

## General correlator Probe — with pulse sequence

From above discussion we know

$$\langle A(t) \rangle = \sum_{n=0}^{\infty} \langle A(t) \rangle_n \equiv \sum_{n=0}^{\infty} (-i)^n \int_0^t ds_1 \dots ds_n \operatorname{tr} \left\{ \rho_0 [A_1(t), [V_1(t_1), \dots, [V_2(t_{n-1}), V_2(t_n)]]] \right\}$$

$$\text{where } \langle A(t) \rangle_0 \equiv \operatorname{tr} \{ \rho_0 A_1(t) \}$$

$$\begin{aligned} \text{def } \Delta \langle A(t) \rangle &\equiv \langle A(t) \rangle - \langle A(t) \rangle_0 \quad \text{we get required dynamics as correlator} \\ &= \sum_{n=1}^{\infty} (-i)^n \int_0^t ds_1 \dots ds_n \left\langle [A_2(t), [V_2(t_1), \dots, [V_2(t_{n-1}), V_2(t_n)]]] \right\rangle \end{aligned}$$

$$\text{Say the driving term is } V(t) = \sum_{\mu} f_{\mu}(t) \hat{B}^{\mu}(t)$$

where  $\mu$  is some "mode" index, say  $\mu = \{x, y, z\}$   $f_{\mu}(t)$  is amplitude  
 Fourier and  $\hat{B}^{\mu}(t)$  is time-dependent driving operator

For example, simply continuous mag field along  $z$ -direction on spin system

$$\text{is } V(t) = \cos(t) \cdot \epsilon \sum_i \sigma_z^{(i)}$$

Next I focus on, and promote, is pulse sequence at discrete times  
 say along 'direction gives



In such a case  $\hat{B}^x(t) \equiv \sum_i Q_z^{(i)}$  a collective  
 total  $z$ -magnetization, But it does not have  
 to be so. For example pulses at different discrete time can be  
 magnetization along different direction; or some times the operator is  
 extending ( $zz$ -type) while others are local as  $\sum_i Q_z^{(i)}$

$$V(t) = \sum_{j=1}^{N\#} f_x(t - \tau_j) \sum_i Q_z^{(i)}$$

We plug the general expr of driving into the expansion-series,  
 separating the "modulation" ( $f$ ) part and "operator" part ( $B(t)$ ),  
 This gives rise to the following expr,

$$\Delta \langle A(t) \rangle = \sum_{n=1}^{\infty} (-i)^n \sum_{\mu_1 \dots \mu_n} \int_0^t dt_n \underbrace{\prod_{j=1}^n f_{\mu_j}(t_j)}_{\text{filter}} \underbrace{\left( [A_I(t), [B_I^{\mu_1}(t_1), \dots [B_I^{\mu_{n-1}}(t_{n-1}), B_I^{\mu_n}(t_n)]] \right]}_{\text{Green's function}} \\ = \sum_{n=1}^{\infty} (-i)^n \underbrace{\mu}_{\text{time-order}} \int_0^t dt_n \underbrace{\prod_{j=1}^n f_{\mu_j}(t_j)}_{\text{filter}} \underbrace{\chi_{(n)}^{(\mu_1 \dots \mu_n)}(t; t_1 \dots t_n)}_{\text{Green's function}} \quad \swarrow \text{the order}$$

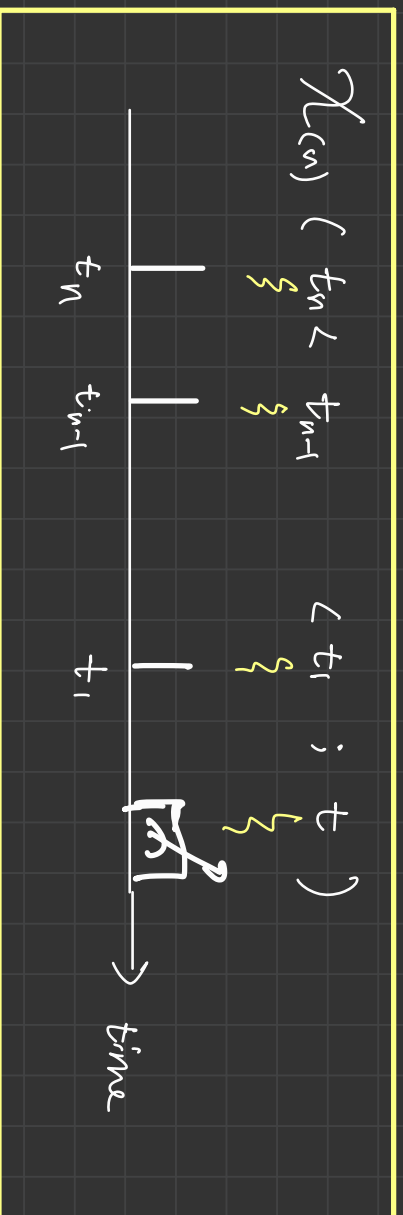
Notice  $t > t_1 > t_2 > \dots > t_n$  due to the causality.

If  $|\mu| \equiv 1$  for whole pulse sequence (for example,  $\mu \equiv z$  as in one-dimension magnetization case), then  $\mu$  can be dropped. We will focus on this case for simplicity.

Question: How to probe Green's function  $\chi_{(n)}(t; t_1 > \dots > t_n)$  for a given  $(t; t_1 > \dots > t_n)$ ? And to separate it from lower order function  $\chi_{(n-1)}(t; t_1 > \dots > t_{n-1})$ ?

Answer: Set the pulse timing to be resonant with  $(t_1 \dots t_n)$

That is



So to probe  $X(n)(t; \vec{t_{en}})$  it's necessary & sufficient to set pulse-sequence at interested time  $\vec{t_{en}}$ .

Under such sequence  $X(n)(t; \vec{t_{en}})$  contribute to  $\Delta\langle A(t) \rangle$  other "times" of  $X(n)$  And any  $X(n|m)$  will not  $\neq$  Their corresponding filter  $\equiv 0$ .

But sum over order  $X(n|n)$  will contribute

What are they (2) For target  $\chi_{(n)}(t; t, \dots, t_n)$

One need to separate  $\chi_{(n)}(t; \tilde{t}_1, \dots, \tilde{t}_n)$  where

$$\{\tilde{t}_1, \dots, \tilde{t}_n\} \subseteq \{t_1, t_2, \dots, t_n\}$$

Thus the central pulsed spectroscopy is

use few pulses (start from one) to estimate  $\chi_{(1)}$  first.

And adding more pulses to probe higher terms.

To probe 'k'-th order, one need k-pulse sequence.

One should vary  $\{t_1, \dots, t_n\}$  for fixed 'n' to get different temporal points @  $\chi_{(n)}$  to get more fine into motion.

Now to think:  $f(t - \tau) \Rightarrow f(t - \tau - \delta_\tau) \cdot (1 + \delta_\tau)$

Control-error mitigation

timing error

amplitude error

