

**NANYANG TECHNOLOGICAL UNIVERSITY**  
**SPMS/DIVISION OF MATHEMATICAL SCIENCES**

2025/26 Semester 2      MH4110 Partial Differential Equations      Tutorial 3, 05 February

**Problem 1** Solve the following first order PDE

$$x^2 v_x + xy v_y = v^2.$$

[Solution:] We define  $u = -1/v$  to change the PDE into  $x^2 u_x + xy u_y = 1$ , and then solve it using the characteristic method. The characteristic curves are given by

$$\frac{dy}{dx} = \frac{xy}{x^2} = \frac{y}{x}.$$

This is a separable ODE, which can be solved to obtain the general solution  $y/x = C$ . Thus, our change of coordinates (no need to be orthogonal but should be non-degenerate) will be

$$\begin{cases} x' = x, \\ y' = \frac{y}{x}. \end{cases}$$

Use the chain rule, we have

$$\begin{cases} \frac{\partial u}{\partial x} = \frac{\partial x'}{\partial x} \frac{\partial u}{\partial x'} + \frac{\partial y'}{\partial x} \frac{\partial u}{\partial y'} = 1 \cdot \frac{\partial u}{\partial x'} - \frac{y}{x^2} \frac{\partial u}{\partial y'} = \frac{\partial u}{\partial x'} - \frac{y'}{x'} \frac{\partial u}{\partial y'}, \\ \frac{\partial u}{\partial y} = \frac{\partial x'}{\partial y} \frac{\partial u}{\partial x'} + \frac{\partial y'}{\partial y} \frac{\partial u}{\partial y'} = 0 \cdot \frac{\partial u}{\partial x'} + \frac{1}{x} \frac{\partial u}{\partial y'} = \frac{1}{x'} \frac{\partial u}{\partial y'}. \end{cases}$$

In these coordinates the equation takes the form

$$x'^2 \left( \frac{\partial u}{\partial x'} - \frac{y'}{x'} \frac{\partial u}{\partial y'} \right) + x'^2 y' \cdot \frac{1}{x'} \frac{\partial u}{\partial y'} = 1,$$

or

$$u_{x'} = \frac{1}{x'^2}.$$

Integrating both sides in  $x'$ , we arrive at

$$u = -\frac{1}{x'} - f(y').$$

Finally, substituting the expressions of  $x'$  and  $y'$  in terms of  $(x, y)$  into the solution, we obtain

$$u(x, y) = -\frac{1}{x} - f\left(\frac{y}{x}\right).$$

or

$$v(x, y) = \frac{x}{1 + x f\left(\frac{y}{x}\right)}.$$

One should again check by substitution that this is indeed a solution to the PDE.

**Problem 2 (Ex. 5 on Page 28)** Consider the equation

$$u_x + y u_y = 0.$$

with the boundary condition  $u(x, 0) = \phi(x)$ .

- (a) For  $\phi(x) \equiv x$ , show that no solution exists.
- (b) For  $\phi(x) \equiv 1$ , show that there are many solutions.

[Solution:] We look for the characteristic curve  $y = y(x)$  satisfying

$$\frac{dy}{dx} = y \quad \Rightarrow \quad y(x) = Ce^x.$$

Hence, the general solution of the problem is

$$u(x, y) = f(e^{-x}y).$$

Given the boundary condition  $u(x, 0) = \phi(x)$ , we have  $u(x, 0) = f(0) = \phi(x)$ .

- (a) If  $\phi(x) \equiv x$ , this contradicts the fact that  $u(x, 0) = f(0)$  is a constant. So no solution exists.
- (b) If  $\phi(x) \equiv 1$ , that means  $f(0) = 1$ . There are many arbitrary functions satisfying  $f(0) = 1$ . For example,  $f(x) = 1 + \sum_{i=1}^n a_i x^{n-i}$  ( $a_i$  are constants and  $n$  is an integer). So there are many solutions.

**Problem 3** Let's consider the PDE

$$a_{11}u_{xx} + 2a_{12}u_{xy} + a_{22}u_{yy} + a_1u_x + a_2u_y + a_0u = 0, \quad (1)$$

which is a linear equation of order two in two variables with six real constant coefficients. We know that

- Equation (1) is of **elliptic type**, if  $a_{11}a_{22} - a_{12}^2 > 0$ . By a linear transform, it can be reduced to

$$u_{xx} + u_{yy} + \{\text{terms of lower order 1 or 0}\} = 0. \quad (2)$$

- Equation (1) is of **hyperbolic type**, if  $a_{11}a_{22} - a_{12}^2 < 0$ . By a linear transform, it can be reduced to

$$u_{xx} - u_{yy} + \{\text{terms of lower order 1 or 0}\} = 0. \quad (3)$$

- Equation (1) is of **parabolic type**, if  $a_{11}a_{22} - a_{12}^2 = 0$ . By a linear transform, it can be reduced to

$$u_{xx} + \{\text{terms of lower order 1 or 0}\} = 0, \quad (4)$$

(unless  $a_{11} = a_{12} = a_{22} = 0$ .)

Please explicitly show how to reduce (1) into the form of (2), (3), or (4). (*Hint: use the method of completing the square.*)

[Solution:] Since the linear transform does not change the order of partial derivatives, we assume that  $a_1 = a_2 = a_0 = 0$ . Without loss of generality, we assume that  $a_{11} = 1$ . Then the

second-order part is

$$\begin{aligned} u_{xx} + 2a_{12}u_{xy} + a_{22}u_{yy} &= (\partial_x^2 + 2a_{12}\partial_{xy} + a_{22}\partial_y^2)u \\ &= \{(\partial_x + a_{12}\partial_y)^2 + (-a_{12}^2 + a_{22})\partial_y^2\}u. \end{aligned}$$

Now, we know the classification of three cases is based on the sign of  $a_{22} - a_{12}^2$ .

We look for a suitable linear transform of the form:

$$x = \alpha\xi + \beta\eta, \quad y = \gamma\xi + \delta\eta. \quad (5)$$

Notice that

$$\begin{aligned} u_\xi &= u_x \frac{\partial x}{\partial \xi} + u_y \frac{\partial y}{\partial \xi} = \alpha u_x + \gamma u_y, \\ u_\eta &= u_x \frac{\partial x}{\partial \eta} + u_y \frac{\partial y}{\partial \eta} = \beta u_x + \delta u_y, \\ \Rightarrow \quad \partial_\xi &= \alpha \partial_x + \gamma \partial_y, \quad \partial_\eta = \beta \partial_x + \delta \partial_y. \end{aligned}$$

We take

$$\alpha = 1, \quad \gamma = a_{12}, \quad \beta = 0, \quad \delta = \sqrt{|a_{22} - a_{12}^2|}, \quad (6)$$

in (5), i.e., the transform:

$$x = \xi, \quad y = a_{12}\xi + \sqrt{|a_{22} - a_{12}^2|}\eta. \quad (7)$$

- If  $a_{22} - a_{12}^2 > 0$  (elliptic case), we can transform the equation to  $u_{\xi\xi} + u_{\eta\eta} = 0$ .
- If  $a_{22} - a_{12}^2 < 0$  (hyperbolic case), we can use the same transform (6) to convert the equation to  $u_{\xi\xi} - u_{\eta\eta} = 0$ .
- If  $a_{22} - a_{12}^2 = 0$  (parabolic case), we cannot use the transform (6) to convert the equation to  $u_{\xi\xi} = 0$  as it is degenerate. The characteristic lines corresponding to the linear operator  $\partial_x + a_{12}\partial_y$  are  $a_{12}x - y = C$ . Thus, we can use the transform:

$$\xi = x, \quad \eta = a_{12}x - y \quad (8)$$

to reduce the equation to  $u_{\xi\xi} = 0$ .

You also can refer to the proof of Theorem 1 on Page 29 of the reference book.

**Problem 4 (Ex. 1 on Page 31)** What is the type of each of the following equations?

- $u_{xx} - u_{xy} + 2u_y + u_{yy} - 3u_{yx} + 4u = 0$ .
- $9u_{xx} + 6u_{xy} + u_{yy} + u_x = 0$ .

[Solution:]

- Assume that the mixed partial derivatives  $u_{xy}$  and  $u_{yx}$  are equal to each other, the PDE  $u_{xx} - u_{xy} + 2u_y + u_{yy} - 3u_{yx} + 4u = 0$  can be simplified into  $u_{xx} - 4u_{xy} + u_{yy} + 2u_y + 4u = 0$ . The coefficient matrix is therefore

$$A = \begin{bmatrix} 1 & -2 \\ -2 & 1 \end{bmatrix}.$$

Its determinant is  $\det(A) = 1 \times 1 - (-2) \times (-2) = -3 < 0$ . The PDE is of hyperbolic type.

(b) The coefficient matrix is

$$A = \begin{bmatrix} 9 & 3 \\ 3 & 1 \end{bmatrix}.$$

Its determinant is  $\det(A) = 9 \times 1 - 3 \times 3 = 0$ . The PDE is of parabolic type.

**Problem 5 (Ex. 2 on Page 31)** Find the regions in the  $xy$  plane where the equation

$$(1+x)u_{xx} + 2xyu_{xy} - y^2u_{yy} = 0$$

is elliptic, hyperbolic, or parabolic. Sketch them.

[Solution:]

The coefficient matrix is

$$A = \begin{bmatrix} 1+x & xy \\ xy & -y^2 \end{bmatrix}.$$

Its determinant is  $\det(A) = (1+x)(-y^2) - (xy)^2 = -(x^2 + x + 1)y^2 = -[(x + \frac{1}{2})^2 + \frac{3}{4}]y^2$ .

We have

- If  $y = 0$ , then  $\det(A) = 0$ . The PDE is of parabolic type.
- If  $y \neq 0$ , then  $\det(A) < 0$ . The PDE is of hyperbolic type.

**Problem 6 (Ex. 6 on Page 32)** Consider the equation  $3u_y + u_{xy} = 0$ .

- (a) What is its type?
- (b) Find the general solution. *Hint: Substitute  $v = u_y$ .*
- (c) With the auxiliary conditions  $u(x, 0) = e^{-3x}$  and  $u_y(x, 0) = 0$ , does a solution exist? Is it unique?

[Solution:]

- (a) We see that  $a_{11} = 0, a_{12} = 1/2, a_{22} = 0$ , so we have  $a_{11}a_{22} - a_{12}^2 = -1/4 < 0$ . The equation is of hyperbolic type.
- (b) Set  $v = u_y$ , the original equation  $3u_y + u_{xy} = 0$  can be rewritten as  $3v + v_x = 0$ . Its general solution is  $v(x, y) = f(y)e^{-3x}$  with  $f$  as an arbitrary function. Integrating over  $y$ , we have the general solution to  $u_y = v$  as  $u(x, y) = F(y)e^{-3x} + G(x)$ . Here  $F$  (a primitive function of  $f$ ) and  $G$  are both arbitrary functions of a single variable.
- (c) Given that  $u(x, 0) = e^{-3x}$ , we have  $u(x, 0) = F(0)e^{-3x} + G(x) = e^{-3x}$ , so  $G(x) = [1 - F(0)]e^{-3x}$ . Additionally,  $u_y(x, 0) = 0$  implies that  $u_y(x, 0) = F'(0)e^{-3x} = 0$  and hence  $F'(0) = 0$ . Therefore, the solution is  $u(x, y) = [F(y) + 1 - F(0)]e^{-3x}$  with the function  $F(y)$  satisfying  $F'(0) = 0$ . There are many possibilities for  $F(y)$ . For example,  $F(y) = cy^2$  or  $F(y) = cy^3$  ( $c$  is an arbitrary constant). The solution exists but is not unique.

**Problem 7** Reduce the elliptic equation

$$u_{xx} + 3u_{yy} - 2u_x + 24u_y + 5u = 0$$

to the form  $v_{xx} + v_{yy} + cv = 0$ .

[Solution:] We look for a suitable linear transform of the form:

$$\xi = ax + by, \quad \eta = cx + dy.$$

The Jacobian determinant of the linear transform is

$$J = \left| \frac{\partial(\xi, \eta)}{\partial(x, y)} \right| = ad - bc.$$

To have a non-degenerate transform,  $ad \neq bc$  is required. The first and second order derivatives in the  $\xi - \eta$  coordinates can be expressed as

$$\begin{aligned} u_x &= u_\xi \frac{\partial \xi}{\partial x} + u_\eta \frac{\partial \eta}{\partial x} = au_\xi + cu_\eta, \\ u_y &= u_\xi \frac{\partial \xi}{\partial y} + u_\eta \frac{\partial \eta}{\partial y} = bu_\xi + du_\eta, \end{aligned}$$

and further

$$\begin{aligned} u_{xx} &= (au_\xi + cu_\eta)_\xi \frac{\partial \xi}{\partial x} + (au_\xi + cu_\eta)_\eta \frac{\partial \eta}{\partial x} = (a^2u_{\xi\xi} + acu_{\xi\eta}) + (acu_{\xi\eta} + c^2u_{\eta\eta}) \\ &= a^2u_{\xi\xi} + 2acu_{\xi\eta} + c^2u_{\eta\eta}, \\ u_{yy} &= (bu_\xi + du_\eta)_\xi \frac{\partial \xi}{\partial y} + (bu_\xi + du_\eta)_\eta \frac{\partial \eta}{\partial y} = (b^2u_{\xi\xi} + bdu_{\xi\eta}) + (bdu_{\xi\eta} + d^2u_{\eta\eta}) \\ &= b^2u_{\xi\xi} + 2bdu_{\xi\eta} + d^2u_{\eta\eta}. \end{aligned}$$

So, the PDE can be rewritten as

$$\begin{aligned} 0 &= u_{xx} + 3u_{yy} - 2u_x + 24u_y + 5u \\ &= (a^2u_{\xi\xi} + 2acu_{\xi\eta} + c^2u_{\eta\eta}) + 3(b^2u_{\xi\xi} + 2bdu_{\xi\eta} + d^2u_{\eta\eta}) \\ &\quad - 2(au_\xi + cu_\eta) + 24(bu_\xi + du_\eta) + 5u \\ &= (a^2 + 3b^2)u_{\xi\xi} + (2ac + 6bd)u_{\xi\eta} + (c^2 + 3d^2)u_{\eta\eta} + (-2a + 24b)u_\xi + (-2c + 24d)u_\eta + 5u \end{aligned}$$

To make the above equation have a form of  $v_{xx} + v_{yy} + cv = 0$ ,  $2ac + 6bd = 0$ ,  $-2a + 24b = 0$ , and  $-2c + 24d = 0$  are required. These imply that  $a = b = 0$  or  $c = d = 0$ . The linear transform would be a degenerate one. We cannot find a linear transform to change the PDE into the desired form.

Alternatively, let  $u = v(x, y)e^{\alpha x + \beta y}$ , then

$$u_x = v_x e^{\alpha x + \beta y} + \alpha v e^{\alpha x + \beta y}, \quad u_y = v_y e^{\alpha x + \beta y} + \beta v e^{\alpha x + \beta y}$$

and

$$\begin{aligned}u_{xx} &= v_{xx}e^{\alpha x + \beta y} + \alpha v_x e^{\alpha x + \beta y} + \alpha v_x e^{\alpha x + \beta y} + \alpha^2 v e^{\alpha x + \beta y}, \\u_{yy} &= v_{yy}e^{\alpha x + \beta y} + \beta v_y e^{\alpha x + \beta y} + \beta v_y e^{\alpha x + \beta y} + \beta^2 v e^{\alpha x + \beta y}.\end{aligned}$$

The original PDE can be rewritten as

$$\begin{aligned}0 &= u_{xx} + 3u_{yy} - 2u_x + 24u_y + 5u \\&= \left( v_{xx}e^{\alpha x + \beta y} + \alpha v_x e^{\alpha x + \beta y} + \alpha v_x e^{\alpha x + \beta y} + \alpha^2 v e^{\alpha x + \beta y} \right) \\&\quad + 3 \left( v_{yy}e^{\alpha x + \beta y} + \beta v_y e^{\alpha x + \beta y} + \beta v_y e^{\alpha x + \beta y} + \beta^2 v e^{\alpha x + \beta y} \right) \\&\quad - 2 \left( v_x e^{\alpha x + \beta y} + \alpha v e^{\alpha x + \beta y} \right) + 24 \left( v_y e^{\alpha x + \beta y} + \beta v e^{\alpha x + \beta y} \right) + 5v e^{\alpha x + \beta y},\end{aligned}$$

which can be simplified as

$$\begin{aligned}0 &= (v_{xx} + 2\alpha v_x + \alpha^2 v) + 3(v_{yy} + 2\beta v_y + \beta^2 v) - 2(v_x + \alpha v) + 24(v_y + \beta v) + 5v \\&= v_{xx} + 3v_{yy} + (2\alpha - 2)v_x + (6\beta + 24)v_y + (\alpha^2 + 3\beta^2 - 2\alpha + 24\beta + 5)v\end{aligned}$$

Assuming that  $\alpha = 1$  and  $\beta = -4$ , the change of variable  $u = ve^{x-4y}$  leads to a new PDE as

$$v_{xx} + 3v_{yy} - 44v = 0.$$

Letting  $y' = \sqrt{1/3}y$ , we have  $v_{yy} = v_{y'y'}(dy'/dy)^2 = (1/3)v_{y'y'}$ . In the  $x - 0 - y'$  coordinates, the above PDE has the form

$$v_{xx} + v_{y'y'} - 44v = 0.$$

**Problem 8 (Ex. 2 on Page 38)** Solve  $u_{tt} = 3u_{xx}$ ,  $u(x, 0) = \ln(1 + x^2)$ ,  $u_t(x, 0) = 4 + x$ .

[Solution:] In this problem the speed of the wave is  $c = \sqrt{3}$ . By the d'Alembert formula,

$$u(x, t) = \frac{\phi(x - ct) + \phi(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds.$$

Since  $\phi(x) = \ln(1 + x^2)$  and  $\psi(x) = 4 + x$ , we have

$$\begin{aligned}u(x, t) &= \frac{\ln[1 + (x - ct)^2] + \ln[1 + (x + ct)^2]}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} (4 + s) ds \\&= \frac{1}{2} \ln[1 + (x - ct)^2 + (x + ct)^2 + (x^2 - c^2 t^2)^2] \\&\quad + \frac{1}{2c} \left[ 4(x + ct) + \frac{1}{2}(x + ct)^2 - 4(x - ct) - \frac{1}{2}(x - ct)^2 \right] \\&= \frac{1}{2} \ln[1 + 2x^2 + 2c^2 t^2 + (x^2 - c^2 t^2)^2] + \frac{1}{2c} [8ct + 2ctx] \\&= \frac{1}{2} \ln[1 + 2x^2 + 2c^2 t^2 + (x^2 - c^2 t^2)^2] + (4 + x)t.\end{aligned}$$

Thus, the solution is

$$\begin{aligned} u(x, t) &= \frac{1}{2} \ln[1 + 2x^2 + 2c^2t^2 + (x^2 - c^2t^2)^2] + (4 + x)t \\ &= \frac{1}{2} \ln[1 + 2x^2 + 6t^2 + (x^2 - 3t^2)^2] + (4 + x)t. \end{aligned}$$

**Problem 9 (Ex. 7 on Page 38)** If both  $\phi$  and  $\psi$  are even functions of  $x$ , show that the solution  $u(x, t)$  of the wave equation is also even in  $x$  for all  $t$ .

[Solution:] By the d'Alembert formula,

$$u(x, t) = \frac{\phi(x - ct) + \phi(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds.$$

Since  $\phi(-s) = \phi(s)$  and  $\psi(-s) = \psi(s)$ , we have

$$u(-x, t) = \frac{\phi(-x - ct) + \phi(-x + ct)}{2} + \frac{1}{2c} \int_{-x-ct}^{-x+ct} \psi(s) ds$$

Make a change of variables as  $s = -y$  and note that  $ds = -dy$ , then

$$\begin{aligned} u(-x, t) &= \frac{\phi(x + ct) + \phi(x - ct)}{2} - \frac{1}{2c} \int_{x+ct}^{x-ct} \psi(-y) dy \\ &= \frac{\phi(x + ct) + \phi(x - ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(y) dy \\ &= u(x, t). \end{aligned}$$

This means  $u(x, t)$  is even in  $x$  for all  $t$ .

**Problem 10 (Ex. 5 on Page 38: The hammer blow)** Let  $\phi(x) \equiv 0$  and  $\psi(x) = 1$  for  $|x| < a$  and  $\psi(x) = 0$  for  $|x| \geq a$ . Sketch the string profile ( $u$  versus  $x$ ) at each of the successive instants  $t = a/2c, a/c, 3a/2c, 2a/c$ , and  $5a/c$ .

[Solution:] By the d'Alembert formula,

$$u(x, t) = \frac{\phi(x - ct) + \phi(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds.$$

Given  $\phi(x) \equiv 0$  and  $\psi(x) = 1$  for  $|x| < a$  and  $\psi(x) = 0$  for  $|x| \geq a$ , we have

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds.$$

Since the initial velocity  $\psi(x)$  is nonzero only in the interval  $[-a, a]$ , the integral must be computed differently according to how the intervals  $[-a, a]$  and  $[x - ct, x + ct]$  intersect.

- For a small value of  $t$  (only if  $ct \leq a$ ), the two intervals  $[-a, a]$  and  $[x - ct, x + ct]$  intersect in the following 5 different ways

- If  $x - ct < x + ct \leq -a < a$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = 0.$$

- If  $x - ct \leq -a < x + ct \leq a$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = \frac{1}{2c} \int_{-a}^{x+ct} 1 ds = \frac{x + ct + a}{2c}.$$

- If  $-a \leq x - ct < x + ct \leq a$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = \frac{1}{2c} \int_{x-ct}^{x+ct} 1 ds = t.$$

- If  $-a \leq x - ct \leq a < x + ct$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = \frac{1}{2c} \int_{x-ct}^a 1 ds = \frac{-x + ct + a}{2c}.$$

- If  $-a < a \leq x - ct < x + ct$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = 0.$$

- For a big value of  $t$  (only if  $ct > a$ ), the two intervals  $[-a, a]$  and  $[x - ct, x + ct]$  also intersect in 5 different ways

- If  $x - ct < x + ct \leq -a < a$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = 0.$$

- If  $x - ct \leq -a < x + ct \leq a$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = \frac{1}{2c} \int_{-a}^{x+ct} 1 ds = \frac{x + ct + a}{2c}.$$

- If  $x - ct \leq -a < a \leq x + ct$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = \frac{1}{2c} \int_{-a}^a 1 ds = \frac{a}{c}.$$

- If  $-a \leq x - ct \leq a < x + ct$ , then

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = \frac{1}{2c} \int_{x-ct}^a 1 ds = \frac{-x + ct + a}{2c}.$$

- If  $-a < a \leq x - ct < x + ct$ , then

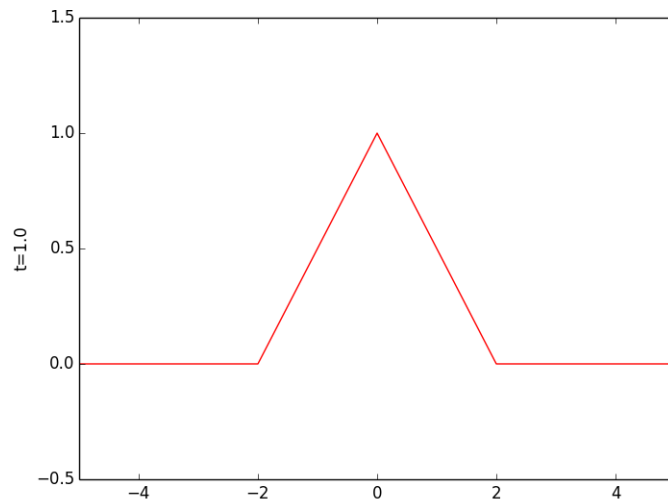
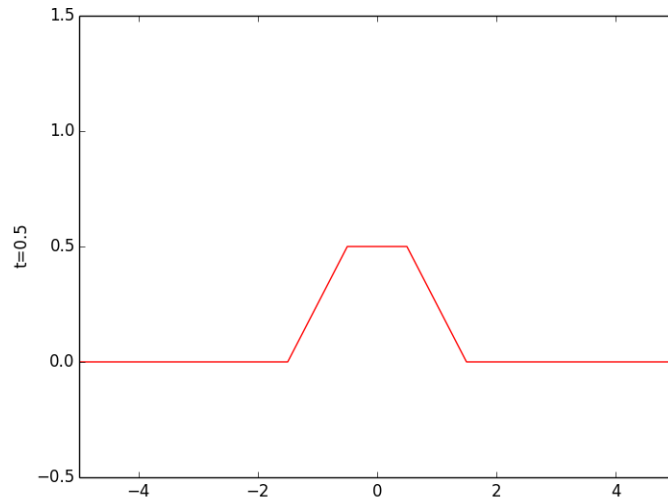
$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds = 0.$$

**Problem 11** Consider the equation

$$au_{tt} + bu_{xt} + cu_{xx} = 0, \quad -\infty < x < \infty, \quad t > 0, \quad (9)$$

$$u(x, 0) = \phi(x), \quad u_t(x, 0) = \psi(x), \quad -\infty < x < \infty, \quad (10)$$





where  $a$ ,  $b$ , and  $c$  are constants such that  $ac < 0$ . Show that the equation is hyperbolic, and derive the solution formula.

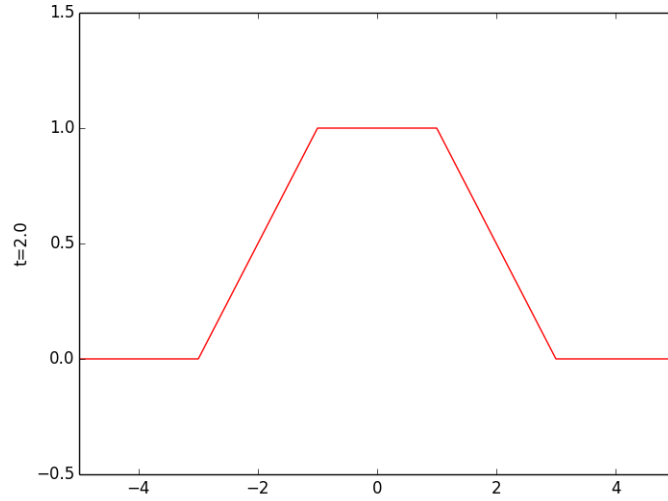
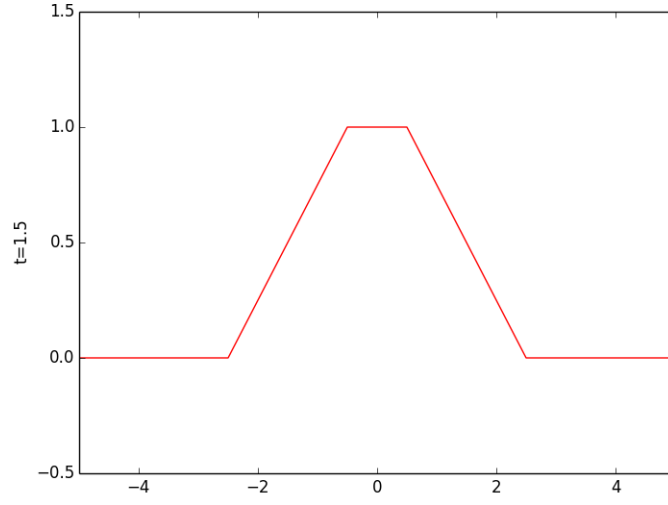
[Solution:] By the definition in Chapter 1, we see that the associated determinant is  $ac - b^2/4 < 0$ , so the equation is hyperbolic.

Letting  $\alpha = (b + \sqrt{b^2 - 4ac})/(2a)$  and  $\beta = (b - \sqrt{b^2 - 4ac})/(2a)$ , we find that  $\alpha + \beta = b/a$  and  $\alpha\beta = c/a$ . We factor the differential operator as

$$(a\partial_t^2 + b\partial_t\partial_x + c\partial_x^2)u = a(\partial_t + \alpha\partial_x)(\partial_t + \beta\partial_x)u = 0.$$

Following the idea of solving the wave equation, we introduce  $v = (\partial_t + \beta\partial_x)u$ , and so we have

$$a(v_t + \alpha v_x) = 0 \Rightarrow v(x, t) = h(x - \alpha t).$$



We now solve

$$u_t + \beta u_x = h(x - \alpha t).$$

Notice that it has a particular solution of the form  $u = f(x - \alpha t)$  with  $f'(s) = h(s)/(\beta - \alpha)$ , and the general solution of  $u_t + \beta u_x = 0$  is  $g(x - \beta t)$ . Therefore, the general solution of (9) is

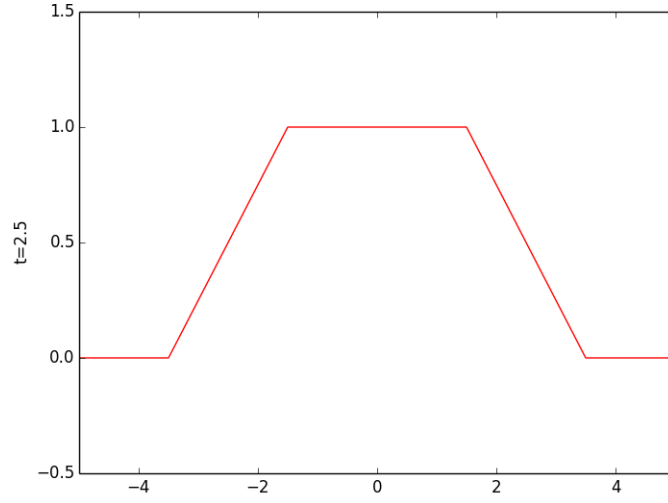
$$u(x, t) = f(x - \alpha t) + g(x - \beta t).$$

Now, we impose the initial conditions (10), and find that

$$f(x) + g(x) = \phi(x), \quad -\alpha f'(x) - \beta g'(x) = \psi(x).$$

Solving the system:

$$f'(x) + g'(x) = \phi'(x), \quad -\alpha f'(x) - \beta g'(x) = \psi(x),$$



leads to

$$f(x) = -\frac{\beta}{\alpha - \beta}\phi(x) - \frac{1}{\alpha - \beta} \int_0^x \psi(s)ds + A,$$

$$g(x) = \frac{\alpha}{\alpha - \beta}\phi(x) + \frac{1}{\alpha - \beta} \int_0^x \psi(s)ds + B,$$

where  $A + B = 0$ .

Finally, we obtain the d'Alembert-type formula

$$u(x, t) = \frac{\alpha\phi(x - \beta t) - \beta\phi(x - \alpha t)}{\alpha - \beta} + \frac{1}{\alpha - \beta} \int_{x-\alpha t}^{x-\beta t} \psi(s)ds,$$

where  $\alpha = (b + \sqrt{b^2 - 4ac})/(2a)$  and  $\beta = (b - \sqrt{b^2 - 4ac})/(2a)$ .