

$$\begin{aligned}
& \frac{1}{3} \frac{t}{2} \cdot \mathcal{A}_{abi} \mathcal{A}^{abi} - \frac{2}{3} \frac{t}{2} \cdot \mathcal{A}_{aib} \mathcal{A}^{abi} + \left(-\frac{r}{2} + \frac{r}{5} \right) \partial_b \mathcal{A}_i^j \partial^i \mathcal{A}^{ab}_a + \\
& \left(-\frac{r}{2} - \frac{r}{5} \right) \partial_b \mathcal{A}_b^j \partial^i \mathcal{A}^{ab}_a - \frac{2}{3} \frac{t}{2} \cdot \mathcal{A}_{abi} \partial^i f^{ab} + \frac{2}{3} \frac{t}{2} \cdot \mathcal{A}_{aib} \partial^i f^{ab} - \frac{2}{3} \frac{t}{2} \cdot \mathcal{A}_{bia} \partial^i f^{ab} + \\
& \frac{1}{3} \frac{t}{2} \cdot \partial_a f_{bi} \partial^i f^{ab} - \frac{1}{6} \frac{t}{2} \cdot \partial_a f_{ib} \partial^i f^{ab} - \frac{1}{6} \frac{t}{2} \cdot \partial_b f_{ai} \partial^i f^{ab} + \frac{1}{6} \frac{t}{2} \cdot \partial_i f_{ab} \partial^i f^{ab} - \frac{1}{6} \frac{t}{2} \cdot \partial_i f_{ba} \partial^i f^{ab} + \\
& \left(-\frac{r}{2} - \frac{r}{5} \right) \partial_a \mathcal{A}^{abi} \partial_j \mathcal{A}_b^j + \left(\frac{r}{3} + 2 \frac{r}{5} \right) \partial^i \mathcal{A}^{ab}_a \partial_j \mathcal{A}_b^j + \left(-\frac{r}{2} + \frac{r}{5} \right) \partial_a \mathcal{A}^{abi} \partial_j \mathcal{A}_i^j + \\
& \left(\frac{r}{3} - 2 \frac{r}{5} \right) \partial^i \mathcal{A}^{ab}_a \partial_j \mathcal{A}_i^j + \frac{4}{3} \frac{r}{2} \cdot \partial_b \mathcal{A}_{aib} \partial^i \mathcal{A}^{abi} - \frac{2}{3} \frac{r}{2} \cdot \partial_b \mathcal{A}_{aji} \partial^i \mathcal{A}^{abi} + \\
& \frac{2}{3} \left(\frac{r}{2} - 6 \frac{r}{3} \right) \partial_b \mathcal{A}_{ija} \partial^i \mathcal{A}^{abi} - \frac{1}{3} \frac{r}{2} \cdot \partial_i \mathcal{A}_{abj} \partial^i \mathcal{A}^{abi} + \frac{1}{3} \frac{r}{2} \cdot \partial_j \mathcal{A}_{abi} \partial^i \mathcal{A}^{abi} - \frac{2}{3} \frac{r}{2} \cdot \partial_j \mathcal{A}_{aib} \partial^i \mathcal{A}^{abi}
\end{aligned} \tag{1}$$

Now we pass this theory into the PSALTer package, which computes the particle spectrum:

The (possibly singular) a -matrices associated

with the Lagrangian, as defined below Eq. (18) of arXiv:1812.02675:

Matrix for spin-0 sector:

$$\begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & k^2 \frac{r}{2} + \frac{t}{2} \end{pmatrix}$$

Matrix for spin-1 sector:

$$\begin{pmatrix} k^2 \left(2 \frac{r}{3} + \frac{r}{5} \right) + \frac{2t}{3} & \frac{\sqrt{2} t}{3} & -\frac{1}{3} i \sqrt{2} k \frac{t}{2} & 0 & 0 & 0 & 0 \\ \frac{\sqrt{2} t}{3} & \frac{t}{3} & -\frac{1}{3} i k \frac{t}{2} & 0 & 0 & 0 & 0 \\ \frac{1}{3} i \sqrt{2} k \frac{t}{2} & \frac{i k t}{3} & \frac{k^2 t}{3} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{2} k^2 \left(\frac{r}{3} + 2 \frac{r}{5} \right) & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

Matrix for spin-2 sector:

$$\begin{pmatrix} -\frac{3k^2 r}{2} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Gauge constraints on source currents:

$$\frac{0}{2} \cdot \tau^\dagger = 0$$

$$\frac{0}{2} \cdot \tau^\parallel = 0$$

$$\frac{0}{2} \cdot \sigma^\parallel = 0$$

$$\frac{1}{2} \cdot \tau^\perp = 0$$