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# PSALTer Calibration for the WGTE cases based on Lin's draft

### About xPlain and formatting

Welcome to the calibration file for the PSALTer package. Commentary is provided in this green text throughout by virtue of the xPlain package.

**Key observation:** Occasionally, more important points will be highlighted in boxes like this.

The xPlain package is not part of PSALTer, so the output from PSALTer itself will contrast with this formatting and be quite distinctive.

#### The structure of this file

The calibration file runs PSALTer on a very long list of theories, whose particle spectra are already known.

The first step is to load the PSALTer package.

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Package xAct`PSALTer` version 1.0.0-developer, {2024, 1, 11}

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Great, so PSALTer is now loaded and we can start to do some science.

# Weyl gauge theory extended (eWGT)

Key observation: We will test the WeylGaugeTheoryExtended module. This is an extension to test eWGT [Lasenby and Hobson 2016].

(2)

**Key observation:** This section is still under development by Zhiyuan.

# **Preamble: setting out the fields**

We present the tetrad,

••• ValidateSymbol: Symbol WeylTetrad is already used as a tensor.

and the inverse tetrad.

••• ValidateSymbol: Symbol WeylInvTetrad is already used as a tensor.

We present the tetrad in terms of perturbation tetrad f,

 $\delta_{\alpha}^{\chi} + f_{\alpha}^{\chi}$ 

and the inverse tetrad in terms of perturbation tetrad f.

 $\delta^{\alpha}_{\chi} + f^{\alpha\beta} f_{\beta\chi} - f_{\chi}^{\alpha}$ 

••• ValidateSymbol: Symbol WeylDaggerA is already used as a tensor.

\cdots General: Further output of ValidateSymbol::used will be suppressed during this calculation. 🕖

We present and expand the eWGT (dagger) field strengths:

T+:

$$\mathcal{T}^{\alpha}_{\beta\chi} - \frac{1}{3} \delta^{\alpha}_{\chi} \mathcal{T}^{\delta}_{\beta\delta} + \frac{1}{3} \delta^{\alpha}_{\beta} \mathcal{T}^{\delta}_{\chi}$$

$$-\frac{1}{3} \partial_{\alpha} \mathcal{T}^{\chi}_{\beta \chi} + \frac{1}{3} \partial_{\beta} \mathcal{T}^{\chi}_{\alpha \chi} + h_{\alpha}^{\chi} h_{\beta}^{\delta} \partial_{\chi} \mathcal{B}_{\delta} - h_{\alpha}^{\chi} h_{\beta}^{\delta} \partial_{\delta} \mathcal{B}_{\chi}$$

CovD+(Phi):

$$\frac{1}{3} \phi \mathcal{T}^{\beta}_{\alpha\beta} b^{\alpha}_{i} - \phi \mathcal{B}_{i} + \partial_{i} \phi$$

$$\mathcal{A}^{\dagger}{}^{\alpha\gamma}{}_{\phi} \ \mathcal{A}^{\dagger}{}^{\beta}{}_{\chi\chi} \ h_{\delta}^{\ \chi} \ h_{\epsilon}^{\ \phi} - \mathcal{A}^{\dagger}{}^{\alpha\gamma}{}_{\chi} \ \mathcal{A}^{\dagger}{}^{\beta}{}_{\gamma\phi} \ h_{\delta}^{\ \chi} \ h_{\epsilon}^{\ \phi} + h_{\delta}^{\ \chi} \ h_{\epsilon}^{\ \phi} \ \partial_{\chi} \mathcal{A}^{\dagger}{}^{\alpha\beta}{}_{\phi} - h_{\delta}^{\ \chi} \ h_{\epsilon}^{\ \phi} \ \partial_{\phi} \mathcal{A}^{\dagger}{}^{\alpha\beta}{}_{\chi}$$

I want to check the outputs for Einstein Gauge expansion

(9)

(16)

Here is the non-linear expansion of A+ to level of perturbation field f:

$$\mathcal{A}^{\alpha\beta}_{\ \theta} - \delta^{\beta}_{\ \theta} \mathcal{B}^{\alpha} - f^{\beta\chi}_{\ \chi\theta} f_{\chi\theta} \mathcal{B}^{\alpha} + f^{\beta}_{\ \theta} \mathcal{B}^{\alpha} + \delta^{\alpha}_{\ \theta} \mathcal{B}^{\beta} + f^{\alpha\chi}_{\ \chi\theta} f_{\chi\theta} \mathcal{B}^{\beta} - f^{\alpha}_{\ \theta} \mathcal{B}^{\beta} + \delta^{\beta}_{\ \theta} f^{\chi\alpha} \mathcal{B}_{\chi} - \delta^{\alpha}_{\ \theta} f^{\chi\beta} \mathcal{B}_{\chi} + f^{\chi\beta}_{\ \theta} f^{\alpha}_{\ \theta} \mathcal{B}_{\chi} - f^{\chi\alpha}_{\ \theta} f^{\beta}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} f^{\beta}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} f^{\beta}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} f^{\beta}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} f^{\beta}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} f^{\beta}_{\ \theta} \mathcal{B}_{\chi} - f^{\alpha}_{\ \theta} \mathcal{B}_{\chi} - f$$

Here is the linearised expansion:

$$\mathcal{A}^{\alpha\beta}_{\phantom{\alpha\beta}\theta} - \delta^{\beta}_{\phantom{\beta}\theta} \mathcal{B}^{\alpha} + \delta^{\alpha}_{\phantom{\alpha}\theta} \mathcal{B}^{\beta}$$

We present and expand A+ and T into PGT field strengths:

A+:

$$\mathcal{A}^{\alpha\beta}_{\phantom{\alpha\beta}\theta} - \delta^{\beta}_{\phantom{\beta}\theta} \mathcal{B}^{\alpha} + \delta^{\alpha}_{\phantom{\alpha}\theta} \mathcal{B}^{\beta}$$

T:

$$\mathcal{A}_{\chi\delta}^{\alpha} h_{\beta}^{\delta} - \mathcal{A}_{\beta\delta}^{\alpha} h_{\chi}^{\delta} + h_{\beta}^{\delta} h_{\chi}^{\epsilon} \partial_{\delta} b_{\epsilon}^{\alpha} - h_{\beta}^{\delta} h_{\chi}^{\epsilon} \partial_{\epsilon} b_{\delta}^{\alpha}$$

Check two expressions for A+ are the same:

Here is the non-linear expansion of A+ to level of perturbation field f:

$$\mathcal{A}^{\alpha\beta}_{\phantom{\alpha\beta}\theta} - \delta^{\beta}_{\phantom{\beta}\theta} \mathcal{B}^{\alpha} - f^{\beta\chi}_{\phantom{\beta}\chi} f_{\phantom{\alpha}\chi\theta} \mathcal{B}^{\alpha} + f^{\beta}_{\phantom{\alpha}\theta} \mathcal{B}^{\alpha} + \delta^{\alpha}_{\phantom{\alpha}\theta} \mathcal{B}^{\beta} + f^{\alpha\chi}_{\phantom{\alpha}\chi} f_{\phantom{\alpha}\chi\theta} \mathcal{B}^{\beta} - f^{\alpha}_{\phantom{\alpha}\theta} \mathcal{B}^{\beta} + \delta^{\beta}_{\phantom{\alpha}\theta} f^{\chi\alpha}_{\phantom{\alpha}\theta} \mathcal{B}_{\chi} - \delta^{\alpha}_{\phantom{\alpha}\theta} f^{\chi\beta}_{\phantom{\alpha}\theta} \mathcal{B}_{\chi} - f^{\chi\alpha}_{\phantom{\alpha}\theta} f^{\beta}_{\phantom{\alpha}\theta} \mathcal{B}_{\chi} - f^{\chi\alpha}_{\phantom{\alpha}\theta} \mathcal{B}_{\chi} - f^{\chi\alpha$$

Here is the linearised expansion:

$$\mathcal{A}^{\alpha\beta}_{\ \theta} - \delta^{\beta}_{\ \theta} \mathcal{B}^{\alpha} + \delta^{\alpha}_{\ \theta} \mathcal{B}^{\beta}$$

Λ

## **Key observation:** Now we have defined all the fields we need.

In eqn 15 of Lin's draft, we check that the T+ contraction = 0. Here we expand T+ to PGT T.

0 (17)

In Eq. (18) this is the non-linear Lagrangian as given in eqn 13 of Lin's draft paper.

$$\frac{1}{2} \stackrel{\mathbf{v}}{\cdot} \mathcal{D}^{\dagger} \phi_{i} \mathcal{D}^{\dagger} \phi^{i} + \xi \stackrel{\mathbf{v}}{\cdot} \mathcal{H}^{\dagger}_{\alpha\beta} \mathcal{H}^{\dagger\alpha\beta} + \lambda \stackrel{\mathbf{v}}{\cdot} \phi^{2} \mathcal{R}^{\dagger\alpha\beta}_{\alpha\beta} + \left(\frac{r}{3} + \frac{r}{6}\right) \mathcal{R}^{\dagger}_{\alpha\beta\chi\delta} \mathcal{R}^{\dagger\alpha\beta\chi\delta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger}_{\alpha\chi\beta\delta} \mathcal{R}^{\dagger\alpha\beta\chi\delta} - c \stackrel{\mathbf{v}}{\cdot} \mathcal{H}^{\dagger\alpha\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\chi\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger}_{\alpha\chi\beta\delta} \mathcal{R}^{\dagger\alpha\beta\chi\delta} - c \stackrel{\mathbf{v}}{\cdot} \mathcal{H}^{\dagger\alpha\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\chi\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger}_{\alpha\chi\beta\delta} \mathcal{R}^{\dagger\alpha\beta\chi\delta} - c \stackrel{\mathbf{v}}{\cdot} \mathcal{H}^{\dagger\alpha\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\chi\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\beta\chi\delta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\beta\chi\delta} - c \stackrel{\mathbf{v}}{\cdot} \mathcal{H}^{\dagger\alpha\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\beta\chi\delta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\beta\gamma\delta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\beta\gamma\delta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\beta\gamma\delta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\beta\gamma\delta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\beta\gamma\delta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} \mathcal{R}^{\dagger\alpha\gamma\beta}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3}\right) \mathcal{R}^{\dagger\alpha\gamma}_{\alpha\beta} + \left(\frac{2r}{3} - \frac{2r}{3$$

$$\left(\begin{matrix} r_{\text{.}} + r_{\text{.}} \\ 4 + 5 \end{matrix}\right) \mathcal{R}^{\dagger}_{\alpha\delta\beta} \mathcal{R}^{\dagger}_{\alpha\delta\beta} \mathcal{R}^{\dagger}_{\lambda} \mathcal{R}^{\dagger}_{\lambda} + \left(\begin{matrix} r_{\text{.}} - r_{\text{.}} \\ 4 - r_{\text{.}} \end{matrix}\right) \mathcal{R}^{\dagger}_{\beta\delta\alpha} \mathcal{R}^{\dagger}_{\lambda} \mathcal{R}^{\dagger}_{\beta\delta\alpha} \mathcal{R}^{\dagger}_{\lambda} \mathcal{R}^{\dagger}_{\beta\delta\alpha} \mathcal{R}^{\dagger}_{\lambda} \mathcal{R}^{\dagger}_{\lambda} \mathcal{R}^{\dagger}_{\beta\delta\alpha} \mathcal{R}^{\dagger}_{\lambda} \mathcal{R}^{\dagger}_{\lambda}$$

Diagnostic: Now the non-linear Lagrangian has been expanded to PGT quantities. This is now stored for linearisation.

## Test case 1: E--H action.

We test the case of the modified Einstein-Hilbert action, and the code will give only 2 propagating graviton modes.

$$\lambda. \phi^2 \mathcal{R}^{\dagger \alpha \beta}_{\alpha \beta}$$

Here, we perform rescalings after application of Einstein Gauge:  $\phi_-0^2*\lambda \rightarrow \lambda$ ,  $\phi_-0^2*v \rightarrow v$ ,  $\phi_-0^2*t_-i \rightarrow t_-i$ . Also

 $\phi_{-}0 \rightarrow 1$ , i.e. here I am making the compensator dimensionless, any possible masses order 1. I do this to prevent any denominators phi/phi0.

Here is the linearised Lagrangian before feeding into ParticleSpectrum[].

$$\lambda. \mathcal{A}_{\alpha\chi\beta} \mathcal{A}^{\alpha\beta\chi} + \lambda. \mathcal{A}_{\alpha\beta}^{\alpha\beta} \mathcal{A}_{\beta\chi}^{\chi} + 4\lambda. \mathcal{A}_{\alpha\beta}^{\beta} \mathcal{B}^{\alpha} - 6\lambda. \mathcal{B}_{\alpha\beta}^{\beta} \mathcal{B}^{\alpha} - 6\lambda. \partial_{\alpha}\mathcal{B}^{\alpha} + 2\lambda. f^{\alpha\beta} \partial_{\beta}\mathcal{A}_{\alpha\chi}^{\chi} - 2\lambda. \partial_{\beta}\mathcal{A}_{\alpha\beta}^{\alpha\beta} - 4\lambda. f^{\alpha\beta} \partial_{\beta}\mathcal{B}_{\alpha} + 4\lambda. f^{\alpha}_{\alpha} \partial_{\beta}\mathcal{B}^{\beta} - 2\lambda. f^{\alpha\beta} \partial_{\chi}\mathcal{A}_{\alpha\beta}^{\chi} + 2\lambda. f^{\alpha}_{\alpha} \partial_{\chi}\mathcal{A}_{\beta\beta}^{\chi}$$

(19)

TaskRemove: A string or a TaskObject is expected instead of CellObject 58808c55-75f1-4216-8386-296cebfad10c

#### **PSALTer results panel**

$$S = \iiint (\phi \, \rho + \, \sigma^{\alpha \beta \chi} \, \, \mathcal{R}_{\alpha \beta \chi} + \, \mathcal{J}^{\alpha \beta} \, \, f_{\alpha \beta} + \, \mathcal{J}^{\alpha} \, \, \mathcal{B}_{\alpha} + \, \lambda \cdot (\, \mathcal{R}_{\alpha \chi \beta} \, \, \mathcal{R}^{\alpha \beta \chi} + \, \mathcal{R}^{\alpha \beta}_{\alpha \beta} \, \, \mathcal{R}^{\chi} + \, 4 \, \, \mathcal{R}^{\beta}_{\alpha \beta} \, \, \mathcal{B}^{\alpha} - \, 6 \, \, \mathcal{B}_{\alpha} \, \, \mathcal{B}^{\alpha} + \, 2 \, \, f^{\alpha \beta}_{\alpha \beta} \, \, \partial_{\beta} \mathcal{R}^{\chi}_{\alpha \beta} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}_{\alpha} + \, 4 \, \, f^{\alpha}_{\alpha \beta} \, \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 2 \, \, f^{\alpha \beta}_{\alpha \beta} \, \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - \, 4 \, \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha}$$

## **Wave operator**

$\begin{array}{cccccccccccccccccccccccccccccccccccc$											
$0^{+}\mathcal{J} + \begin{array}{ c c c c c c c c c c c c c c c c c c c$											
0° \rho + 0 0 0 0 0 0 0											
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$											
•											
$0^+\tau^{\parallel} + \begin{vmatrix} \frac{i\sqrt{3}}{7k\lambda} & 0 & -\frac{i}{7\sqrt{2}k\lambda} & \frac{1}{2k^2\lambda} & 0 \end{vmatrix} = 0$											
$0.7$ $\tau^{\perp}$ $\uparrow$ 0 0 0 0 0 0											
$[\cdot, \sigma] + \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & -\frac{1}{\lambda} \end{bmatrix}$	1. σ <sup>  </sup> αβ	$^{1^+}\sigma^{\scriptscriptstyle\perp}_{\alpha\beta}$	$^{1^+}\tau^{\parallel}_{\alpha\beta}$	$^{1}\mathcal{J}_{lpha}$	$^{1.}\sigma^{\parallel}{}_{\alpha}$	$^{1}\sigma_{\alpha}^{\perp}$	$^{1.}\tau^{\parallel}_{\alpha}$	$^{1.}\tau^{\perp}_{\alpha}$			
	0	$-\frac{\sqrt{2}}{\lambda.+k^2\lambda.}$	$-\frac{i\sqrt{2}k}{\lambda + k^2\lambda}.$	0	0	0	0	0			
1. σ₁ 1	$\frac{1}{\lambda + k^2 \lambda}$	$\frac{1}{(1+k^2)^2 \lambda}.$	$\frac{ik}{(1+k^2)^2 \lambda}.$	0	0	0	0	0			
1.° r " †	$\frac{i \sqrt{2} k}{\lambda + k^2 \lambda}.$	$-\frac{ik}{(1+k^2)^2\lambda}.$	$\frac{k^2}{(1+k^2)^2 \lambda}.$	0	0	0	0	0			
1: <i>g</i>	- <sup>α</sup> 0	0	0	$-\frac{2(3+8k^2)}{(7+10k^2)^2\lambda}.$	$\frac{2(-1+6k^2)}{(7+10k^2)^2\lambda}.$	$\frac{\sqrt{2} (1+20 k^2)}{(7+10 k^2)^2 \lambda}.$	0	$\frac{2 i (k+20 k^3)}{(7+10 k^2)^2 \lambda}.$			
1. ol	-α 0	0	0	$\frac{2(-1+6k^2)}{(7+10k^2)^2\lambda}.$	$\frac{16 (2+k^2)}{(7+10 k^2)^2 \lambda}.$	$\frac{\sqrt{2} (33+10 k^2)}{(7+10 k^2)^2 \lambda}$	0	$\frac{2 i k (33+10 k^2)}{(7+10 k^2)^2 \lambda}.$			
1. σ <sup>L</sup>	0	0	0	$\frac{\sqrt{2} (1+20 k^2)}{(7+10 k^2)^2 \lambda}$	$\frac{\sqrt{2} (33+10 k^2)}{(7+10 k^2)^2 \lambda}$	$\frac{65}{(7+10 k^2)^2 \lambda}$	0	$\frac{65 i \sqrt{2} k}{(7+10 k^2)^2 \lambda}.$			
1'τ"	-α 0	0	0			0		0			
1,1,1	-α 0	0	0	$-\frac{2 i (k+20 k^3)}{(7+10 k^2)^2 \lambda}$	$-\frac{2ik(33+10k^2)}{(7+10k^2)^2\lambda}.$	$-\frac{65 i \sqrt{2} k}{(7+10 k^2)^2 \lambda}.$	0	$\frac{130 k^2}{(7+10 k^2)^2 \lambda}.$	$^{2^{+}}\sigma^{\parallel}_{\alpha\beta}$ $^{2^{+}}\tau^{\parallel}$	αβ 2	$\sigma^{\parallel}_{\alpha\beta\chi}$
			_					$^{2^{+}}\sigma^{\parallel}$ † $^{\alpha\beta}$	$0 - \frac{i}{k}$	√ <u>2</u>	0
								$2^+\tau^{\parallel} \uparrow^{\alpha\beta}$	$\frac{i\sqrt{2}}{k\lambda}$ $-\frac{1}{k^2}$	<u>1</u> λ.	0
								$\mathcal{E}^{\sigma} \sigma^{\parallel} \uparrow^{\alpha\beta\chi}$	0 0	)	<u>2</u> λ.

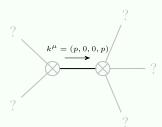
### Source constraints

Spin-parity form	Covariant form	Multiplicitie			
$0^+\tau^{\perp}=0$	$\partial_{\beta}\partial_{\alpha}\mathcal{J}^{\alpha\beta}$ == 0	1			
$2^{0^{+}}\sigma^{\parallel} + {}^{0^{+}}\mathcal{J} == 0$	$\partial_{\alpha} \mathcal{J}^{\alpha} == 2  \partial_{\beta} \sigma^{\alpha}_{\alpha}^{\beta}$	1			
0÷ρ == 0	$\rho == 0$	1			
$2ik!\sigma^{\parallel}^{\alpha} + 1\tau^{\perp}^{\alpha} - ik!\sigma^{\perp}$		3			
$1.\tau^{\alpha} = 0$	$\partial_{\chi}\partial_{\beta}\partial^{\alpha}\mathcal{J}^{\beta\chi} == \partial_{\chi}\partial^{\chi}\partial_{\beta}\mathcal{J}^{\beta\alpha}$	3			
$\frac{1}{2} \cdot \int_{0}^{\alpha} d^{\alpha} = 2 \cdot \int_{0}^{\alpha} d^{\alpha} + \int_{0}^{\alpha} \int_{0}^{\alpha} d^{\alpha}$	$\partial_{\beta}\partial^{\alpha}\mathcal{J}^{\beta} == \partial_{\beta}\partial^{\beta}\mathcal{J}^{\alpha} + 2\left(\partial_{\chi}\partial^{\alpha}\sigma^{\beta}_{\beta}^{\chi} + \partial_{\chi}\partial^{\chi}\sigma^{\beta\alpha}_{\beta}\right)$	3			
$i k \cdot \dot{\tau}^{+} \sigma^{\perp}^{\alpha\beta} + \dot{\tau}^{+} \tau^{\parallel}^{\alpha\beta} == 0$	$\partial_{\chi}\partial^{\alpha}\mathcal{J}^{\beta\chi} + \partial_{\chi}\partial^{\beta}\mathcal{J}^{\chi\alpha} + \partial_{\chi}\partial^{\chi}\mathcal{J}^{\alpha\beta} + 2\partial_{\delta}\partial_{\chi}\partial^{\alpha}\sigma^{\chi\beta\delta} + 2\partial_{\delta}\partial^{\delta}\partial_{\chi}\sigma^{\chi\alpha\beta} = \partial_{\chi}\partial^{\alpha}\mathcal{J}^{\chi\beta} + \partial_{\chi}\partial^{\beta}\mathcal{J}^{\alpha\chi} + \partial_{\chi}\partial^{\chi}\mathcal{J}^{\beta\alpha} + 2\partial_{\delta}\partial_{\chi}\partial^{\beta}\sigma^{\chi\alpha\delta}$	3			
Total expected gauge generators:					

### **Massive spectrum**

(No particles)

### **Massless spectrum**



Massless particle

Pole residue:  $\left| -\frac{1}{\lambda} \right| > 0$ Polarisations: 2

## **Unitarity conditions**

λ. < 0

**Key observation:** This marks the completion of the particle spectrum analysis for the modified E--H action.

## Test case 2: E--H action with propagating compensator.

We test the case of the modified Einstein-Hilbert action, with propagating compensator.

$$\frac{1}{2} v. \mathcal{D}^{\dagger} \phi_{\alpha} \mathcal{D}^{\dagger} \phi^{\alpha} + \lambda. \phi^{2} \mathcal{R}^{\dagger \alpha \beta}_{\alpha \beta}$$

Here, we perform rescalings after application of Einstein Gauge:  $\phi_-0^{\circ}2*\lambda \rightarrow \lambda$ ,  $\phi_-0^{\circ}2*\nu \rightarrow \nu$ ,  $\phi_-0^{\circ}2*t_-i \rightarrow t_-i$ . Also

 $\phi_{-0}$  -> 1, i.e. here I am making the compensator dimensionless, any possible masses order 1. I do this to prevent any denominators phi/phi0.

Here is the linearised Lagrangian before feeding into ParticleSpectrum[].

$$\lambda. \mathcal{A}_{\alpha\chi\beta} \mathcal{A}^{\alpha\beta\chi} + \left(\lambda. - \frac{v.}{18}\right) \mathcal{A}_{\alpha}^{\alpha\beta} \mathcal{A}_{\beta\chi}^{\chi} + \left(4\lambda. - \frac{v.}{3}\right) \mathcal{A}_{\alpha\beta}^{\beta} \mathcal{B}^{\alpha} + \left(-6\lambda. + \frac{v.}{2}\right) \mathcal{B}_{\alpha} \mathcal{B}^{\alpha} + \frac{1}{3} v. \mathcal{B}^{\alpha} \partial_{\alpha}f^{\beta}_{\beta} - 6\lambda. \partial_{\alpha}\mathcal{B}^{\alpha} + 2\lambda. f^{\alpha\beta} \partial_{\beta}\mathcal{A}_{\alpha\chi}^{\chi} - 2\lambda. \partial_{\beta}\mathcal{A}^{\alpha\beta}_{\alpha} - \frac{1}{2} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} - \frac{1}{2} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} - \frac{1}{2} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} - \frac{1}{2} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} - 2\lambda. f^{\alpha\beta} \partial_{\chi}\mathcal{A}_{\alpha\beta}^{\chi} + 2\lambda. f^{\alpha\beta} \partial_{\chi}\mathcal{A}_{\beta\beta}^{\chi} + \frac{1}{18} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} \partial_{\chi}f^{\chi}_{\beta} - \frac{1}{2} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} \partial_{\chi}f^{\chi}_{\beta} - \frac{1}{2} v. \partial_{\beta}f^{\alpha\beta}_{\alpha} \partial_{\chi}f^{\gamma}_{\beta} - \frac{1}{2} v. \partial_{\gamma}f^{\gamma}_{\beta} \partial_{\chi}f^{\gamma}_{\alpha} - \frac{1}{2$$

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## **PSALTer results panel**

 ${}^{2^{+}}\mathcal{H}^{\parallel} \dagger^{\alpha\beta}$ 

 $^{2^{+}}f^{\parallel}\dagger^{\alpha\beta}$ 

 $^{2}\mathcal{H}^{\parallel}$  †  $^{\alpha\beta\chi}$ 

0

## Wave operator

	$\overset{0^{+}}{\cdot}\mathcal{B}$	0÷ <b>φ</b>	${}^{0^{\scriptscriptstyle +}}_{\cdot}\mathcal{A}^{\scriptscriptstyle \parallel}$	$0^{+}f^{\parallel}$	$\overset{0^+}{\cdot}f^{\scriptscriptstyle \perp}$	${}^{0}\mathcal{A}^{\parallel}$	_								
0⁺. ₿†	$-6 \lambda. + \frac{v}{2}$	0	12 λν. 2 √6	$-\frac{i k (12 \lambda - v.)}{2 \sqrt{3}}$	0	0									
0 <sup>+</sup> <b>φ</b> †	0	0	0	0	0	0									
<sup>0⁺</sup> Æ <sup>  </sup> †	12 λν. 2 √6	0	$-\lambda$ . $+\frac{v}{12}$	$\frac{i k (12 \lambdav.)}{6 \sqrt{2}}$	0	0									
0 <sup>+</sup> <i>f</i> <sup>∥</sup> †		0	$-\frac{i k (12 \lambdav.)}{6 \sqrt{2}}$	$\frac{k^2 v}{6}$	0	0									
0 <sup>+</sup> f <sup>⊥</sup> †	0	0	0	0	0	0									
º.º A∥ †	0	0	0	0	0	-λ.	$^{1.}\mathcal{A}^{\parallel}_{lphaeta}$	${}^{1,\dagger}\mathcal{H}^{\perp}{}_{lphaeta}$	$f^{\dagger}f^{\parallel}_{\alpha\beta}$	$^{1}\mathcal{B}_{lpha}$	${}^{1}\mathcal{A}^{\parallel}{}_{\alpha}$	<sup>1</sup> . A <sup>±</sup> α	$f^{\parallel}_{\alpha}$	$f_{\alpha}^{\perp}$	
						$^{1}\mathcal{A}^{\parallel}\dagger^{lphaeta}$	λ. - <u>-</u> 2	$-\frac{\lambda}{\sqrt{2}}$	$-\frac{i k \lambda}{\sqrt{2}}$	0	0	0	0	0	
						$^{1^{+}}\mathcal{A}^{\perp}\dagger^{\alpha\beta}$	$-\frac{\lambda}{\sqrt{2}}$	0	0	0	0	0	0	0	
						$\dot{f}^{\dagger}f^{\parallel} \dagger^{\alpha\beta}$	$\frac{i k \lambda}{\sqrt{2}}$	0	0	0	0	0	0	0	
						¹.'Β† <sup>α</sup>	0	0	0	$-6 \lambda. + \frac{v}{2}$	$-2\lambda.+\frac{v.}{6}$	$\frac{12 \lambda - v}{6 \sqrt{2}}$	0	$\frac{1}{6} ik (12 \lambda v.)$	
						${}^{1}\mathcal{A}^{\parallel}\dagger^{\alpha}$	0	0	0	$-2\lambda.+\frac{v}{6}$	$\frac{1}{18} \left( -9  \lambda_{\cdot} + v_{\cdot} \right)$	$\frac{18 \lambdav.}{18 \sqrt{2}}$	0	$\frac{1}{18}  i  k  (18  \lambda v.)$	
						¹. ℋ† <sup>α</sup>	0	0	0	$\frac{12 \lambda \cdot v}{6 \sqrt{2}}$	$\frac{18 \lambdav.}{18 \sqrt{2}}$	v. 36	0	$\frac{ikv.}{18\sqrt{2}}$	
						$f^{\parallel} \uparrow^{\alpha}$	0	0	0	0	0	0	0	0	
						$f^{\perp}f^{\perp}$	0	0	0	$k\left(-2i\lambda.+\frac{i\nu.}{6}\right)$	$k\left(-i\lambda.+\frac{i\nu.}{18}\right)$	$-\frac{i k v}{18 \sqrt{2}}$	0	$\frac{k^2 v}{18}$	<sup>2†</sup> .Æ

Saturated propagator

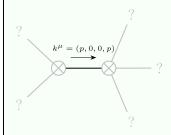
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	τ <sup>⊥</sup> 0. σ									
6 v. √6 v. : √3	0									
√6 v v. ;	0									
$0^{+}\sigma^{\parallel} + \frac{7}{588\lambda^{2}-49\lambda \cdot \nu}  0  \frac{7}{49\lambda \cdot (-12\lambda \cdot + \nu)}  \frac{7}{7} \sqrt{2} k\lambda.$	0									
$0^+\tau^{\parallel} + \frac{i\sqrt{3}}{7k\lambda} \qquad 0 \qquad -\frac{i}{7\sqrt{2}k\lambda} \qquad \frac{1}{2k^2\lambda}$	0									
	0									
°. σ † 0 0 0 0	1	1⁺a∥ .	$\dot{\sigma}^{\perp}_{\alpha\beta}$	1, ,	$^{1}\mathcal{J}_{lpha}$	1. σ <sup>  </sup> α	$1.\sigma_{\alpha}$	$^{1} \tau^{\parallel}_{\alpha}$	$^{1}$ $\tau^{\perp}_{\alpha}$	
	•		$-\frac{\sqrt{2}}{\lambda \cdot + k^2 \lambda \cdot}$							Ì
	$^{1^{+}}\sigma^{\parallel}$ † $^{\alpha\beta}$				0	0	0	0	0	
	$1.^{+}\sigma^{\perp} \uparrow^{\alpha\beta}$	$-\frac{\sqrt{2}}{\lambda.+k^2\lambda.}$	$\frac{1}{(1+k^2)^2 \lambda}.$	$\frac{i k}{(1+k^2)^2 \lambda}.$	0	0	0	0	0	
	$1.\tau^{\parallel} + \alpha\beta$	$\frac{i \sqrt{2} k}{\lambda . + k^2 \lambda .}$	$-\frac{ik}{(1+k^2)^2\lambda}.$	$\frac{k^2}{(1+k^2)^2 \lambda}.$	0	0	0	0	0	
	1. σ. ι α			0	8 (9 (3+8 k²) λ.+4 k⁴ v.)	8 ((9-54 k²) λ.+k² (7+2 k²) v.)	4 √2 (9 (1+20 k²) λ14 k² v.)	0	8 i (9 k λ.+2 k³ (90 λ7 v.))	
	! <i>J</i> † <sup>a</sup>	0	0	0	$3(7+10 k^2)^2 \lambda. (12 \lambdav.)$	$3 (7+10 k^2)^2 \lambda. (12 \lambdav.)$	3 $(7+10 k^2)^2 \lambda$ . $(12 \lambdav.)$	0	$3 (7+10 k^2)^2 \lambda. (12 \lambdav.)$	
	$\frac{1}{2}\sigma^{\parallel} + \alpha$	0	0	0	$-\frac{8 \left(\left(9-54  k^2\right) \lambda.+k^2 \left(7+2  k^2\right) v.\right)}{3 \left(7+10  k^2\right)^2 \lambda. \left(12  \lambdav.\right)}$	$\frac{576 (2+k^2) \lambda \cdot -2 (7+2 k^2)^2 v \cdot}{3 (7+10 k^2)^2 \lambda \cdot (12 \lambda \cdot -v \cdot)}$	$\frac{2\sqrt{2}(18(33+10k^2)\lambda7(7+2k^2)v.}{3(7+10k^2)^2\lambda.(12\lambdav.)}$	) - 0	$\frac{4 i k (18 (33+10 k^2) \lambda7 (7+2 k^2) v.)}{3 (7+10 k^2)^2 \lambda. (12 \lambdav.)}$	
						$2 \sqrt{2} (18 (33+10 k^2) \lambda7 (7+2 k^2) v.)$	4 (585 λ49 v.)		4 i √2 k (585 λ49 v.)	
	$!\sigma^{\perp}\uparrow^{\alpha}$	0	0	0	$3(7+10 k^2)^2 \lambda. (12 \lambdav.)$	$3 (7+10 k^2)^2 \lambda. (12 \lambdav.)$	$3(7+10 k^2)^2 \lambda. (12 \lambdav.)$	0	$\frac{1}{3(7+10 k^2)^2 \lambda. (12 \lambdav.)}$	
	$1.\tau^{\parallel} + \alpha$	0	0	0	0	0	0	0	0	
	1. τ <sup>⊥</sup> † α	0	0	0	8 i (9 k λ.+2 k³ (90 λ7 v.))	4 i k (18 (33+10 k²) λ7 (7+2 k²) v.)	4 i √2 k (585 λ49 v.)	0	8 k <sup>2</sup> (585 λ49 v.)	
	:τ Τ	U	0	U	$3(7+10 k^2)^2 \lambda. (12 \lambdav.)$	3 (7+10 k <sup>2</sup> ) <sup>2</sup> λ. (12 λν.)	$\frac{1}{3} (7+10 k^2)^2 \lambda. (12 \lambdav.)$	0	$3 (7+10 k^2)^2 \lambda. (12 \lambdav.)$	
									$\overset{2^{+}}{\cdot}\sigma^{\parallel}\uparrow^{lphaeta}$	
									2 <sup>*</sup> τ <sup>  </sup> † <sup>αβ</sup>	
									$\frac{2}{3}\sigma^{\parallel} + \frac{\alpha\beta\chi}{3}$	
										L
Source constraints										

Spin-parity form	Covariant form	Multiplicities				
0 <sup>+</sup> τ <sup>⊥</sup> == 0	$\partial_{eta}\partial_{lpha}\mathcal{J}^{lphaeta}=0$	1				
$\frac{1}{2 \cdot 0^{\cdot} \sigma^{\parallel} + 0^{\cdot} \mathcal{J}} = 0$	$\partial_{\alpha} \mathcal{J}^{\alpha} = 2 \partial_{\beta} \sigma^{\alpha}_{\alpha}{}^{\beta}$	1				
0 <sup>+</sup> ρ == 0	ρ == 0	1				
$2ik!\sigma^{\parallel}^{\alpha} + i\tau^{\perp}^{\alpha} - ik!\sigma^{\alpha}$	$==0 \left[ \partial_{\chi} \partial^{\chi} \partial_{\beta} \mathcal{J}^{\alpha\beta} + \partial_{\chi} \partial^{\chi} \partial_{\beta} \partial^{\alpha} \mathcal{J}^{\beta} + 2 \partial_{\delta} \partial^{\delta} \partial_{\chi} \partial_{\beta} \sigma^{\beta\alpha\chi} = \partial_{\chi} \partial_{\beta} \partial^{\alpha} \mathcal{J}^{\beta\chi} + \partial_{\chi} \partial^{\chi} \partial_{\beta} \partial^{\beta} \mathcal{J}^{\alpha} + 2 \left( \partial_{\delta} \partial^{\delta} \partial_{\chi} \partial^{\alpha} \sigma^{\beta}_{\beta}^{\chi} + \partial_{\delta} \partial^{\delta} \partial_{\chi} \partial^{\chi} \sigma^{\beta\alpha}_{\beta} \right) \right]$	3				
1. T    a == 0	$\partial_{\chi}\partial_{\beta}\partial^{\alpha}\mathcal{J}^{\beta\chi} == \partial_{\chi}\partial^{\chi}\partial_{\beta}\mathcal{J}^{\beta\alpha}$	3				
$\frac{1}{2} \int_{0}^{1} d^{\alpha} = 2 \int_{0}^{1} \sigma^{\perp} + \int_{0}^{1} \sigma^{\alpha}$	$\partial_{\beta}\partial^{\alpha}\mathcal{J}^{\beta} = \partial_{\beta}\partial^{\beta}\mathcal{J}^{\alpha} + 2\left(\partial_{\chi}\partial^{\alpha}\sigma^{\beta}_{\beta}{}^{\chi} + \partial_{\chi}\partial^{\chi}\sigma^{\beta\alpha}_{\beta}\right)$	3				
$i k \stackrel{1^+}{\cdot} \sigma^{\perp}^{\alpha\beta} + 1^+_{\cdot} \tau^{\parallel}^{\alpha\beta} == 0$	$\partial_{\chi}\partial^{\alpha}\mathcal{J}^{\beta\chi} + \partial_{\chi}\partial^{\beta}\mathcal{J}^{\chi\alpha} + \partial_{\chi}\partial^{\chi}\mathcal{J}^{\alpha\beta} + 2\partial_{\delta}\partial_{\chi}\partial^{\alpha}\sigma^{\chi\beta\delta} + 2\partial_{\delta}\partial^{\delta}\partial_{\chi}\sigma^{\chi\alpha\beta} = \partial_{\chi}\partial^{\alpha}\mathcal{J}^{\chi\beta} + \partial_{\chi}\partial^{\beta}\mathcal{J}^{\alpha\chi} + \partial_{\chi}\partial^{\chi}\mathcal{J}^{\beta\alpha} + 2\partial_{\delta}\partial_{\chi}\partial^{\beta}\sigma^{\chi\alpha\delta}$	3				
Total expected gauge generators:						

### **Massive spectrum**

(No particles)

## **Massless spectrum**



Massless particle

Pole residue: Polarisations: 2

## **Unitarity conditions**

λ. < 0

**Key observation:** This marks the completion of the particle spectrum analysis for the modified E--H action with propagating compensator.

# Test case 3: E--H action with propagating compensator and vector.

We test the case of the modified Einstein-Hilbert action, with propagating compensator and vector B.

$$\frac{1}{2} v. \mathcal{D}^{\dagger} \phi_{\alpha} \mathcal{D}^{\dagger} \phi^{\alpha} + \xi. \mathcal{H}^{\dagger}_{\alpha\beta} \mathcal{H}^{\dagger^{\alpha\beta}} - c. \mathcal{H}^{\dagger^{\alpha\beta}} \mathcal{R}^{\dagger^{\alpha\beta}}_{\alpha\beta\chi} + \lambda. \phi^{2} \mathcal{R}^{\dagger^{\alpha\beta}}_{\alpha\beta}$$

Here is the linearised Lagrangian before feeding into ParticleSpectrum[].

$$\lambda. \mathcal{A}_{\alpha\chi\beta} \mathcal{A}^{\alpha\beta\chi} + \left(\lambda. - \frac{v.}{18}\right) \mathcal{A}^{\alpha\beta}_{\alpha} \mathcal{A}^{\chi}_{\beta\chi} + \left(4\lambda. - \frac{v.}{3}\right) \mathcal{A}_{\alpha\beta}^{\beta}_{\beta} \mathcal{B}^{\alpha} + \left(-6\lambda. + \frac{v.}{2}\right) \mathcal{B}_{\alpha} \mathcal{B}^{\alpha}_{\beta} + \frac{1}{3}v. \mathcal{B}^{\alpha}_{\alpha}\partial_{\beta}\mathcal{B}^{\beta}_{\beta} - 6\lambda. \partial_{\alpha}\mathcal{B}^{\alpha}_{\beta} + 2\lambda. f^{\alpha\beta}_{\beta}\partial_{\beta}\mathcal{A}^{\chi}_{\alpha} - 2\lambda. \partial_{\beta}\mathcal{A}^{\alpha\beta}_{\alpha} - \frac{1}{3}v. \mathcal{B}^{\alpha}_{\beta}\partial_{\beta}\mathcal{B}^{\alpha}_{\alpha} + 4\lambda. f^{\alpha\beta}_{\alpha}\partial_{\beta}\mathcal{B}^{\alpha}_{\alpha} + 4\lambda. f^{\alpha\beta}_{\alpha}\partial_{\beta}\mathcal{B}^{\beta}_{\alpha} + 4\lambda. f^{\alpha\beta$$

TaskRemove: A string or a TaskObject is expected instead of CellObject

General: Further output of TaskRemove::taskid will be suppressed during this calculation.

### **PSALTer results panel**

$$S = \iiint (\phi \, \rho + \, \sigma^{\alpha \beta \chi} \, \mathcal{A}_{\alpha \beta \chi} + \, \mathcal{J}^{\alpha \beta} \, f_{\alpha \beta} + \, \mathcal{J}^{\alpha} \, \mathcal{B}_{\alpha} + \lambda . \, (\mathcal{A}_{\alpha \chi \beta} \, \mathcal{A}^{\alpha \beta \chi} + \, \mathcal{A}^{\alpha \beta}_{\alpha} \, \mathcal{A}^{\chi}_{\beta \chi} + 4 \, \mathcal{A}^{\beta}_{\alpha \beta} \, \mathcal{B}^{\alpha} - 6 \, \mathcal{B}_{\alpha} \, \mathcal{B}^{\alpha} - 6 \, \partial_{\alpha} \mathcal{B}^{\alpha} + 2 \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{A}^{\chi}_{\alpha \chi} - 2 \, \partial_{\beta} \mathcal{A}^{\alpha \beta}_{\alpha \chi} - 4 \, f^{\alpha \beta}_{\alpha \beta} \, \partial_{\beta} \mathcal{B}^{\alpha} - 4 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\beta} \mathcal{B}^{\beta} - 2 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\gamma} \mathcal{A}^{\chi}_{\alpha \beta} + 2 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{A}^{\alpha \chi}_{\alpha \chi} - 2 \, \partial_{\beta} \mathcal{A}^{\alpha \beta}_{\alpha \chi} - 4 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\beta} \mathcal{B}^{\alpha} - 4 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\beta} \mathcal{B}^{\beta} - 2 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\gamma} \mathcal{A}^{\chi}_{\alpha \beta} + 2 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{A}^{\alpha \chi}_{\alpha \chi} - 2 \, \partial_{\beta} \mathcal{A}^{\alpha \beta}_{\alpha \chi} - 4 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha} - 4 \, f^{\alpha \beta}_{\alpha \alpha} \, \partial_{\beta} \mathcal{B}^{\beta}_{\alpha \gamma} - 2 \, f^{\alpha \beta}_{\alpha \alpha \beta} \, \partial_{\gamma} \mathcal{A}^{\alpha \chi}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 4 \, f^{\alpha \beta}_{\alpha \alpha \chi} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} \, \partial_{\beta} \mathcal{B}^{\alpha \gamma}_{\alpha \gamma} - 2 \, \partial$$

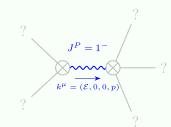
## Wave operator

$0^{\circ}\mathcal{B}$ $0^{\circ}\phi$ $0^{\circ}\mathcal{A}^{\parallel}$ $0^{\circ}f^{\parallel}$ $0^{\circ}f^{\perp}$ $0^{\circ}\mathcal{A}^{\parallel}$									
$^{+}\mathcal{B}\dagger -6\lambda . + \frac{v}{2}  0  \frac{12\lambdav}{2\sqrt{6}}  -\frac{ik(12\lambdav)}{2\sqrt{3}}  0$									
$ \phi \uparrow $ 0 0 0 0 0 0									
$\mathcal{A}^{\parallel} + \begin{vmatrix} \frac{12 \lambda - \nu}{2 \sqrt{6}} & 0 & -\lambda \cdot + \frac{\nu}{12} & \frac{i k (12 \lambda - \nu)}{6 \sqrt{2}} & 0 \end{vmatrix} = 0$									
$ik(12\lambda-\nu)$ $ik(12\lambda-\nu)$ $k^2\nu$ .									
$f^{\perp} + \begin{bmatrix} 2\sqrt{3} & 0 & -\frac{1}{6\sqrt{2}} & \frac{1}{6} \\ 0 & 0 & 0 & 0 \end{bmatrix}$									
$\mathcal{A}^{\parallel} \dagger \begin{array}{ccccccccccccccccccccccccccccccccccc$	1. A αβ	$^{1^{+}}\mathcal{H}^{\perp}{}_{lphaeta}{}^{1}$	$\dot{f}^{\parallel}_{\alpha\beta}$	$^{1}\mathcal{B}_{lpha}$	${}^{!}\mathcal{A}^{\parallel}{}_{lpha}$	${}^{1}\mathcal{A}^{\perp}{}_{lpha}$	$f^{\parallel}_{\alpha}$	$f^{\perp}_{\alpha}$	
1.* All + c			ikλ. - √2	0	0	0	0	0	
¹˙ <i>'</i> ⁄⁄⁄⁄⁄⁄⁄//////////////////////////////	$\beta = \frac{\lambda}{\sqrt{2}}$	0	0	0	0	0	0	0	
1.* <i>f</i>    † <sup>c</sup>		0	0	0	0	0	0	0	
! <i>B</i> †		0	0	$-6 \lambda_{\cdot} + \frac{\sqrt{2}}{2} + 2 k^2 (-c_{\cdot} + \xi_{\cdot})$	$\frac{1}{6} \left( -12 \lambda_{\cdot} + v_{\cdot} + k^2 \left( -5 c_{\cdot} + 4 \xi_{\cdot} \right) \right)$	$\frac{12 \lambda - v + 2 k^2 (c - 2 \xi)}{6 \sqrt{2}}$	0	$\frac{1}{6} ik (12 \lambda v.)$	
!' <i>Я</i> " †	α 0	0	0		)) $\frac{1}{18} \left( -9 \lambda. + v. + k^2 \left( -6 c. + 4 \xi. \right) \right)$				
!' <i>Я</i> ² †	α 0	0	0	$\frac{12 \lambda - v + 2 k^2 (c - 2 \xi)}{6 \sqrt{2}}$	$\frac{18 \lambda - \nu + k^2 (3 c - 4 \xi)}{18 \sqrt{2}}$	$\frac{1}{36} (v. + 4 k^2 \xi.)$	0	ikv. 18 √2	
¹.´f <sup>∥</sup> †	0	0	0	0	0	0	0	0	
¹.'f* †	σ 0	0	0	$k\left(-2i\lambda.+\frac{i\nu.}{6}\right)$	$k\left(-i\lambda.+\frac{i\nu}{18}\right)$	$-\frac{i k v}{18 \sqrt{2}}$	0	$\frac{1}{18} k^2 (v_1 + 4 k^2 \xi_2)$	${}^{2^{+}}\mathcal{A}^{\parallel}_{\alpha\beta}  {}^{2^{+}}f^{\parallel}_{\alpha\beta}  {}^{2^{-}}\mathcal{A}^{\parallel}_{\alpha\beta\chi}$
								²∹ <i>Я</i> <sup>∥</sup> † <sup>αβ</sup>	$\frac{\lambda}{2} - \frac{ik\lambda}{\sqrt{2}}$ 0
								$^{2^{+}}f^{\parallel}$ † $^{lphaeta}$	$\frac{ik\lambda}{\sqrt{2}}$ 0 0
								$^{2}\mathcal{H}^{\parallel}$ † $^{lphaeta\chi}$	$0  0  \frac{\lambda}{2}$
aturated propagator									

#### Source constraints

Spin-parity form	Covariant form	Multiplicities			
$0^+_{\cdot} \tau^{\perp} == 0$	$\partial_{\beta}\partial_{\alpha}\mathcal{J}^{\alpha\beta} == 0$	1			
$2^{0^{+}}\sigma^{\parallel} + {}^{0^{+}}\mathcal{J} == 0$	$\partial_{\alpha} \mathcal{J}^{\alpha} == 2  \partial_{\beta} \sigma^{\alpha \beta}_{\alpha}$	1			
$0^+ \rho == 0$	$\rho = 0$	1			
$1 \tau^{\parallel^{\alpha}} == 0$	$\partial_{\chi}\partial_{\beta}\partial^{\alpha}\mathcal{J}^{\beta\chi} == \partial_{\chi}\partial^{\chi}\partial_{\beta}\mathcal{J}^{\beta\alpha}$	3			
$2 \cdot \left[ \sigma^{\parallel} \right]^{\alpha} = 2 \cdot \left[ \sigma^{\perp} \right]^{\alpha} + \left[ \sigma^{\perp} \right]^{\alpha}$	$\partial_{\beta}\partial^{\alpha}\mathcal{J}^{\beta} = \partial_{\beta}\partial^{\beta}\mathcal{J}^{\alpha} + 2\left(\partial_{\chi}\partial^{\alpha}\sigma^{\beta}_{\beta}^{\chi} + \partial_{\chi}\partial^{\chi}\sigma^{\beta\alpha}_{\beta}\right)$	3			
$i k \cdot 1^+ \sigma^{\perp} \alpha^{\beta} + 1^+ \tau^{\parallel} \alpha^{\beta} == 0$	$\partial_{\chi}\partial^{\alpha}\mathcal{J}^{\beta\chi} + \partial_{\chi}\partial^{\beta}\mathcal{J}^{\chi\alpha} + \partial_{\chi}\partial^{\chi}\mathcal{J}^{\alpha\beta} + 2\partial_{\delta}\partial_{\chi}\partial^{\alpha}\sigma^{\chi\beta\delta} + 2\partial_{\delta}\partial^{\delta}\partial_{\chi}\sigma^{\chi\alpha\beta} = \partial_{\chi}\partial^{\alpha}\mathcal{J}^{\chi\beta} + \partial_{\chi}\partial^{\beta}\mathcal{J}^{\alpha\chi} + \partial_{\chi}\partial^{\chi}\mathcal{J}^{\beta\alpha} + 2\partial_{\delta}\partial_{\chi}\partial^{\beta}\sigma^{\chi\alpha\delta}$	3			
Total expected gauge generators:					

## Massive spectrum



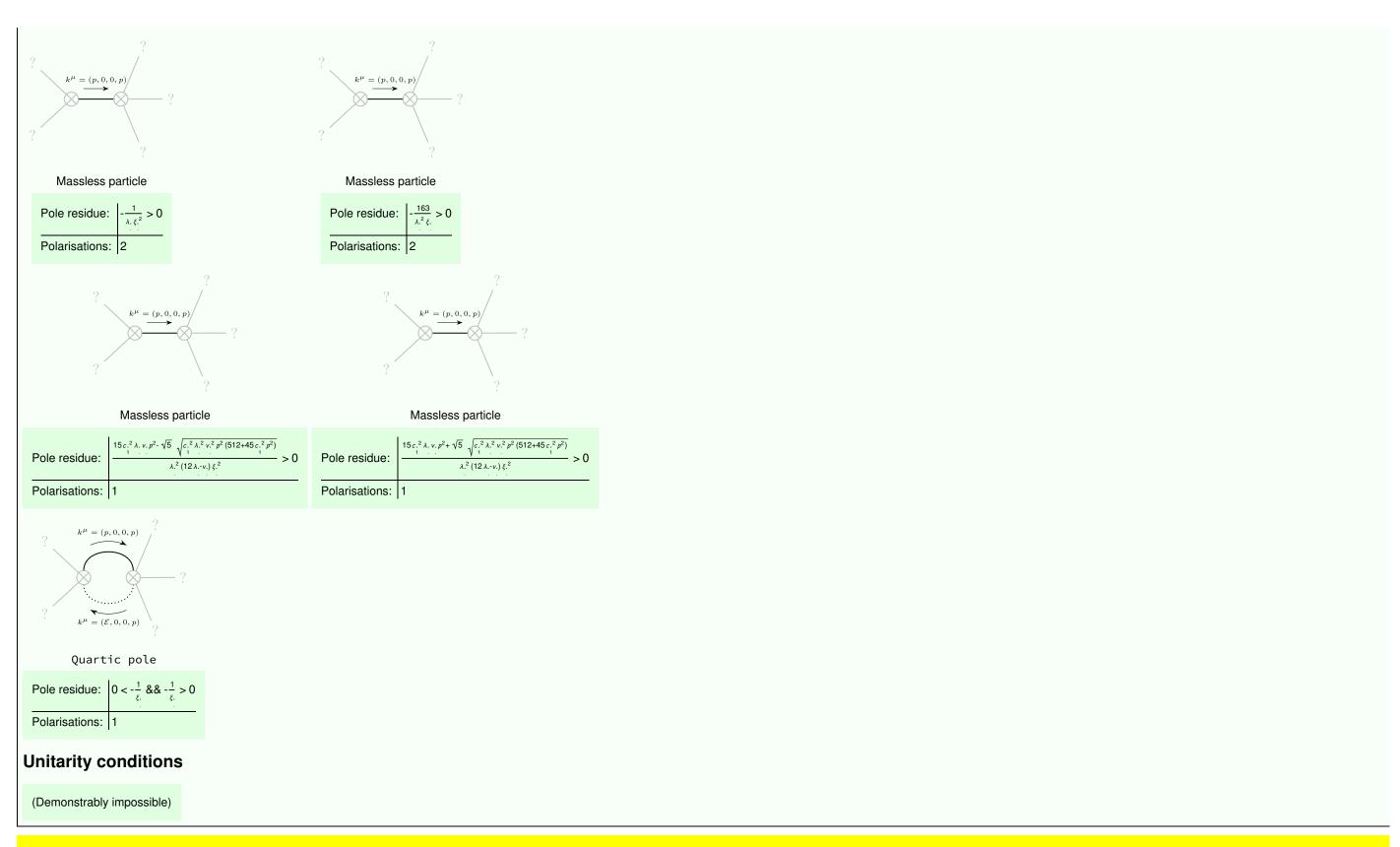
Massive particle

Pole residue:	$ \left(6\left(21c_{.}^{8}v_{.}^{2}+1008c_{.}^{.7}\lambda_{.}v_{.}\xi_{.}-4320c_{.}^{.6}\lambda_{.}^{2}v_{.}\xi_{.}-672c_{.}^{.6}\lambda_{.}v_{.}^{2}\xi_{.}-28c_{.}^{.6}v_{.}^{.3}\xi_{.}^{.2}+207360c_{.}^{.5}\lambda_{.}^{.3}\xi_{.}^{.2}+1008c_{.}^{.6}\lambda_{.}v_{.}^{2}\xi_{.}^{.2}-2688c_{.}^{.5}\lambda_{.}v_{.}^{2}\xi_{.}^{.2}+1755648c_{.}^{.4}\lambda_{.}^{.3}\xi_{.}^{.3}-19584c_{.}^{.4}\lambda_{.}^{.2}v_{.}\xi_{.}^{.3}-1344c_{.}^{.4}\lambda_{.}v_{.}^{.2}\xi_{.}^{.3}+1658880c_{.}^{.3}\lambda_{.}^{.3}k_{.}^{.2}$
	$3096576 \lambda.^{3} \xi.^{7} + 21 c.^{6} v. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 4320 c.^{4} \lambda.^{2} \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 672 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi 48 \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v. + 48 c. \lambda. \xi.^{2})^{2}} - 28 c.^{4} \lambda. v. \xi. \sqrt{384 c.^{2} \lambda. (-12 \lambda$
	$1344  c.^3  \lambda.  v.  \xi.^2  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} + 146304  c.^2  \lambda.^2  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} + 2688  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi 48  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi.^2)^2} - 19356  c.^2  \lambda.  v.  v.  \xi.^3  \sqrt{384  c.^2  \lambda.  (-12  \lambda. + v.)  \xi.^2 + (c.^2  v. + 48  c.  \lambda.  \xi.^2)^2} - 19356 $
	$(7c.^{2}(c.^{6}v.^{3} + 144c.^{5}\lambda.v.^{2}\xi. + 2304c.^{4}\lambda.^{2}v.\xi.^{2} + 240c.^{4}\lambda.v.^{2}\xi.^{2} + 240c.^{4}\lambda.v.^{2}\xi.^{2} - 110592c.^{3}\lambda.^{3}\xi.^{3} + 4608c.^{3}\lambda.^{2}v.\xi.^{3} - 110592c.^{2}\lambda.^{3}\xi.^{4} - 11520c.^{2}\lambda.^{2}v.\xi.^{4} + 331776c.\lambda.^{3}\xi.^{5} - 110592\lambda.^{3}\xi.^{6} + c.^{4}v.^{2}\sqrt{384c.^{2}\lambda.(-12\lambda.+v.)}\xi.^{2} + (c.^{2}v. + 48c.\lambda.\xi.^{-2})^{4}$
	$96c.^2\lambda.v.\xi.^2\sqrt{384c.^2\lambda.(-12\lambda.+v.)\xi.^2+(c.^2v.+48c.\lambda.\xi.-48\lambda.\xi.^2)^2}-4608c.\lambda.^2\xi.^3\sqrt{384c.^2\lambda.(-12\lambda.+v.)\xi.^2+(c.^2v.+48c.\lambda.\xi.-48\lambda.\xi.^2)^2}+2304\lambda.^2\xi.^4\sqrt{384c.^2\lambda.(-12\lambda.+v.)\xi.^2+(c.^2v.+48c.\lambda.\xi.-48\lambda.\xi.^2)^2}))>0$
Square mass:	$\frac{c.^{2} v.+48 c. \lambda. \xi48 \lambda. \xi.^{2} + \sqrt{384 c.^{2} \lambda. (-12 \lambda. + v.) \xi.^{2} + (c.^{2} v.+48 c. \lambda. \xi48 \lambda. \xi.^{2})^{2}}{1 \cdot 1 \cdot$

## Massless spectrum

Odd

Spin: Parity:



**Key observation:** This marks the completion of the particle spectrum analysis for the modified E--H action with propagating compensator and vector.

# Test case 4: Only propagating compensator and vector.

We test the case of preserving only the propagating compensator and vector B.

$$\frac{1}{2} v. \mathcal{D}^{\dagger} \phi_{\alpha} \mathcal{D}^{\dagger} \phi^{\alpha} + \xi. \mathcal{H}^{\dagger}_{\alpha\beta} \mathcal{H}^{\dagger^{\alpha\beta}}$$
(25)

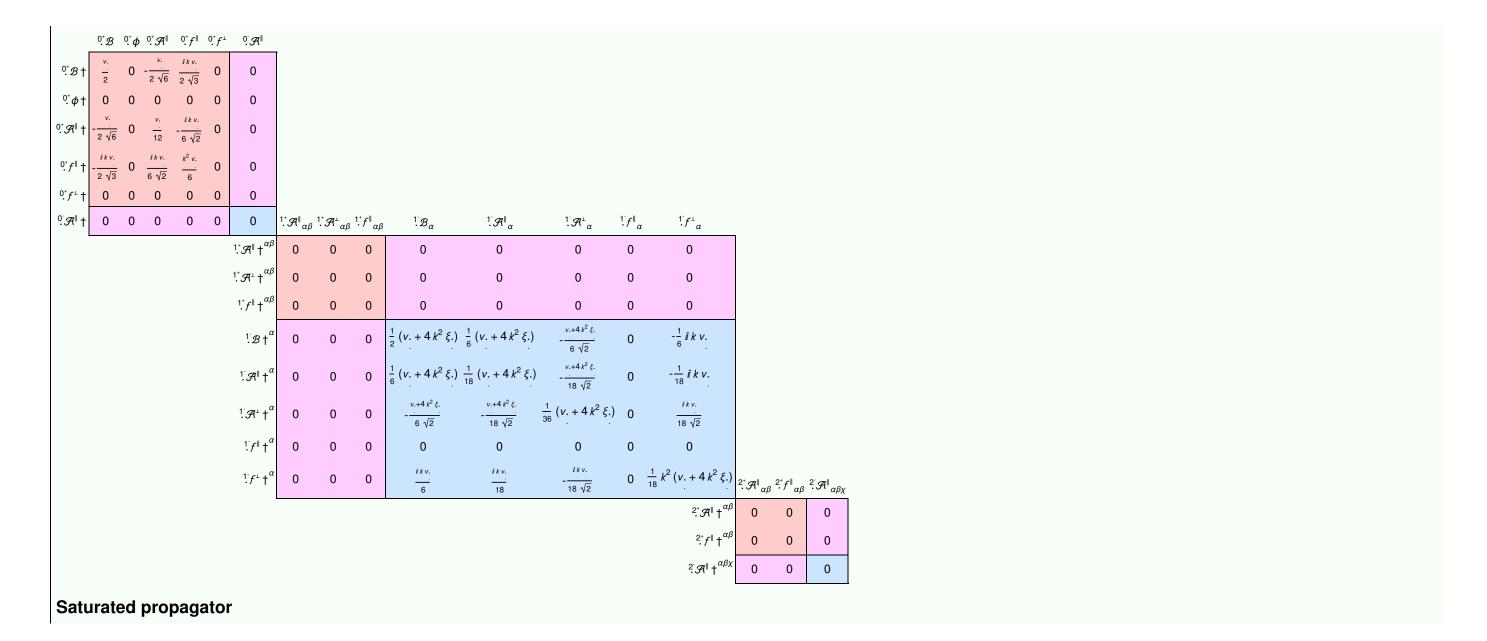
Here, we perform rescalings after application of Einstein Gauge:  $\phi_-0^2 \times \lambda - \lambda$ ,  $\phi_-0^2 \times \nu - \nu$ ,  $\phi_-0^2 \times t_-i$  -> t\_i. Also  $\phi_-0$  -> 1, i.e. here I am making the compensator dimensionless, any possible masses order 1. I do this to prevent any denominators phi/phi0.

Here is the linearised Lagrangian before feeding into ParticleSpectrum[].

$$-\frac{1}{18} v. \mathcal{A}^{\alpha\beta}_{\alpha} \mathcal{A}^{\chi}_{\beta} - \frac{1}{3} v. \mathcal{A}^{\beta}_{\alpha} \mathcal{B}^{\alpha} + \frac{1}{2} v. \mathcal{B}_{\alpha} \mathcal{B}^{\alpha} + \frac{1}{3} v. \mathcal{B}^{\alpha}_{\alpha} \partial_{\alpha} f^{\beta}_{\beta} - \frac{1}{3} v. \mathcal{B}^{\alpha}_{\alpha} \partial_{\beta} f^{\beta}_{\alpha} + \frac{1}{9} v. \mathcal{A}^{\chi}_{\alpha} \partial_{\beta} f^{\alpha\beta} - \frac{1}{9} v. \mathcal{A}^{\chi}_{\beta} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} + \frac{1}{18} v. \partial_{\beta} f^{\chi}_{\alpha} \partial^{\beta}_{\beta} f^{\alpha}_{\alpha} + \frac{4}{3} \xi. \partial_{\alpha} \mathcal{A}^{\chi}_{\beta} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} + \frac{4}{3} \xi. \partial_{\alpha} \mathcal{A}^{\chi}_{\beta} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} + \frac{4}{3} \xi. \partial^{\beta}_{\alpha} \mathcal{A}^{\alpha}_{\beta} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} + \frac{4}{3} \xi. \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} - \frac{2}{9} \xi. \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} - \frac{4}{3} \xi. \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} f^{\alpha}_{\alpha} - \frac{2}{9} \xi. \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_{\alpha} \partial^{\beta}_$$

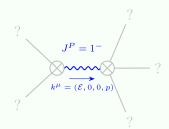
#### **PSALTer results panel**

## Wave operator



Total expected gauge generators:

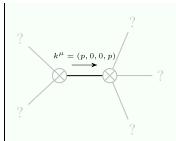
## **Massive spectrum**



Massive particle

Pole residue:	$\left  \frac{9}{2 v} - \frac{3}{14 \xi} \right  > 0$				
Square mass:	$-\frac{v}{2\xi} > 0$				
Spin:	1				
Parity:	Odd				

## **Massless spectrum**



Massless particle

Pole residue: Polarisations: 2

## **Unitarity conditions**

 $\xi$ . < 0 && v. > 0

**Key observation:** This marks the completion of the particle spectrum analysis for action with only propagating compensator and vector.

# Killing off the quartic pole

We will kill the quartic pole.

$$\frac{1}{2} \stackrel{\mathbf{v}}{\cdot} \mathcal{D}^{\dagger} \phi_{\alpha} \mathcal{D}^{\dagger} \phi^{\alpha} - c_{1} \mathcal{H}^{\dagger \alpha \beta} \mathcal{R}^{\dagger \chi}_{\alpha \beta \chi} + \lambda \stackrel{\mathbf{v}}{\cdot} \phi^{2} \mathcal{R}^{\dagger \alpha \beta}_{\alpha \beta} + \frac{1}{6} \left( 2 \stackrel{\mathbf{r}}{\cdot} + \stackrel{\mathbf{r}}{\cdot} \right) \mathcal{R}^{\dagger}_{\alpha \beta \chi \delta} \mathcal{R}^{\dagger \alpha \beta \chi \delta} + \frac{2}{3} \left( \stackrel{\mathbf{r}}{\cdot} - \stackrel{\mathbf{r}}{\cdot} \right) \mathcal{R}^{\dagger \alpha \beta \chi \delta} + \left( \stackrel{\mathbf{r}}{\cdot} \stackrel{\mathbf{r}}{\cdot} - \stackrel{\mathbf{r}}{\cdot} \right) \mathcal{R}^{\dagger \alpha \beta \chi \delta} \mathcal{R}^{\dagger \alpha \beta \chi \delta} + \left( \stackrel{\mathbf{r}}{\cdot} - \stackrel{\mathbf{r}}{\cdot} \right) \mathcal{R}^{\dagger \alpha \beta \chi \delta} \mathcal{R}^{\dagger \alpha \beta \chi \delta} + \frac{1}{6} \left( 3 \stackrel{\mathbf{r}}{\cdot} - \stackrel{\mathbf{r}}{\cdot} \right) \mathcal{R}^{\dagger \alpha \beta \chi \delta} \mathcal{R}^{\dagger \alpha \beta \chi \delta} + \left( \stackrel{\mathbf{r}}{\cdot} - \stackrel{\mathbf{r}}{\cdot} \right) \mathcal{R}^{\dagger \alpha \beta \chi \delta} \mathcal{R}^{\dagger \alpha \delta \chi \delta} + \frac{1}{12} \left( 3 \stackrel{\mathbf{r}}{\cdot} - \stackrel{\mathbf{r}}{\cdot} \right) \phi^{2} \mathcal{T}^{\dagger \alpha \beta \chi} \mathcal{T}^{\dagger \alpha \gamma \chi} \mathcal{T}^{\dagger \alpha \beta \chi} \mathcal{T}^{\dagger \alpha \gamma \chi} \mathcal{T}^$$

Out[4]= \$Aborted