

# On Longitudinal Emittance

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## 1 Emittance Conversions

The standard formula for an upright ellipse in phase-space  $\Delta\phi \otimes w$  is:

$$\frac{\Delta\phi^2}{\Delta\phi_0^2} + \frac{w^2}{w_0^2} = 1 \quad (1)$$

with  $\Delta\phi = \phi - \phi_s$  and  $w \equiv \delta\gamma = \Delta W/mc^2$ .

$\phi_s$  being the synchronous phase,  $mc^2$  the rest energy,  $W$  the total energy and  $\gamma$  the relativistic factor. It has the emittance

$$\epsilon_w = \Delta\phi_0 w_0 \quad (2)$$

and units [rad].

The ellipse intersects the  $\Delta\phi$ -axis at  $\Delta\phi_0$  and the  $w$ -axis at  $w_0$ .

Let's change to new coordinates, for instance the pair of canonical variables  $\Delta z \otimes \Delta p/p$ , as it is used internally in Trace 3D. The transformation from old to new coordinates is:  $|\Delta z| = \kappa|\Delta\phi| = \frac{\beta\lambda}{2\pi}|\Delta\phi|$  and  $\Delta p/p = \tau w = \gamma/(\gamma^2 - 1)w = (\gamma\beta^2)^{-1}w$ . This gives the modified ellipse equation:

$$\frac{\Delta z^2}{(\kappa\Delta\phi_0)^2} + \frac{(\Delta p/p)^2}{(\tau w_0)^2} = 1 \quad (3)$$

which has the transformed emittance

$$\epsilon_z = \kappa\Delta\phi_0\tau w_0 = \kappa\tau\epsilon_w = \frac{\beta\lambda}{2\pi}\gamma/(\gamma^2 - 1)\epsilon_w = \frac{\lambda}{2\pi\gamma\beta}\epsilon_w, \quad (4)$$

with units [m].

For the  $\Delta\phi \otimes \Delta W$  phase space, because  $\Delta W = mc^2 w$ , we have  $\kappa = 1$  and  $\tau = mc^2$ . So that

$$\epsilon_W = mc^2\epsilon_w \quad (5)$$

with units  $[rad * eV]$ .

Finally for the  $\Delta z \otimes \Delta W$  phase space we get the emittance

$$\epsilon_{zW} = \frac{\beta\lambda}{2\pi} mc^2 \epsilon_w \quad (6)$$

with units  $[m * eV]$ .

## 2 Twiss Parameter Values

To simplify we assume the twiss parameter  $\alpha = 0$ . The twiss parameter  $\gamma$  then reduces to  $1/\beta$  and only two free parameters  $\epsilon$  and  $\beta$  remain to describe the ellipse in phase space completely.

For the longitudinal dynamics in the passage of an rf-gap the intersection  $w_0$  on the w-axis is given by

$$w_0 = \frac{\Delta W}{mc^2} = \sqrt{qE_0 T \beta_s^3 \gamma_s^3 \lambda \sin(-\phi_s) \Delta\phi_0^2 / 2\pi mc^2} \quad (7)$$

$$= \Delta\phi_0 \sqrt{qE_0 T \beta_s^3 \gamma_s^3 \lambda \sin(-\phi_s) / 2\pi mc^2} \quad (8)$$

If  $w_0$  is given  $\Delta\phi_0$  follows from (8) *and vice versa*. Putting  $w_0 = \epsilon_w / \Delta\phi_0$  into (8) we get

$$\Delta\phi_0 = \sqrt{\epsilon_w / \sqrt{qE_0 T \beta_s^3 \gamma_s^3 \lambda \sin(-\phi_s) / 2\pi mc^2}} \quad (9)$$

and from (9) we get finally

$$\gamma_0 = \epsilon_w / \Delta\phi_0^2 = \sqrt{qE_0 T \beta_s^3 \gamma_s^3 \lambda \sin(-\phi_s) / 2\pi mc^2} \quad (10)$$

and

$$\beta_0 = 1/\gamma_0 \quad (11)$$

NOTE: the two twiss parameters  $\gamma_0$  and  $\beta_0$  are independent from the emittance  $\epsilon_w$  and completely defined by the cavity field  $E_0$ , rf-phase  $\phi_s$ , rf-wavelength  $\lambda$  and particle impuls  $\sim \gamma\beta$ .

### 3 Appendix

Table 1: SIMULAC variable names

$\epsilon_w = \text{emitw}$	$\Delta\phi = \text{Dphi}$	$\Delta\phi_0 = \text{Dphi0}$	$w = \text{w}$	$w_0 = \text{w0}$
$\epsilon_W = \text{emitW}$	$\Delta z = \text{z}$	$\Delta W = \text{DW}$	$\Delta p/p = \text{Dp2p}$	$\Delta p/p_0 = \text{Dp2p0}$
$\epsilon_z = \text{emitz}$	$\beta_z = \text{betaz}$	$\gamma_z = \text{gammaz}$	$\alpha_z = \text{alphaz}$	$\lambda = \text{lamb}$
$Ez_{avg} = \text{EzAvg}$	$Ez_{peak} = \text{EzPeak}$	$\phi_+ = \text{phi\_1}$	$\phi_- = \text{phi\_2}$	$\psi = \text{psi}$
$\gamma = \text{gamma}$	$\gamma\beta = \text{gb}$	$\beta = \text{beta}$	$E_0T = \text{E0T}$	$mc^3 = \text{m0c3}$
$mc^2 = \text{m0c2}$	$\epsilon_{xi} = \text{emitx\_i}$	$\epsilon_{yi} = \text{emity\_i}$	$\epsilon_{zi} = \text{emitz\_i}$	$\beta_{xi} = \text{betax\_i}$
$\beta_{yi} = \text{betay\_i}$	$\alpha_{xi} = \text{alfax\_i}$	$\alpha_{yi} = \text{alfay\_i}$	$\gamma_{xi} = \text{gamax\_i}$	$\gamma_{yi} = \text{gamay\_i}$
$\omega = \text{omg}$	$\phi = \text{phi}$	$\phi_s = \text{phis}$		