The Harmonic Oscillator

Department of Mathematics

Salt Lake Community College

(Slides by Adam Wilson)

Newton's Dot Notation

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$$\dot{x} = \frac{dx}{dt}$$
 and $\ddot{x} = \frac{dx^2}{d^2t}$ and $\ddot{x} = \frac{dx^3}{d^3t}$ and $\ddot{x} = \frac{dx^4}{d^4t}$

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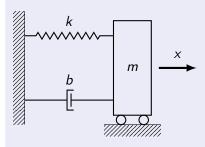
Definition

A very important DE is the second-order homogeneous equation

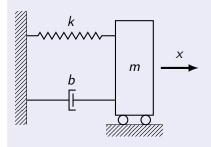
$$m\ddot{x} + b\dot{x} + kx = 0$$

where m > 0, b, and k are constants.

This models a class of phenomena called damped harmonic oscillation.

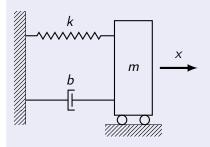


We will model the Mass-Spring system using Newton's Second Law of Motion, $F = m\ddot{x}$, where F is the sum of the following forces:



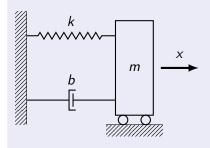
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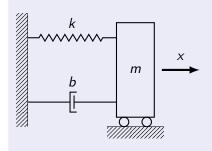
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Summing these forces gives:

mass × acceleration =
$$F_{\text{restoring}}$$
 + F_{damping} + F_{external}
 $m\ddot{x} = -kx$ - $b\dot{x}$ + $f(t)$

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- When b = 0, the motion is called undamped; otherwise it is damped.
- If f(t) = 0 for all t, then the equation is homogeneous:

$$m\ddot{x} + b\dot{x} + kx = 0$$

and the motion is called **unforced**, **undriven**, or **free**; otherwise it is called **forced** or **driven**.

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We also measure the damping force of the object sliding on the table to be 0.5 newtons when the velocity is 0.25 meters per second.

$$b = \frac{0.5 \text{ newton}}{0.25 \frac{\text{meter}}{\text{second}}} = 2 \frac{\text{newton second}}{\text{meter}}$$

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Notice that a second-order DE requires two initial conditions.

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Gravity (Earth)	9.8 $\frac{m}{s^2}$	980.665 $\frac{\text{cm}}{\text{s}^2}$	$32 \frac{ft}{s^2}$

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Comparing

$$(\sin(\omega_0 t))'' = -\omega_0^2 \sin(\omega_0 t)$$
 and $\ddot{x} = -\frac{k}{m}x$

we see that if $\omega_0 = \sqrt{\frac{k}{m}}$, then $x(t) = \sin(\omega_0 t)$ is a solution.

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Note

Another solutions is $x(t) = \cos(\omega_0 t)$.

Solution of the Undamped Unforced Oscillator

For the undamped unforced oscillator

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The Superposition Principle tells us that any linear combination of these two solutions is itself a solution. Thus, for $c_1, c_2 \in \mathbb{R}$, the family of solutions is

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Note

We will see next section that these cover all solutions.

Alternate Form of the Undamped Unforced Oscillator Solution

Solutions to the undamped unforced oscillator may also be expressed as

$$x(t) = A\cos(\omega_0 t - \delta)$$

- *A* is the **amplitude**
- ullet δ is the **phase angle**, measured in radians
- The motion has **circular frequency** $\omega_o = \sqrt{\frac{k}{m}}$, measured in $\frac{\text{radians}}{\text{second}}$
- The natural frequency is $f_0 = \frac{\omega_0}{2\pi}$.
- The **period** is $\frac{1}{f_0} = \frac{2\pi}{\omega_0} = 2\pi\sqrt{\frac{m}{k}}$, measured in seconds
- ullet This is a horizontal translation of $A\cos{(\omega_o t)}$ with **phase shift** $rac{\delta}{\omega_0}$

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Converting Between the Two Forms

The translation is given by

$$A=\sqrt{c_1^2+c_2^2}, \qquad an\left(\delta
ight)=rac{c_2}{c_1}$$

and

$$c_1 = A\cos(\delta), \qquad c_2 = A\sin(\delta)$$

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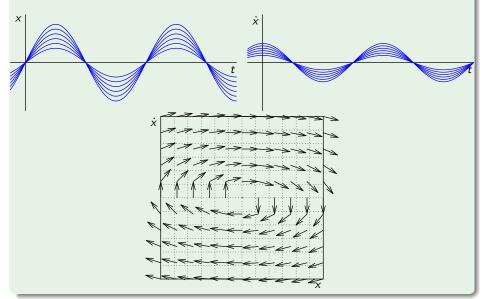
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Substituting t = 0, x(0) = 0, and $\dot{x}(0) = 1$ into this system gives the solution $c_1 = 0$ and $c_2 = 1$.

Let us look at some plots concerning $\ddot{x} + 0.25x = 0$:



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$$\ddot{x} = F(x, \dot{x})$$

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The phase plane has a **vector field** specified by the DE, which at any point in the phase plane gives a direction vector with

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Note

Phase portraits can be graphed without solving the DE.

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is equivalent to the system of first-order equations:

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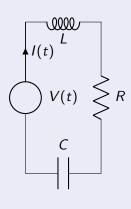
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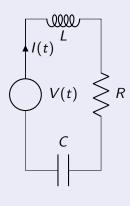
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Then, pplane may be used to plot the phase portrait.



The current I in a wire, measured in amperes, is the flow of charge Q. That is, the current is the rate of change of the charge

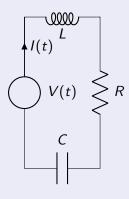
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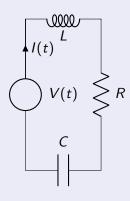
Kirchoff's Voltage Law tell us that the input voltage V(t) is the sum of voltage drops around the circuit. In our circuit, we have three such voltage drops.



Drop across a Resistor: By **Ohm's Law**, the voltage drop across a resistor is proportional to the current passing through it.

$$V_R(t) = R \cdot I(t)$$

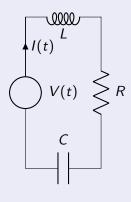
Where *R* is the **resistance** of the resistor and is measured in *ohms*.



Drop across an Inductor: By Faraday's Law, the voltage drop across an inductor is proportional to the time rate of change of the current passing through it.

$$V_L(t) = L \cdot \dot{I}(t)$$

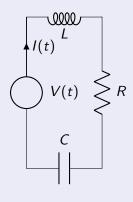
where *L* is the **inductance** and is measured in *henries*.



Drop across a Capacitor: The voltage drop across a capacitor is proportional to the charge Q(t) on the capacitor.

$$V_C(t) = rac{1}{C}Q(t) = rac{1}{C}\int I(t)dt$$

where *C* is the **capacitance** of the capacitor and is measured in *farads*.



Thus, the voltage drop across the circuit is

$$V(t) = R \cdot I + L \cdot \dot{I} + \frac{1}{C} \int I(t) dt$$

This is called an **integro-differential equation** because it contains both a derivative and an integral.

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Series Circuit Equation (Charge)

$$L\ddot{Q} + R\dot{Q} + \frac{1}{C}Q = V(t)$$

If there is no voltage source (V(t) = 0), then

$$L\ddot{Q} + R\dot{Q} + \frac{1}{C}Q = 0$$

Using the fact that $I(t) = \dot{Q}(t)$ we can build the following equations.

Series Circuit Equation (Charge)

$$L\ddot{Q} + R\dot{Q} + \frac{1}{C}Q = V(t)$$

If there is no voltage source (V(t) = 0), then

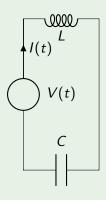
$$L\ddot{Q} + R\dot{Q} + \frac{1}{C}Q = 0$$

Series Circuit Equation (Current)

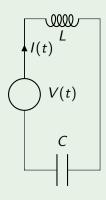
$$L\ddot{I} + R\dot{I} + \frac{1}{C}I = \dot{V}(t)$$

If there is no voltage source (V(t) = 0), then

$$L\ddot{I} + R\dot{I} + \frac{1}{C}I = 0$$



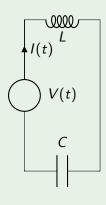
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Thus, the solution is

$$Q(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t)$$

where

$$\omega_0 = \sqrt{\frac{1}{LC}}$$