

• 1 i) Show $\tilde{F} = \langle \tilde{\psi} | \hat{F} | \tilde{\psi} \rangle = \text{Tr}(\underline{f} \underline{P})$

$$\hat{F} = \hat{f}(x) + \hat{f}(y).$$

$\tilde{\psi}$ is our Slater determinant of basis functions.

Suppose we perform a unitary transformation U , $\chi'_a = \sum_b \chi_b U_{ba}$.

Since our $\tilde{\psi}$ is defined to be $\det(\underline{A})$, $\underline{A} = \begin{pmatrix} \chi_1(x) & \chi_2(x) \\ \chi_1(y) & \chi_2(y) \end{pmatrix}$, then $\underline{A}' = \underline{A} U$ will have $\tilde{\psi}' = \det(\underline{A}' U) = \det(\underline{A}) \det(U)$. $\therefore U$ is unitary, the new antisymmetric wavefunction differs just by a rotational factor, of which vanishes in matrix elements.

\therefore It is always possible to find such U that diagonalises ε (Szabo 122), then we have $\langle \chi'_a | \hat{f} | \chi'_b \rangle = \varepsilon_b \delta_{ab}$. We can just drop the prime since it's arbitrary:

$$\Rightarrow \langle \tilde{\psi} | \hat{F} | \tilde{\psi} \rangle = \sum_{ab} \langle \chi_a | \hat{f} | \chi_b \rangle \text{ by orthonormality, where the little } \hat{f} \text{ is an one electron operator } \chi_a = \sum_{v=1}^k c_{va} \phi_v(\tau), \hat{f} | \chi_a \rangle = \varepsilon_a | \chi_a \rangle.$$

$$\begin{aligned} \text{Then we have: } & \sum_a \langle \chi_a | \hat{f} | \chi_a \rangle \\ &= \sum_a \int d\tau \chi_a^*(\tau) \hat{f}(\tau) \chi_a(\tau) \end{aligned}$$

For $\chi_a = \sum_v^k c_{va} \phi_v(\tau)$, we get:

$$\begin{aligned} & \sum_a \left(\int d\tau \sum_v^k \sum_v^k c_{va}^* c_{va} \phi_{va}^*(\tau) \phi_{va}(\tau) \right) \\ &= \sum_a \left(\sum_{v,v}^{k,k} c_{va}^* c_{va} \int d\tau \phi_{va}^*(\tau) \phi_{va}(\tau) \right) \\ &= \sum_a \left(\sum_{v,v}^{k,k} c_{va} c_{va}^* F_{vv} \right) \\ &= \sum_{v,v}^{k,k} F_{vv} \sum_a c_{va} c_{va}^* \end{aligned}$$

$$= \text{Tr}(\underline{f} \underline{P}), \quad \text{where } \underline{f}_{vv} = F_{vv}, \quad \underline{P}_{vv} = \sum_a c_{va} c_{va}^*.$$