

Forced Oscillations

Department of Mathematics

Salt Lake Community College

(Slides by Adam Wilson)

Example 1

Let us look at the mass-spring system described by

$$\ddot{x} + x = 2 \cos(3t)$$

Example 1

Let us look at the mass-spring system described by

$$\ddot{x} + x = 2 \cos(3t)$$

The characteristic equation is $r^2 + 1 = 0$ and thus

$$x_h = c_1 \cos(t) + c_2 \sin(t)$$

Example 1

Let us look at the mass-spring system described by

$$\ddot{x} + x = 2 \cos(3t)$$

The characteristic equation is $r^2 + 1 = 0$ and thus

$$x_h = c_1 \cos(t) + c_2 \sin(t)$$

Using the method of undetermined coefficients

$$x_p = A \cos(3t) + B \sin(3t)$$

Example 1

Let us look at the mass-spring system described by

$$\ddot{x} + x = 2 \cos(3t)$$

The characteristic equation is $r^2 + 1 = 0$ and thus

$$x_h = c_1 \cos(t) + c_2 \sin(t)$$

Using the method of undetermined coefficients

$$x_p = A \cos(3t) + B \sin(3t)$$

Which leads to

$$x_p = -\frac{1}{4} \cos(3t)$$

Example 1

Let us look at the mass-spring system described by

$$\ddot{x} + x = 2 \cos(3t)$$

The characteristic equation is $r^2 + 1 = 0$ and thus

$$x_h = c_1 \cos(t) + c_2 \sin(t)$$

Using the method of undetermined coefficients

$$x_p = A \cos(3t) + B \sin(3t)$$

Which leads to

$$x_p = -\frac{1}{4} \cos(3t)$$

and

$$x = c_1 \cos(t) + c_2 \sin(t) - \frac{1}{4} \cos(3t)$$

General Solution

We can now look at the general solution for the undamped system

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

Where ω_f is the **forcing frequency** and F_0 is the **forcing amplitude**.

General Solution

We can now look at the general solution for the undamped system

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

Where ω_f is the **forcing frequency** and F_0 is the **forcing amplitude**.

We know that

$$x_h = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) \quad \text{where} \quad \omega_0 = \sqrt{\frac{k}{m}}$$

General Solution

We can now look at the general solution for the undamped system

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

Where ω_f is the **forcing frequency** and F_0 is the **forcing amplitude**.

We know that

$$x_h = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) \quad \text{where} \quad \omega_0 = \sqrt{\frac{k}{m}}$$

This leaves two separate cases for x_p :

- 1 The frequencies ω_f and ω_0 are different.
- 2 The frequencies ω_f and ω_0 are the same.

Unmatched Frequencies ($\omega_f \neq \omega_0$)

This means we want to look for

$$x_p = A \cos(\omega_f t) + B \sin(\omega_f t)$$

Unmatched Frequencies ($\omega_f \neq \omega_0$)

This means we want to look for

$$x_p = A \cos(\omega_f t) + B \sin(\omega_f t)$$

and so, we find that

$$A = \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \quad \text{and} \quad B = 0$$

Unmatched Frequencies ($\omega_f \neq \omega_0$)

This means we want to look for

$$x_p = A \cos(\omega_f t) + B \sin(\omega_f t)$$

and so, we find that

$$A = \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \quad \text{and} \quad B = 0$$

So, where c_1 and c_2 are determined by initial conditions, we have

$$x(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

Unmatched Frequencies ($\omega_f \neq \omega_0$)

This means we want to look for

$$x_p = A \cos(\omega_f t) + B \sin(\omega_f t)$$

and so, we find that

$$A = \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \quad \text{and} \quad B = 0$$

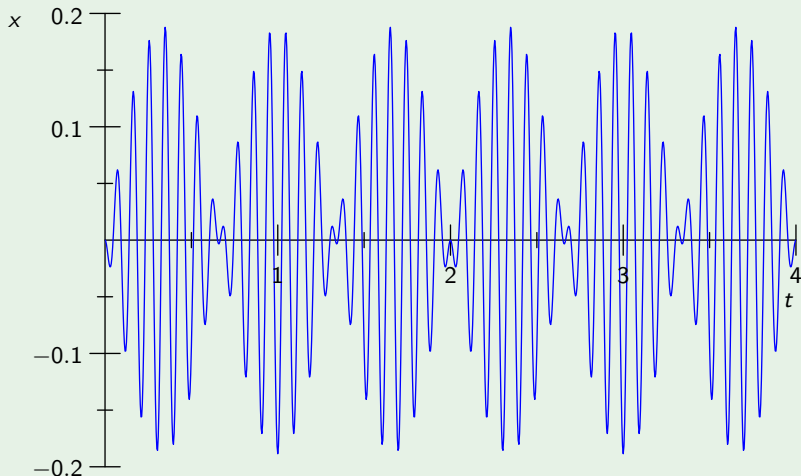
So, where c_1 and c_2 are determined by initial conditions, we have

$$\begin{aligned} x(t) &= c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \\ &= C \cos(\omega_0 t - \delta) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \end{aligned}$$

where $C = \sqrt{c_1^2 + c_2^2}$ and $\tan(\delta) = \frac{c_2}{c_1}$.

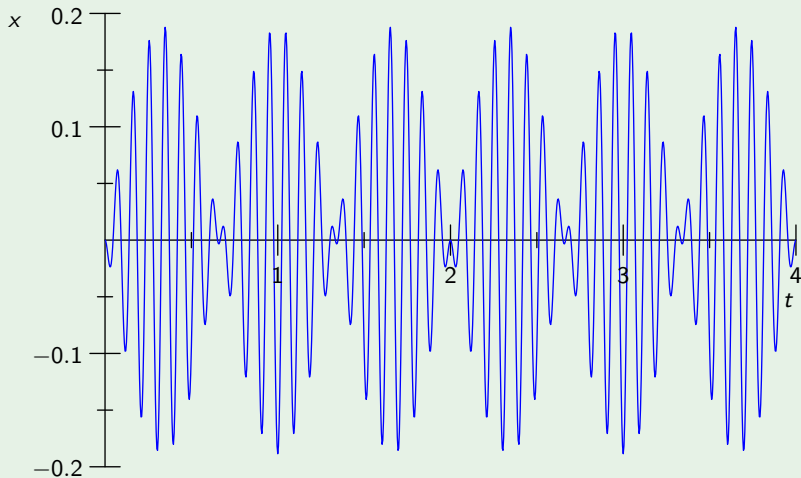
Example 2

An example of $\omega_f \neq \omega_0$ being different, but close, is



Example 2

An example of $\omega_f \neq \omega_0$ being different, but close, is



The regular periodic patterns are called **beats**.

Resonance ($\omega_f = \omega_0$)

This means we want to look at

$$x_p = At \cos(\omega_0 t) + Bt \sin(\omega_0 t)$$

Resonance ($\omega_f = \omega_0$)

This means we want to look at

$$x_p = At \cos(\omega_0 t) + Bt \sin(\omega_0 t)$$

and so, we find that

$$A = 0 \quad \text{and} \quad B = \frac{F_0}{2m\omega_0}$$

Resonance ($\omega_f = \omega_0$)

This means we want to look at

$$x_p = At \cos(\omega_0 t) + Bt \sin(\omega_0 t)$$

and so, we find that

$$A = 0 \quad \text{and} \quad B = \frac{F_0}{2m\omega_0}$$

So, where c_1 and c_2 are determined by initial conditions, we have

$$x(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + \frac{F_0}{2m\omega_0} t \sin(\omega_0 t)$$

Example 3

We can see an example of resonance

$$\ddot{x} + x = 2 \cos(t), \quad x(0) = 0, \quad \dot{x}(0) = 0$$

Example 3

We can see an example of resonance

$$\ddot{x} + x = 2 \cos(t), \quad x(0) = 0, \quad \dot{x}(0) = 0$$

Since $\frac{F_0}{2m\omega_0} = \frac{2}{2 \cdot 1 \cdot 1} = 1$, we have $x_p = t \sin(t)$.

Example 3

We can see an example of resonance

$$\ddot{x} + x = 2 \cos(t), \quad x(0) = 0, \quad \dot{x}(0) = 0$$

Since $\frac{F_0}{2m\omega_0} = \frac{2}{2 \cdot 1 \cdot 1} = 1$, we have $x_p = t \sin(t)$.

So, given that $x(0) = 0$ and $\dot{x}(0) = 0$ we have $x = t \sin(t)$.

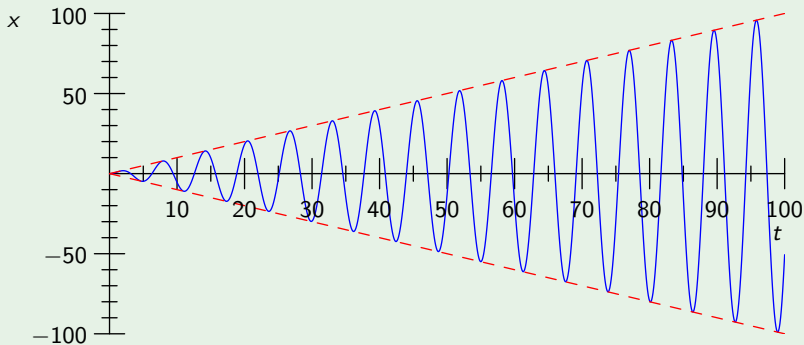
Example 3

We can see an example of resonance

$$\ddot{x} + x = 2 \cos(t), \quad x(0) = 0, \quad \dot{x}(0) = 0$$

Since $\frac{F_0}{2m\omega_0} = \frac{2}{2 \cdot 1 \cdot 1} = 1$, we have $x_p = t \sin(t)$.

So, given that $x(0) = 0$ and $\dot{x}(0) = 0$ we have $x = t \sin(t)$.



Beats

As we saw earlier when ω_f is close to ω_0 , though still not equal, there is clearly periodic behavior.

Beats

As we saw earlier when ω_f is close to ω_0 , though still not equal, there is clearly periodic behavior.

This periodic pattern has a frequency, called the **beat frequency**, that is lower than both ω_f and ω_0 .

Beats

As we saw earlier when ω_f is close to ω_0 , though still not equal, there is clearly periodic behavior.

This periodic pattern has a frequency, called the **beat frequency**, that is lower than both ω_f and ω_0 .

If the system is initially at rest ($x(0) = 0$ and $\dot{x}(0) = 0$) then the solution is

$$x(t) = -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

Beats

As we saw earlier when ω_f is close to ω_0 , though still not equal, there is clearly periodic behavior.

This periodic pattern has a frequency, called the **beat frequency**, that is lower than both ω_f and ω_0 .

If the system is initially at rest ($x(0) = 0$ and $\dot{x}(0) = 0$) then the solution is

$$x(t) = -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

So, we can simplify using the trigonometric identity

$$\cos(u) - \cos(v) = -2 \sin\left(\frac{u-v}{2}\right) \sin\left(\frac{u+v}{2}\right)$$

Beats

$$x(t) = -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

Beats

$$\begin{aligned}x(t) &= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \\&= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} (\cos(\omega_0 t) - \cos(\omega_f t))\end{aligned}$$

Beats

$$\begin{aligned}x(t) &= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \\&= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} (\cos(\omega_0 t) - \cos(\omega_f t)) \\&= \frac{2F_0}{m(\omega_0^2 - \omega_f^2)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right) \sin\left(\frac{\omega_0 + \omega_f}{2}t\right)\end{aligned}$$

Beats

$$\begin{aligned}x(t) &= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \\&= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} (\cos(\omega_0 t) - \cos(\omega_f t)) \\&= \frac{2F_0}{m(\omega_0^2 - \omega_f^2)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right) \sin\left(\frac{\omega_0 + \omega_f}{2}t\right)\end{aligned}$$

When the difference between ω_f and ω_0 is small, then $\omega_0 - \omega_f$ is much smaller than $\omega_0 + \omega_f$.

Beats

$$\begin{aligned}x(t) &= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \\&= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} (\cos(\omega_0 t) - \cos(\omega_f t)) \\&= \frac{2F_0}{m(\omega_0^2 - \omega_f^2)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right) \sin\left(\frac{\omega_0 + \omega_f}{2}t\right)\end{aligned}$$

When the difference between ω_f and ω_0 is small, then $\omega_0 - \omega_f$ is much smaller than $\omega_0 + \omega_f$.

Thus $\sin\left(\frac{\omega_0 - \omega_f}{2}t\right)$ oscillates much slower than $\sin\left(\frac{\omega_0 + \omega_f}{2}t\right)$.

Beats

$$\begin{aligned}x(t) &= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t) \\&= -\frac{F_0}{m(\omega_0^2 - \omega_f^2)} (\cos(\omega_0 t) - \cos(\omega_f t)) \\&= \frac{2F_0}{m(\omega_0^2 - \omega_f^2)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right) \sin\left(\frac{\omega_0 + \omega_f}{2}t\right)\end{aligned}$$

When the difference between ω_f and ω_0 is small, then $\omega_0 - \omega_f$ is much smaller than $\omega_0 + \omega_f$.

Thus $\sin\left(\frac{\omega_0 - \omega_f}{2}t\right)$ oscillates much slower than $\sin\left(\frac{\omega_0 + \omega_f}{2}t\right)$.

The two curves

$$\pm \frac{2F_0}{m(\omega_0^2 - \omega_f^2)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right)$$

form an envelope of the more rapid oscillation and is called the **sinusoidal amplitude**.

Solutions to the Undamped Forced Oscillator ($\omega_f \neq \omega_0$)

For

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

Solutions to the Undamped Forced Oscillator ($\omega_f \neq \omega_0$)

For

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

The general solution is

$$x(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

where c_1 and c_2 are determined by initial conditions and $\omega_0 = \sqrt{\frac{k}{m}}$.

Solutions to the Undamped Forced Oscillator ($\omega_f \neq \omega_0$)

For

$$m\ddot{x} + kx = F_0 \cos(\omega_f t)$$

The general solution is

$$x(t) = c_1 \cos(\omega_0 t) + c_2 \sin(\omega_0 t) + \frac{F_0}{m(\omega_0^2 - \omega_f^2)} \cos(\omega_f t)$$

where c_1 and c_2 are determined by initial conditions and $\omega_0 = \sqrt{\frac{k}{m}}$.

If the system starts from rest ($x(0) = 0$ and $\dot{x}(0) = 0$), the solution can be written as

$$x(t) = \underbrace{\frac{2F_0}{m(\omega_0^2 - \omega_f^2)} \sin\left(\frac{\omega_0 - \omega_f}{2}t\right)}_{\text{sinusoidal amplitude}} \underbrace{\sin\left(\frac{\omega_0 + \omega_f}{2}t\right)}_{\text{rapid oscillation within beats}}$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}}$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}} \quad \rightarrow \quad \omega_0^2 = \frac{k}{m}$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}} \quad \rightarrow \quad \omega_0^2 = \frac{k}{m} \quad \rightarrow \quad k = m\omega_0^2$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}} \rightarrow \omega_0^2 = \frac{k}{m} \rightarrow k = m\omega_0^2 = 484\pi^2$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}} \rightarrow \omega_0^2 = \frac{k}{m} \rightarrow k = m\omega_0^2 = 484\pi^2$$

Since

$$\frac{\omega_0 - \omega_f}{2} = \pi \quad \text{and} \quad \frac{\omega_0 + \omega_f}{2} = 21\pi$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}} \rightarrow \omega_0^2 = \frac{k}{m} \rightarrow k = m\omega_0^2 = 484\pi^2$$

Since

$$\frac{\omega_0 - \omega_f}{2} = \pi \quad \text{and} \quad \frac{\omega_0 + \omega_f}{2} = 21\pi$$

The solution is

$$x(t) = 1 \cdot \sin(\pi t) \sin(21\pi t)$$

Example 4

Given an oscillator with $\omega_0 = 22\pi$, $\omega_f = 20\pi$, $m = 1$, and $F_0 = 42\pi^2$.

We can then calculate k :

$$\omega_0 = \sqrt{\frac{k}{m}} \rightarrow \omega_0^2 = \frac{k}{m} \rightarrow k = m\omega_0^2 = 484\pi^2$$

Since

$$\frac{\omega_0 - \omega_f}{2} = \pi \quad \text{and} \quad \frac{\omega_0 + \omega_f}{2} = 21\pi$$

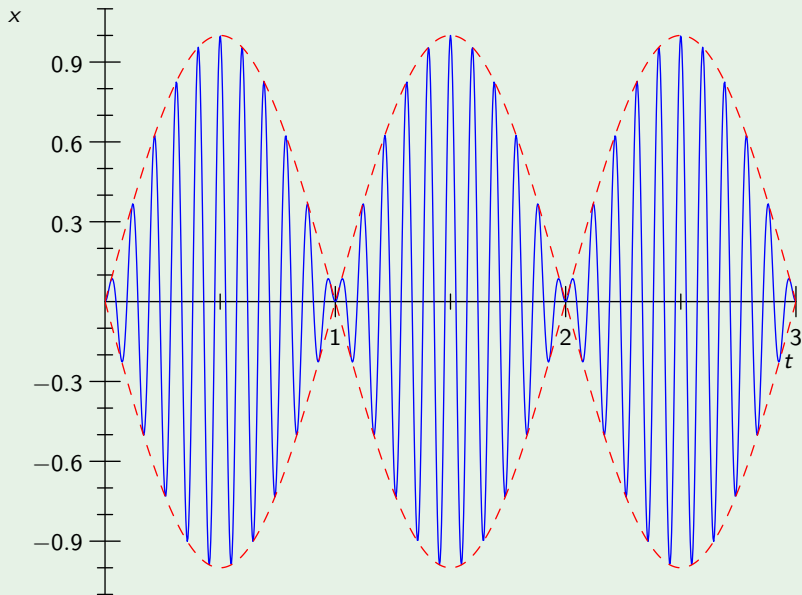
The solution is

$$x(t) = 1 \cdot \sin(\pi t) \sin(21\pi t)$$

Thus, envelope curves are

$$y = \pm 1 \cdot \sin(\pi t)$$

Example 4



Example 5

Let us consider the damped forced oscillator

$$\ddot{x} + 4\dot{x} + 5x = 10 \cos(3t), \quad x(0) = \dot{x}(0) = 0$$

Example 5

Let us consider the damped forced oscillator

$$\ddot{x} + 4\dot{x} + 5x = 10 \cos(3t), \quad x(0) = \dot{x}(0) = 0$$

First, we need to find the solution to the associated homogeneous equation.

Example 5

Let us consider the damped forced oscillator

$$\ddot{x} + 4\dot{x} + 5x = 10 \cos(3t), \quad x(0) = \dot{x}(0) = 0$$

First, we need to find the solution to the associated homogeneous equation.

The characteristic equation $r^2 + 4r + 5 = 0$ has solutions $r = -2 \pm i$, so we have

$$x_h = e^{-2t} (c_1 \cos(t) + c_2 \sin(t))$$

Example 5

Let us consider the damped forced oscillator

$$\ddot{x} + 4\dot{x} + 5x = 10 \cos(3t), \quad x(0) = \dot{x}(0) = 0$$

First, we need to find the solution to the associated homogeneous equation.

The characteristic equation $r^2 + 4r + 5 = 0$ has solutions $r = -2 \pm i$, so we have

$$x_h = e^{-2t} (c_1 \cos(t) + c_2 \sin(t))$$

Next we need to find a particular solution.

Example 5

Let us consider the damped forced oscillator

$$\ddot{x} + 4\dot{x} + 5x = 10 \cos(3t), \quad x(0) = \dot{x}(0) = 0$$

First, we need to find the solution to the associated homogeneous equation.

The characteristic equation $r^2 + 4r + 5 = 0$ has solutions $r = -2 \pm i$, so we have

$$x_h = e^{-2t} (c_1 \cos(t) + c_2 \sin(t))$$

Next we need to find a particular solution.

We can use the method of undetermined coefficients:

$$\begin{aligned} x(t) &= A \cos(3t) + B \sin(3t) \\ \dot{x}(t) &= -3A \sin(3t) + 3B \cos(3t) \\ \ddot{x}(t) &= -9A \cos(3t) - 9B \sin(3t) \end{aligned}$$

Example 5

Substituting into the DE gives

$$\begin{aligned} &(-9A \cos(3t) - 9B \sin(3t)) \\ &\quad + 4(-3A \sin(3t) + 3B \cos(3t)) \\ &\quad + 5(A \cos(3t) + B \sin(3t)) = 10 \cos(3t) \end{aligned}$$

Example 5

Substituting into the DE gives

$$\begin{aligned} &(-9A \cos(3t) - 9B \sin(3t)) \\ &\quad + 4(-3A \sin(3t) + 3B \cos(3t)) \\ &\quad + 5(A \cos(3t) + B \sin(3t)) = 10 \cos(3t) \end{aligned}$$

Which gives the system

$$\begin{aligned} -9A + 12B + 5A &= 10 \\ -9B - 12A + 5B &= 0 \end{aligned}$$

Example 5

Substituting into the DE gives

$$\begin{aligned} &(-9A \cos(3t) - 9B \sin(3t)) \\ &\quad + 4(-3A \sin(3t) + 3B \cos(3t)) \\ &\quad + 5(A \cos(3t) + B \sin(3t)) = 10 \cos(3t) \end{aligned}$$

Which gives the system

$$\begin{aligned} -9A + 12B + 5A &= 10 \\ -9B - 12A + 5B &= 0 \end{aligned}$$

Which has solution

$$A = -\frac{1}{4} \quad \text{and} \quad B = \frac{3}{4}$$

Example 5

Substituting into the DE gives

$$\begin{aligned} &(-9A \cos(3t) - 9B \sin(3t)) \\ &\quad + 4(-3A \sin(3t) + 3B \cos(3t)) \\ &\quad + 5(A \cos(3t) + B \sin(3t)) = 10 \cos(3t) \end{aligned}$$

Which gives the system

$$\begin{aligned} -9A + 12B + 5A &= 10 \\ -9B - 12A + 5B &= 0 \end{aligned}$$

Which has solution

$$A = -\frac{1}{4} \quad \text{and} \quad B = \frac{3}{4}$$

Thus,

$$x_p = -\frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

Example 5

The general solution is

$$x = e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

Example 5

The general solution is

$$x = e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

To solve the IVP, we need to calculate

$$\begin{aligned} \dot{x} = & -2e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) \\ & + e^{-2t} (-c_1 \sin(t) + c_2 \cos(t)) \\ & + \frac{1}{4} \sin(3t) - \frac{3}{4} \cos(3t) \end{aligned}$$

Example 5

The general solution is

$$x = e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

To solve the IVP, we need to calculate

$$\begin{aligned}\dot{x} &= -2e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) \\ &\quad + e^{-2t} (-c_1 \sin(t) + c_2 \cos(t)) \\ &\quad + \frac{1}{4} \sin(3t) - \frac{3}{4} \cos(3t)\end{aligned}$$

Next, we need to plug in the initial conditions

$$x(0) = 0$$

$$\dot{x}(0) = 0$$

Example 5

The general solution is

$$x = e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

To solve the IVP, we need to calculate

$$\begin{aligned}\dot{x} &= -2e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) \\ &\quad + e^{-2t} (-c_1 \sin(t) + c_2 \cos(t)) \\ &\quad + \frac{1}{4} \sin(3t) - \frac{3}{4} \cos(3t)\end{aligned}$$

Next, we need to plug in the initial conditions

$$\begin{aligned}x(0) = 0 &\Rightarrow c_1 - \frac{1}{4} = 0 \\ \dot{x}(0) = 0 &\Rightarrow c_2 + \frac{7}{4} = 0\end{aligned}$$

Example 5

The general solution is

$$x = e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

To solve the IVP, we need to calculate

$$\begin{aligned}\dot{x} &= -2e^{-2t} (c_1 \cos(t) + c_2 \sin(t)) \\ &\quad + e^{-2t} (-c_1 \sin(t) + c_2 \cos(t)) \\ &\quad + \frac{1}{4} \sin(3t) - \frac{3}{4} \cos(3t)\end{aligned}$$

Next, we need to plug in the initial conditions

$$\begin{aligned}x(0) = 0 &\Rightarrow c_1 - \frac{1}{4} = 0 &\Rightarrow c_1 = \frac{1}{4} \\ \dot{x}(0) = 0 &\Rightarrow c_2 + \frac{7}{4} = 0 &\Rightarrow c_2 = -\frac{7}{4}\end{aligned}$$

Example 5

The solution to the IVP is

$$x = e^{-2t} \left(\frac{1}{4} \cos(t) - \frac{7}{4} \sin(t) \right) - \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)$$

Example 5

The solution to the IVP is

$$x = e^{-2t} \underbrace{\left(\frac{1}{4} \cos(t) - \frac{7}{4} \sin(t) \right)}_{\text{Transient}} \underbrace{- \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)}_{\text{Steady-State}}$$

We call x_h **transient**, because for $b > 0$ the solution tends towards zero.

Example 5

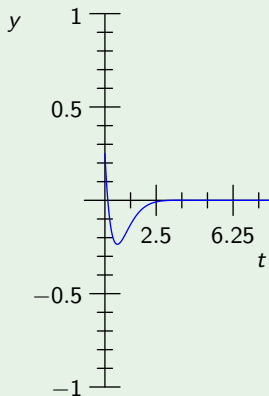
The solution to the IVP is

$$x = e^{-2t} \underbrace{\left(\frac{1}{4} \cos(t) - \frac{7}{4} \sin(t) \right)}_{\text{Transient}} \underbrace{- \frac{1}{4} \cos(3t) + \frac{3}{4} \sin(3t)}_{\text{Steady-State}}$$

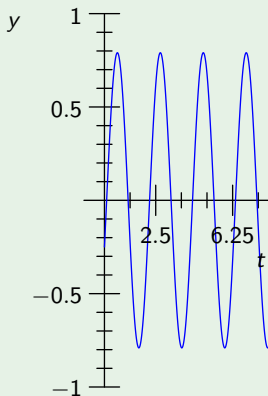
We call x_h **transient**, because for $b > 0$ the solution tends towards zero.

The particular solution x_p may either be constant or a periodic **steady-state** solution.

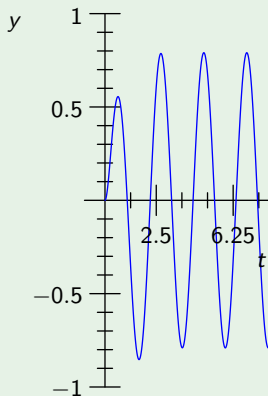
Example 5



(a) Transient Solution



(b) Steady-State



(c) IVP Solution

Particular Solution x_p of a Damped Mass-Spring System

The damped mass-spring system

$$m\ddot{x} + b\dot{x} + kx = F_0 \cos(\omega_f t)$$

has particular solution

$$x_p = A \cos(\omega_f t) + B \sin(\omega_f t)$$

with

$$A = \frac{m(\omega_0^2 - \omega_f^2) F_0}{m^2(\omega_0^2 - \omega_f^2)^2 + (b\omega_f)^2} \quad \text{and} \quad B = \frac{b\omega_f F_0}{m^2(\omega_0^2 - \omega_f^2)^2 + (b\omega_f)^2}$$

with natural circular frequency $\omega_0 = \sqrt{\frac{k}{m}}$.

Particular Solution x_p of a Damped Mass-Spring System

The damped mass-spring system

$$m\ddot{x} + b\dot{x} + kx = F_0 \cos(\omega_f t)$$

has particular solution

$$x_p = A \cos(\omega_f t) + B \sin(\omega_f t)$$

with

$$A = \frac{m(\omega_0^2 - \omega_f^2) F_0}{m^2(\omega_0^2 - \omega_f^2)^2 + (b\omega_f)^2} \quad \text{and} \quad B = \frac{b\omega_f F_0}{m^2(\omega_0^2 - \omega_f^2)^2 + (b\omega_f)^2}$$

with natural circular frequency $\omega_0 = \sqrt{\frac{k}{m}}$.

Note

You will verify this in the homework.