Standard Model is best model (working title)

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STANDARD MODEL IS BEST MODEL (WORKING TITLE)

Acknowledgements

I'd like to thanks the Ghosts of Penn Students Past for providing me with such an amazing thesis template.

ABSTRACT

STANDARD MODEL IS BEST MODEL (WORKING TITLE)

William Kennedy DiClemente

J. Kroll

This is the abstract text.

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Preface

This is the preface. It’s optional, but it’s nice to give some context for the reader and stuff.

Will K. DiClemente
Philadelphia, February 2019

Chapter 1

Introduction

The Standard Model (SM)[1] has been remarkably successful...

[1] Here's a footnote.

Chapter 2

Theoretical Framework

(Some example introductory text for this chapter)...

2.1 Introduction to the Standard Model

Modern particle physics is generally interpreted in terms of the Standard Model (SM). This is a quantum field theory which encapsulates our understanding of the electromagnetic, weak, and strong interactions...

2.2 Electroweak Mixing and the Higgs Field

When the theory of the electroweak interaction was first developed [2, 3], the W and Z bosons were predicted to be massless (a typical mass term in the Lagrangian would violate the SU(2) symmetry). However, these were experimentally observed to have masses...

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3.2.1.3 Transition Radiation Tracker

The Transition Radiation Tracker is the outermost component of the ID...

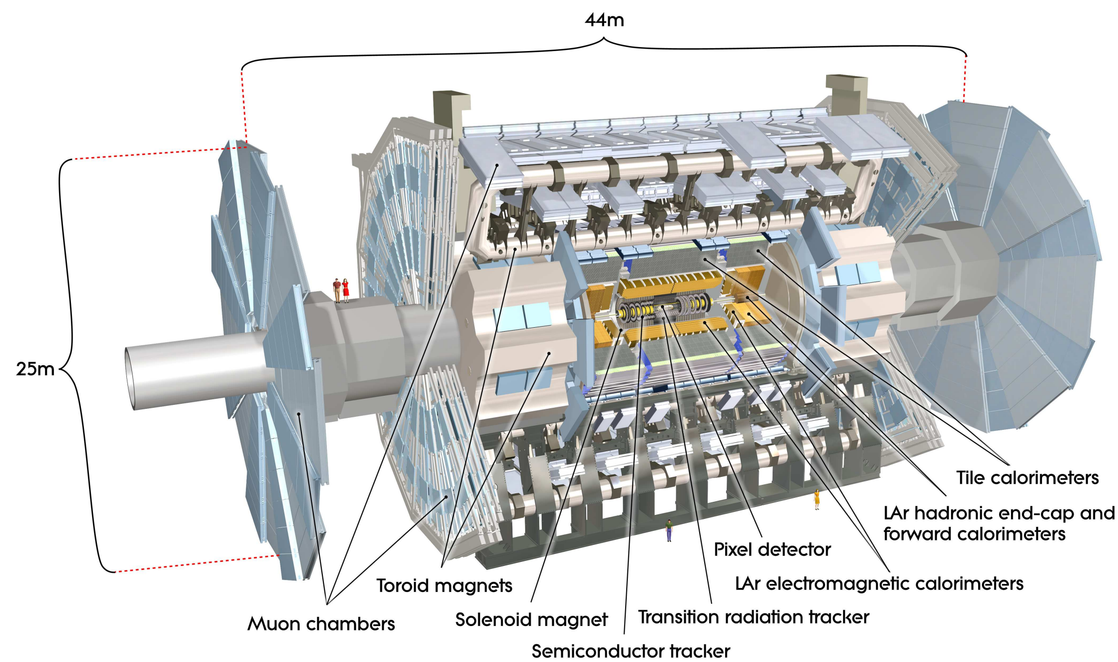

Figure 3.1: General cut-away view of the ATLAS detector [5].

3.2.2 The Calorimeters

ATLAS includes two types of calorimeter system for measuring electromagnetic and hadronic showers. These are the Liquid Argon (LAr) calorimeters and the Tile calorimeters. Together, these cover the region with $|\eta| < 4.9$...

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The Liquid Argon system consists of...

3.2.2.2 Tile Calorimeters

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In order for the subdetectors of the ID to operate at their designed precisions, it is essential that the locations of the sensors be known as precisely as possible. Differences between the expected and actual positions of a sensor can result in displaced particle hits and degrade track reconstruction quality. These misalignments can occur for any number of reasons, including but not limited to elemnts shifting during maintenance periods or cycles in ATLAS's magnetic field, or simply small movements during normal detector operations. Since it is not practical to physically realign hundreds of thousands of detector elements to μm precision by hand, an iterative track-based alignment algorithm is used to determine the physical positions and orientations of these elements [6]. The effects of misalignments and the steps taken to correct and monitor them are detailed in this chapter.

4.1 Effects of Misalignment

Hello world!

4.2 The Alignment Method

Hello world!

4.3 Momentum Bias Corrections

Hello world!

4.4 Alignment of the IBL

Hello world!

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Hello world!

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5.1 Theoretical motivation

Hello world!

5.2 Signal definition

Hello world!

5.3 Background estimations

Hello world!

5.4 Cross section measurement

Hello world!

Chapter 6

Prospects for same-sign WW at the High Luminosity LHC

On December 3, 2018, Run 2 of the LHC officialy ended, and the collider was shut down to begin the first of two scheduled extended maintenance periods [7]. During these two long shutdowns, the Phase-I and Phase-II upgrades of the LHC and ATLAS will occur in order to prepare for the High-Luminosity LHC (HL-LHC) which is scheduled to begin operation in 2026 [8].

The HL-LHC is planned to run at a center-of-mass energy of $\sqrt{s} = 14$ TeV with an instantaneous luminosity of $\mathcal{L} = 5 \times 10^{34}$ cm^{-2}s^{-1} with up to 200 collisions per beam-crossing. Over the course of operation, the HL-LHC is expected to collect a total integrated luminosity of $\mathcal{L} = 3000$ fb^{-1} by 2035 [9].

These run conditions are much harsher than what ATLAS has experienced so far, and as a result there are several planned upgrades to the detector. Most notably, the entire ID will be replaced with an all-silicon tracker which will extend the coverage from $|\eta| \leq 2.7$ up to $|\eta| \leq 4.0$. This will allow for reconstruction of charged particle tracks which can in turn be matched to clusters in the calorimeters for electron identification or forward jet tagging [10].

The upgraded detector combined with the higher beam energy and the considerable increase in integrated luminosity means that many analyses with low signal statistics in Run 2 have the potential to be greatly improved with the HL-LHC. While the ATLAS 13 TeV $W^{\pm}W^{\pm}jj$ cross section measurement certainly did not suffer greatly from low statistics TODO: –reword–, the accuracy of the measurement can still be improved at the HL-LHC. Of particular interest is the longitudinal polarization of the W bosons due to its sensitivity to electroweak symmetry breaking [11].

The analysis detailed in this chapter is based off of the 2018 public ATLAS $W^{\pm}W^{\pm}jj$ prospects study [12] which is itself an extension of the 2017 ATLAS study [13]. TODO: mention CMS's study + yellow report?

6.0.1 Analysis Overview

The experimental signature of interest here is identical to the 13 TeV analysis detailed in Chapter 5: two prompt leptons (electrons or muons) with the same charge, missing transverse energy, and two jets. Once again the two leading jets are required to have a large angular separation and a high combined invariant mass to preferentially select EWK VBS production over QCD $W^{\pm}W^{\pm}jj$ events.

Background processes that can mimic the signal are again similar to the 13 TeV analysis. The dominant source of prompt background from WZ+jets events where both bosons decay leptonically. If the lepton from the Z-decay with opposite charge from the W falls outside of the detector acceptance or is not identified, the remainder could appear to be a $W^{\pm}W^{\pm}jj$ signal event. To a lesser extent, ZZ+jets events can enter the signal region in much the same way provided two leptons are "lost". Other prompt sources include $t\bar{t}+V$ and and multiple parton interactions, however these processes do not contribute much. The upgrades to the ATLAS detector are expected to reduce the size of these prompt contributions due in large part to the increased detector acceptance from the forward tracking. Jets mis-reconstructed as leptons or leptons from hacronic decays (such as $t\bar{t}$ and W+jets production) comprise the non-prompt lepton background. Lastly, events with two prompt, opposite-charge electrons can contribute provided one of the electrons is mis-reconstructed as the wrong charge.

In this analysis, the EWK production of $W^{\pm}W^{\pm}jj$ is studied in the context of the planned HL-LHC run conditions and upgraded ATLAS detector. An optimized event selection (referred to as the *optimized selection*) is also explored in an effort to gain increased signal significance over the *default selection*. The cross section of the inclusive EWK production is measured for both the default and optimized selections, and the extraction of the longitudinal scattering significance is measured with the optimized selection.

6.1 Theoretical motivation

The theoretical motivation for studying the ssWW process is detailed in Section 5.1. The particular interest in polarization is the potential for the scattering amplitude of longitudinally polarized weak bosons to diverge linearly as the center of mass energy increases, ultimately violating unitarity around 1 TeV [14]. In the Standard Model, the Higgs boson cancels these divergences. However, as the Higgs is recently discovered it is still extremely to study the mechanism of electroweak symmetry breaking (EWSB), and the longitudinal scattering of W bosons is expected to be one of the most sensitive tests of EWSB [11].

6.1.1 Experimental sensitivity to longitudinal polarization

There are three possible polarization states for a massive vector boson: two transverse ($+$ or $-$) and one longitudinal (0). Therefore, in a system with two W bosons, the overall polarization can be purely longitudinal (00), purely transverse ($++$, $--$, and $+-$), or mixed ($+0$ and -0). The three combinations will be referred to as LL, TT, and LT respectively.

In order extract the longitudinal scattering component, it is necessary to find variables that distinguish the LL from the TT and LT. Several variables were studied, and those with the best discriminating power between the polarizations were the leading and subleading lepton p_{T} as well as the azimuthal separation ($|\Delta\phi_{jj}|$) of the two VBS jets. The LL events preferred lower p_{T} for both signal leptons (see Figure 6.1), which motivates keeping these two cuts as low as possible in the event selection in order to preserve as much longitudinal polarization as possible. In the case of $|\Delta\phi_{jj}|$, the LL events generally had a larger dijet separation (see Figure 6.2), and this variable is used in a binned likelihood fit to extract the longitudinal scattering significance.

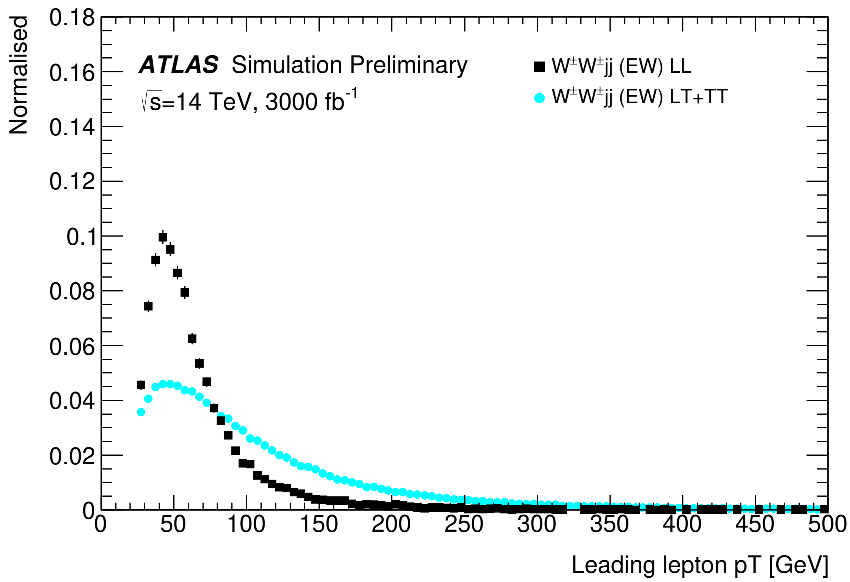
6.2 Monte Carlo samples

As no real HL-LHC data will be available for many years, all processes in this prospects study must be simulated using Monte Carlo (MC) generators. Signal and background processes were generated at $\sqrt{s} = 14$ TeV, and the event yields scaled to the anticipated HL-LHC integrated luminosity of $\mathcal{L} = 3000\ \mathrm{fb}^{-1}$.

TODO: Consider putting all this in a table

The signal sample consists of both VBS and non-VBS electroweak (EWK) $W^{\pm}W^{\pm}jj$ production, and it is sumulated with the `Madgraph5_aMC@NLO` generator [15] using the `NNPDF3.0` PDF set [16] and interfaced with `PYTHIA v8` [17] for hadronization and parton showering. To study the longitudinal polarization more directly, two additional `Madgraph5_aMC@NLO` $W^{\pm}W^{\pm}jj$ samples are used: one containing only the longitudinal contribution (LL) and a second containing the transverse (TT) and mixed (LT) contributions.

There are many other processes that can produce the same final state as the $W^{\pm}W^{\pm}jj$ and must also be accounted for using MC simulations. WZ events are generated using `SHERPA v2.2.0` [18, 19,



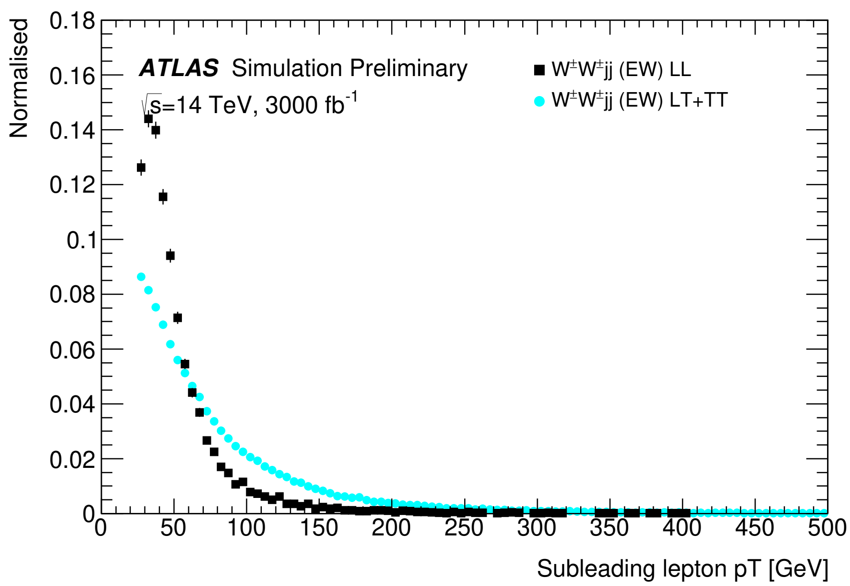


Figure 6.1: Comparison of the leading (top) and subleading (bottom) lepton p_{T} distributions for purely longitudinal (LL, black) and mixed polarization (LT+TT, cyan) $W^{\pm}W^{\pm}jj$ events. Plots from [1].

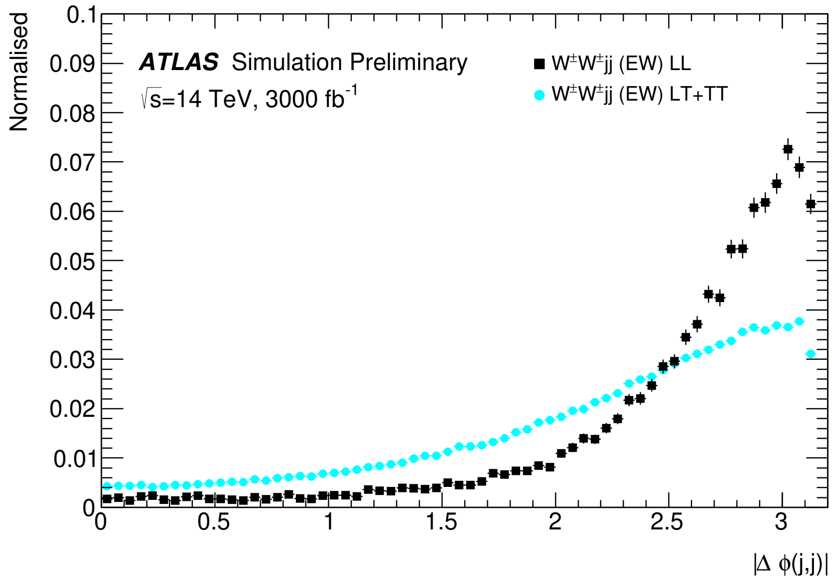


Figure 6.2: Comparison of the azimuthal dijet separation ($|\Delta\phi_{jj}|$) for purely longitudinal (LL, black) and mixed polarization (LT+TT, cyan) $W^{\pm}W^{\pm}jj$ events. Plot from [1].

20], which includes up to one parton at next-to-leading order (NLO) in the strong coupling constant α_s and up to three additional partons at leading order (LO). Both EWK and QCD production are included in these samples. ZZ events are generated using `SHERPA v2.2.2` with up to two additional partons in the final state. Triboson backgrounds $VVV, V = W, Z$ where the bosons can decay leptonically or hadronically are simulated with `SHERPA v2.2.2` with up to two additional partons in the final state. W+jets backgrounds are generated for electron, muon, and tau final states are generated at LO with `Madgraph5_aMC@NLO` and the `NNPDF3.0` set with showering from `PYTHIA v8` . Z+jets events are generated using `POWHEG-BOX` [21] and the `CT10` PDF set [22] interfaced with `PYTHIA v8` . Finally, $t\bar{t}$ and single-top events are generated using `POWHEG-BOX` with showering from `PYTHIA v6` .

Since the MC samples used in the analysis are generated at particle-level and have not been run through the typical full simulation of the ATLAS detector, smearing functions are instead used to estimate detector effects. These are derived from a `GEANT4` simulation of the upgraded ATLAS detector [23]. In addition, pileup events are fully simulated.

6.3 Background estimations

In this analysis, all background contributinos are estimated using MC simulations. Backgrounds such as electron charge misidentification and fake electrons from jets–which are traditionally estimated using data-driven techniques–are estimated using a set of parameterization functions applied to the MC. These functions calculate the probability that an electron is assigned the wrong charge or a jet is mis-reconstructed as an electron parameterized by the p_{T} and η of the electron or jet. The probabilites are derived from studies on expected electron performance with the upgraded ATLAS detector [24].

Processes involving two W and Z bosons are grouped together as *diboson* backgrounds, with the exception of $W^{\pm}W^{\pm}jj$ events produced via QCD interactions, which are kept separate. Similarly, all backgrounds with three vector bosons are combined and labeled as *triboson*. Any W+jets or top events that pass selection and do not contain a fake electron, as well as any Z+jets events without an electron identifed as having its charge misidentified are combined as *other non-prompt* backgrounds.

6.3.1 Truth-based isolation

Since the MC samples used in this analysis have not been run through a full detector simulation, they lack any kind of particle isolation variables (since they require, for example, information on the calorimeter response). Generally, this is not a large concern, as at truth-level, high p_{T} signal leptons tend to be well isolated to begin with. However, isolation is one of the most powerful tools for rejecting leptons from non-prompt sources such as top events, which are produced in association with additional nearby particles from b and c quark decays. In the absence of any sort of isolation requirement, contributions from top backgrounds (including single top, $t\bar{t}$ and $t\bar{t}+V$) were more than an order of magnitude higher than expected.

As a result, it was necessary to create an analogue to the isolation information that is available in fully-simulated samples. Track- and calorimeter-based isolation variables were constructed by summing the momentum and energy, respectively, of stable truth particles with $p_{\mathrm{T}} > 1$ GeV within a specified radius of each signal lepton. For the track-based isolation, only charged truth particles were used; both charged and neutral particles (excluding neutrinos) were included for the calorimeter-based isolation. Ultimately, a set of isolation cuts were chosen that are similar to those recommended by ATLAS for Run 2 analyses. The truth-based isolation requirements are listed in Table 6.1.

The truth-based isolation requirement reduced the top background by over 99% and reduced the contribution of top events to the total background from 83% to 2%. Additional details on the

	Electron Isolation	Muon Isolation
Track-based isolation cone size	$\Delta R < 0.2$	$\Delta R < 0.3$
Track-based isolation requirement	$\sum p_{\mathrm{T}}/p_{\mathrm{T}}^{e} < 0.06$	$\sum p_{\mathrm{T}}/p_{\mathrm{T}}^{\mu} < 0.04$
Calorimeter-based isolation cone size	$\Delta R < 0.2$	$\Delta R < 0.2$
Calorimeter-based isolation requirement	$\sum E_{\mathrm{T}}/p_{\mathrm{T}}^{e} < 0.06$	$\sum E_{\mathrm{T}}/p_{\mathrm{T}}^{\mu} < 0.15$

Table 6.1: Truth-based isolation requirements for electrons and muons.

truth-based isolation studies are presented in Appendix A.

6.4 Object and Event selection

6.4.1 Object selection

Electrons and muons are preselected to have $p_{\mathrm{T}} > 7$ and 6 GeV, respectively, and $|\eta| \leq 4.0$. The likelihood of a given lepton to pass the trigger or identification requirements is estimated by estimating an efficiency dependent on the p_{T} and η of the lepton. The leptons are also required to pass the isolation criteria detailed in Table 6.1. Jets that have been tagged as a fake electron by the functions described earlier in Section 6.3 are treated as electrons for the purpose of the object selection and are subject to the same criteria. In order to be considered a signal lepton, an additional requirement of $p_{\mathrm{T}} > 25$ GeV is applied on top of the preselection. The two highest p_{T} leptons passing this selection are chosen to be the leading and subleading signal leptons.

Jets are clustered using the anti-k_t algorithm [25] from final-state particles within a radius of $\Delta R = 0.4$ (excluding muons and neutrinos). Jets are required to have $p_{\mathrm{T}} > 30$ GeV and lie within $|\eta| < 4.5$, with an additional cut of $p_{\mathrm{T}} > 70$ GeV for jets above $|\eta| \geq 3.8$ in order to suppress jets from pileup interactions. Jets overlapping with a preselected electron within $\Delta R_{e,j} < 0.05$ are removed in order to prevent double counting. The two highest p_{T} jets are defined as the leading and subleading *tag jets*.

6.4.2 Event selection

The default event selection is summarized in Table 6.2 and described here. Exactly two signal leptons are required with the same electric charge and separated from each other by 0.3 in ΔR. In order to suppress contributions from Drell-Yan backgrounds, the two signal leptons must have an invariant mass m_{ll} greater than 20 GeV. Additionally, if both signal leptons are electrons, their mass must be at least 10 GeV from the Z-boson mass in order to reduce background from Z-boson decays[2]. The event is required to have at least 40 GeV of missing transverse energy ($E_{\mathrm{T}}^{\mathrm{miss}}$) to account for the two neutrinos from the W decays. Events with additional preselected leptons are vetoed, which greatly reduces WZ and ZZ backgrounds. Both tag jets are required to not overlap with the signal leptons, and there is a veto on events with one or more b-jets. In order to preferentially select VBS production, the tag jets are also required to have a large separation between them and a large invariant mass. Finally, a cut on the lepton centrality, ζ, defined in Equation 6.1 enhances the EWK $W^{\pm}W^{\pm}jj$ signal.

$$\zeta = \min[\min(\eta_{\ell 1}, \eta_{\ell 2}) - \min(\eta_{j1}, \eta_{j2}), \max(\eta_{j1}, \eta_{j2}) - \max(\eta_{\ell 1}, \eta_{\ell 2})] \tag{6.1}$$

Selection requirement	Selection value
Lepton kinematics	$p_{\mathrm{T}} > 25$ GeV $\lvert\eta\rvert \leq 4.0$
Jet kinematics	$p_{\mathrm{T}} > 30$ GeV for $\lvert\eta\rvert \leq 4.5$ $p_{\mathrm{T}} > 70$ GeV for $\lvert\eta\rvert > 3.8$
Dilepton charge	Exactly two signal leptons with same charge
Dilepton separation	$\Delta R_{l,l} \geq 0.3$
Dilepton mass	$m_{ll} > 20$ GeV
Z boson veto	$\lvert m_{ee} - m_Z\rvert > 10$ GeV (ee-channel only)
$E_{\mathrm{T}}^{\mathrm{miss}}$	$E_{\mathrm{T}}^{\mathrm{miss}} > 40$ GeV
Jet selection	At least two jets with $\Delta R_{l,j} > 0.3$
b jet veto	$N_{\text{b-jet}} = 0$
Dijet separation	$\Delta\eta_{jj} > 2.5$
Trilepton veto	No additional preselected leptons
Dijet mass	$m_{jj} > 500$ GeV
Lepton-jet centrality	$\zeta > 0$

Table 6.2: Summary of the signal event selection.

6.5 Selection optimization

As mentioned earlier, the HL-LHC will feature forward tracking, an increase in center of mass energy, and a higher integrated luminosity. Therefore, this study is an excellent time to see if there are new optimizations to the signal event selection that can improve the signal to background ratio.

[2]The electron charge mis-ID rate is high enough that contributions from $Z \to ee$ backgrounds are non-negligible.

6.5.1 Random grid search algorithm

The chosen method for optimizing the event selection is a cut-based algorithm known as the Random Grid Search (RGS) [26]. Consider a simple case of two variables x and y chosen to differentiate the signal from the background. In order to be considered a signal event, a given event would be required to pass a *cut point* $c = \{x > x_c, y > y_c\}$. A simple method to choose the optimal cut point (i.e. the "best" values of the cuts x_c and y_c) would be to construct an $n \times m$ rectangular grid in x and y consisting of points $(x_0, y_0), (x_1, y_1), ..., (x_n, y_m)$, as in Figure 6.3. One can then choose a cut point $c_k = \{x > x_i, y > y_j\}$ that maximizes the signal significance as measured by a chosen metric. This would be considered a *regular* or *rectangular* grid search.

While effective in principle, this rectangular grid search comes with two major drawbacks:

1. The algorithm does not scale well as the number of variables to be optimized–the dimensionality of the grid–increases. In the case of a square grid with N bins per variable v, the number of cut points to be evaluated grows as N^v.

2. Signal and background samples are rarely evenly distributed over the entire grid, resulting in many cut points being sub-optimal and evaluating them would be a waste of computing resources.

To combat these limitations, the RGS algorithm constructs a grid of cut points directly from the signal sample itself. In the two-dimensional example, this means that the variables x_i and y_j making up the cut point $c_k = \{x > x_i, y > y_j\}$ take their values directly from a given signal event. This has the benefit of creating a *random grid* of cut points that is by construction biased towards regions of high signal concentration. This reduces the need for exponentially increasing numbers of cut points while ensuring that computing resources are not wasted in regions with few to no signal events. An example of the the two-dimensional random grid is shown in Figure 6.4.

Once the random grid of cut points is constructed, the optimal cut point can be chosen using whatever metric the analyzer chooses, such as signal to background ratio. For the purpose of the $W^{\pm}W^{\pm}jj$ upgrade study, the optimal cut point is the one that mazimizes the signal significance Z defined as in Equation 6.2 [27].

$$Z = \sqrt{2\left[(s+b)\ln\left(\frac{s+b}{b_0}\right) + b_0 - s - b\right] + \frac{(b-b_0)^2}{\sigma_b^2}} \tag{6.2}$$

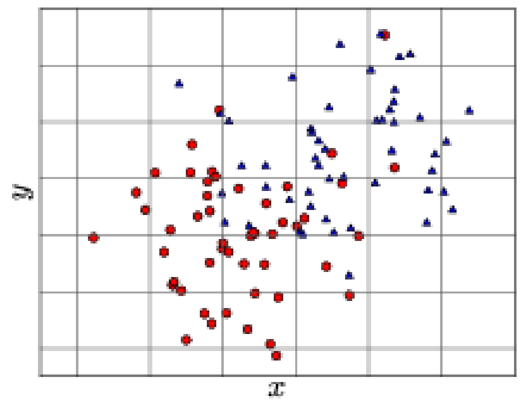


Figure 6.3: A visual representation of a rectangular grid search algorithm. The signal events are the blue triangles, and the red circles are the background events. TODO: replace with own figure

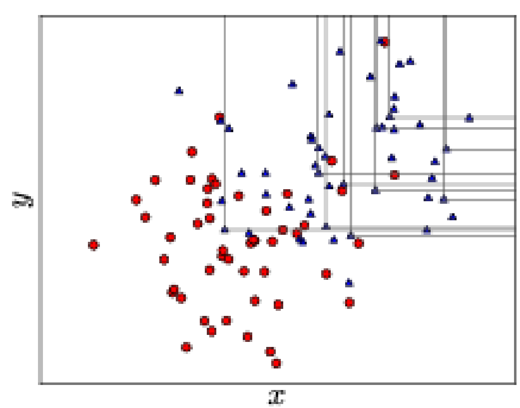


Figure 6.4: A visual representation of a random grid search algorithm. The signal events are the blue triangles, and the red circles are the background events. TODO: replace with own figure

where s and b are the number of signal and background events, respectively, σ_b is the total uncertainty on the background, and b_0 is defined as:

$$b_0 = \frac{1}{2}\left(b - \sigma_b^2 + \sqrt{(b - \sigma_b^2)^2 + 4(s+b)\sigma_b^2}\right) \tag{6.3}$$

In the case where the backround is known precisely (i.e. $\sigma_b = 0$), Equation 6.2 simplifies to

$$Z = \sqrt{2\left(b\left[(1 + s/b)\ln(1 + s/b) - s/b\right]\right)} \tag{6.4}$$

which further reduces to the familiar $Z = s/\sqrt{b}$ for the case when $s << b$.

6.5.2 Inputs to the optimization

In order to train the RGS, signal and background samples were prepared from events passing the event selection outlined in Table 6.2 up through the b-jet veto. The signal sample was chosen to be the longitudinally polarized $W^\pm W^\pm jj$ EWK events, and the transverse and mixed polarizations were treated as background along with $W^\pm W^\pm jj$ events from QCD interactions and the traditional backgrounds listed in Section 6.3. Splitting the inclusive $W^\pm W^\pm jj$ EWK events by polarization allows the optimization to favor the longitunally polarized events as much as possible, even though they both contribute to the EWK signal.

The following variables were chosen for optimization:

- Leading lepton p_{T}
- Dilepton invariant mass (m_{ll})
- Leading and subleading jet p_{T}
- Dijet invariant mass (m_{jj})
- Lepton-jet centrality (ζ)

Subleading lepton p_{T} was omitted as it is desirable to keep the cut value as low as possible due to its sensitivity to the longitudinal polarization (as discussed in Section 6.1.1). Additionally, the dijet separation $\Delta\eta_{jj}$ was included in the optimization originally, however it was dropped from the list due to the cut value being motivated by differences between EWK and QCD produced $W^\pm W^\pm jj$ events.

Two additional constraints were imposed when selecting the optimal cut point:

1. At least 1000 signal events must survive in order to prevent the optimization from being too aggressive and unnecssarily reducing signal statistics.

2. The dijet invariant mass may only vary within a 50 GeV range of the default value (from 450-550 GeV) due to the cut being physically motivated by the VBS event topology (TODO: reference where this is discussed in the 13TeV section).

Lastly, the decision was made to use calculate the signal significance without taking into account the uncertainty of the background using Equation 6.4. This was due to the fact that the statistical uncertainties of the fake electron and charge-misID backgrounds were quite large, and if Equation 6.2 were used instead, the optimization would cut unreasonably hard against these backgrounds. Since Monte Carlo statistics is not expected to be a limiting factor when this analysis is performed at the HL-LHC, it is more realistic to simply ignore these large statistical uncertainties for the purpose of the selection optimization.

6.5.3 Results of the optimization

Ultimately, the random grid was constructed from over 38,000 LL-polarized $W^{\pm}W^{\pm}jj$ events in the variables listed above. After applying the constraints, an optimal cut point was chosen which reduced the total background from 9900 to 2310 while reducing the signal from 3489 to 2958. This corresponds to an increase in signal significance from $Z = 33.26$ to $Z = 52.63$ as calculated by Equation 6.4. The updates to the event selection are listed in Table 6.3.

The large reduction in the background is primarily a result of the increase in the leading and subleading jet p_{T} from 30 GeV to 90 GeV and 45 GeV, respectively. As can be seen in Figure 6.7, this increase removes a significant portion of the backgrounds from jets faking electrons and charge mis-ID. Additionally, the loosening of the lepton-jet centrality cut ζ allows more signal events to survive the event selection (see Figure 6.9). Other changes to the event selection are minor and do not individually have a large impact on the signal or background yields.

The full event yields after optimization as well as the cross section measurement are detailed alongside those using the default selection in Section 6.6.

TODO: It's a bit awkward to reference the results of the default/optimized before they're properly presented. Maybe move the sections around? not sure...

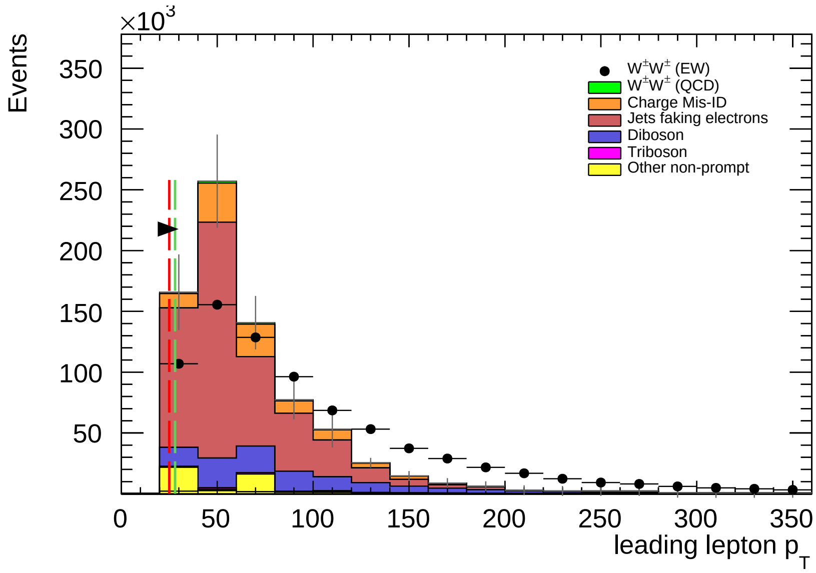


Figure 6.5: Leading lepton p_{T} distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^{\pm}W^{\pm}jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram). TODO: Move to appendix or omit

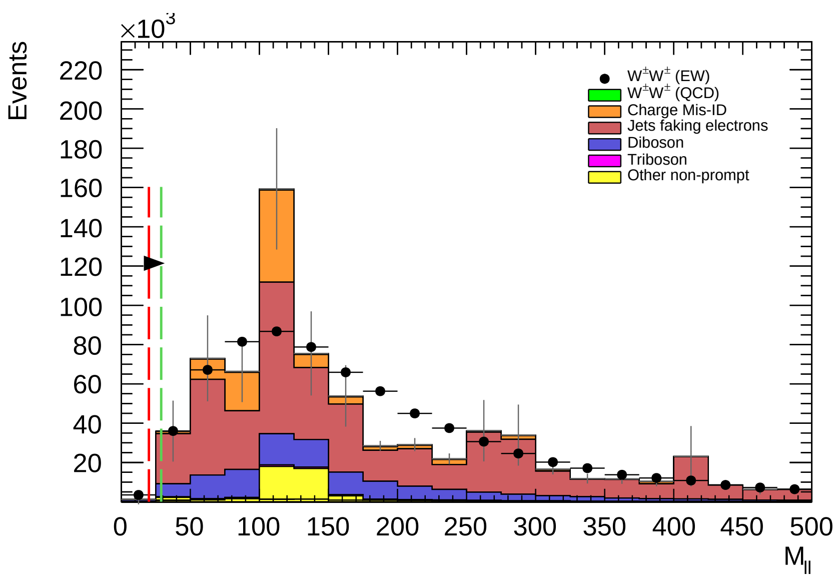
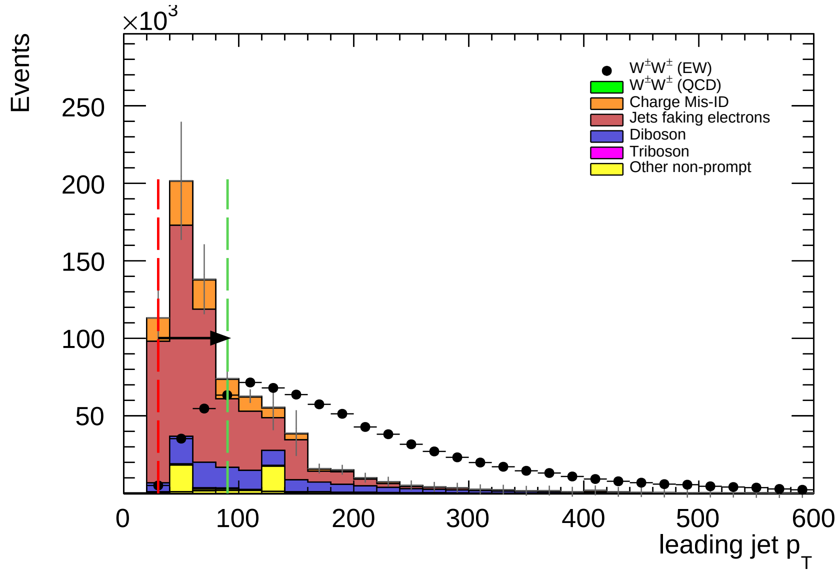


Figure 6.6: Dilepton invariant mass distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^{\pm}W^{\pm}jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram). TODO: Move to appendix or omit



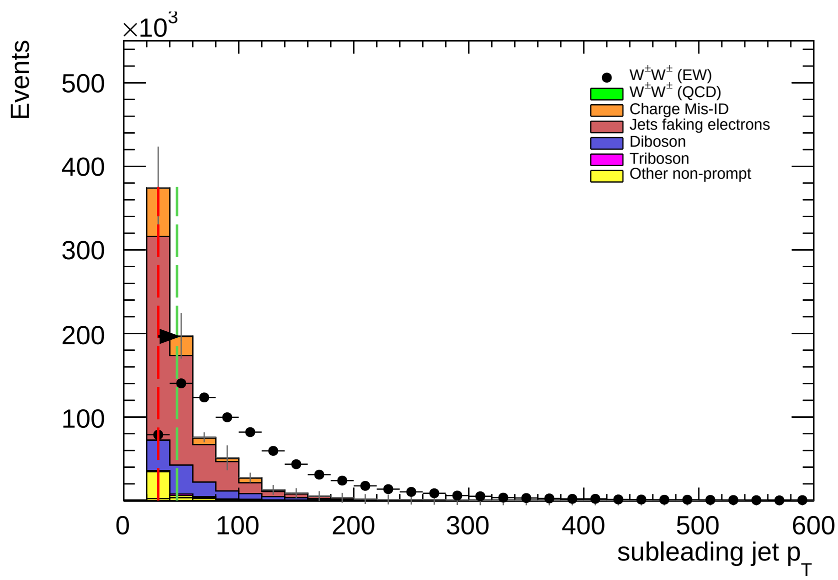


Figure 6.7: Leading (top) and subleading (bottom) jet p_T distributions. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^{\pm}W^{\pm}jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

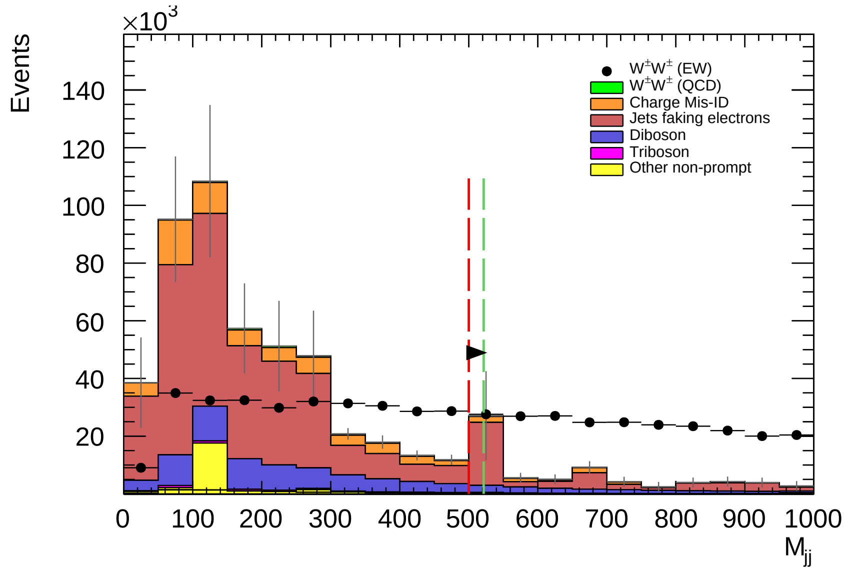


Figure 6.8: Dijet invariant mass distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^{\pm}W^{\pm}jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram). TODO: Move to appendix or omit

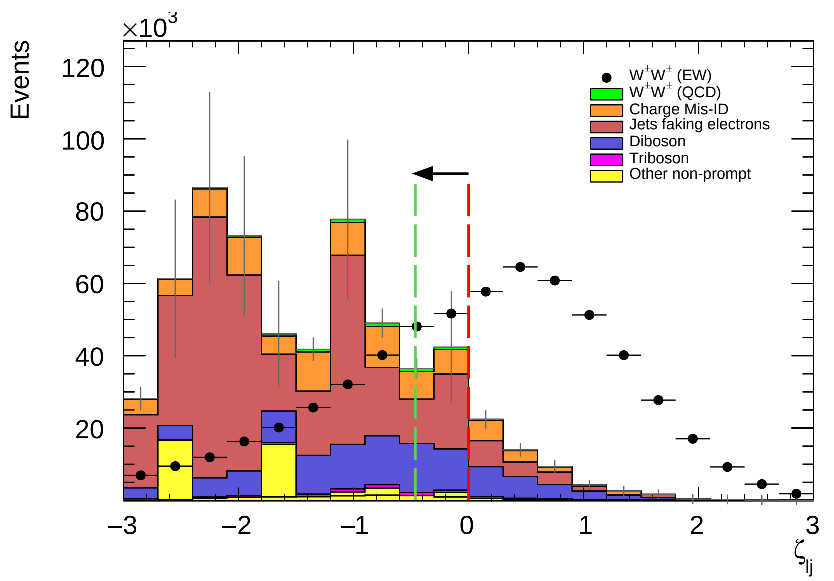


Figure 6.9: Lepton-jet centrality distribution. The default and optimized cuts are represented by the red and green dashed lines, respectively. The $W^{\pm}W^{\pm}jj$ EWK signal (black points) is normalized to the same area as the sum of the backgrounds (colored histogram).

Selection requirement	Selection value
Lepton kinematics	$p_{\mathrm{T}} > 28$ GeV (leading lepton only)
Jet kinematics	$p_{\mathrm{T}} > 90$ GeV (leading jet) $p_{\mathrm{T}} > 45$ GeV (subleading jet)
Dilepton mass	$m_{ll} > 28$ GeV
Dijet mass	$m_{jj} > 520$ GeV
Lepton-jet centrality	$\zeta > -0.5$

Table 6.3: Updates to the $W^{\pm}W^{\pm}jj$ event selection criteria after optimization. Cuts not listed remain unchanged from the default selection in Table 6.2.

6.6 Results

6.6.1 Event yields

6.6.2 Uncertainties

6.6.3 Cross section measurement

CHAPTER 7

Conclusion

Here's where you wrap it up.

Looking Ahead

Here's an example of how to have an "informal subsection".

447 # Appendix A

448 # Additional material on truth isolation

yields by type	all channels	$\mu\mu$	ee	μe	$e\mu$
signal	4011	1583.2	531.7	793.1	1103.1
ww qcd	252.6	105.8	30.4	48	68.4
charge flip	2528.4	0.0	2075.4	255.1	197.8
fakes	7135.4	0.0	4675.1	1904.3	555.9
diboson	2370.4	581.2	491.8	517.9	779.6
triboson	125.5	49.1	17.8	24.6	34.1
top	90150.5	26618	15301.6	25277.9	22953.1
z+jets	241.2	0.0	0.0	0.0	241.2
w+jets	31.4	3.9	7.6	13.2	6.7
total bkg	102803.9	27354	22592	28027.8	24830.1
signal	4011	1583.2	531.7	793.1	1103.1

Table A.1: Event yields prior to applying any form of truth-based isolation criteria.

yields by type	all channels	$\mu\mu$	ee	μe	$e\mu$
signal	3470.5	1427.3	428.8	675.8	938.7
ww qcd	205.8	90.8	22.7	38.3	54
charge flip	2398.3	0.0	2104.6	95.8	197.9
fakes	4309.7	0.0	3390.6	750.8	168.3
diboson	1552.4	311.3	355.6	346.8	538.7
triboson	115	46.8	15.4	21.6	31.2
top	156.9	42.3	14.8	76.6	23.3
z+jets	0.0	0.0	0.0	0.0	0.0
w+jets	0.3	0.0	0.0	0.3	0.0
total bkg	8738.1	491.3	5903.7	1329.8	1013.4
signal	3470.5	1427.3	428.8	675.8	938.7

Table A.2: Event yields after applying a test version of the truth-based isolation.

449 TODO: Add tables for tight vs loose working point, information on the necessity of TRUTH1++

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