## Forced Oscillations

Colby Community College

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and

$$x = c_1 \cos(t) + c_2 \sin(t) - \frac{1}{4} \cos(3t)$$

#### General Solution

We can now look at the general solution for the undamped system

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

Where  $\omega_f$  is the forcing frequency and  $F_0$  is the forcing amplitute.

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This leaves two separate cases for  $x_p$ :

- **1** The frequencies  $\omega_f$  and  $\omega_0$  are different.
- 2 The frequencies  $\omega_f$  and  $\omega_0$  are the same.

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So, where  $c_1$  and  $c_2$  are determined by initial conditions, we have

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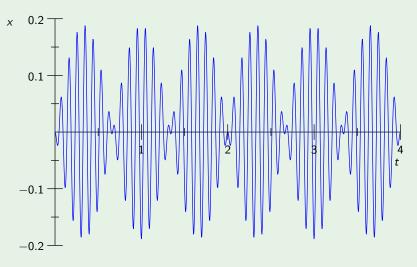
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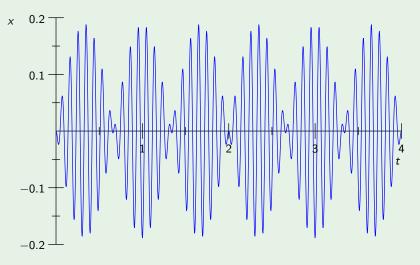
$$x(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$
$$= C \cos(\omega_0 t - \delta) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

where 
$$C = \sqrt{c_1^2 + c_2^2}$$
 and  $\tan(\delta) = \frac{c_2}{c_1}$ .

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The regular periodic patterns are called beats.

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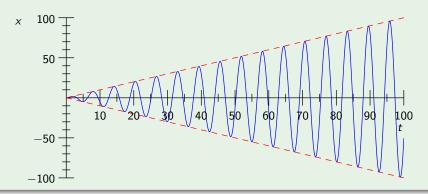
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If the system is initially at rest  $(x(0) = 0 \text{ and } \dot{x}(0) = 0)$  then the solution is

$$x(t) = -\frac{F_0}{m\left(\omega_0^2 - \omega_f^2\right)}\cos\left(\omega_0 t\right) + \frac{F_0}{m\left(\omega_0^2 - \omega_f^2\right)}\cos\left(\omega_f t\right)$$

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So, we can simplify using the trigonometric identity

$$\cos(u) - \cos(v) = -2\sin\left(\frac{u-v}{2}\right)\sin\left(\frac{u+v}{2}\right)$$

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$$= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} (\cos(\omega_0 t) - \cos(\omega_f t))$$

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Thus  $\sin\left(\frac{\omega_0-\omega_f}{2}t\right)$  oscillates much slower than  $\sin\left(\frac{\omega_0+\omega_f}{2}t\right)$ .

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The two curves

$$\pm \frac{2F_0}{m\left(\omega_0^2 - \omega_f^2\right)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right)$$

form an envelope of the more rapid oscillation and is called the **sinusoidal** amplitude.

# Solutions to the Undamped Forced Oscillator ( $\omega_{\it f} \neq \omega_{\it 0}$ )

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where  $c_1$  and  $c_2$  are determined by initial conditions and  $\omega_0 = \sqrt{rac{k}{m}}.$ 

If the system starts from rest (x(0) = 0) and  $\dot{x}(0) = 0$ , the solution can be written as

$$x(t) = \underbrace{\frac{2F_0}{m\left(\omega_0^2 - \omega_f^2\right)}\sin\left(\frac{\omega_0 - \omega_f}{2}t\right)}_{\text{sinusoidal amplitude}} \underbrace{\sin\left(\frac{\omega_0 + \omega_f}{2}t\right)}_{\text{rapid oscillation within beats}}$$

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We can then calculate k:

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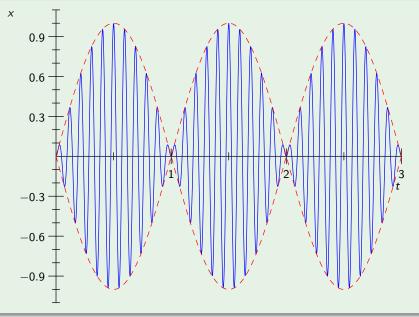
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Thus, envelope curves are

$$y = \pm 1 \cdot \sin(\pi t)$$



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$$x(t) = A\cos(3t) + B\sin(3t)$$
  
 $\dot{x}(t) = -3A\sin(3t) + 3B\cos(3t)$   
 $\ddot{x}(t) = -9A\cos(3t) - 9B\sin(3t)$ 

Substituting into the DE gives

$$(-9A\cos(3t) - 9B\sin(3t)) + 4(-3A\sin(3t) + 3B\cos(3t)) + 5(A\cos(3t) + B\sin(3t)) = 10\cos(3t)$$

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Thus,

$$x_p = -\frac{1}{4}\cos(3t) + \frac{3}{4}\sin(3t)$$

The general solution is

$$x = e^{-2t} \left( c_1 \cos(t) + c_2 \sin(t) \right) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

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To solve the IVP, we need to calculate

$$\dot{x} = -2e^{-2t} \left( c_1 \cos(t) + c_2 \sin(t) \right) + e^{-2t} \left( -c_1 \sin(t) + c_2 \cos(t) \right) + \frac{1}{4} \sin(3t) - \frac{3}{4} \cos(3t)$$

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$$x(0) = 0$$
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 $\dot{x}(0) = 0$   $\Rightarrow$   $c_2 + \frac{7}{4} = 0$ 

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 $\dot{x}(0) = 0$   $\Rightarrow$   $c_2 + \frac{7}{4} = 0$   $\Rightarrow$   $c_2 = -\frac{7}{4}$ 

The solution to the IVP is

$$x = e^{-2t} \left( \frac{1}{4} \cos(t) - \frac{7}{4} \sin(t) \right) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

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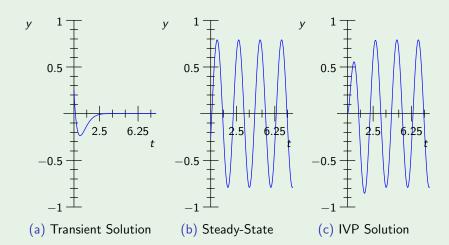
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The particular solution  $x_p$  may either be constant or a periodic **steady-state** solution.



## Particular Solution $x_p$ of a Damped Mass-Spring System

The damped mass-spring system

$$m\ddot{x} + b\dot{x} + kx = F_0 \cos(\omega_f t)$$

has particular solution

$$x_p = A\cos(\omega_f t) + B\sin(\omega_f t)$$

with

$$A = \frac{m(\omega_0^2 - \omega_f^2) F_0}{m^2(\omega_0^2 - \omega_f^2)^2 + (b\omega_f)^2} \quad \text{and} \quad B = \frac{b\omega_f F_0}{m^2(\omega_0^2 - \omega_f^2)^2 + (b\omega_f)^2}$$

with natural circular frequency  $\omega_0 = \sqrt{\frac{k}{m}}.$ 

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#### Note

You will verify this in the homework.