

A SELECTION OF FORMULAE AND DATA USEFUL
FOR THE DESIGN OF A.G. SYNCHROTRONS

C. Bovet, R. Gouiran, I. Gumowski, K.H. Reich

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European Organization for Nuclear Research
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Switzerland

CERN/MPS-SI/Int. DL/70/4
23 April, 1970

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LIST OF FREQUENTLY OCCURRING SYMBOLS, THEIR MEANINGS AND UNITS *)

A	RF "bucket" area (in longitudinal phase plane **) [see page 31]
A_H	acceptance in horizontal phase plane ** [see page 18] (= area of largest
A_V	acceptance in vertical phase plane ** [see page 18] acceptable ellipse/ π)
B, B_0	magnetic flux density, in teslas[T], nominal value
C, C_0	length of orbit [m], nominal value
c	velocity of light [m/s]
e	electronic charge [C]
eV	maximum energy gain per turn [keV]
E	total energy of particle [Gev]
E_0	rest energy of particle [Gev]
f_a	accelerating frequency [Hz]
f	revolution frequency [Hz]
f_∞	asymptotic value of f_a reached at $\beta = 1$
g	gradient of magnetic field, in teslas per metre [Tm^{-1}]
h	harmonic number = f_a/f
K	focal constant [m^{-2}]
m	mass of particle [GeV/c^2]
m_p	mass of proton [GeV/c^2]
n	field index = $(-\rho_0/B_0)(\partial B/\partial x)$
p, p_0	momentum of particle [GeV/c], nominal value
Q	number of betatron oscillations per revolution
R, R_0	mean orbit radius (= $C/2\pi$), nominal value, [m] unless stated otherwise
T	kinetic energy of particle [GeV]
V	peak accelerating voltage per turn [kV]

*) In square brackets

**) With these definitions the available six-dimensional hypervolume
is $A_6 = \pi^2 A_H A_V RA$ where A is in $(\Delta p/m_0 c) - \varphi$ coordinates.

α_p	momentum compaction factor
$\alpha(s)$	Twiss parameter
β	ratio of particle velocity to that of light ($= v/c$)
$\beta(s), \beta_{H,V}$	betatron amplitude function
Γ	$= \sin \varphi_s$ where φ_s refers to the synchronous particle
$\gamma(s)$	Twiss parameter
ϑ	deflection angle
ϵ	emittance in transverse plane [see page 18] ($=$ area of ellipse/ π)
ϵ_H	horizontal beam emittance* [see page 18] occupied by beam in
ϵ_V	vertical beam emittance* [see page 18] respective plane)
Θ	azimuthal angle
μ	phase shift of betatron oscillation for one focusing period
ρ	bending radius [m], positive from centre towards outside
φ	"phase angle" between particle and zero crossing of RF voltage
φ_s	"phase angle" for synchronous (phase stationary) particle
$\psi(s)$	phase advance of the betatron oscillation

Other symbols are defined as they occur.

Coordinate system of particle: (Definitions of s, x, y, z as in Courant and Snyder[1.3**])

s	distance along beam axis
x	horizontal transverse coordinate, same sign as ρ
z	vertical transverse coordinate, positive towards sky
y	general transverse coordinate
\bar{x}	arithmetic mean of x
$\langle x \rangle$	r.m.s. value of x

A prime denotes differentiation with respect to s .

A dot denotes differentiation with respect to time.

*) With these definitions the six-dimensional invariant hypervolume occupied by the beam is $V_6 = (\pi \beta \gamma)^2 \epsilon_H \epsilon_V R S_\varphi$ where S_φ is the area [in $(\Delta p/m_e c) - \varphi$ coordinates] occupied by one bunch in the longitudinal phase plane

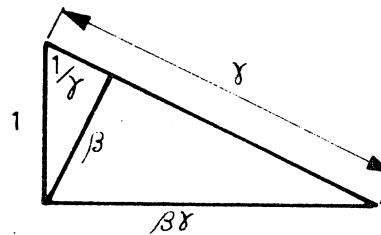
**) See page 43 for references

P A R T I

BASIC RELATIONS

1. PARTICLE VELOCITY, MOMENTUM AND ENERGY

1.1 Relations between β , cp , E_0 , T , E , γ



In terms of wanted	β	cp	T	E	γ
$\beta =$	β	$[(E_0/cp)^2 + 1]^{-\frac{1}{2}}$	$[1 - (1 + T/E_0)^{-2}]^{\frac{1}{2}}$	$[1 - (E_0/E)^2]^{\frac{1}{2}}$	$(1 - \gamma^{-2})^{\frac{1}{2}}$
		cp/E		cp/E	
$cp =$	$E_0(\beta^{-2} - 1)^{-\frac{1}{2}}$	cp	$[T(2E_0 + T)]^{\frac{1}{2}}$	$(E^2 - E_0^2)^{\frac{1}{2}}$	$E_0(\gamma^2 - 1)^{\frac{1}{2}}$
	$E\beta$		$T[(\gamma + 1)/(\gamma - 1)]^{\frac{1}{2}}$	$E\beta$	
$E_0 =$	$cp/\beta\gamma$	$cp(\gamma^2 - 1)^{-\frac{1}{2}}$	$T/(\gamma - 1)$	$(E^2 - c^2 p^2)^{\frac{1}{2}}$	E/γ
	$E(1 - \beta^2)^{\frac{1}{2}}$				
$T =$	$[(1 - \beta^2)^{-\frac{1}{2}} - 1]E_0$	$[E_0^2 + c^2 p^2]^{\frac{1}{2}} - E_0$	T	$E - E_0$	$E_0(\gamma - 1)$
		$cp[(\gamma - 1)/(\gamma + 1)]^{\frac{1}{2}}$			
$\gamma =$	$(1 - \beta^2)^{-\frac{1}{2}}$	$cp/E_0\beta$	$1 + T/E_0$	E/E_0	γ
		$[1 + (cp/E_0)^2]^{\frac{1}{2}}$			

In a synchrotron:

$$\beta = 2\pi Rf/c \quad (f = f_a/h)$$

$$\text{and } p[\text{GeV}/c] = 0.2997925 B\rho [\text{Tm}], \quad \text{or} \quad p [\text{VAs}^2 \text{m}^{-1}] = eB\rho [\text{As Tm}].$$

1.2 First Derivatives

In terms of wanted	$d\beta$	$d(cp)$	$dY = dE/E_0 = dT/T_0$
$d\beta =$	$d\beta$	$[1 + (cp/E_0)^2]^{-\frac{3}{2}} d(cp)/E_0$	$\gamma^{-2}(\gamma^2 - 1)^{-\frac{1}{2}} dY$
		$\gamma^{-3} d(cp)/E_0$	$\beta^{-1} \gamma^{-3} dY$
$d(cp) =$	$E_0(1 - \beta^2)^{-\frac{3}{2}} d\beta$	$d(cp)$	$E_0 \gamma (\gamma^2 - 1)^{-\frac{1}{2}} dY$
	$E_0 \gamma^3 d\beta$		$E_0 \beta^{-1} dY$
$dY =$ $= dE/E_0 =$ $= dT/T_0 =$	$\beta(1 - \beta^2)^{-\frac{3}{2}} d\beta$	$[1 + (E_0/cp)^2]^{-\frac{1}{2}} d(cp)/E_0$	dY
	$\beta \gamma^3 d\beta$	$\beta d(cp)/E_0$	

1.5 Logarithmic first derivatives

In terms of wanted	$d\beta/\beta$	dp/p	dT/T	$dE/E = dY/Y$
$d\beta/\beta =$	$d\beta/\beta$	$\gamma^{-2} dp/p$	$[\gamma(\gamma + 1)]^{-1} dT/T$	$(\gamma^2 - 1)^{-1} dY/Y$
		$dp/p - dY/Y$		$(\beta\gamma)^{-2} dY/Y$
$dp/p =$	$\gamma^2 d\beta/\beta$	dp/p	$[\gamma/(\gamma + 1)] dT/T$	$\beta^{-2} dY/Y$
$dT/T =$	$\gamma(\gamma + 1) d\beta/\beta$	$(1 + \gamma^{-1}) dp/p$	dT/T	$\gamma(\gamma - 1)^{-1} dY/Y$
$dE/E =$	$(\beta\gamma)^2 d\beta/\beta$	$\beta^2 dp/p$	$(1 - \gamma^{-1}) dT/T$	dY/Y
	$(\gamma^2 - 1) d\beta/\beta$	$dp/p - d\beta/\beta$		

See page 1 for meaning of symbols.

2. ENERGY AVAILABLE IN COLLISION BETWEEN TWO PARTICLES

β , γ , and θ_c measured in the laboratory frame.

2.1 General two-body collision along the same line

$$E_{c.m.} = [m_1^2 + m_2^2 + 2m_1m_2\gamma_1\gamma_2(1 - \beta_1\beta_2)]^{1/2},$$

where β is counted algebraically, and $E_{c.m.}$ is the total energy in the centre-of-mass frame, i.e. the maximum available energy.

2.2 Two identical particles

i) One particle at rest: $\beta_1 = 0$, $\gamma_1 = 1$

$$E_{c.m.} = m(2 + 2\gamma_2)^{1/2} \approx m(2\gamma_2)^{1/2} \quad \text{for } \gamma_2 \gg 1.$$

ii) Two particles having velocities of the same magnitude but of opposite sign: $\gamma_1 = \gamma_2 = \gamma$; $\beta_1 = -\beta_2$:

$$E_{c.m.} = 2E = 2m\gamma.$$

If a proton colliding with another proton at rest can liberate the same energy as a collision between two protons with opposite velocities, its energy is defined by

$$\gamma_{eq} = 2\gamma^2 - 1 \approx 2\gamma^2 \quad \text{for } \gamma \gg 1.$$

iii) Two particles having velocities of the same magnitude but making a small angle θ_c :

$$\begin{aligned} E_{c.m.} &= 2E(1 - \beta^2 \sin^2(\theta_c/2))^{1/2} \\ &\approx 2E \cos(\theta_c/2) \quad \text{for } \beta \approx 1. \end{aligned}$$

$$\begin{aligned} \gamma_{eq} &= 2\gamma^2 \cos^2(\theta_c/2) - 1 \\ &\approx 2\gamma^2 \cos^2(\theta_c/2) \quad \text{for } \gamma \gg 1. \end{aligned}$$

3. MAGNETIC AND ELECTRIC DEFLECTION

3.1 Magnetic deflection

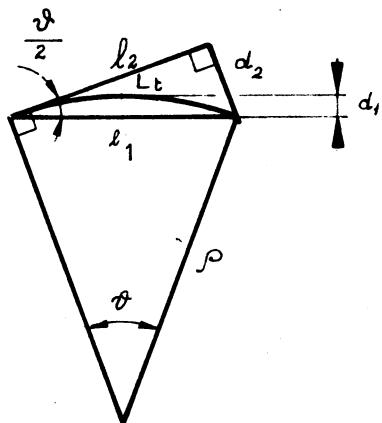
a) Deflection angle ϑ [rad] = $BL_t/(B\rho) = 0.2997925 BL_t/p$ [Tm/GeV/c)]

b) Beam rigidity (magnetic bending radius)

$$B\rho \text{ [Tm]} = 3.5356 p \text{ [GeV/c]}$$

= $3.1297 \beta\gamma$ for protons (refer to Table 1.1 for other expressions)

c) Sagitta

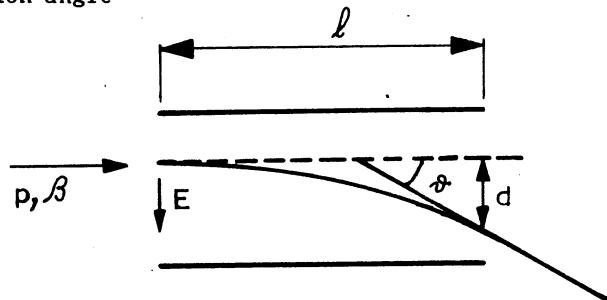


$$d_1 = \frac{1}{2} L_t \tan(\vartheta/4) = 2\rho \sin^2(\vartheta/4)$$
$$\approx \frac{\rho\vartheta^2}{8} = \frac{L_t^2}{8\rho}$$

$$d_2 = L_t \tan(\vartheta/2) = \rho(1 - \cos \vartheta)$$
$$\approx \frac{\rho\vartheta^2}{2} = \frac{L_t^2}{2\rho}$$

3.2 Electric deflection

a) Deflection angle



$$\vartheta \text{ [rad]} = \arctan(E\ell/p\beta) [10^9 \text{ V/(GeV/c)}]$$

b) Sagitta

$$d[\text{m}] = E\ell^2/2\beta p [10^9 \text{ Vm/(GeV/c)}].$$

5.5 Comparison of electric and magnetic deflection

For small ϑ ,

$$B[T] \approx E/(300 \beta) [MV/m] \text{ for the same deflection.}$$

Equivalent deflection for high fields, $B = 2T$, $E = 10 MV/m$ corresponds to $\beta = 1/60$; and, for protons, $p = 16 MeV/c$, $T = 0.13 MeV$.

4. SOME FORMULAE FOR QUANTITIES RELATED TO SYNCHROTRONS

4.1 Mean machine radius

$$R_0 = C_0/2\pi = (1 + k)\rho_0, \text{ where } k \text{ is the circumference factor.}$$

4.2 Relations between p , R , B , f , β and their derivatives

4.2.1 p, R, B^*)

a) Fundamental equation

$$p = e\rho_0 (R/R_0)^{1/\alpha_p} B$$

b) Definition of α_p

$$\alpha_p = \frac{p}{R} \left(\frac{\partial R}{\partial p} \right)_B \left(\approx \frac{1}{Q^2} \right)$$

4.2.2 f, β, R

$$f = \beta c / 2\pi R.$$

4.2.3 Definition of transition energy $E_{tr} = \gamma_{tr} E_0$

$$\frac{p}{f} \left(\frac{\partial f}{\partial p} \right)_B = \frac{1}{\gamma_{tr}^2} - \alpha_p = 0$$

$$\gamma_{tr} = 1/\sqrt{\alpha_p} (\approx Q).$$

*) B is defined on the nominal orbit C_0 .

4.2.4 Differential relations

Variables	Equations
B, p, R	$\frac{dp}{p} = \gamma_{tr}^2 \frac{dR}{R} + \frac{dB}{B}$
f, p, R	$\frac{dp}{p} = \gamma^2 \frac{df}{f} + \gamma^2 \frac{dR}{R}$
B, f, p	$\frac{dB}{B} = \gamma_{tr}^2 \frac{df}{f} + \frac{\gamma^2 - \gamma_{tr}^2}{\gamma^2} \frac{dp}{p}$
B, f, R	$\frac{dB}{B} = \gamma^2 \frac{df}{f} + (\gamma^2 - \gamma_{tr}^2) \frac{dR}{R}$

4.3 Relation between currents and number of particles

- a) Number of injected particles in terms of linac current:

In the case of multturn injection

$$N = 1.5082 \times 10^{11} n_t \epsilon_a \epsilon_{tp} R I_L / \beta \quad [\text{Am}]$$

where

N is the number of trapped particles in the synchrotron;

n_t is the number of injected turns;

I_L is the linac current [A];

ϵ_a is the mean transverse phase space injection efficiency;

ϵ_{tp} is the longitudinal trapping efficiency.

- b) Circulating current:

$$I \quad [\text{A}] = (ec/2\pi)(N\beta/R) = 7.6441 \times 10^{-12} N\beta/R \quad [\text{mA}] .$$

Number of charges passing per microsecond = $6.2418 \times 10^9 I \quad [\text{mA}]$,

or $I \quad [\text{mA}] = 1.6021 \times 10^{-10} \times \text{number of charges passing per microsecond.}$

ADDITIONAL FORMULAE

- 11 -

P A R T II

TRANSVERSE PHASE SPACE

1. MATRIX FORMULATION OF BEAM DYNAMICS

1.1 General form of matrices with dispersive terms

The general matrix M is defined by

[2,52]*
[3,52]

$$\begin{pmatrix} y \\ y' \\ \Delta p/p \end{pmatrix}_B = M(B|A) \begin{pmatrix} y \\ y' \\ \Delta p/p \end{pmatrix}_A$$

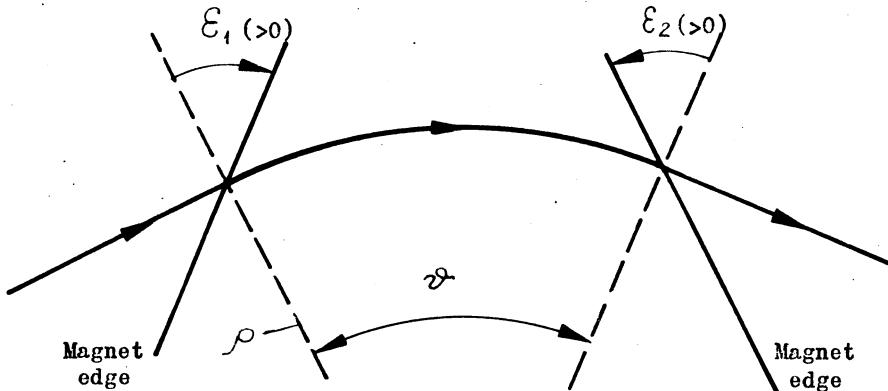
1.2 Drift length ℓ

$$M_\ell = \begin{pmatrix} 1 & \ell & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

1.3 Dipole magnet

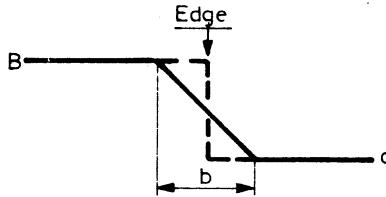
a) Notation and definitions

[3,154]



* See page 43 for references

Fringe field, linear approximation:



b) Pure sector magnet

$$(\epsilon_1 = \epsilon_2 = 0, b = 0)$$

$$M_H^S = \begin{pmatrix} \cos \vartheta & \rho \sin \vartheta & \rho(1 - \cos \vartheta) \\ -\frac{\sin \vartheta}{\rho} & \cos \vartheta & \sin \vartheta \\ 0 & 0 & 1 \end{pmatrix} \quad \text{in horizontal plane (plane of deflection)}$$

$$M_V^S = \begin{pmatrix} 1 & \rho \vartheta & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \text{in vertical plane}$$

c) Magnet with parallel faces

$$M_H^R = \begin{pmatrix} 1 & \rho \sin \phi & 0 \\ 0 & 1 & 0 \end{pmatrix} \quad M_V^R = M_V^E M_V^S M_V^E$$

d) Edge effect with linear fringe field

$$[4,100]^* \quad M_H^E = \begin{pmatrix} 1 & 0 & 0 \\ \frac{\tan \epsilon}{\rho} & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \begin{array}{l} \epsilon = \epsilon_1 \text{ for entrance} \\ \epsilon = \epsilon_2 \text{ for exit} \end{array}$$

$$M_V^E = \begin{pmatrix} 1 & 0 & 0 \\ \frac{1}{\rho} \left(\frac{b}{6\rho \cos \epsilon} - \tan \epsilon \right) & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \begin{array}{l} \epsilon = \epsilon_1 \text{ for entrance} \\ \epsilon = \epsilon_2 \text{ for exit} \end{array}$$

*) See page 43 for references

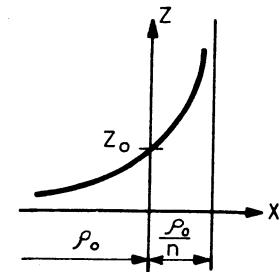
1.4 Gradient sector magnet

a) Equation of profile

$$\left(\frac{\rho_0}{n} - x \right) z = \frac{\rho_0}{n} z_0$$

z_0 is the half-aperture where $\rho = \rho_0$.

$$\begin{aligned} \text{Also } g[\text{Tm}^{-1}] &= - n B_0 / \rho_0 & [\text{Tm}^{-1}] \\ &= - 3.3356 n p_0 / \rho_0^2 & [(\text{GeV}/c)\text{m}^{-2}] \end{aligned}$$



b) Focusing plane

$$[3,53] \quad M_F^* = \begin{pmatrix} \cos \zeta & \frac{1}{\sqrt{K}} \sin \zeta & \frac{1}{\rho K} (1 - \cos \zeta) \\ -\sqrt{K} \sin \zeta & \cos \zeta & \frac{1}{\rho \sqrt{K}} \sin \zeta \\ 0 & 0 & 1 \end{pmatrix},$$

where $K[\text{m}^{-2}] = (|n| + 1)/\rho_0^2$ (horizontal plane)

$= |n|/\rho_0^2$ (vertical plane)

$\zeta = \ell_m \sqrt{K}$, ℓ_m is the magnetic length.

c) Defocusing plane

$$M_D^* = \begin{pmatrix} \cosh \zeta & \frac{1}{\sqrt{K}} \sinh \zeta & \frac{1}{\rho K} (\cosh \zeta - 1) \\ \sqrt{K} \sinh \zeta & \cosh \zeta & \frac{1}{\rho \sqrt{K}} \sinh \zeta \\ 0 & 0 & 1 \end{pmatrix},$$

where $K[\text{m}^{-2}] = (|n| - 1)/\rho_0^2$ (horizontal plane)

$= |n|/\rho_0^2$ (vertical plane)

$\zeta = \ell_m \sqrt{K}$, ℓ_m is the magnetic length

d) Pure quadrupole lens

The same matrices as in points (b) and (c) are used, with $1/\rho = 0$

and $K[\text{m}^{-2}] = 0.2997925 g/p_0$ $[\text{Tm}^{-1}/(\text{GeV}/c)]$

$= - n/\rho_0^2 = g/(B\rho)$ $[\text{m}^{-2}]$

$= 0.31952 g/\beta_Y$ $[\text{Tm}^{-1}]$ for protons.

*) Valid for coordinate system defined on page 2.

2. DESCRIPTION OF SINGLE PARTICLE MOTION IN A SYNCHROTRON

2.1 Equation of motion

$$[2,33] \quad d^2x/ds^2 + K_x(s)x = (\Delta p/p)/\rho(s)$$

$$[5,55] \quad d^2z/ds^2 + K_z(s)z = 0.$$

[1,3]

2.2 Solution of the equation of motion

$$[1,11] \quad y(s) = \sqrt{\epsilon/\beta(s)} \cos[\psi(s) + \delta]$$

$$[2,79] \quad \left\{ \begin{array}{l} y'(s) = -\sqrt{\epsilon/\beta(s)} \{ \alpha(s) \cos[\psi(s) + \delta] + \sin[\psi(s) + \delta] \} = \sqrt{\epsilon\gamma(s)} \cos[\chi(s) + \delta] \\ \psi(s) = \int_0^s ds/\beta(s) \end{array} \right.$$

$$[6,4] \quad \beta(s)\gamma(s) = 1 + \alpha^2(s)$$

$$\beta'(s) = -2\alpha(s)$$

$$\tan[\psi(s) - \chi(s)] = 1/\alpha(s)$$

$$\sin[\psi(s) - \chi(s)] = -[\beta(s)\gamma(s)]^{-1/2}.$$

Envelope equation:

$$\sqrt{\beta''} + K(s)\sqrt{\beta} - \beta^{-3/2} = 0$$

Initial conditions:

$$\epsilon = \gamma(0)y^2(0) + \beta(0)y'^2(0) + 2\alpha(0)y(0)y'(0)$$

$$\cos \delta = y(0)/\sqrt{\epsilon\beta(0)}$$

$$\tan \delta = -[\alpha(0) + \beta(0)y'(0)/y(0)].$$

For $y(0) = 0$:

$$y(s) = y'(0)\sqrt{\beta(0)\beta(s)} \sin \psi(s)$$

$$y'(s) = -y'(0)\sqrt{\beta(0)/\beta(s)} [\alpha(s)\sin \psi(s) - \cos \psi(s)].$$

$$Q = \psi(0)/(2\pi)$$

Betatron wavelength

$$\lambda = 2\pi R/Q$$

Form factor

$$F = \beta_{\max} Q/R.$$

2.3 Sinusoidal approximation

$$\psi(s) \approx 2\pi s/\lambda = Q 2\pi s/C; \quad \beta(s) \approx R/Q$$

$$y(s) \approx \sqrt{\epsilon R/Q} \cos(Qs/R + \delta).$$

2.4 Motion through one period (cell) of length L_p

$\kappa(s)$, $\alpha(s)$, $\beta(s)$, $\gamma(s)$ are periodic with period L_p ,
and $\psi(s + L_p) = \psi(s) + \mu$, $\chi(s + L_p) = \chi(s) + \mu$.

The transfer matrix through one period may be written as

$$[1,6] \quad M(s + L_p | s) = \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} = \begin{pmatrix} \cos \mu + \alpha(s) \sin \mu & \beta(s) \sin \mu \\ -\gamma(s) \sin \mu & \cos \mu - \alpha(s) \sin \mu \end{pmatrix}$$

$$\cos \mu = \frac{1}{2} (a_{11} + a_{22}) \quad \beta(s) = a_{12} / \sin \mu$$

$$\gamma(s) = -a_{21} / \sin \mu \quad \alpha(s) = \frac{1}{2} (a_{11} - a_{22}) / \sin \mu.$$

2.5 Transfer matrix through any section

- a) If the Twiss parameters at points s_1 , s_2 are $(\beta_1, \alpha_1, \gamma_1)$ and $(\beta_2, \alpha_2, \gamma_2)$, respectively, the 2×2 transfer matrix from s_1 to s_2 can be written as

$$[1,9] \quad M(s_2 | s_1) = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} =$$

$$= \begin{pmatrix} \sqrt{\frac{\beta_2}{\beta_1}} (\cos \Delta\psi + \alpha_1 \sin \Delta\psi) & \sqrt{\beta_1 \beta_2} \sin \Delta\psi \\ -\left[\frac{(1 + \alpha_1 \alpha_2) \sin \Delta\psi + (\alpha_2 - \alpha_1) \cos \Delta\psi}{\sqrt{\beta_1 \beta_2}} \right] & \sqrt{\frac{\beta_1}{\beta_2}} (\cos \Delta\psi - \alpha_2 \sin \Delta\psi) \end{pmatrix},$$

where $\Delta\psi = \psi(s_2) - \psi(s_1)$.

- b) Transformation of the Twiss parameters through a beam transfer section:

$$[7,100] \quad \begin{pmatrix} \beta_2 \\ \alpha_2 \\ \gamma_2 \end{pmatrix} = \begin{pmatrix} m_{11}^2 & -2m_{11}m_{12} & m_{12}^2 \\ -m_{21}m_{11} & 1 + 2m_{12}m_{21} & -m_{12}m_{22} \\ m_{21}^2 & -2m_{22}m_{21} & m_{22}^2 \end{pmatrix} \begin{pmatrix} \beta_1 \\ \alpha_1 \\ \gamma_1 \end{pmatrix}.$$

$$\tan \Delta\psi = m_{12} / [m_{11}\beta(s_1) - m_{12}\alpha(s_1)]$$

Example: drift length ℓ ,

$$[8,2] \quad \begin{aligned} \beta(s_2) &= \beta(s_1) - 2\alpha(s_1)\ell + \gamma(s_1)\ell^2 \\ \alpha(s_2) &= \alpha(s_1) - \gamma(s_1)\ell \\ \gamma(s_2) &= \gamma(s_1), \end{aligned}$$

and $\psi(s_2) = \psi(s_1) + \arctan \{\ell / [\beta(s_1) - \alpha(s_1)\ell]\}$.

2.6 Normalization

After a normalisation transformation $(\frac{\eta}{\eta'}, \frac{y}{y'}) = M(\frac{y}{y}, \frac{y'}{y})$ (with $\eta' = d\eta/d\psi$), the transverse phase space trajectories have the form of circles on which a phase advance $\psi(s)$ produces simply a rotation by ψ . Possible transformations matrices:

$$M = \begin{pmatrix} \frac{1}{\sqrt{\beta(s)}} & 0 \\ \frac{\alpha(s)}{\sqrt{\beta(s)}} & \sqrt{\beta(s)} \end{pmatrix}; \text{ or, when } \alpha = 0, M = \begin{pmatrix} 1 & 0 \\ 0 & \beta(s) \end{pmatrix}$$

3. ELLIPSE REPRESENTATION IN TRANSVERSE PHASE SPACE (see page 18 for units)

3.1 The Courant-Snyder invariant

$$\gamma(s)y^2 + 2\alpha(s)yy' + \beta(s)y'^2 = \epsilon = \frac{\text{area}}{\pi} .$$

The largest area contained in the synchrotron is given by the acceptance $A_{H,V} = r^2/\beta_{\max}$, where r is the half-aperture of the vacuum chamber at β_{\max} .

3.2 Ellipse parameters

An ellipse centred at the origin of the phase plane is determined by three independent parameters. Depending on the particular problem, one can choose one of the following sets:

- i) Twiss parameters α , β , γ and ϵ giving the emittance (see Section 3.1). These parameters transform with the matrix given in Section 2.5 (b);
- ii) the elements c_i of a 2×2 matrix which transforms by multiplication with $M(s_2 | s_1)$; $(c_3 y - c_1 y')^2 + (c_4 y - c_2 y')^2 = \epsilon^2$
- iii) L , S , and ϵ where L is the ratio of the ellipse axes a/b at the waist, S is the distance of the waist along the beam (>0 if waist upstream).

The optical transformations from s_1 to s_2 are:

$$L(s_2) = \frac{L(s_1)}{[m_{21} L(s_1)]^2 + [m_{21} S(s_1) + m_{22}]^2}$$

$$S(s_2) = \frac{m_{11}m_{21} L^2(s_1) + [m_{11} S(s_1) + m_{12}][m_{21} S(s_1) + m_{22}]}{[m_{21} L(s_1)]^2 + [m_{21} S(s_1) + m_{22}]^2} .$$

Conversion from one set to the other is given in Table 3.3.

3.3 Conversion of ellipse parameters

given	$\alpha, \beta, \gamma, \epsilon$	$\begin{pmatrix} c_1 & c_2 \\ c_3 & c_4 \end{pmatrix}$	L, S, ϵ
wanted	$\beta\gamma - \alpha^2 = 1$	$\epsilon = c_1c_4 - c_2c_3$	
α	α	$-(c_1c_3 + c_2c_4)/\epsilon$	$-S/L$
β	β	$(c_1^2 + c_2^2)/\epsilon$	$L + S^2/L$
γ	γ	$(c_3^2 + c_4^2)/\epsilon$	$1/L$
c_1	$\sqrt{\epsilon\beta}$	c_1	$\sqrt{L\epsilon}$
c_2	0	c_2	$S\sqrt{\epsilon/L}$
c_3	$-\alpha\sqrt{\epsilon/\beta}$	c_3	0
c_4	$\sqrt{\epsilon/\beta}$	c_4	$\sqrt{\epsilon/L}$
L	$1/\gamma$	$\epsilon/(c_3^2 + c_4^2)$	L
S	$-\alpha/\gamma$	$(c_1c_3 + c_2c_4)/(c_3^2 + c_4^2)$	S

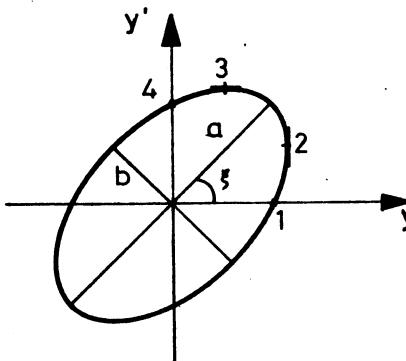
Comments on units:

The equations of Sections 3.1 to 3.4 are valid for either of the two following sets of units:

- i) All lengths in metres, all angles in radians, the emittances in rad m, L ($= y_1/y_{1'}$) in m/rad, S ($= y_3/y_{3'}$) in m/rad.
- ii) All phase plane dimensions in mm and mrad, emittances in mrad mm, $\beta = [\text{mm}/\text{mrad}]$ if defined as $y_2/y_{2'}$ or $\beta = [m]$ if defined as reduced betatron wavelength (similarly $\gamma = [\text{mm}/\text{mrad}]$, or $[m^{-1}]$, L = [mm/mrad], S = [mm/mrad] if defined as $y_3/y_{3'}$ or S = [m] if defined as distance from waist).

N.B.: the values of α, β, γ, L and S do not depend on the choice between i) and ii).

5.4 Geometrical properties of the ellipse



$\alpha, \beta, \gamma, \epsilon$	$c_1 \quad c_2$ $c_3 \quad c_4$	L, S, ϵ	
$\beta\gamma - \alpha^2 = 1$ $H = \frac{1}{2}(\beta + \gamma)$	$\epsilon = c_1c_4 - c_2c_3$ $H = \frac{1}{2}(c_1^2 + c_2^2 + c_3^2 + c_4^2)/\epsilon$	$H = \frac{1}{2L}(L^2 + S^2 + 1)$	
y_1 y_2 y_3 y_4	$\sqrt{\epsilon/\gamma}$ $\sqrt{\epsilon\beta}$ $-\alpha\sqrt{\epsilon/\beta}$ $-\alpha\sqrt{\epsilon/\gamma}$ $\sqrt{\epsilon\gamma}$ $\sqrt{\epsilon/\beta}$	$\epsilon/\sqrt{c_3^2 + c_4^2}$ $\sqrt{c_1^2 + c_2^2}$ $(c_1c_3 + c_2c_4)/\sqrt{c_1^2 + c_2^2}$ $(c_1c_3 + c_2c_4)/\sqrt{c_3^2 + c_4^2}$ $\sqrt{c_3^2 + c_4^2}$ $\epsilon/\sqrt{c_1^2 + c_2^2}$	$\sqrt{\epsilon L}$ $\sqrt{\epsilon/L}\sqrt{S^2 + L^2}$ $S\sqrt{\epsilon/L}/\sqrt{S^2 + L^2}$ $S\sqrt{\epsilon/L}$ $\sqrt{\epsilon/L}$ $\sqrt{\epsilon L}/\sqrt{S^2 + L^2}$
a b $a/b > 1$		$\sqrt{\epsilon/2}(\sqrt{H+1} + \sqrt{H-1})$ $\sqrt{2\epsilon}/(\sqrt{H+1} + \sqrt{H-1}) = \sqrt{\epsilon/2}(\sqrt{H+1} - \sqrt{H-1})$ $H + \sqrt{H^2 - 1}$	
$\tan \xi$ $\sin 2\xi$ $\cos 2\xi$ $\tan 2\xi$	$[-\alpha(H + \sqrt{H^2 - 1})]/[\beta(H + \sqrt{H^2 - 1}) - 1]$ $-\alpha/\sqrt{H^2 - 1}$ $(\beta - \gamma)/2\sqrt{H^2 - 1}$ $-2\alpha/(\beta - \gamma)$	$[c_2 + c_3(H + \sqrt{H^2 - 1})]/[c_1(H + \sqrt{H^2 - 1}) - c_4]$ $(c_1c_3 + c_2c_4)/\epsilon\sqrt{H^2 - 1}$ $(c_1^2 + c_2^2 - c_3^2 - c_4^2)/2\epsilon\sqrt{H^2 - 1}$ $2(c_1c_3 + c_2c_4)/(c_1^2 + c_2^2 - c_3^2 - c_4^2)$	$S/[L(H + \sqrt{H^2 - 1}) - 1]$ $S/L\sqrt{H^2 - 1}$ $(L^2 + S^2 - 1)/2L\sqrt{H^2 - 1}$ $2S/(L^2 + S^2 - 1)$

3.4.4 Two ellipses S_1, S_2 with same area S and centre

Common area S_c is given by

$$\frac{S_c}{S} = \frac{4}{\pi} \operatorname{arc} \tan [D - \sqrt{D^2 - 1}]^{1/2}$$

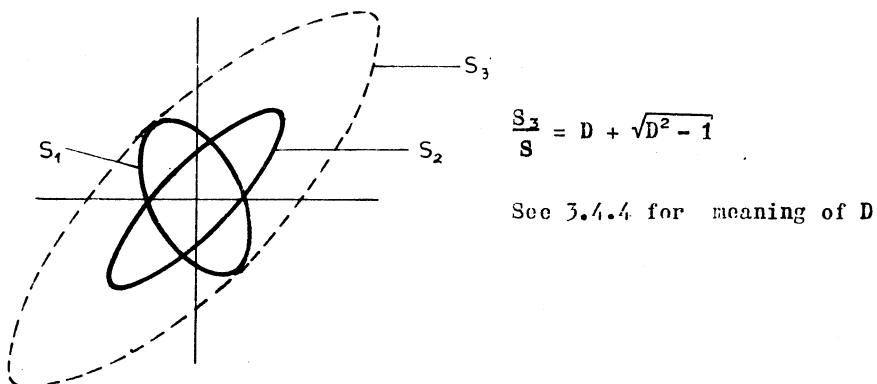
where $D = \frac{1}{2} (\beta_2 \gamma_1 + \gamma_2 \beta_1 - 2\alpha_1 \alpha_2)$

$$= 1 + \frac{(L_2 - L_1)^2 + (S_2 - S_1)^2}{2L_1 L_2} .$$

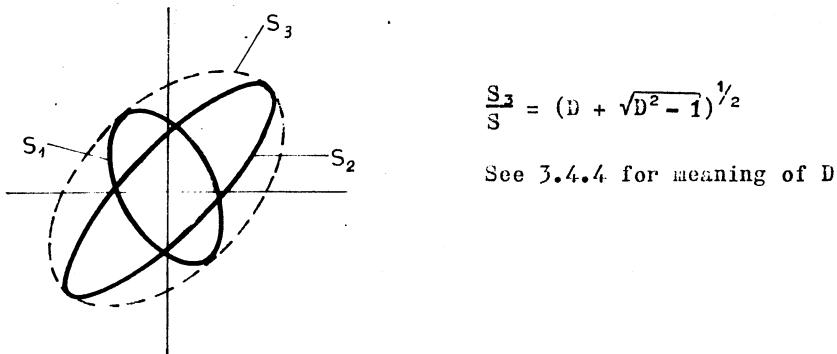
For meaning of $\alpha, \beta, \gamma, L, S$ see 3.2.

3.4.5 Three ellipses

- a) Area of ellipse S_3 , similar to S_2 , such that S_3 circumscribes S_1 : (area S_1 = area $S_2 = S$)

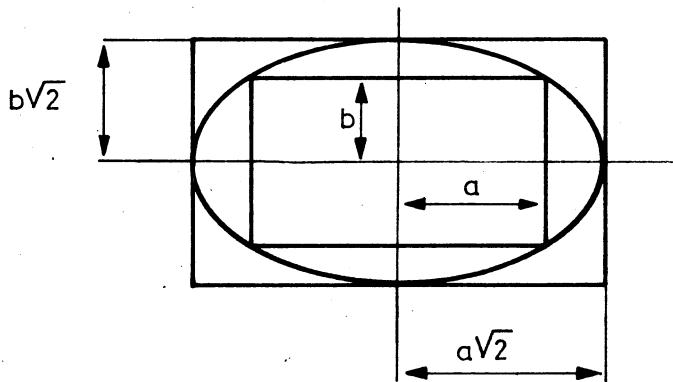


- b) Area of ellipse S_3 circumscribing two ellipses S_1, S_2 of same area S :



3.5 Relations between areas of rectangles, ellipses and circles

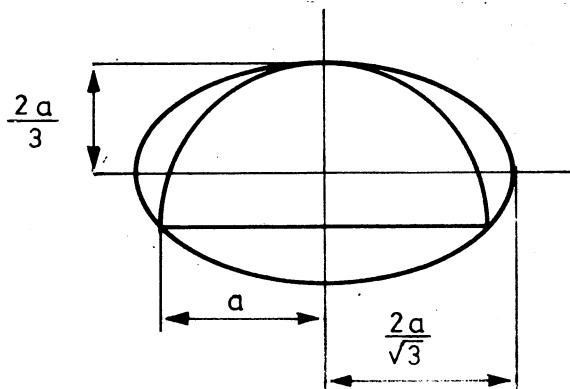
3.5.1 Rectangle and ellipse



$$\frac{S_{\text{min ellipse}}}{S_{\text{inscribed rect.}}} = \frac{\pi}{2}$$

$$\frac{S_{\text{circumscribed rect.}}}{S_{\text{ellipse}}} = \frac{4}{\pi}$$

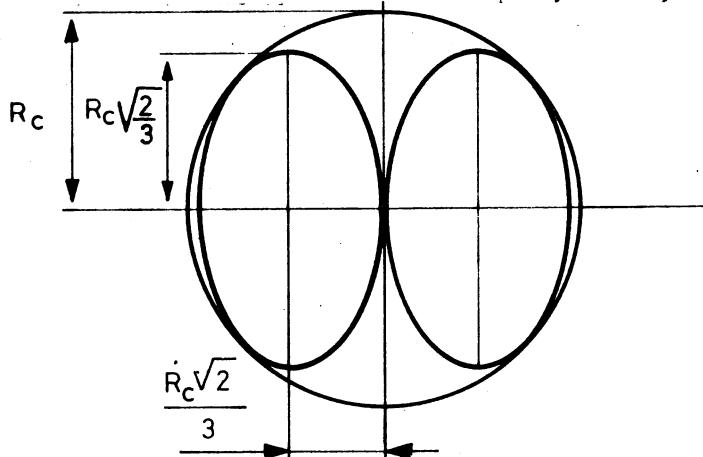
3.5.2 Semi-circle and ellipse



$$\frac{S_{\text{min ellipse}}}{S_{\text{semi-circle}}} = \frac{8\sqrt{3}}{9} \approx 1.54$$

3.5.3 Two maximum ellipses circumscribed by a circle

For a non-zero septum, see [9, Fig. 9]



$$\frac{S_{\text{circle}}}{S_{\text{max ellipse}}} = \frac{3\sqrt{3}}{2} \approx 2.60$$

4. CLOSED ORBIT

4.1 Closed orbit for a momentum deviation $\Delta p/p$

The closed orbit for a given $\Delta p/p$ is given by its phase space coordinates:

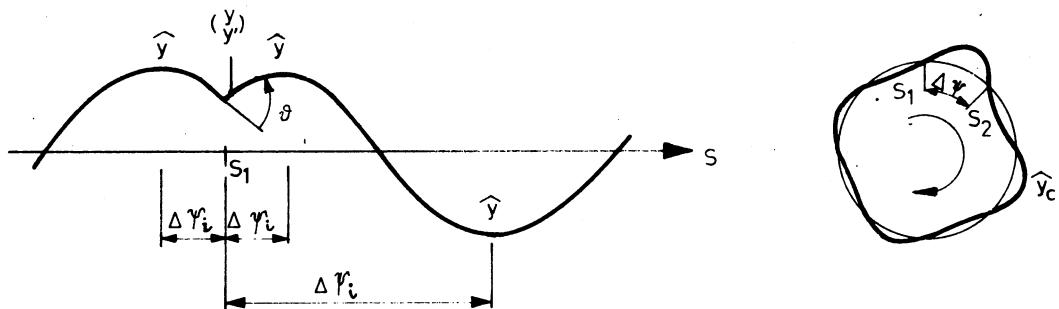
$$[6,19] \quad \begin{pmatrix} e(s) \\ e'(s) \end{pmatrix} = \frac{\Delta p/p}{2(1 - \cos \mu)} \begin{pmatrix} m_{13} + m_{12}m_{23} - m_{22}m_{13} \\ m_{23} + m_{21}m_{13} - m_{11}m_{23} \end{pmatrix},$$

the m_{ij} being the elements of the 3×3 transfer matrix through one period.

The transformation of the closed orbit vector through the machine is as in Section II.1.1., page 12.

4.2 Orbit deformations *)

4.2.1 One and two dipoles



a) One dipole producing a deflection angle ϑ

$$y(s_1) = 0.5\beta(s_1)\vartheta(s_1)\cot(\pi Q)$$

$$y'(s_1) = 0.5\vartheta(s_1)[1 - \alpha(s_1)\cot(\pi Q)].$$

The maximum orbit deviations are

$$\hat{y}(s) = 0.5\vartheta(s_1)[\beta(s_1)\beta(s)]^{1/2}/\sin(\pi Q)$$

and occur approximately at distances (in betatron oscillation phase) of

$$\Delta\psi_i \approx \pm \pi(Q - m), \quad m = 1, 2, 3, \dots < Q.$$

*) Complete decoupling between betatron and synchrotron oscillations, i.e. $\nu_{\text{betatron}} \gg \nu_{\text{synchrotron}}$, is assumed throughout.

b) Reduction of the deformation by a second dipole

For best reduction between s_2 and s_1 of a deformation described in a), a second dipole positioned at s_2 and spaced by $\Delta\psi$ should provide a deflection

$$\vartheta(s_2) = - \cos \Delta\psi [\beta(s_1)\beta(s_2)]^{1/2} \vartheta(s_1).$$

The remaining relative deformations of the corrected orbit are:

between s_2 and s_1 ("outside" dipoles):

$$\hat{y}_c/\hat{y} = \sin \Delta\psi$$

between s_1 and s_2 ("between" dipoles):

$$\hat{y}_c/\hat{y} = [\cos \Delta\psi + 2\sin^2 \pi Q - \cos 2(\pi Q - \Delta\psi)]^{1/2}.$$

4.2.2 Distortions due to random errors

The number of magnet units m is assumed to be

[10,6]*

$$m > 3Q.$$

a) r.m.s. value of $y(s)$

$$\langle y \rangle = \frac{\pi}{\sqrt{2} |\sin \pi Q|} \frac{R}{Q} \frac{|n|}{\rho} \frac{\langle \delta \rangle}{\sqrt{m}},$$

where δ is the position error of one of the m gradient magnets.

More generally, one has

$$\langle y \rangle = \frac{1}{2\sqrt{2} |\sin \pi Q|} \sqrt{\bar{\beta}} \sqrt{\sum_i m_i \beta_i \Psi_i^2}$$

where $\bar{\beta} = (1/C) \int_0^C \beta(s) ds$ and the equivalent kicks Ψ_i are for the various cases of interest:

*) See page 43 for references

Type of element	Source of kick	r.m.s. value	ψ_i	Directions
Gradient element	Displacement	$\langle \Delta y \rangle$	$K_i \ell_i \langle \Delta y \rangle$	x and z
Bending elements	Tilt	$\langle \Delta \theta_e \rangle$	$\vartheta_i \langle \Delta \theta_e \rangle$	z only
Bending elements	Field error	$\langle \Delta B/B \rangle$	$\vartheta_i \langle \Delta B/B \rangle$	x only
Straight sections	Stray field	$\langle \Delta B_s \rangle$	$\ell_i \langle \Delta B_s \rangle / \rho B_{inj}$	x and z
Gradient and bending elements	Displacement and field error	$\langle \Delta y \rangle$ and $\langle \Delta B/B \rangle$	$K_i \ell_i [\langle \Delta y^2 \rangle + \rho^2/n^2 \langle (\Delta B/B)^2 \rangle]^{1/2}$	x and z

b) Value \hat{y} not exceeded with a probability P

$$\hat{y}_P(s) = k(P) \left[1 + \frac{|\sin \pi Q|}{3} \right]^* \sqrt{\frac{\beta(s)}{\bar{\beta}}} \sqrt{2} \langle y \rangle$$

with $k(P)$ given by

[12, Fig.1]

k	P	50%	75%	90%	98%
rectangular vacuum chamber	1.11	1.41	1.72	2.14	
elliptical vacuum chamber	1.28	1.63	1.95	2.40	

[11, Fig.2] *) This bracket takes into account the mean influence of higher harmonics.

5. EFFECTS OF VARIOUS FOCUSING PERTURBATIONS ON THE FREQUENCY AND AMPLITUDE OF BETATRON OSCILLATIONS

5.1 Change in frequency due to tuning of quadrupole lenses or gradient errors

The frequency shift is given by

$$[1,25]* \quad \cos(2\pi Q) - \cos(2\pi Q_0) = 0.5 \sin(2\pi Q_0) \int_0^C \beta(s)k(s)ds$$

where $2\pi Q_0$ is the unperturbed phase shift around the orbit of length C and $k(s)$ the focal constant of the perturbation(s).

In the case of small frequency shifts this becomes

$$\Delta Q = (1/4\pi) \int_0^C \beta(s)k(s)ds .$$

For $m \gg Q$, random errors in m elements produce a shift

$$[1,27] \quad \Delta Q = \frac{1}{4\pi} \sqrt{\sum_i^m (\beta_i k_i e_i)^2} < \frac{\Delta K}{K} >$$

where the symbols are the same as in Section 4.2.2, pages 23 and 24.

5.2 Tuning of momentum dependent frequency shifts by means of sextupoles

To compensate a shift caused by $k(s) = -K(s)\Delta p/p$ one needs (in the case of an ideal closed orbit) a sextupole field $\partial^2 B_z / \partial x^2$ such that

$$\int_0^C \beta(s) \left[\frac{\partial^2 B_z}{\partial x^2}(s)e(s) - B\beta k(s) \right] ds = 0$$

where $e(s)$ is defined in Section 4.1 on page 22.

5.3 "Beating" of amplitudes

The beat factor characterising the amplitude function modified by gradient errors is

$$[1,25] \quad G = [\beta(\text{actual})/\beta(\text{ideal})]_{\max} .$$

*) See page 43 for references

In practice one is more interested in $(\hat{\Delta y}/\hat{y})_P = 0.5(G-1)$. Similarly as for the orbit (Section 4.2.2. page 23) one has:

$$(\hat{\Delta y}/\hat{y})_P = \frac{k(P)}{4} \left[\frac{1}{3} + \frac{1}{|\sin 2\pi Q|} \right] \sqrt{\sum_i m_i (\beta_i K_i \ell_i)^2} < \frac{\Delta K}{K} >$$

5.4 Stopbands due to random gradient errors

The total width of the stopband is

$$\delta Q = 2\Delta Q$$

where the ΔQ is given in Section 5.1 on page 25.

6. SPACE CHARGE LIMIT

Symbols:

N	:	limit of the number of particles in the synchrotron
B_f	:	bunching factor (< 1)
$b[m]$:	mean semi-minor beam axis (vertical)
$a[m]$:	mean semi-major beam axis (horizontal)
$\Delta(Q^2)$:	$Q_0^2 - Q^2 \approx 2Q_0\Delta Q$
r	:	classical particle radius ($= e/\{4\pi \epsilon_0 m c^2 [eV]\}$, see p. 45)
$2h[m]$:	vertical aperture of the vacuum chamber
$2w[m]$:	horizontal aperture of the vacuum chamber
$2v[m]$:	height of the magnet gap.

a) Individual particle limit (without neutralization)

[13,331]*

$$N_{ind} = -0.5 \pi b(a+b)(R r F)^{-1} \beta^2 \gamma^3 B_f \Delta(Q_i^2)$$
$$\approx -(\pi \epsilon_V \beta \gamma)(1 + \sqrt{\epsilon_H/\epsilon_V})(rF)^{-1} \beta \gamma^2 B_f \Delta Q_i$$

where

$$F = 1 + [b(a+b)/h^2] \{ \epsilon_1 [1 + B_f(\gamma^2 - 1)] + \epsilon_2 B_f(\gamma^2 - 1)(h^2/v^2) \}$$

with ϵ_1, ϵ_2 , the image force coefficients given on page 27, and ϵ_H, V in rad m.

*) See page 43 for references

w/h	1(circle)	5/4	4/3	3/2	2/1	∞ (parallel plates)
ε_1	0	0.090	0.107	0.134	0.172	0.206
ξ_1	0.5	0.553	0.559	0.575	0.599	0.617

For parallel straight pole pieces, and to good approximation for wedged-shaped poles, the magnetostatic image coefficients have the values $\varepsilon_2 = 0.411$, $\xi_2 = 0.617$.

b) Coherent particle limit (without neutralization)

[13,342]*

$$N_{coh} = -\pi Q_0 h^2 (R r F)^{-1} \beta^2 \gamma^3 B_f \Delta Q_c$$

where near an integral resonance

$$F = \xi_1 [1 + B_f (\gamma^2 - 1)] + \xi_2 B_f (\gamma^2 - 1) h^2 / v^2$$

and near a half-integral resonance

[13a,150]

$$F = \xi_1 + \varepsilon_1 B_f (\gamma^2 - 1) + \varepsilon_2 B_f (\gamma^2 - 1) h^2 / v^2$$

with ε_1 , ε_2 , ξ_1 and ξ_2 as given above.

For $B_f \gamma^2 \gg 1$, one has near an integral resonance

$$N_{coh} \approx -\pi Q_0 h^2 [R r (\xi_1 + \xi_2 h^2 / v^2)]^{-1} \gamma \Delta Q_c$$

*) See page 43 for references

ADDITIONAL FORMULAE

P A R T III

LONGITUDINAL PHASE SPACE

1. ACCELERATING VOLTAGE

$$V \sin \phi_s [\text{kV}] = 10^{-3} [C(\rho \dot{B} + B \dot{\rho}) - \dot{B} S_F] [\text{m}^2 \text{T/s}]$$

where the index s refers to the synchronous particle and S_F is an equivalent area such that $B S_F$ is the total flux enclosed by C .

For $\dot{\rho} = S_F = 0$, one has

$$V \sin \phi_s [\text{kV}] = 0.020958 R_p [\text{m(GeV/c)/s}]$$

(see Table I.1.1 for other expressions).

2. ACCELERATING FREQUENCY

$$\begin{aligned} f_a [\text{Hz}] &= h f = h e \beta/C [\text{s}^{-1}] \\ &= h e C^{-1} B \left[B^2 [\text{T}^2] + \left(\frac{E_0 [\text{MeV}]}{e \rho [\text{km}^2/\text{s}]} \right)^2 \right]^{-1/2} \end{aligned}$$

(see Table I.1.1 for other expressions).

See Section I.4.2.4, page 9, for differential relations.

3. SYNCHROTRON OSCILLATIONS

3.1 Equation of motion

(Above transition ϕ should be replaced by $\phi + \pi - 2\phi_s$; for transition see Section I.4.2.3, page 8.)

$$d/dt[(\beta_s^2 E_s / \eta_s \Omega_s^2) \dot{\phi}] = (h e V / 2\pi)(\sin \phi - \sin \phi_s)$$

where E_s is in keV, $\eta = \gamma_{tr}^{-2} - \gamma^{-2}$, and Ω is the particle angular velocity on the orbit of length C . To transform to other co-ordinates one has in the absence of perturbations

$$\begin{aligned} \dot{\phi} &= h (\Omega_s - \Omega) \\ &= (h \Omega_s \eta_s \gamma_{tr}^{-2} / R) \Delta R \\ &= (h \Omega_s \eta_s / \beta_s \gamma_s) (\Delta p / m_0 c) \\ &= (h \Omega_s \eta_s / \beta_s^2 E_s) \Delta E \end{aligned}$$

[14] 3.2 Bucket size

3.2.1 Without space charge effects

a) Bucket area, (half) height and coordinates

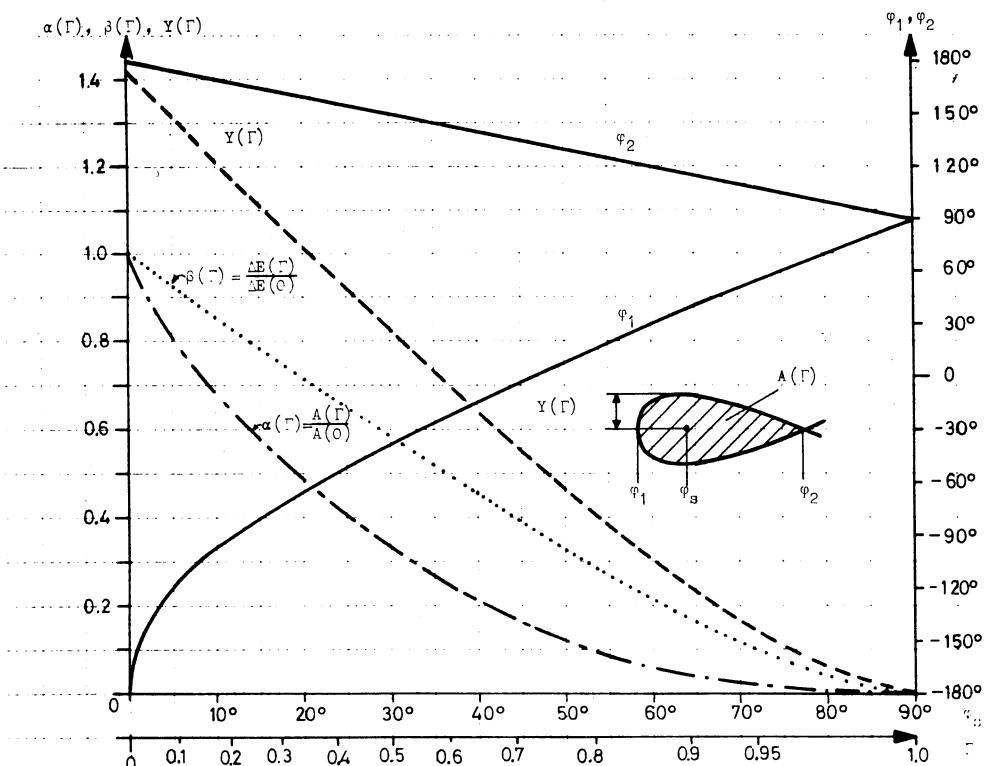
Bucket area	Bucket (half) height	Coordinates *
$(\text{heV})^{\frac{1}{2}} \alpha(\Gamma) (16\gamma/h) (2\pi E \eta)^{-\frac{1}{2}}$	$(\text{heV})^{\frac{1}{2}} Y(\Gamma)(\gamma/h) (\pi E \eta)^{-\frac{1}{2}}$	$(\Delta p/m_c) - \varphi$
$(\text{heV})^{\frac{1}{2}} \alpha(\Gamma) (16\beta/h) [E/(2\pi \eta)]^{\frac{1}{2}}$	$(\text{heV})^{\frac{1}{2}} Y(\Gamma)(\beta/h)[E/(\pi \eta)]^{\frac{1}{2}}$	$(\Delta E) - \varphi$
$(\text{heV})^{\frac{1}{2}} \alpha(\Gamma) [16\alpha_p R/(h\beta)] (2\pi \eta)^{-\frac{1}{2}}$	$(\text{heV})^{\frac{1}{2}} Y(\Gamma)[R/(\gamma_{tr}^2 h\beta)] (\pi E \eta)^{-\frac{1}{2}}$	$(\Delta R) - \varphi$
$(\text{heV})^{\frac{1}{2}} \alpha(\Gamma) [16\beta/(h^2\Omega)] [E/(2\pi \eta)]^{\frac{1}{2}}$	$(\text{heV})^{\frac{1}{2}} Y(\Gamma)[\beta/(h^2\Omega)][E/(\pi \eta)]^{\frac{1}{2}}$	$(\Delta E/h\Omega) - \varphi$

For $\alpha(\Gamma)$ see below and Appendix C; for η see Section 3.1 on preceding page.

$$Y = Y(\Gamma) = \dot{\varphi}_{\max}/(\sqrt{2} 2\pi v_0)_{\varphi_s=0} = \dot{\varphi}_{\max} (\text{heV})^{-\frac{1}{2}} (\beta/\Omega) (\pi E|\eta|)^{\frac{1}{2}} * [\text{E and eV in keV} \\ \varphi \text{ in rad, } \Delta R \text{ in cm}]$$

Ideal adiabatic trapping of a linac beam with $\pm \Delta E_L$ leads to a minimum bucket (half) height $\Delta E = (\pi/2)\Delta E_L$

b) Bucket width, normalised (half) height and area (see Appendix C for numbers)



3.2.2 Reduction of bucket area due to space charge effects (below transition)

This reduction can be obtained from Fig. III.3.2.2, where

$$\Delta A_{\text{sp.c.}} = 4\pi h g_c E_0 r_p N / (\text{ReV} \gamma^2)$$

with

[15,3]

N = number of accelerated particles

$g_c = 1 + 2 \ln (\text{vacuum chamber diameter}/\text{beam diameter})$

r_p = classical proton radius

and E_0 and eV are in the same units (as are r_p and R).

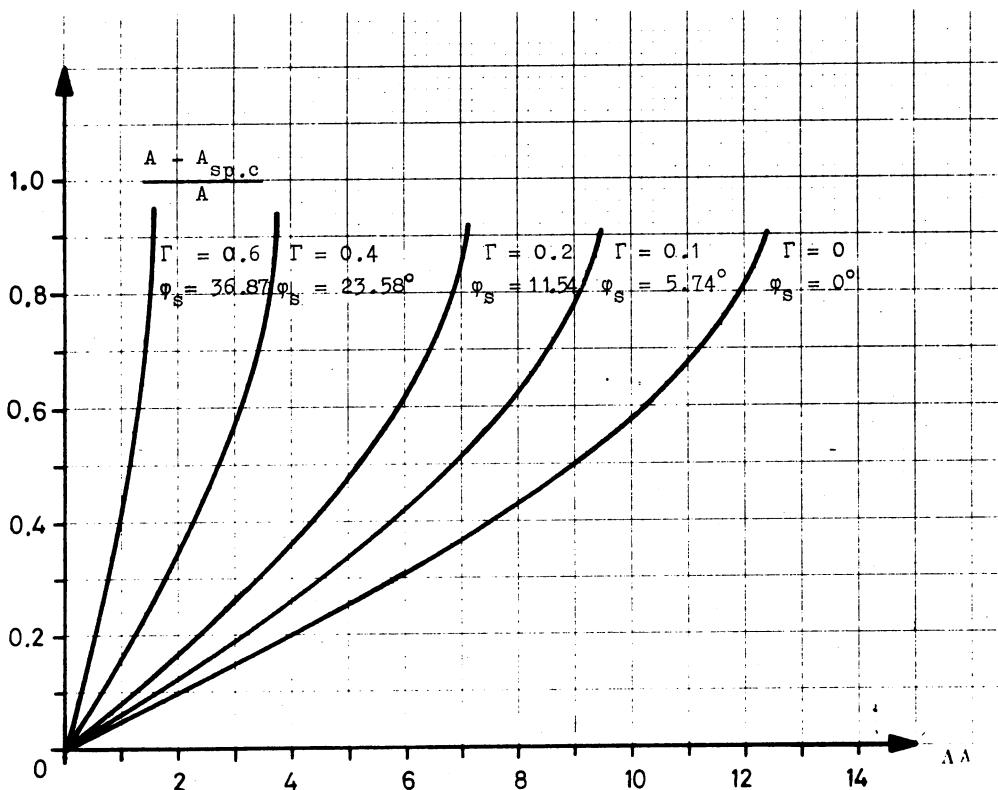


Fig. III.3.2.2 $(A - A_{\text{sp.c.}})/A = f(\Delta A_{\text{sp.c.}})$ (for constant density in phase space)

For $\varphi_s = 0^\circ$ (and a \cos^2 distribution in real space) one has

[18,
Appendix IV]

$$A_{\text{sp.c.}}/A = [1 - g_c e h N / (4\pi \epsilon_0 \gamma^2 R V)]^{1/2}$$

where V is in volts.

3.3 Frequency of synchrotron oscillations

In the case of small phase oscillation amplitudes around the stable phase φ_s , the equation of motion is

$$[17,10]* \quad (\ddot{\Delta\varphi}) = (\delta^2 - 1)\delta^{-3} C' [\sin(\varphi_s + \Delta\varphi) - \sin \varphi_s] = (\delta^2 - 1)\delta^{-3} C' \cos \varphi_s \Delta\varphi$$

where $\delta^2 = \eta \gamma^2 + 1 = \alpha_p \gamma^2$, $C' = 2\pi f_\infty^2 \gamma_{tr}^{-3} eV/(hE_0)$ and eV and E₀ are in the same units. (See Section I.4.2.3, page 8, for transition.)

The frequency of these oscillations is [eV and E_s in same units]

$$[17,11] \quad \nu_0 = [(1 - \delta^2)\delta^{-3} C' \cos \varphi_s]^{1/2}/(2\pi) = [f_\infty^2 |\eta| eV \cos \varphi_s / (2\pi E_s h)]^{1/2}.$$

In terms of the bucket area A [in $\Delta p/(mc)$ - φ coordinates]:

$$\nu_0 = [A E_0 f_\infty |\eta| / (16 E_s \alpha(\Gamma))] (|\cos \varphi_s|)^{1/2}$$

or

$$\nu_0 = [A eV |\eta| / (32\pi R_s \gamma_s \alpha(\Gamma))] (|\cos \varphi_s|)^{1/2}.$$

Alternatively

$$[16,4] \quad \nu_0 = \{[\cos \varphi_s / (4\pi^2)][c^2 / (2\pi R^2 E_0)](heV)(|1 - \gamma_{tr}^{-2} \gamma^2|) / \gamma^3\}^{1/2}.$$

3.4 Adiabatic damping of small-amplitude oscillations

$$\gamma(t) = \{1 + [B(t)/B(t_0)]^2\}^{1/2}.$$

3.4.1 Phase amplitude [eV and E_s in same units]

$$[16,5] \quad \Delta\varphi(t) = D [eV(t) \cos \varphi_s]^{-1/4} [|1 - \gamma_{tr}^{-2} \gamma^2(t)| / \gamma^3(t)]^{1/4}$$

where

$$D = \text{constant} = \Delta\varphi_i [2\pi R^2 E_0 / (h c^2)]^{1/4}$$

with i denoting the initial values.

$$\Delta\varphi(t) / \Delta\varphi_i = [V_i/V(t)]^{1/4} [|1 - \gamma_{tr}^{-2} \gamma^2(t) / \gamma^3(t)|]^{1/4} (|1 - \gamma_{tr}^{-2} \gamma_i^2| / \gamma_i^3)^{-1/4}.$$

*) See page 43 for references

3.4.2 Energy amplitude

$$\Delta E(t) = G [eV(t) \cos \varphi_s]^{1/4} \beta(t) [\gamma^3(t) / |1 - \gamma_{tr}^{-2} \gamma^2(t)|]^{1/4}$$

where

$$G = \text{constant} = \Delta\varphi_i [E_0 R / (hc)] [h c^2 / (2\pi R^2 E_0)]^{1/4}$$

$$\Delta E(t) / \Delta E_i = [V(t)/V_i]^{1/4} [\beta(t)/\beta_i] [|1 - \gamma_{tr}^{-2} \gamma^2(t)|/\gamma^3(t)]^{-1/4} (|1 - \gamma_{tr}^{-2} \gamma^2|/\gamma_i^3)^{1/4}.$$

3.4.3 Radial amplitude

$$\Delta R(t) = H [eV(t) \cos \varphi_s]^{1/4} [\beta(t) \gamma(t)]^{-1} [\gamma^3(t) / |1 - \gamma_{tr}^{-2} \gamma^2(t)|]^{1/4}$$

where

$$H = \Delta\varphi_i [\gamma_{tr}^{-2} R^2 / (hc)] [h c^2 / (2\pi R^2 E_0)]^{1/4}$$

$$\Delta R / \Delta R_i = [\beta_i \gamma_i / \beta(t) \gamma(t)] [\Delta\varphi_i / \Delta\varphi(t)].$$

4. DEBUNCHING

4.1 Debunching time

In the absence of RF fields the beam "debunches" *) itself in the synchrotron (i.e. front end of one bunch reaches the tail of the next bunch ahead) after a time **)

$$t_{db} \approx (\pi - \Delta\varphi) [2\pi f_a |\gamma^{-2} - \gamma_{tr}^{-2}| \Delta p/p]^{-1}$$

where $2\Delta\varphi$ and $2\Delta p$ are the total phase and momentum spreads.

Special cases:

a) Low energy

$$\gamma^2 \ll \gamma_{tr}^2 \quad \Delta\varphi \ll \pi \text{ (strong damping in linac)}$$

$$\begin{aligned} t_{db} &= \gamma^2 (2f_a \Delta p/p)^{-1} \\ &= \gamma(\gamma+1) (2f_a \Delta T/T)^{-1}. \end{aligned}$$

*) This azimuthal spreading does not involve any reduction of $\Delta n/p$.

**) Strictly valid for "rectangular" bunches. For "oval" bunches more time may be required in practice.

b) Well above transition

$$\gamma^2 \gg \gamma_{tr}^2 \quad \Delta\phi \ll \pi \text{ (damping in synchrotron)}$$

$$t_{db} = \gamma_{tr}^2 (2f_a \Delta p/p)^{-1}$$

Complete overlapping (front end reaches centre of next bunch ahead) requires $2t_{db}$.

4.2 Travelling distance required for debunching

$$S_{db} = c \beta t_{db} .$$

For $\gamma^2 \ll \gamma_{tr}^2$ and $\Delta\phi \ll \pi$, this becomes with $2D$ = distance between centre of bunches

$$S_{db} = D/(\Delta\beta/\beta).$$

See Table I.1.1 for other expressions.

ADDITIONAL FORMULAE

P A R T IV

ROUGH EVALUATION OF MAJOR ACCELERATOR SYSTEMS

1. MAGNETS (non-saturated)

1.1 Bending magnet

a) Excitation current

$$N_B I [\text{ampere-turns}] = B h_B / \mu_0 [\text{mT}/(\text{H m}^{-1})]$$

where h_B is the (mean) gap height

$$N_B I / (B h_B) \approx 800 \text{ ampere-turns} / (0.1\text{T} \times 0.01 \text{ m gap height}).$$

b) Inductance

$$L_B [\text{H}] \approx N_B^2 \mu_0 w \ell_B / h_B$$

$$w = w_a + \frac{2}{3} w_c \quad (\text{for window frame magnet})$$

$$w = w_p + \frac{1}{2} h_B \quad (\text{for magnet with poles})$$

where

w_a = aperture between coils, w_c = coil width

w_p = pole width

ℓ_B = the total magnetic length

c) Stored energy

$$W_B [\text{Ws}] \approx B^2 h_B w \ell_B / (2\mu_0) \quad [\text{T}^2 \text{ m}^3 / (\text{H m}^{-1})]$$

where w is as in b).

1.2 Quadrupole lens

a) Excitation current per pole

$$N_Q I [\text{ampere-turns}] = g r_Q^2 / (2\mu_0) \quad [\text{Tm}/\text{H m}^{-1}]$$

$$N_Q I / (g r_Q^2) \approx 400 \text{ ampere-turns} / [10 \text{ T/m} \cdot (0.01 \text{ m bore radius})^2]$$

b) Inductance

$$L_Q [\text{H}] \approx 8\mu_0 N_Q^2 y_{\max} (y_{\max} + 2/3 w_c) \ell_Q / r_Q^2$$

where y_{\max} is the distance from the lens centre to the coil face
and ℓ_Q is the total magnetic length.

c) Stored energy

$$W_Q [Ws] \approx g^2 r_Q^2 y_{max} (y_{max} + 2/3 w_c) \ell_Q / \mu_0 .$$

1.3 Bending magnet and quadrupole lens excited in series

$$B \approx N_B r_Q^2 \beta \gamma K / (0.639 h_B N_Q) , \text{ (for protons)}$$

$$\text{or } K \approx (0.6/p)(N_Q/N_B)(h_B/r_Q^2) B$$

where r_Q and h_B in m.

See Section II.1.4.d) for other expressions.

1.4 Cooling water requirements

To cool a conductor heated by a power loss $N[kW]$, one needs a water flow of

$$G_w [\ell/s] \approx N/(4.2 \Delta t) = 10^{-3} v A_F = 10^{-3} v F_s d_h^2 ,$$

where

- $\Delta t_w [^\circ C]$ is the allowed temperature increase
- $v[m/s]$ is the velocity of the cooling water
- $A_F [mm^2]$ is the flow area
- $F_s = A_F/d_h^2$ is the shape factor ($= \pi/4$ for round holes)
- $d_h [mm] = 4A_F/\text{perimeter}$ is the hydraulic diameter.

For turbulent flow the required pressure drop may be obtained from

$$\Delta P_w [kg/cm^2] = 0.18 L_c v^{1.75} / (F_s^{1.75} d_h^{1.25})$$

where $L_c [m] =$ length of conductor, and it is noted that $0.18(\pi/4)^{-1.75} = 0.28$.

If d_h is in metres, this becomes for round holes

$$[4,67]* \quad \Delta P_w [kg/cm^2] = 5 \cdot 10^{-5} L v^{1.75} / d_h^{1.25} .$$

Alternatively, one has (with somewhat more pessimistic assumptions about the pressure loss)

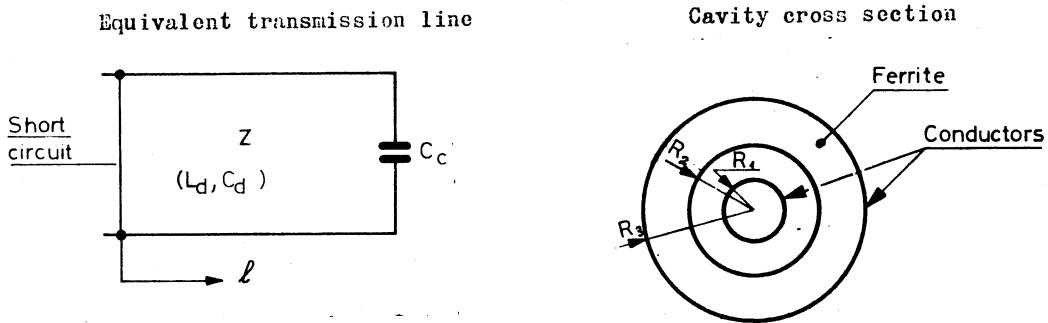
$$G_w [\ell/s] = 1.25 \cdot 10^{-3} (1 + 0.009 t_w) F_s d_h^{2.71} (\Delta P_w / L_c)^{0.57}$$

where $t_w [^\circ C]$ is the water temperature.

*) See page 43 for references

2. FERRITE-LOADED RF ACCELERATING CAVITY

2.1 Basic quantities



a) Effective permittivity and permeability

$$\epsilon_{\text{eff}} = \epsilon / [x + (1-x)\epsilon]$$

$$\mu_{\text{eff}} = 1 + x(\mu - 1) \approx \mu x$$

where $x = \ln(R_3/R_2)/\ln(R_3/R_1)$.

b) Inductance of ferrite cylinder of length ℓ_c

$$L_d = [\mu_{\text{eff}} \mu_0 \ell_c \ln(R_3/R_1)] / 2\pi$$

c) Capacitance

$$C_d = 2\pi \epsilon_{\text{eff}} \epsilon_0 \ell_c / \ln(R_3/R_1)$$

d) Characteristic impedance

$$Z = (L_d/C_d)^{1/2} = 60(\mu_{\text{eff}}/\epsilon_{\text{eff}})^{1/2} \ln(R_3/R_1)$$

e) Wavelength

$$\lambda = v/f_a = c/[f_a (\mu_{\text{eff}} \epsilon_{\text{eff}})^{1/2}] = \ell_c / [f_a (L_d C_d)^{1/2}]$$

2.2 Length of cavity

For resonance (assuming negligible losses)

$$1/(w_0 C_c) = Z \tan(2\pi \ell_c / \lambda) \quad \text{i.e.}$$

$$\tan(2\pi \ell_c / \lambda) = (w_0 C_c Z)^{-1},$$

which becomes for small arguments with $\omega_0 \approx (L_d C_e)^{-1/2}$

$$\ell_e \approx \lambda (C_d/C_e)^{1/2} / (2\pi).$$

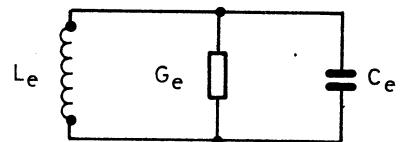
2.3 Equivalent resonant circuit

If one knows the complex cavity admittance

$$Y = G + jB = (G^2 + B^2)^{1/2} e^{j\Phi} \text{ as a function of } \omega,$$

one can find C_e from

$$C_e = (G_e \tan \Phi) / (2\Delta\omega)$$



where the phase angle Φ pertains to the frequency

$\omega = \omega_0 \pm \Delta\omega$, and hence

$$L_e = (\omega_0^2 C_e)^{-1}.$$

The relation between the equivalent and the "real" quantities is as follows (on the basis of $B = B_e = 0$; $dB/d\omega = dB_e/d\omega$ for $\omega = \omega_0$)

$$C_e = 0.5 C_c + 0.5 C_d [\sin^2(2\pi \ell_e/\lambda)]^{-1}$$

$$L_e = 2L_d \left((2\pi \ell_e/\lambda)^2 \left\{ (C_e/C_d) + [\sin^2(2\pi \ell_e/\lambda)]^{-1} \right\} \right)^{-1}.$$

2.4 Longitudinal variation of power loss

$$P(\ell) = P_{\max} \cos^2(2\pi \ell/\lambda)$$

$$\bar{P} = P_{\max} \{ 0.5 + [0.25 \sin(4\pi \ell_e/\lambda)] / (2\pi \ell_e/\lambda) \}$$

where P_{\max} is the maximum loss (occurring at the short circuit) and the power loss per unit volume is assumed to be constant.

3. VACUUM PRESSURE REQUIRED

The natural growth of the beam emittance in either transverse plane is given by

$$[19,8]* \quad \Delta(\varepsilon\beta\gamma) [\text{rad m}] = 0.32 \pi P |\ln(1 - \eta)| \int_{t_0}^{t_1} \beta^{-2} \gamma^{-1} dt [\text{m Torr s}]$$

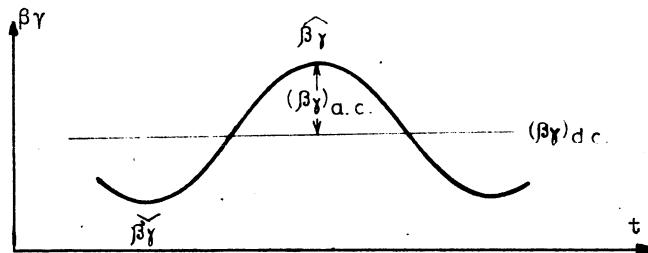
where $\pi = R/Q$, P = nitrogen equivalent pressure and η is the fraction of the particles contained in the emittance

η	0.5	0.8	0.9	0.95	0.97
$ \ln(1 - \eta) $	0.694	1.61	2.3	3.0	3.5

a) For $(\beta\gamma) = \text{const}$, one has

$$\int_{t_0}^{t_1} \beta^{-2} \gamma^{-1} dt = (\beta\gamma)^{-1} \left(\frac{1}{\beta(t_0)} - \frac{1}{\beta(t_1)} \right) + \ln \left\{ \gamma(t_1)[1 + \beta(t_1)] / \gamma(t_0)[1 + \beta(t_0)] \right\} .$$

b) In the case of sinusoidal excitation of the magnet field



$$\begin{aligned} \beta(t)\gamma(t) &= (\beta\gamma)_{\text{d.c.}} - (\beta\gamma)_{\text{a.c.}} \cos \omega_m t \\ &= 0.5 \left\{ (\hat{\beta}\gamma) + (\check{\beta}\gamma) - [(\hat{\beta}\gamma) - (\check{\beta}\gamma)] \cos \omega_m t \right\} \end{aligned}$$

one has to good approximation (for $\hat{\beta} \approx 1$ and $\check{\beta} \ll 1$)

$$\int_{t_0}^{t_1} \beta^{-2} \gamma^{-1} dt \approx \pi (1/\hat{\beta} + 1/\check{\beta}) / \{ 2\omega_m [(\hat{\beta}\gamma)(\check{\beta}\gamma)]^{1/2} \} .$$

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APPENDIX A

T A B L E O F C O N S T A N T S

Symbols	Meaning	Value	Units
c	velocity of light	2.997925×10^8	ms^{-1}
e	electronic charge	1.6021×10^{-19}	C
e/m_e	charge to mass ratio for an electron	1.75880×10^8	Cg^{-1}
e/m_p	charge to mass ratio for a proton	9.57896×10^4	Cg^{-1}
h	Planck's constant	6.6256×10^{-34}	Js
\hbar	Planck's constant/ 2π	1.0545×10^{-34}	Js
h/e	quantum charge ratio	4.1355×10^{-15}	Js C^{-1}
k	Boltzmann's constant	(1.3805×10^{-23} 8.6171×10^{-14})	(J/ $^{\circ}\text{K}$ GeV/ $^{\circ}\text{K}$)
m_d	rest mass of deuteron	(3.3433×10^{-27} 1.87558)	(kg GeV/c ²)
m_e	rest mass of electron	(9.1091×10^{-31} 5.11006×10^{-4})	(kg GeV/c ²)
m_p	rest mass of proton	(1.6725×10^{-27} 0.93826)	(kg GeV/c ²)
m_p/m_e	ratio of proton and electron masses	1.83610×10^3	
r_e	classical electron radius	2.8178×10^{-15}	m
r_d	classical deuteron radius	0.76774×10^{-18}	m
r_p	classical proton radius	1.5347×10^{-18}	m
μ_0	permeability of free space	($= 4\pi \times 10^{-7}$ $= 1.25664 \times 10^{-6}$)	Hm ⁻¹
ϵ_0	permittivity of free space	($= (\mu_0 c^2)^{-1}$ $= 8.8542 \times 10^{-12}$)	Fm ⁻¹

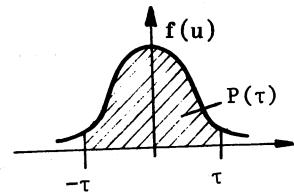
These values are taken from the 48th edition of the CRC Handbook of Chemistry and Physics. (Still within the limits of errors of the values given in the 50th edition.)

APPENDIX B

GAUSSIAN DISTRIBUTIONS

Definitions: $f(u) = (2\pi)^{-1/2} \exp(-u^2/2)$; $\overline{u^2} = 1$

$$P(\tau) = \int_{-\infty}^{\tau} f(u) du$$



One-dimensional density distribution:

$$\rho(x) = (2\pi)^{-1/2} \sigma^{-1} \exp[-x^2/(2\sigma^2)] = \sigma^{-1} f(x/\sigma)$$

$$\text{Normalization: } \int_{-\infty}^{\infty} \rho dx = 1; \int_{-\infty}^{x} \rho(x) dx = P(x/\sigma)$$

$$\text{Variance: } \overline{x^2} = \int_{-\infty}^{\infty} x^2 \rho dx = \sigma^2$$

$$\text{Standard deviation: } \sqrt{\overline{x^2}} = \langle x \rangle = \sigma$$

Two-dimensional density distribution (x, z uncorrelated):

$$\rho(x, z) = (2\pi\sigma_1\sigma_2)^{-1} \exp[-x^2/(2\sigma_1^2) - z^2/(2\sigma_2^2)] = (\sigma_1\sigma_2)^{-1} f(x/\sigma_1) f(z/\sigma_2)$$

Normalization: over the whole plane $\iint \rho dxdz = 1$

Variances: $\overline{x^2} = \sigma_1^2$ for any z or for all z ; $\overline{z^2} = \sigma_2^2$ similarly

The ellipses $x^2/\sigma_1^2 + z^2/\sigma_2^2 = v^2$ are lines of constant ρ with

$$\rho = (2\pi\sigma_1\sigma_2)^{-1} \exp(-v^2/2); \quad \overline{v^2} = 2.$$

Over this ellipse $\iint \rho dxdz = 1 - \exp(-v^2/2)$

u, τ, v	$f(u)$	$P(\tau)$	$1 - \exp(-v^2/2)$	$v/\sqrt{2}$	$v^2/2$
0	0.399	0	0	0	0
0.5	0.352	0.383	0.118	0.354	0.125
0.674	0.318	0.500	0.203	0.477	0.227
0.707	0.311	0.520	0.221	0.500	0.250
1	0.242	0.683	0.393	0.707	0.500
1.414	0.147	0.843	0.632	1.000	1.000
1.5	0.130	0.866	0.675	1.061	1.125
2	0.0540	0.954	0.865	1.414	2.000
2.5	0.0175	0.9876	0.9561	1.768	3.125
2.797	0.0080	0.9948	0.980	1.978	3.912
3	0.0044	0.9973	0.9889	2.121	4.500
∞	0	1.000	1.000	∞	∞

$E_0 = 938.26 \text{ MeV}$

T [MeV]	c_p [MeV]	βp [fm]	β	β^2	γ	γ^2	$\beta\gamma$	$\beta^2\gamma$	$\beta\gamma^2$	$\beta^2\gamma^3$
20.00	194.7573	6.4964E-01	2.0324E-01	4.1307E-02	1.0213E+00	1.0431E+00	2.0757E-01	4.2187F-02	2.1200E-01	4.4005E-02
50.00	310.3643	1.0353E+00	3.1405E-01	9.8628E-02	1.0533E+00	1.1094E+00	3.3079E-01	1.0388E-01	3.4844E-01	1.1525E-01
200.00	644.4408	2.1496E+00	5.6616E-01	3.2054E-01	1.2132E+00	1.4718E+00	6.8685E-01	3.8887E-01	8.3326E-01	5.7232E-01
400.00	954.2573	3.1831E+00	7.1306E-01	5.0845E-01	1.4263E+00	2.0344E+00	1.0171E+00	7.2522E-01	1.4506E+00	1.4754E+00
500.00	1090.3734	3.6361E+00	7.5791E-01	5.7443E-01	1.5329E+00	2.3498E+00	1.1618E+00	8.8054E-01	1.7809E+00	2.0691E+00
600.00	1218.9799	4.0661E+00	7.9244E-01	6.2796E-01	1.6395E+00	2.6879E+00	1.2992E+00	1.0295E+00	2.1306E+00	2.7673E+00
700.00	1342.9684	4.4797E+00	8.1975E-01	6.7199E-01	1.7461E+00	3.0487E+00	1.4313E+00	1.1733E+00	2.4992E+00	3.5772E+00
750.00	1403.5277	4.9817E+00	8.3135E-01	6.9114E-01	1.7994E+00	3.2377E+00	1.4959E+00	1.2436E+00	2.6916E+00	4.0264E+00
800.00	1403.2894	4.8810E+00	8.4181E-01	7.0865E-01	1.8526E+00	3.4323E+00	1.5596E+00	1.3129E+00	2.8893E+00	4.5061E+00
850.00	1522.3475	5.0780E+00	8.5130E-01	7.2471E-01	1.9059E+00	3.6326E+00	1.6225E+00	1.3813E+00	3.0924E+00	5.0175E+00
900.00	1580.7808	5.2729E+00	8.5993E-01	7.3949E-01	1.9592E+00	3.8386E+00	1.6848E+00	1.4488E+00	3.3009E+00	5.5614E+00
950.00	1638.6562	5.4660E+00	8.6781E-01	7.5310E-01	2.0125E+00	4.0502E+00	1.7465E+00	1.5156E+00	3.5148E+00	6.1386E+00
		[GeV]								
1.00	1.6960	5.6574E+00	8.7503E-01	7.6567E-01	2.0658E+00	4.2675E+00	1.8076E+00	1.5817E+00	3.7342E+00	6.7501E+00
3.00	3.2249	1.2758E+01	9.7121E-01	9.4324E-01	4.1974E+00	1.7618E+01	4.0765E+00	3.9592E+00	1.7111E+01	6.9754E+01
6.00	6.8745	2.2931E+01	9.9081E-01	9.8171E-01	7.3948E+00	5.46683E+01	7.3269E+00	7.2596E+00	5.4181E+01	3.9698E+02
8.00	E.8889	2.9650E+01	9.9448E-01	9.8898E-01	9.5264E+00	9.0753E+01	9.4738E+00	9.4215E+00	9.0251E+01	8.5502L+02
10.00	1C.8979	3.6352E+01	9.9631E-01	9.9264E-01	1.1658E+01	1.3591E-02	1.1615E+01	1.1572E+01	1.3541E+02	1.5728E+03
12.00	12.9042	4.3044E+01	9.9737E-01	9.9474E-01	1.3790E+01	1.9015E+02	1.3753E+01	1.3717E+01	1.8965E+02	2.6084E+03
18.00	15.9150	6.3094E+01	9.9877E-01	9.9755E-01	2.0184E+01	4.0741E+02	2.0160E+01	2.0135E+01	4.0691E+02	8.2032E+03
21.00	21.9182	7.3111E+01	9.9909E-01	9.9817E-01	2.3332E+01	5.4671L+02	2.3360E+01	2.3339L+01	5.4621E+02	1.2760E+04
24.00	24.9206	8.3126E+01	9.9929E-01	9.9858E-01	2.6579E+01	7.6646E+02	2.6560E+01	2.6542E+01	7.0596E+02	1.8751E+04
27.00	27.9225	9.3139E+01	9.9944E-01	9.9867E-01	2.9777E+01	6.8665E+02	2.9760E+01	2.9743E+01	6.8615E+02	2.6372E+04
33.00	33.9253	1.1316E+02	9.9962E-01	9.9924E-01	3.6171L+01	1.3084E+03	3.6158E+01	3.6144E+01	1.3079E+03	4.7290E+04
76.00	76.9325	2.5662E+02	9.9993E-01	9.9985E-01	8.2001E+01	6.7242E+03	8.1995E+01	8.1989E+01	6.7237E+03	5.5131E+05
200.00	400.00	6.7025E+02	9.9999E-01	9.9998E-01	2.1416E+02	4.5865E+04	2.1416E+02	2.1416E+02	4.5864E+04	9.8222E+06
300.00	300.00	1.0038E+03	1.00000E+00	9.9999E-01	3.2074L+02	1.0287L+05	3.2074E+02	1.0287E+02	3.2996E+07	3.2996E+07
400.00	400.00	1.3374E+03	1.00000E+00	9.9999E-01	4.2732E+02	1.8260E+05	4.2732E+02	1.8260E+05	7.8030L+07	7.8030L+07

*) See page 48 for values useful for PSB and CPS operation.

RF HBUCKETS WIDTH, NORMALISED (HALF) HEIGHT AND AREA

(From Ref. 14; note that in this reference $Y(0) = \varphi_{\max}(0)[\sqrt{2} 2\pi v_0(0)] = \sqrt{2}$
 - rather than 2 - for computational convenience.
 See page 31 for other definitions and a graphical representation.)

All values of φ_s , φ_1 , φ_2 , $\Delta\varphi$ in degrees

Stable phase	Γ	'Bucket' width			Half height	Height	Area
		φ_1	φ_2	$\Delta\varphi$			
φ_s	Γ	φ_1	φ_2	$\Delta\varphi$	$Y(\Gamma)$	$\beta = \frac{\Delta E(\Gamma)}{\Delta E(0)}$	$\alpha = \frac{A(\Gamma)}{A(0)}$
0	0.000000	-180.0	180	360	1.414214	1.000000	1.000000
1	• 0.17452	-154.0	179	333	1.394803	• 986275	• 954105
2	• 0.34859	-143.5	178	321	1.375347	• 972517	• 917558
3	• 0.52336	-135.5	177	312	1.355847	• 958129	• 884511
4	• 0.69756	-128.8	176	305	1.336309	• 944913	• 853747
5	• 0.87156	-122.9	175	298	1.316736	• 931073	• 824676
6	• 1.04528	-117.6	174	292	1.297132	• 917211	• 796983
7	• 1.21869	-112.6	173	286	1.277500	• 903329	• 770443
8	• 1.39173	-108.1	172	280	1.257846	• 889431	• 744906
9	• 1.56434	-103.7	171	275	1.238171	• 875519	• 720257
10	• 1.73648	-99.6	170	270	1.218482	• 861597	• 696413
11	• 1.908C9	-95.7	169	265	1.198781	• 847666	• 673303
12	• 2.07912	-92.0	168	260	1.179072	• 833730	• 650875
13	• 2.24951	-88.4	167	255	1.159360	• 819791	• 629092
14	• 2.41922	-84.9	166	251	1.139648	• 805653	• 607888
15	• 2.58819	-81.5	165	247	1.119940	• 791917	• 587261
16	• 2.75637	-78.2	164	242	1.100240	• 771987	• 567174
17	• 2.92372	-75.0	163	238	1.080552	• 760666	• 547603
18	• 3.09017	-71.9	162	234	1.060881	• 750156	• 528529
19	• 3.25568	-68.9	161	230	1.041230	• 736261	• 509933
20	• 3.42020	-65.9	160	226	1.021603	• 722382	• 491799
21	• 3.58368	-63.0	159	222	1.002004	• 708524	• 474114
22	• 3.74607	-60.1	158	218	• 694688	• 456865	• 456865
23	• 3.90731	-57.3	157	214	• 662907	• 680878	• 440040
24	• 4.06737	-54.5	156	210	• 943418	• 667097	• 423630
25	• 4.22618	-51.0	155	207	• 923972	• 653347	• 407624
26	• 4.38371	-49.1	154	203	• 904576	• 639632	• 392016
27	• 4.53990	-46.4	153	199	• 885232	• 625954	• 376714
28	• 4.69472	-43.8	152	196	• 865945	• 612316	• 361955
29	• 4.84810	-41.2	151	192	• 846719	• 598721	• 347489
30	• 500000	-38.7	150	189	• 827559	• 585172	• 333392
31	• 515038	-36.2	149	185	• 808467	• 571673	• 319673
32	• 529919	-33.7	148	182	• 789440	• 558225	• 306279
33	• 544629	-31.2	147	178	• 770510	• 544833	• 293232
34	• 559193	-28.8	146	175	• 751653	• 531499	• 280571
35	• 573576	-26.3	145	171	• 732882	• 518226	• 268630
36	• 587785	-23.9	144	168	• 714202	• 505021	• 256229
37	• 601815	-21.6	143	165	• 695618	• 491876	• 244560
38	• 615661	-19.2	142	161	• 677132	• 478805	• 233218
39	• 629326	-16.9	141	158	• 658751	• 465807	• 222202

Stable phase	Γ	'Bucket' width			Half height	Height	Area
		φ_1	φ_2	$\Delta\varphi$			
φ_s	Γ	φ_1	φ_2	$\Delta\varphi$	$Y(\Gamma)$	$\beta = \frac{\Delta E(\Gamma)}{\Delta E(0)}$	$\alpha = \frac{A(\Gamma)}{A(0)}$
4.0	• 642788	-14.6	140	155	• 640479	• 452887	• 211505
4.1	• 650559	-12.3	139	151	• 622319	• 440046	• 201125
4.2	• 669331	-10.0	138	148	• 604277	• 427288	• 191058
4.3	• 686198	-7.7	137	145	• 586357	• 414617	• 181304
4.4	• 697658	-5.4	136	141	• 568564	• 402035	• 171848
4.5	• 7011C7	-3.2	135	138	• 550902	• 389546	• 162698
4.6	• 719340	-1.0	134	135	• 533376	• 377154	• 153845
4.7	• 711354	1.2	133	132	• 515991	• 364861	• 145288
4.8	• 743145	3.5	132	129	• 498752	• 352671	• 137022
4.9	• 754710	5.6	131	125	• 481664	• 340588	• 129044

Stable phase		'Bucket' width				Half height		Area		Stable phase		'Bucket' width				Half height		Area					
Γ	Φ_S	Φ_1	Φ_2	$\Delta\varphi$	$\Upsilon(\Gamma)$	$\beta = \frac{\Delta E(\Gamma)}{\Delta E(0)}$	$\alpha = \frac{A(\Gamma)}{A(0)}$	Γ	Φ_S	Φ_1	Φ_2	$\Delta\varphi$	$\Upsilon(\Gamma)$	$\beta = \frac{\Delta E(\Gamma)}{\Delta E(0)}$	$\alpha = \frac{A(\Gamma)}{A(0)}$	Γ	Φ_S	Φ_1	Φ_2	$\Delta\varphi$	$\Upsilon(\Gamma)$	$\beta = \frac{\Delta E(\Gamma)}{\Delta E(0)}$	$\alpha = \frac{A(\Gamma)}{A(0)}$
0.000	0.000	-180.0	180.0	360.0	1.414214	1.000000	1.000000	500	30.0000	-38.7	150.0	188.7	-827559	.585172	.333392	500	30.664	-37.0	149.3	186.3	-827559	.585172	.333392
-0.010	.573	-160.2	179.4	339.6	1.403098	1.000000	1.000000	510	30.664	-37.0	149.3	186.3	-827559	.585172	.333392	510	31.332	-35.3	148.7	184.0	-802140	.567199	.315173
-0.020	1.146	-152.2	178.9	331.0	1.391966	1.000000	1.000000	520	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	520	32.684	-32.0	147.3	179.3	-776693	.558152	.324234
-0.030	1.719	-146.1	178.3	324.4	1.380816	1.000000	1.000000	530	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	530	33.637	-30.3	147.6	176.6	-763280	.558152	.324234
-0.040	2.292	-141.0	177.7	318.7	1.369816	1.000000	1.000000	540	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	540	34.056	-28.6	145.9	174.6	-750603	.539932	.297335
-0.050	2.866	-136.5	177.1	313.6	1.358463	1.000000	1.000000	550	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	550	34.570	-26.9	145.2	172.2	-750757	.539932	.297335
-0.060	3.440	-132.4	176.6	309.0	1.347259	1.000000	1.000000	560	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	560	35.470	-25.3	144.5	172.2	-750757	.539932	.297335
-0.070	4.014	-128.7	176.0	304.7	1.336035	1.000000	1.000000	570	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	570	36.157	-23.6	143.8	167.4	-711278	.512267	.271282
-0.080	4.589	-125.2	175.4	300.7	1.324793	1.000000	1.000000	580	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	580	36.157	-23.6	143.8	167.4	-711278	.512267	.271282
-0.090	5.164	-122.0	174.8	296.8	1.313520	1.000000	1.000000	590	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	590	36.157	-23.6	143.8	167.4	-711278	.512267	.271282
-0.100	5.739	-118.9	174.3	293.2	1.302248	1.000000	1.000000	600	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	600	36.870	-21.9	143.1	165.0	-698030	.493582	.246059
-0.110	6.315	-116.0	173.7	289.7	1.290944	1.000000	1.000000	610	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	610	37.590	-20.2	142.4	162.6	-684708	.484162	.237834
-0.120	6.892	-113.2	173.1	286.3	1.279620	1.000000	1.000000	620	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	620	38.316	-18.5	141.7	160.2	-671310	.474688	.229701
-0.130	7.470	-110.5	172.5	283.0	1.262737	1.000000	1.000000	630	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	630	39.050	-16.8	140.9	160.2	-664273	.465158	.221658
-0.140	8.048	-107.8	172.0	279.8	1.256905	1.000000	1.000000	640	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	640	40.542	-15.0	140.2	155.2	-630629	.445922	.205845
-0.150	8.627	-105.3	171.4	276.7	1.245153	1.000000	1.000000	650	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	650	41.300	-13.3	139.5	152.8	-636211	.436211	.198074
-0.160	9.207	-102.9	170.8	273.7	1.234099	1.000000	1.000000	660	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	660	42.067	-9.8	138.7	150.3	-646436	.436211	.198074
-0.170	9.788	-100.5	170.2	270.7	1.222661	1.000000	1.000000	670	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	670	42.844	-8.1	137.2	145.2	-639151	.416592	.180394
-0.180	10.370	-98.2	169.6	267.8	1.211998	1.000000	1.000000	680	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	680	43.630	-6.3	136.4	142.6	-606678	.406678	.175309
-0.190	10.953	-95.9	169.0	265.0	1.201711	1.000000	1.000000	690	32.065	-33.7	148.0	181.7	-789346	.558152	.324234	690	43.900	-6.3	136.4	142.6	-575130	.302848	.175309
-0.200	11.537	-93.7	168.5	262.2	1.188199	1.000000	1.000000	700	44.4277	-4.5	135.6	140.1	-396691	.561006	.396691	700	45.235	-2.7	134.8	137.4	-546772	.532425	.160591
-0.210	12.122	-91.5	167.9	259.4	1.176660	1.000000	1.000000	710	46.054	-8	133.9	134.8	-386626	.532425	.160591	710	46.886	-1.0	133.9	134.8	-532425	.532425	.153372
-0.220	12.709	-89.4	167.3	256.7	1.165096	1.000000	1.000000	720	46.054	-8	133.9	134.8	-386626	.532425	.160591	720	47.731	-2.9	132.3	129.4	-532425	.532425	.153372
-0.230	13.297	-87.3	166.7	254.0	1.153504	1.000000	1.000000	730	48.590	-4.7	131.4	126.7	-355935	.532425	.153372	730	49.590	-4.7	131.4	126.7	-488645	.355935	.142245
-0.240	13.887	-85.3	166.1	251.4	1.141884	1.000000	1.000000	740	49.590	-4.7	131.4	126.7	-355935	.532425	.153372	740	50.354	-8.6	129.5	123.9	-473784	.355935	.142245
-0.250	14.478	-83.3	165.5	249.8	1.130236	1.000000	1.000000	750	50.354	-8.6	129.5	123.9	-355935	.532425	.153372	750	51.261	-10.6	128.7	118.2	-443616	.313684	.142245
-0.260	15.070	-81.3	164.9	246.2	1.118559	1.000000	1.000000	760	51.261	-10.6	127.8	115.3	-428292	.512650	.142245	760	52.186	-12.6	127.8	115.3	-428292	.512650	.142245
-0.270	15.664	-79.3	164.3	243.7	1.106853	1.000000	1.000000	770	52.186	-12.6	127.8	115.3	-428292	.512650	.142245	770	53.130	-14.6	126.9	112.3	-412793	.291889	.099065
-0.280	16.260	-77.4	163.7	241.1	1.095116	1.000000	1.000000	780	54.096	-16.7	125.9	109.2	-397110	.546772	.099065	780	55.096	-16.7	125.9	109.2	-381228	.546772	.099065
-0.290	16.858	-75.5	163.1	238.6	1.083348	1.000000	1.000000	790	55.096	-16.7	124.9	106.2	-397110	.546772	.099065	790	56.099	-20.9	123.9	103.0	-365135	.546772	.099065
-0.300	17.458	-73.6	162.5	236.1	1.071549	1.000000	1.000000	800	55.096	-16.7	124.9	106.2	-397110	.546772	.099065	800	57.140	-23.1	122.9	99.7	-348813	.246648	.080375
-0.310	18.059	-71.7	161.9	233.7	1.059717	1.000000	1.000000	810	55.096	-16.7	124.9	106.2	-397110	.546772	.099065	810	58.212	-25.4	121.8	96.4	-332245	.246648	.080375
-0.320	18.663	-69.9	161.3	231.2	1.047851	1.000000	1.000000	820	56.435	-16.7	124.9	106.2	-397110	.546772	.099065	820	59.317	-27.7	120.7	93.0	-315408	.234933	.084847
-0.330	19.269	-68.1	160.7	228.8	1.035952	1.000000	1.000000	830	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	830	60.435	-30.1	119.5	89.5	-298278	.223027	.082690
-0.340	19.877	-66.2	160.1	226.4	1.024018	1.000000	1.000000	840	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	840	61.642	-32.6	118.4	85.8	-280826	.198574	.081506
-0.350	20.487	-64.4	159.5	224.0	1.012048	1.000000	1.000000	850	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	850	62.877	-35.1	117.1	82.0	-263017	.185981	.081506
-0.360	21.100	-62.7	158.9	221.6	1.000042	1.000000	1.000000	860	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	860	64.172	-35.1	117.1	82.0	-263017	.185981	.081506
-0.370	21.716	-60.9	158.3	219.2	987998	1.000000	1.000000	870	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	870	65.505	-37.8	115.8	78.1	-244809	.173106	.040856
-0.380	22.334	-59.1	157.7	216.8	975916	1.000000	1.000000	880	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	880	66.926	-40.5	114.5	73.9	-244809	.173106	.040856
-0.390	22.954	-57.4	157.0	214.4	963795	1.000000	1.000000	890	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	890	68.435	-43.5	113.1	69.6	-244809	.173106	.040856
-0.400	23.578	-55.7	156.4	212.1	951634	1.000000	1.000000	900	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	900	69.9	-43.5	113.1	69.6	-244809	.173106	.040856
-0.410	24.205	-53.9	155.8	209.7	939431	1.000000	1.000000	910	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	910	70.502	-46.6	111.6	65.0	-187705	.143355	.030797
-0.420	24.835	-52.2	155.2	207.4	927186	1.000000	1.000000	920	57.140	-16.7	124.9	106.2	-397110	.546772	.099065	920	71.805	-53.4	108.2	65.0	-187705	.143355	.030797
-0.430	25.468	-50.5	154.5	205.0	914897	1.000000	1.000000	930	57.140	-16.7	124.9	106.2	-397110	.546772	.099065</								