veto	No additional preselected leptons
Dijet mass	$m_{jj} > 500 \text{ GeV}$
Lepton-jet centrality	$\zeta > 0$

Table 6.3: Summary of the signal event selection.

2457 to account for the two final-state neutrinos. Events with additional preselected leptons are vetoed,
 2458 which greatly reduces WZ and ZZ backgrounds.

2459 Each event must have at least two jets, and both tag jets are required to not overlap with the
 2460 signal leptons. Events with one or more b -jets are vetoed to suppress backgrounds from heavy-flavor
 2461 decays. In order to preferentially select EWK production, the tag jets are required to have a large
 2462 separation between them and a large invariant mass. Finally, a cut on the lepton centrality ζ^{22} ,
 2463 defined in Equation 6.1, further enhances the EWK $W^\pm W^\pm jj$ signal.

$$\zeta = \min[\min(\eta_{\ell 1}, \eta_{\ell 2}) - \min(\eta_{j 1}, \eta_{j 2}), \max(\eta_{j 1}, \eta_{j 2}) - \max(\eta_{\ell 1}, \eta_{\ell 2})] \quad (6.1)$$

2464

2465 6.5 Selection optimization

2466 The default event selection is optimized in order to improve the strength of the $W^\pm W^\pm jj$ EWK
 2467 signal. The expectation is that the increased detector acceptance from the forward tracking combined
 2468 with an increase in center of mass energy and much higher integrated luminosity will allow tighter
 2469 selection cuts without jeopardizing signal statistics.

²² ζ is a measurement of whether the two signal leptons lie between the two tagging jets in η , as is preferred by the VBS topology.

2470 **6.5.1 Random grid search algorithm**

2471 The chosen method for optimizing the event selection is a cut-based algorithm known as the Random
 2472 Grid Search (RGS) [128]. Consider a simple case of two variables x and y chosen to differentiate
 2473 signal from background. In order to be considered a signal event, a given event would be required to
 2474 pass a set of selection criteria, called a *cut point*: $c = \{x > x_c, y > y_c\}$. A simple method to choose
 2475 the optimal cut point (the “best” values of the cuts x_c and y_c) would be to construct an $n \times m$
 2476 rectangular grid in x and y consisting of points $(x_0, y_0), (x_1, y_1), \dots, (x_n, y_m)$, as in Figure 6.3a. One
 2477 can then choose a cut point $c_k = \{x > x_i, y > y_j\}$ that maximizes the signal significance as given
 2478 by a chosen metric. This would be considered a *rectangular grid search*.

2479 While effective in principle, a rectangular grid search comes with two major drawbacks:

- 2480 1. The algorithm scales exponentially as the number of variables to be optimized increases, as
 2481 this is effectively increasing the dimensionality of the grid. In the simple case of a square grid
 2482 with N bins per variable v , the number of cut points to be evaluated grows as N^v .
- 2483 2. Signal and background samples are rarely evenly distributed over the entire grid, resulting
 2484 in many cut points being sub-optimal, and evaluating them would be a waste of computing
 2485 resources.

2486 To combat these limitations, the RGS algorithm constructs a grid of cut points directly from
 2487 the signal sample itself. In the two-dimensional example, this means that the variables x_i and y_j
 2488 making up the cut point $c_k = \{x > x_i, y > y_j\}$ take their values directly from a given signal event.
 2489 This creates a *random grid* of cut points that has the benefit of being biased towards regions of high
 2490 signal concentration by construction. This reduces the need for exponentially increasing numbers of
 2491 cut points while ensuring that computing resources are not wasted in regions with few to no signal
 2492 events. An example of a two-dimensional random grid is shown in Figure 6.3b.

2493 Once the random grid of cut points is constructed, the optimal cut point can be chosen using any
 2494 number of metrics, such as signal to background ratio. For the purpose of the $W^\pm W^\pm jj$ upgrade
 2495 study, the optimal cut point is chosen to be the one that maximizes the signal significance Z , defined
 2496 as [129]:

$$Z = \sqrt{2 \left[(s + b) \ln \left(\frac{s + b}{b_0} \right) + b_0 - s - b \right] + \frac{(b - b_0)^2}{\sigma_b^2}} \quad (6.2)$$

2497 where s and b are the number of signal and background events, respectively, σ_b is the total uncertainty

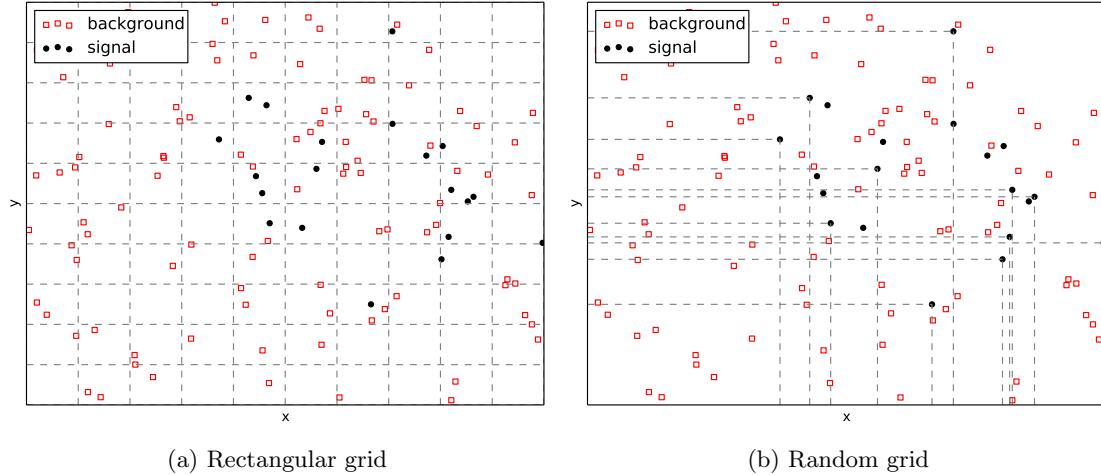


Figure 6.3: A visual representation of a two-dimensional rectangular grid (left) and a random grid (right) in variables x and y . The signal events are the black circles, and the red squares are the background events. Each intersection of gray dashed lines represents a cut point to be evaluated by the optimization.

on the background, and b_0 is defined as:

$$b_0 = \frac{1}{2} \left(b - \sigma_b^2 + \sqrt{(b - \sigma_b^2)^2 + 4(s + b)\sigma_b^2} \right) \quad (6.3)$$

In the case where the background is known precisely (i.e. $\sigma_b = 0$), Equation 6.2 simplifies to

$$Z = \sqrt{2 \left(b \left[(1 + s/b) \ln(1 + s/b) - s/b \right] \right)} \quad (6.4)$$

which further reduces to the familiar $Z = s/\sqrt{b}$ for the case when $s \ll b$.

2501 6.5.2 Inputs to the optimization

In order to train the RGS, signal and background samples are prepared from events passing the event selection outlined in Table 6.3 up through the b -jet veto. The signal sample is chosen to be the longitudinally polarized $W^\pm W^\pm jj$ EWK events, and the transverse and mixed polarizations are treated as background along with $W^\pm W^\pm jj$ events from QCD interactions and the traditional backgrounds listed in Section 6.3. Splitting the inclusive $W^\pm W^\pm jj$ EWK events by polarization allows the optimization to favor the longitudinally polarized events as much as possible, even though they both contribute to the EWK signal.

2509 The following variables are chosen for optimization:

- Leading lepton p_T

2511 • Dilepton invariant mass (m_{ll})

2512 • Leading and subleading jet p_T

2513 • Dijet invariant mass (m_{jj})

2514 • Lepton-jet centrality (ζ)

2515 Subleading lepton p_T is omitted as it is desirable to keep the cut value as low as possible due
 2516 to its sensitivity to the longitudinal polarization (as discussed in Section 6.1.1); despite this, the
 2517 leading lepton p_T is still allowed to be optimized as it can have strong background rejection power.
 2518 Additionally, the dijet separation $\Delta\eta_{jj}$ was included in early studies of the optimization, however it
 2519 was dropped due to the cut value being motivated by well-studied differences between EWK- and
 2520 QCD-produced $W^\pm W^\pm jj$ events (as in Figure 5.6b).

2521 Two additional constraints are imposed on the optimal cut point:

2522 1. At least 1000 signal events must survive in order to prevent the optimization from being too
 2523 aggressive and unnecessarily reducing signal statistics.

2524 2. The dijet invariant mass may only vary within a 50 GeV range of the default value (from
 2525 450–550 GeV) due to the cut being physically motivated by the VBS event topology described
 2526 in Section 5.0.2.

2527 Lastly, the signal significance is calculated without taking into account the uncertainty of the
 2528 background using Equation 6.4. This is due to the fact that the statistical uncertainties of the fake
 2529 electron and charge misidentification backgrounds are quite large, owing to poor MC statisces in a
 2530 few of the samples. If Equation 6.2 were used instead, the optimization would cut unreasonably
 2531 hard against these backgrounds. Since Monte Carlo statistics is not expected to be a limiting factor
 2532 when this analysis is performed at the HL-LHC, it is more realistic to simply ignore these large
 2533 statistical uncertainties for the purpose of the optimization.

2534 6.5.3 Results of the optimization

2535 Ultimately, the random grid is constructed from over 38,000 LL-polarized $W^\pm W^\pm jj$ events in the
 2536 six variables listed above. After applying the constraints, the optimal cut point reduces the total
 2537 background from 9900 to 2310 while reducing the signal from 3489 to 2958. This corresponds to
 2538 an increase in signal significance from $Z = 33.26$ to $Z = 52.63$ as calculated by Equation 6.4. The
 2539 updates to the event selection are listed in Table 6.4.

2540 The large reduction in the background is primarily a result of increasing the leading and sub-
 2541 leading jet p_T from 30 GeV to 90 GeV and 45 GeV, respectively. As can be seen in Figure 6.4,
 2542 this increase removes a significant portion of the backgrounds from jets faking electrons and charge
 2543 mis-ID. Additionally, the loosening of the lepton-jet centrality cut ζ allows more signal events to
 2544 survive the event selection (see Figure 6.5). Other changes to the event selection are minor and do
 2545 not individually have a large impact on the signal or background yields; similar distributions for
 2546 these variables are shown in Appendix B.2.

2547 The full event yields after optimization as well as the cross section measurement are detailed
 2548 alongside those using the default selection in Section 6.6.

Selection requirement	Selection value
Lepton kinematics	$p_T > 28$ GeV (leading lepton only)
Jet kinematics	$p_T > 90$ GeV (leading jet) $p_T > 45$ GeV (subleading jet)
Dilepton mass	$m_{ll} > 28$ GeV
Dijet mass	$m_{jj} > 520$ GeV
Lepton-jet centrality	$\zeta > -0.5$

Table 6.4: Updates to the $W^\pm W^\pm jj$ event selection criteria after optimization. Cuts not listed remain unchanged from the default selection in Table 6.3.

2549 6.6 Results

2550 6.6.1 Event yields

2551 After applying the full event selection, the analysis is broken down into four channels based off of
 2552 the flavor of the signal leptons: $\mu\mu$, ee , μe , and $e\mu$. The full signal and background event yields are
 2553 shown in Table 6.5 for each channel separately and combined using the default event selection. 3489
 2554 EWK $W^\pm W^\pm jj$ events are expected compared to 9900 background events. The dominant sources
 2555 of background are jets faking electrons followed by charge misidentification and diboson processes.
 2556 Triboson events, QCD $W^\pm W^\pm jj$, and other non-prompt sources make up approximately 5% of the
 2557 total background combined.

2558 The event yields for the optimized selection detailed in Section 6.5.3 are listed in Table 6.6. After
 2559 optimization, 2958 signal events and just 2310 background events are expected. Diboson events are
 2560 now the primary source of background, as the optimization greatly reduces the fake and charge
 2561 mis-ID contributions. As discussed earlier, the increase in the leading and subleading jet p_T cuts
 2562 as well as the loosening of the centrality cut are most responsible for the changes in the signal

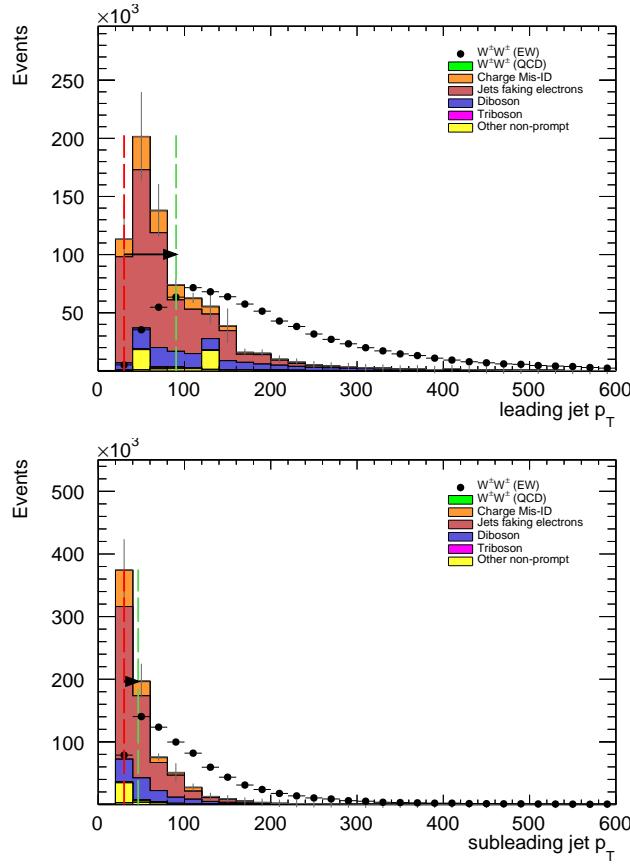


Figure 6.4: Leading (top) and subleading (bottom) jet p_T distributions. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^\pm W^\pm jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

	All channels	$\mu\mu$	ee	μe	$e\mu$
$W^\pm W^\pm jj$ (QCD)	206.4	91.1	22.8	38.4	54.1
Charge Misidentification	2300	0.0	2100	90	160
Jets faking electrons	5000	0.0	3400	1200	340
$WZ + ZZ$	2040	500	438	423	680
Tribosons	115	47	15.4	21.6	31.2
Other non-prompt	210	110	20	60	27
Total Background	9900	750	6000	1900	1290
Signal $W^\pm W^\pm jj$ (EWK)	3489	1435	432	679	944

Table 6.5: Signal and background event yields using the default event selection for an integrated luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$. Events containing a fake or charge-flipped electron are removed from their respective sources and combined into a single entry each.

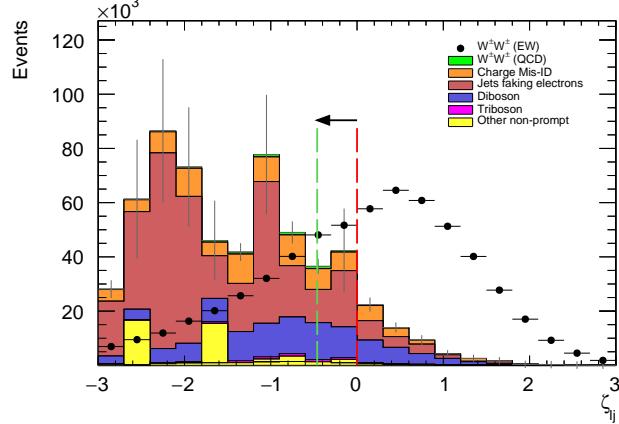


Figure 6.5: Lepton-jet centrality distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^\pm W^\pm jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

and background yields; distributions of these quantities using the default and the optimized event selections can be found in Figures 6.6, 6.7, and 6.8, respectively.

	All channels	$\mu\mu$	ee	μe	$e\mu$
$W^\pm W^\pm jj$ (QCD)	168.7	74.6	19.7	32.2	42.2
Charge Misidentification	200	0.0	11	30	160
Jets faking electrons	460	0.0	130	260	70
$WZ + ZZ$	1286	322	289	271	404
Tribosons	76	30.1	9.6	15.1	21.6
Other non-prompt	120	29	16.6	50	19
Total Background	2310	455	480	660	710
Signal $W^\pm W^\pm jj$ (EWK)	2958	1228	380	589	761

Table 6.6: Signal and background event yields using the optimized event selection for an integrated luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$. Events containing a fake or charge-flipped electron are removed from their respective sources and combined into a single entry each.

It is important to note, however, that the MC sample used to estimate $Z+\text{jets}$ events suffers from poor statistics which results in large per-event weights once scaled to $\mathcal{L} = 3000 \text{ fb}^{-1}$. This sample contributes heavily to the fake and charge misidentification backgrounds, and a handful of these events being cut out by the optimization is largely responsible for the dramatic reduction of the corresponding backgrounds. As a result, the optimized results presented here are likely overly optimistic. However, given proper MC statistics, it is still expected that this optimization will outperform the default selection.

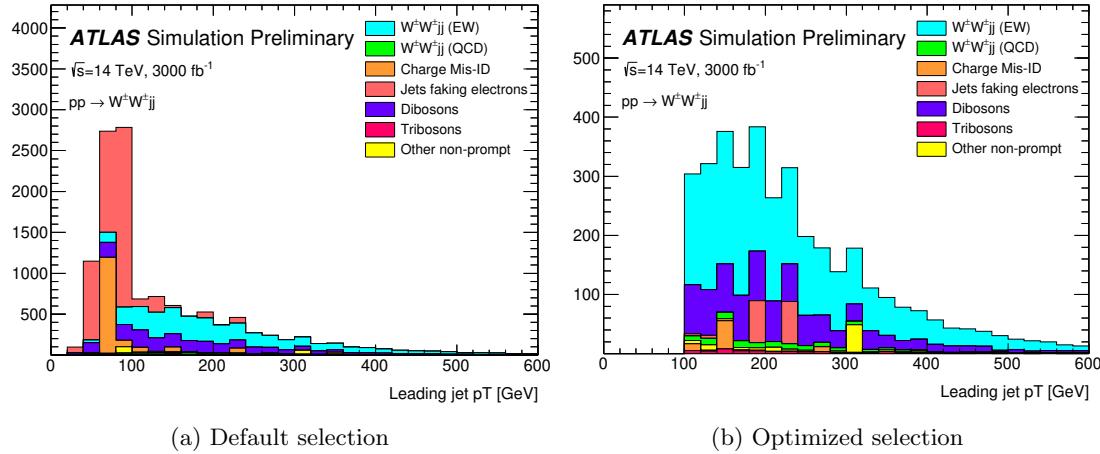


Figure 6.6: p_T distributions for the leading jet using the default (left) and optimized (right) event selections for all channels combined.

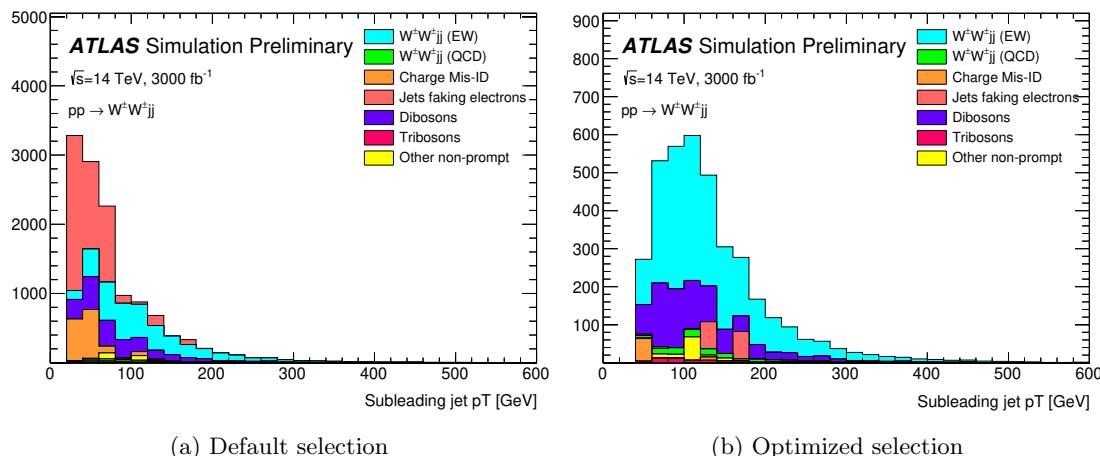


Figure 6.7: p_T distributions for the subleading jet using the default (left) and optimized (right) event selections for all channels combined.

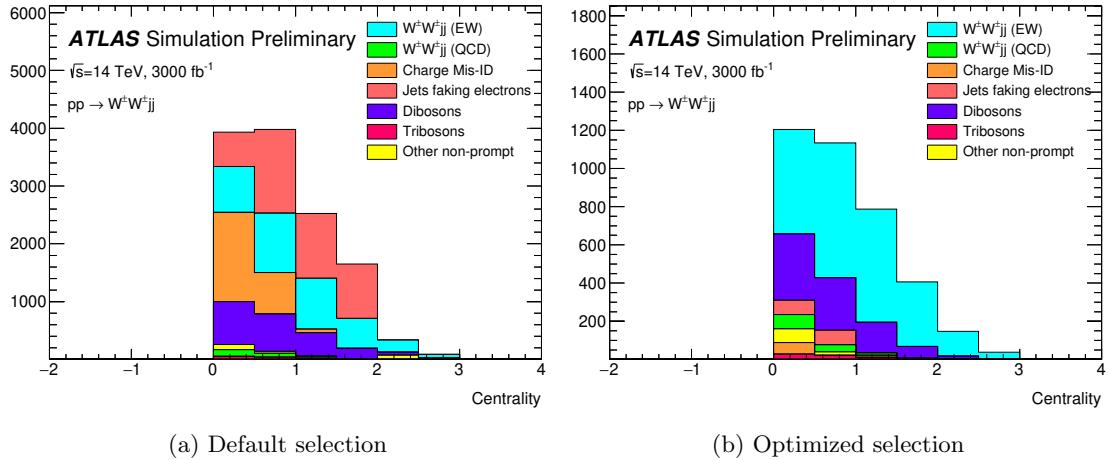


Figure 6.8: p_T distributions for lepton-jet centrality ζ using the default (left) and optimized (right) event selections for all channels combined.

2572 6.6.2 Uncertainties

The uncertainties considered for the analysis are summarized in Table 6.7. Values for experimental systematics on the trigger efficiency, lepton and jet reconstruction, and flavor tagging are taken directly from the 13 TeV analysis [81]. The rate uncertainties for the background processes are halved from the 13 TeV values according to ATLAS recommendations. The uncertainty on the fake electron estimation is also halved from the 13 TeV analysis. Finally, a conservative estimate of the uncertainty on the charge flip background is used as the electron charge mis-ID rate due to material interactions is difficult to predict at this stage.

Source	Uncertainty (%)
$W^\pm W^\pm jj$ (EWK)	3
Luminosity	1
Trigger efficiency	0.5
Lepton reconstruction and identification	1.8
Jets	2.3
Flavor tagging	1.8
Jets faking electrons	20
Charge misidentification	25
$W^\pm W^\pm jj$ (QCD)	20
Top	15
Diboson	10
Triboson	15

Table 6.7: Summary of estimated experimental and rate uncertainties.

2580 **6.6.3 Cross section measurement**

2581 The cross section is calculated using the same method as in the 13 TeV analysis, detailed in Sec-
 2582 tion 5.4. Unlike the previous analysis, however, eight lepton channels are used here instead of six.
 2583 The μe and $e\mu$ channels remain separated in addition to the $\mu\mu$ and ee channels, and each lepton fla-
 2584 vor channel is further split by charge, as this increases the sensitivity of the analysis. Each channel's
 2585 m_{jj} distribution is combined in a profile likelihood fit to extract the EWK $W^\pm W^\pm jj$ production
 2586 cross section. Using the default event selection, the expected cross section calculated to be:

$$\sigma_{W^\pm W^\pm jj}^{\text{expected}} = 16.89 \pm 0.36 \text{ (stat)} \pm 0.53 \text{ (theory)} \pm 0.84 \text{ (syst)} \text{ fb} \quad (6.5)$$

2587 and with the optimized event selection:

$$\sigma_{W^\pm W^\pm jj}^{\text{expected}} = 16.94 \pm 0.36 \text{ (stat)} \pm 0.53 \text{ (theory)} \pm 0.78 \text{ (syst)} \text{ fb} \quad (6.6)$$

2588 The optimized selection should not change the measured value of the cross section, and indeed both
 2589 are consistent with within uncertainties. The systematic uncertainty is reduced by about 7% with
 2590 the optimized selection. The total uncertainty on the cross section measurement is approximately
 2591 6%, compared to the 20% uncertainty on the measured fiducial cross section of the 13 TeV analysis
 2592 reported in Equation 5.30.

2593 Projections of each uncertainty type and the total uncertainty on the cross section as a function of
 2594 integrated luminosity are shown in Figure 6.9. The predictions are made by scaling the event yields
 2595 by different luminosity values and re-running the fitting procedure. As the integrated luminosity
 2596 increases past $\mathcal{L} > 3000 \text{ fb}^{-1}$, the statistical uncertainty continues to reduce; however, the total
 2597 uncertainty becomes limited by the systematics. Additionally, the total uncertainty is expected to
 2598 reduce by less than a percent as the integrated luminosity increases past the planned 3000 fb^{-1} .
 2599 This implies that the precision on the cross section measurement will not improve by much even
 2600 with additional data.

2601 **6.6.4 Longitudinal scattering significance**

2602 The longitudinal scattering significance is extracted in much the same way as the cross section, this
 2603 time using a binned likelihood fit on the $|\Delta\phi_{jj}|$ distribution. In order to increase sensitivity, the
 2604 $|\Delta\phi_{jj}|$ distribution is split into two bins in m_{jj} , and an additional cut on the pseudorapidity of the
 2605 subleading lepton is applied ($|\eta| < 2.5$) to reduce background contributions from fake electrons and
 2606 charge flip. The $|\Delta\phi_{jj}|$ distributions used in the fit are shown in Figure 6.10. Due to limited statistics

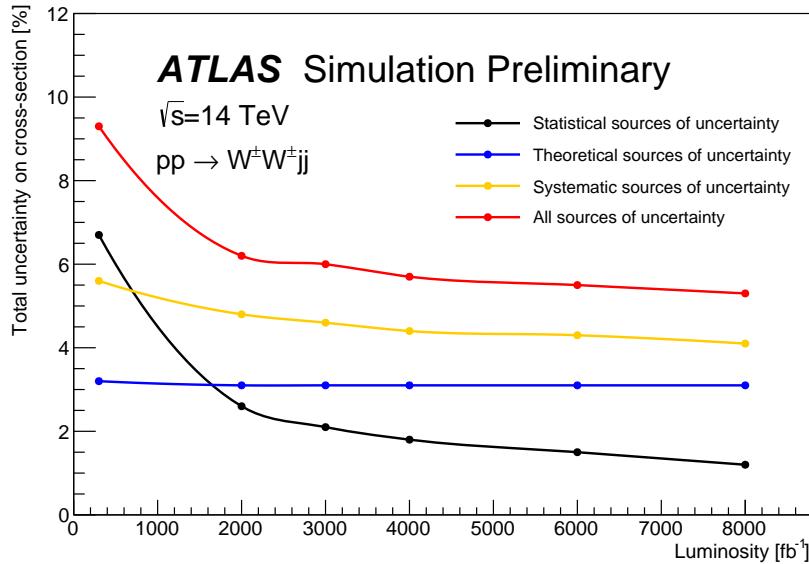


Figure 6.9: Projections of the statistical (black), theoretical (blue), systematic (yellow), and total (red) uncertainties on the measured cross section as a function of integrated luminosity using the optimized event selection.

2607 in the LL events, the four lepton flavor channels are not split by charge. The expected significance
 2608 of the $W_L^\pm W_L^\pm jj$ process is 1.8σ with a precision of 47% on the measurement. Projections of the
 2609 expected significance as a function of integrated luminosity is shown in Figure 6.11.

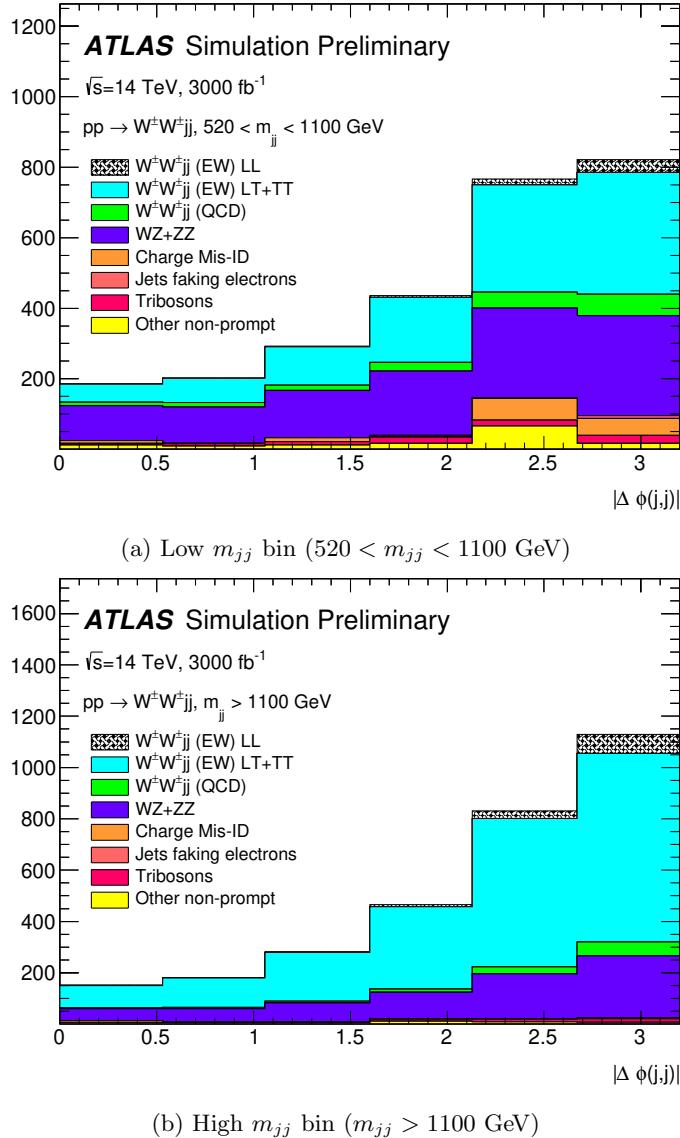


Figure 6.10: Dijet azimuthal separation ($|\Delta\phi_{jj}|$) for the low m_{jj} region ($520 < m_{jj} < 1100$ GeV, top) and the high m_{jj} region ($m_{jj} > 1100$ GeV, bottom). The purely longitudinal (LL, gray) is plotted separately from the mixed and transverse (LT+TT, cyan) polarizations.

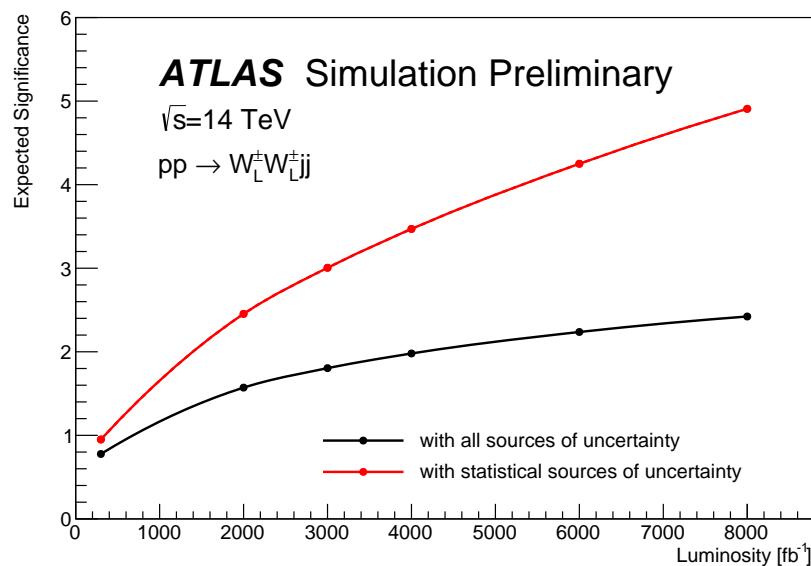


Figure 6.11: Projections of the expected longitudinal scattering significance as a function of integrated luminosity when considering all sources of uncertainties (black) or only statistical uncertainties (red).

2610

CHAPTER 7

2611

Conclusion

2612 Here's where you wrap it up.

2613

APPENDIX A

2614

2615

Additional material on $W^\pm W^\pm jj$ measurement at $\sqrt{s} = 13$ TeV

2616

A.1 Impact of experimental uncertainty on MC background estimations

2617 Tables A.1-A.6 contain the impact of experimental systematic uncertainties for the remaining back-
 2618 grounds estimated from MC simulation. The $W^\pm W^\pm jj$ EWK signal and WZ background sys-
 2619 tematics are listed in the main body of the document, in Tables 5.21 and 5.22, respectively. While the
 2620 percentage of the contributions for some systematics appear large, the size of these backgrounds are
 2621 quite small compared to the total background.

$W^\pm W^\pm jj$ QCD	ee % Yield	$e\mu$ % Yield	$\mu\mu$ % Yield
Jet-related Uncertainties	3.41	3.04	2.85
b-tagging efficiency	2.56	2.48	2.48
Pile-up	4.99	0.45	0.33
Trigger efficiency	0.02	0.08	0.41
Lepton reconstruction/ID	1.62	1.19	1.89
MET reconstruction	0.41	0.22	0.34

Table A.1: Impact of experimental uncertainties for the $W^\pm W^\pm jj$ QCD processes in all channels.

Triboson	ee % Yield	$e\mu$ % Yield	$\mu\mu$ % Yield
Jet-related Uncertainties	13.09	13.39	16.85
b-tagging efficiency	2.96	3.77	4.95
Pile-up	19.37	24.66	6.87
Trigger efficiency	0.02	0.07	0.47
Lepton reconstruction/ID	1.66	1.27	2.48
MET reconstruction	0.00	0.46	0.00

Table A.2: Impact of experimental uncertainties for triboson process in all channels.

$t\bar{t}V$	ee % Yield	$e\mu$ % Yield	$\mu\mu$ % Yield
Jet-related Uncertainties	17.65	11.97	14.27
b-tagging efficiency	15.02	9.04	13.83
Pile-up	8.73	10.69	4.18
Trigger efficiency	0.03	0.08	0.39
Lepton reconstruction/ID	2.57	3.27	2.66
MET reconstruction	1.75	4.16	1.62

Table A.3: Impact of experimental uncertainties for $t\bar{t}V$ processes in all channels.

$W\gamma$	ee % Yield	$e\mu$ % Yield	$\mu\mu$ % Yield
Jet-related Uncertainties	7.05	33.36	—
b-tagging efficiency	1.97	2.94	—
Pile-up	4.11	14.17	—
Trigger efficiency	0.01	0.14	—
Lepton reconstruction/ID	1.40	1.13	—
MET reconstruction	0.00	0.00	—

Table A.4: Impact of experimental uncertainties for the $W\gamma$ process in all channels.

$Z\gamma$	ee % Yield	$e\mu$ % Yield	$\mu\mu$ % Yield
Jet-related Uncertainties	16.22	370.44	—
b-tagging efficiency	1.08	3.10	—
Pile-up	12.57	11.51	—
Trigger efficiency	0.02	0.07	—
Lepton reconstruction/ID	1.26	22.01	—
MET reconstruction	0.00	0.00	—

Table A.5: Impact of experimental uncertainties for the $Z\gamma$ process in all channels.

ZZ	ee % Yield	$e\mu$ % Yield	$\mu\mu$ % Yield
Jet-related Uncertainties	15.71	15.76	35.18
b-tagging efficiency	2.23	2.35	2.89
Pile-up	1.22	3.20	4.58
Trigger efficiency	0.03	0.10	0.36
Lepton reconstruction/ID	3.59	3.10	5.70
MET reconstruction	4.84	3.26	3.24

Table A.6: Impact of experimental uncertainties for the ZZ process in all channels.

2622 A.2 Pre-fit event yields

2623 Table A.7 contains the event yields for each source and channel before the fit.

	e^+e^+	e^-e^-	μ^+e^+	μ^-e^-	$\mu^+\mu^+$	$\mu^-\mu^-$	combined
WZ	1.9 ± 0.6	1.3 ± 0.4	14 ± 4	8.9 ± 2.6	5.5 ± 1.6	3.6 ± 1.1	35 ± 10
Non-prompt	4.1 ± 2.3	2.3 ± 1.7	9 ± 5	6 ± 4	0.57 ± 0.15	0.67 ± 0.25	23 ± 10
e/γ conversions	1.74 ± 0.29	1.8 ± 0.4	6.1 ± 1.6	3.7 ± 0.8	—	—	13.4 ± 2.5
Other prompt	0.17 ± 0.05	0.14 ± 0.04	0.90 ± 0.19	0.60 ± 0.14	0.36 ± 0.10	0.19 ± 0.05	2.4 ± 0.5
$W^\pm W^\pm jj$ QCD	0.38 ± 0.13	0.16 ± 0.05	3.0 ± 1.0	1.2 ± 0.4	1.8 ± 0.6	0.76 ± 0.25	7.3 ± 2.5
Expected background	8.2 ± 2.4	5.7 ± 1.8	33 ± 7	21 ± 5	8.2 ± 1.8	5.3 ± 1.2	81 ± 14
$W^\pm W^\pm jj$ EWK	3.8 ± 0.6	1.49 ± 0.22	16.5 ± 2.5	6.5 ± 1.0	9.1 ± 1.4	3.5 ± 0.5	41 ± 6
Data	10	4	44	28	25	11	122

Table A.7: Table of the data and prediction event yields in the signal region before the fit. Numbers are shown for the six lepton flavor and charge channels and for all channels combined. Here the WZ background yields are normalized to the data in the WZ control region. The background estimations from the fake-factor are included in the “Non-prompt” category, and backgrounds from $V\gamma$ production and electron charge misidentification are combined in the “ e/γ conversions” category. Finally, ZZ , VVV , and $t\bar{t} + V$ backgrounds are combined in the “Other prompt” category.

APPENDIX B

Additional material on $W^\pm W^\pm jj$ prospects at the HL-LHC

B.1 Truth isolation

As mentioned in Section 6.3.1, the size of the background contribution from top processes are much larger than expected when no isolation is applied. The event yields using an earlier version of the event selection with no truth-based isolation requirement are listed in Table B.1. Here, top events make up nearly 90% of the total background, and the contributions from fake and charge-flipped electrons are also large. The event yields using the same event selection with the truth-based isolation included are shown in Figure B.2. When comparing the two tables, the considerable reduction in the top background can be clearly seen.

yields by type	all channels	$\mu\mu$	ee	μe	$e\mu$
signal	4011	1583.2	531.7	793.1	1103.1
ww qcd	252.6	105.8	30.4	48	68.4
charge flip	2528.4	0.0	2075.4	255.1	197.8
fakes	7135.4	0.0	4675.1	1904.3	555.9
diboson	2370.4	581.2	491.8	517.9	779.6
triboson	125.5	49.1	17.8	24.6	34.1
top	90150.5	26618	15301.6	25277.9	22953.1
z+jets	241.2	0.0	0.0	0.0	241.2
w+jets	31.4	3.9	7.6	13.2	6.7
total bkg	102803.9	27354	22592	28027.8	24830.1
signal	4011	1583.2	531.7	793.1	1103.1

Table B.1: Event yields prior to applying any form of truth-based isolation criteria.

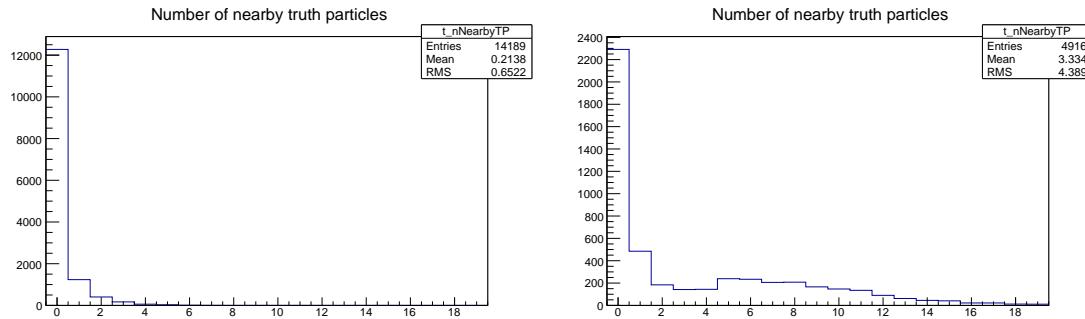
TODO: Add tables for tight vs loose working point, information on the necessity of TRUTH1++

yields by type	all channels	$\mu\mu$	ee	μe	$e\mu$
signal	3470.5	1427.3	428.8	675.8	938.7
ww qcd	205.8	90.8	22.7	38.3	54
charge flip	2398.3	0.0	2104.6	95.8	197.9
fakes	4309.7	0.0	3390.6	750.8	168.3
diboson	1552.4	311.3	355.6	346.8	538.7
triboson	115	46.8	15.4	21.6	31.2
top	156.9	42.3	14.8	76.6	23.3
$z+jets$	0.0	0.0	0.0	0.0	0.0
$w+jets$	0.3	0.0	0.0	0.3	0.0
total bkg	8738.1	491.3	5903.7	1329.8	1013.4
signal	3470.5	1427.3	428.8	675.8	938.7

Table B.2: Event yields after applying a test version of the truth-based isolation.

2636 **B.1.1 TRUTH1++ derivations**

2637 The ATLAS standard TRUTH1 derivations used for this analysis contain a slimmed truth particle
 2638 container in order to reduce the file size. As a result, many of the truth particles that would be
 2639 included in the isolation variables are missing, and the truth-based isolation will not accurately
 2640 model the reconstruction-level isolation variables. In order to recover the performance of the truth-
 2641 based isolation in the top MC samples (where it is most needed), a custom derivation was produced
 2642 privately that duplicated the default TRUTH1 data structure but includes the full truth particle
 2643 record. The reduced size of the truth particle information in the TRUTH1 derivation compared to the
 2644 TRUTH1++ derivation is shown in Figure B.1.

Figure B.1: Number of truth particles within $\Delta R < 0.4$ of a selected muon or electron using the ATLAS standard TRUTH1 (left) and the custom TRUTH1++ (right) derivations in $t\bar{t}$ simulation. The complete truth record is stored in the TRUTH1++ derivation, and this is best seen in the first bin, where the lepton has no nearby truth particles.

2645 **B.1.2 Check of truth-based isolation**

2646 Since the isolation variables are constructed from truth particles, there is an expectation that the
 2647 efficiency of the isolation selection will be higher than what would be seen in the full simulation.
 2648 In order to test this, a truth-level 13 TeV $t\bar{t}$ MC sample was run through a selection altered to
 2649 mimic the 13 TeV $W^\pm W^\pm jj$ measurement as closely as possible. The results were compared to the
 2650 $t\bar{t}$ background in the 13 TeV analysis extrapolated to 14 TeV and 3000 fb^{-1} , and the truth-based
 2651 isolation reduces the expected events by a factor of approximately 4. However, the statistics in the
 2652 13 TeV truth-level sample are low, and it is therefore difficult to measure precisely how much the
 2653 truth-based isolation overperforms.

2654 **B.1.3 Loose isolation working point**

2655 As another check on the truth-based isolation, a second isolation working point was constructed
 2656 to match the official ATLAS Fixed Cut Loose isolation working point. The definition of this loose
 2657 isolation are found in Table B.3.

2658 The primary impact of loosening the isolation is a substantial increase in the non-prompt back-
 2659 round from top processes, and a moderate increase in the charge mis-ID and fake backgrounds.
 2660 Backgrounds from prompt leptons only did not see major changes. As a result, the tight working
 2661 point is chosen for the analysis. The event yields by sample and by background type using the
 2662 loose working point are in Table B.4, and Table B.5 has the numbers using the tight working point
 2663 (defined in Table 6.2) for comparison.

	Electron Isolation	Muon Isolation
Track-based isolation cone size	$\Delta R < 0.2$	$\Delta R < 0.3$
Track-based isolation requirement	$\sum p_T/p_T^e < 0.15$	$\sum p_T/p_T^\mu < 0.15$
Calorimeter-based isolation cone size	$\Delta R < 0.2$	$\Delta R < 0.2$
Calorimeter-based isolation requirement	$\sum E_T/p_T^e < 0.2$	$\sum E_T/p_T^\mu < 0.3$

Table B.3: Electron and muon isolation requirements for the loose working point.

run number	all channels			mm			ee			me			em		
	events	stat	sys	events	stat	sys	events	stat	sys	events	stat	sys	events	stat	sys
signal	3783.21	22.08	0.00	1524.99	15.00	0.00	484.74	7.01	0.00	740.76	9.01	0.00	1032.72	11.50	0.00
ww qcd	223.95	3.54	44.79	97.17	2.51	19.43	25.51	1.03	5.10	42.23	1.40	8.45	59.04	1.80	11.81
charge flip	3025.40	1276.74	0.00	0.00	0.00	0.00	2615.30	1267.89	0.00	197.20	87.94	0.00	212.90	121.63	0.00
fakes	5315.55	1775.87	0.00	0.00	0.00	0.00	3524.24	1694.39	0.00	1356.74	450.60	0.00	434.57	282.33	0.00
diboson	2195.61	38.10	219.58	548.72	18.54	54.87	451.27	18.29	45.14	470.61	15.71	47.07	725.01	22.95	72.50
triboson	117.43	5.90	17.62	47.55	4.32	7.13	15.83	1.94	2.37	22.11	2.18	3.32	31.94	2.76	4.80
top	554.63	218.75	83.21	229.26	135.53	34.40	61.15	38.23	9.18	232.30	167.28	34.85	31.92	6.43	4.78
z+jets	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
w+jets	1.21	0.87	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.02	0.00	0.00
total bkg	11433.78	2198.44	239.70	922.70	136.88	67.99	6693.30	2116.67	46.41	2322.38	488.89	59.27	1495.40	308.36	73.77
signal	3783.21	22.08	0.00	1524.99	15.00	0.00	484.74	7.01	0.00	740.76	9.01	0.00	1032.72	11.50	0.00

Table B.4: Event yields broken down by sample and by background type using the loose isolation workingpoint. Events containing a fake or charge-flipped electron are removed from their respective sample and added to the ‘fakes’ and ‘charge flip’ rows, respectively. Errors include statistical uncertainty and estimated systematic rate uncertainty based on the background process.

run number	all channels			mm			ee			me			em		
	events	stat	sys	events	stat	sys	events	stat	sys	events	stat	sys	events	stat	sys
signal	3489.49	21.23	0.00	1434.85	14.55	0.00	431.75	6.61	0.00	679.09	8.63	0.00	943.8	11.00	0.00
ww qcd	206.42	3.41	41.28	91.12	2.43	18.22	22.84	0.98	4.57	38.37	1.34	7.67	54.09	1.72	10.82
charge flip	2335.73	1163.47	0.00	0.00	0.00	0.00	2087.78	1159.5	0.00	90.37	33.32	0.00	157.58	90.02	0.00
fakes	4979.27	1756.47	0.00	0.00	0.00	0.00	3406.20	1705.03	0.00	1230.80	362.15	0.00	342.27	216.54	0.00
diboson	2039.94	36.93	204.00	499.69	18.04	49.97	437.60	14.12	43.76	422.90	14.18	42.29	679.75	25.25	67.98
triboson	115.03	5.87	17.29	46.84	4.31	7.03	15.40	1.94	2.32	21.55	2.17	3.24	31.24	2.74	4.70
top	211.74	84.14	31.76	107.96	71.12	16.20	19.58	3.76	2.93	57.21	44.47	8.58	26.99	5.40	4.05
z+jets	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.00
w+jets	0.30	0.28	0.00	0.00	0.00	0.00	0.00	0.00	0.00	0.28	0.28	0.00	0.02	0.02	0.00
total bkg	9888.43	2108.87	211.25	745.61	73.54	56.04	5898.40	2061.99	44.16	1861.48	366.67	43.95	1291.94	235.95	69.11
signal	3489.49	21.23	0.00	1434.85	14.55	0.00	431.75	6.61	0.00	679.09	8.63	0.00	943.80	11.00	0.00

Table B.5: Event yields broken down by background type using the tight isolation workingpoint. Events containing a fake or charge-flipped electron are removed from their respective sample and added to the “fakes” and “charge flip” rows, respectively. Errors include statistical uncertainty and estimated systematic rate uncertainty based on the background process.

2664 **B.2 Plots of other optimization variables**

2665 Plots of the remaining optimization variables not shown in Section 6.5.3 are presented here for
 2666 reference. Figures B.2, B.3, and B.4 compare signal and background distributions for the default
 2667 and optimized cuts. None of these cuts change by much in the optimized selection and their impacts
 2668 on the overall event selection is minimal.

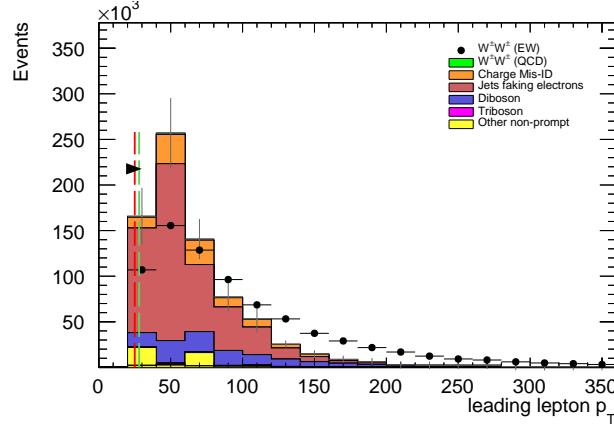


Figure B.2: Leading lepton p_T distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^\pm W^\pm jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

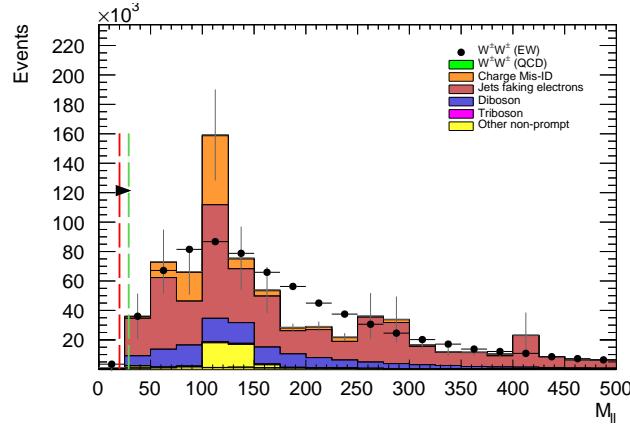


Figure B.3: Dilepton invariant mass distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^\pm W^\pm jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

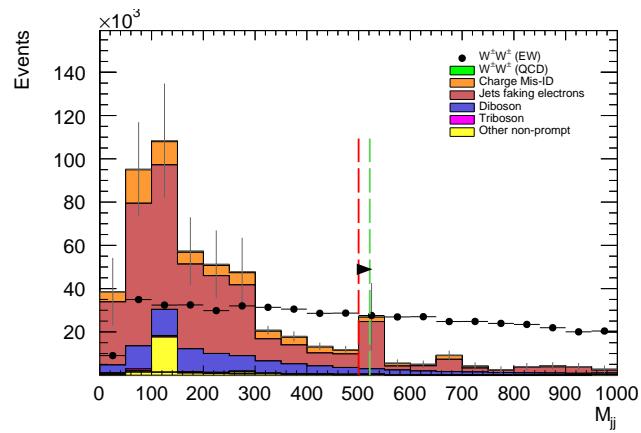


Figure B.4: Dijet invariant mass distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^\pm W^\pm jj$ signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

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