

General correlator Probe – with pulse sequence

From above discussion we know

$$\langle A(t) \rangle = \sum_{n=0}^{\infty} \langle A(t) \rangle_n \equiv \sum_{n=0}^{\infty} (-i)^n \int_0^t d\tau t_{\text{un}} \text{tr} \left\{ \rho_0 [A_I(t), [V_1(t_1), \dots, [V_L(t_{n-1}), V_L(t_n)]]] \right\}$$

where $\langle A(t) \rangle_0 \equiv \text{tr} \{ \rho_0 A_I(t) \}$

$$\begin{aligned} \Delta \langle A(t) \rangle &= \langle A(t) \rangle - \langle A(t) \rangle_0 \quad \text{we get response dynamics as corr} \\ &= \sum_{n=1}^{\infty} (-i)^n \int_0^t d\tau t_{\text{un}} \left\langle \left[[A_I(t), [V_1(t_1), \dots, [V_L(t_{n-1}), V_L(t_n)]]] \right] \right\rangle \end{aligned}$$

Say the driving term is $V(t) = \sum_{\mu} f_{\mu}(t) \hat{B}^{\mu}(t)$

where μ is some "mode" index, say $\mu = \{x, y, z\}$ $|f_{\mu}|$ is amplitude function and $\hat{B}^{\mu}(t)$ is time-dependent driving operator

For example, apply curving mag field along z -direction on spin system

$$iS V(t) = \cos(\omega t) \cdot \omega \sum_i \sigma_z^{(i)}$$

What I focus on, and promote, is pulse sequence at discrete times say along ' X' direction gives

$$V(t) = \sum_{j=1}^{N^{\#}} f_X(t - \tau_j) \sum_i O_z^{(i)}$$



In such a case $\hat{B}^X(t) \equiv \sum_i O_z^{(i)}$ a collective total \hat{x} -magnetization, but it does not have to be so. for example pulses at different discrete time can be magnetization along different direction; or some times the operator is extending (Z type) while others are local as $\sum_i O_z^{(i)}$

We plug the general expr of driving into the expansion-series, separating the "modulation" (f) part and "operator" part ($B(t)$), this gives rise to the following expr,

$$\Delta \langle A(t) \rangle = \sum_{n=1}^{\infty} (-i)^n \sum_{\mu_1 \dots \mu_n} \int_0^t d\tau_1 t^{(\alpha)} \underbrace{\int_0^{\tau_1} f_{\mu_1}(t_j)}_{\text{filter}} \left\langle \left[A_1(t), \overline{B_x^{\mu_1}(t_1)}, \dots [B_z^{\mu_{n-1}}(t_{n-1}), B_z^{\mu_n}(t_n)] \right] \right\rangle$$

$$= \sum_{n=1}^{\infty} (-i)^n \underbrace{\mu}_{\text{Green's function}} \int_0^t d\tau_1 t^{(\alpha)} \underbrace{\sum_{j=1}^n f_{\mu_j}(\tau_j)}_{\text{time-order}} \underbrace{\mathcal{N}_{(\alpha)}^{\mu_1 \dots \mu_n}(t; t_1 > \dots > t_n)}_{\text{time-order}} \text{ & the order}$$

Notice $t > t_1 > t_2 > \dots$ t_n due to the consulating.

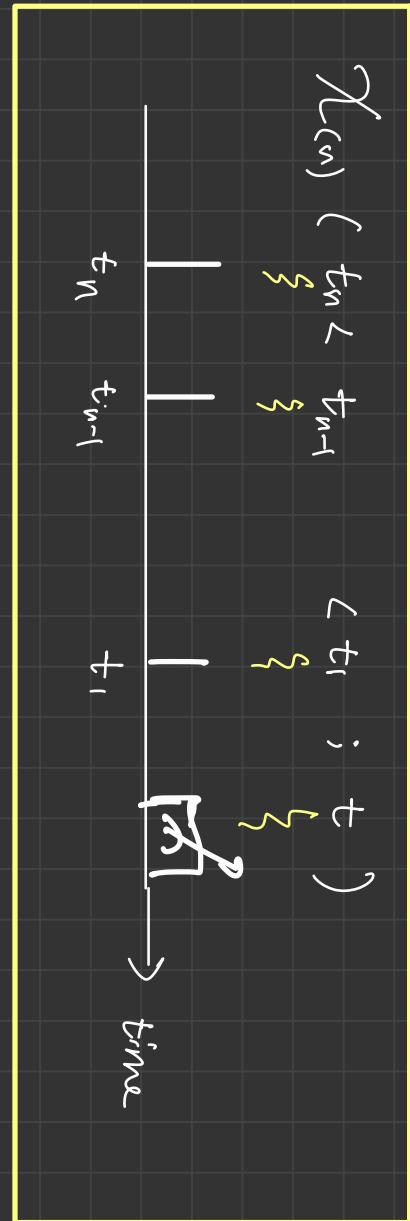
If $|\mu| \equiv 1$ for whole pulse sequence (for example, $\mu \equiv 2$ as in one-direction magnetization case), then μ can be dropped. We will focus on this case for simplifying.

Question:

How to probe Green's function $\mathcal{N}_{(\alpha)}(t; t_1 > \dots > t_n)$ for a given $(t; t_1 > \dots > t_n)$? And to separate it from lower order function $\mathcal{N}_{(n+1)}(t; t_1 > \dots > t_n)$?

Answer: Set the pulse timing to be resonant with $(t_1 \dots t_n)$

That is



So to prove $X(n) (t; t_n)$ it's necessary & sufficient to test pulse-sequence at interested time t_n .

Under such sequence $X(n) (t; t_n)$ contribute to $\Delta \langle f(t) \rangle$ other "times" of $X(n)$. And any $X(n)$ will not. Their corresponding filter $\equiv 0$.

But some lower order $X(n')$ will contribute

What are they (?) for target

$$\chi_{(n)}(t_1, t_2, \dots, t_n)$$

One need to separate

$$\chi_{(n \leq n)}(t_1, t_2, \dots, t_n)$$

where

$$\{\tilde{t}_1, \dots, \tilde{t}_n\} \subseteq \{t_1, t_2, \dots, t_n\}$$

thus the

General

pulse

spreading is

use few pulses (start from one) to estimate $\chi_{(1)}$ first.

And adding more pulses to probe higher terms.

To probe " k "-th order, one need k -pulse sequence.

One should vary $\{t_1, \dots, t_n\}$ for fixed " n " to get different temporal points @ $\chi_{(n)}$ to get more fine information.

Now to think:
 $f(t - \tau) \Rightarrow f(t - \tau - \delta_t) \cdot (\text{if } \delta_t)$
control error mitigation
timing error

