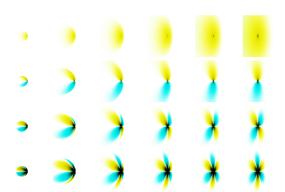
# Linear gravity from lightlike sources

#### W E V Barker

Cavendish Astrophysics Group Supervised by Anthony Lasenby, Mike Hobson & Will Handley



Null waves (e.g. light, GWs) carry energy and spin, so must act as sources of curvature and torsion: therefore what is the gravity of light?

- Null waves (e.g. light, GWs) carry energy and spin, so must act as sources of curvature and torsion: therefore what is the gravity of light?
- Tolman, Ehrenfest & Podolsky 1933 considered 'thin light pencils' (today: 'laser pulses'):

- Null waves (e.g. light, GWs) carry energy and spin, so must act as sources of curvature and torsion: therefore what is the gravity of light?
- Tolman, Ehrenfest & Podolsky 1933 considered 'thin light pencils' (today: 'laser pulses'):

"Parallel light pencils feel no mutual attraction..."

- Null waves (e.g. light, GWs) carry energy and spin, so must act as sources of curvature and torsion: therefore what is the gravity of light?
- Tolman, Ehrenfest & Podolsky 1933 considered 'thin light pencils' (today: 'laser pulses'):
  - "Parallel light pencils feel no mutual attraction..."
- Scully 1979 considered slow light in a dielectric using the Einstein-Maxwell equations:

- Null waves (e.g. light, GWs) carry energy and spin, so must act as sources of curvature and torsion: therefore what is the gravity of light?
- Tolman, Ehrenfest & Podolsky 1933 considered 'thin light pencils' (today: 'laser pulses'):

"Parallel light pencils feel no mutual attraction..."

Scully 1979 considered slow light in a dielectric using the Einstein-Maxwell equations:

"... but slow light pulses do."



■ Tempting to use linear gravity, i.e.

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + \mathcal{O}(h^2) \tag{1}$$

Tempting to use linear gravity, i.e.

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + \mathcal{O}(h^2) \tag{1}$$

The retarded Liénard Wiechert potential

$$h^{\mu\nu} - \frac{1}{2}h\eta^{\mu\nu} = -4\int \frac{T^{\mu\nu}(\overline{x})\delta(\overline{t} - t_{\rm r})}{|\overline{\mathbf{x}} - \mathbf{x}|} d^{4}\overline{x}$$
 (2)

Tempting to use linear gravity, i.e.

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu} + \mathcal{O}(h^2) \tag{1}$$

The retarded Liénard Wiechert potential

$$h^{\mu\nu} - \frac{1}{2}h\eta^{\mu\nu} = -4\int \frac{T^{\mu\nu}(\overline{x})\delta(\overline{t} - t_{\rm r})}{|\overline{\mathbf{x}} - \mathbf{x}|} d^{4}\overline{x}$$
 (2)

 However the notion of retarded time is complicated for sources moving at the fundamenetal velocity



 Very fast timelike sources seem to approach superluminally and slow down only as they near the observer

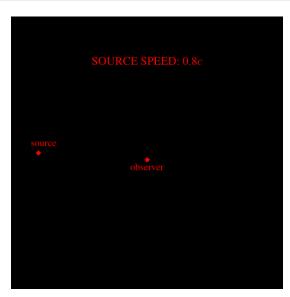
- Very fast timelike sources seem to approach superluminally and slow down only as they near the observer
- In the limiting case, lightlike sources have no retarded position before they pass by

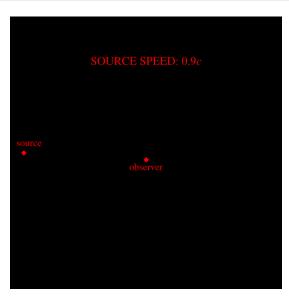
- Very fast timelike sources seem to approach superluminally and slow down only as they near the observer
- In the limiting case, lightlike sources have no retarded position before they pass by
- In traditional calculations at v = 1, advanced position implicitly used for half of all coordinate time...

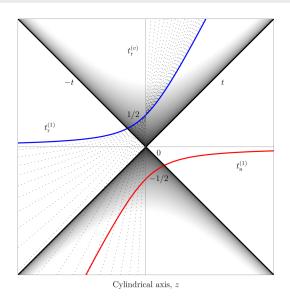
- Very fast timelike sources seem to approach superluminally and slow down only as they near the observer
- In the limiting case, lightlike sources have no retarded position before they pass by
- In traditional calculations at v = 1, advanced position implicitly used for half of all coordinate time...
- ... end up with gauge waves but no physical fields:

$$h_{\mu\nu} \propto k_{\mu} k_{\nu} f\left(k_{\alpha} x^{\alpha}\right) \tag{3}$$









 If a source falls eternally on a timelike geodesic its fields can be found via Lorentz transformation

- If a source falls eternally on a timelike geodesic its fields can be found via Lorentz transformation
- Aichelburg & Sexl 1971 found the exact fields of a massless point particle by ultraboosting the Schwarzschild metric

- If a source falls eternally on a timelike geodesic its fields can be found via Lorentz transformation
- Aichelburg & Sexl 1971 found the exact fields of a massless point particle by ultraboosting the Schwarzschild metric
- After a careful use of analysis, the perturbation in cylindrical coordinates is

$$h_{\mu\nu} \propto k_{\mu} k_{\nu} \ln \left( \rho \right) \delta \left( z - t \right) \tag{4}$$



- If a source falls eternally on a timelike geodesic its fields can be found via Lorentz transformation
- Aichelburg & Sexl 1971 found the exact fields of a massless point particle by ultraboosting the Schwarzschild metric
- After a careful use of analysis, the perturbation in cylindrical coordinates is

$$h_{\mu\nu} \propto k_{\mu} k_{\nu} \ln \left( \rho \right) \delta \left( z - t \right) \tag{4}$$

...i.e. a physical wave abreast of the particle



- If a source falls eternally on a timelike geodesic its fields can be found via Lorentz transformation
- Aichelburg & Sexl 1971 found the exact fields of a massless point particle by ultraboosting the Schwarzschild metric
- After a careful use of analysis, the perturbation in cylindrical coordinates is

$$h_{\mu\nu} \propto k_{\mu} k_{\nu} \ln \left( \rho \right) \delta \left( z - t \right) \tag{4}$$

- ...i.e. a physical wave abreast of the particle
- But we established this portion of spacetime was not associated with any finite retarded time!



We would like to extend the picture to general null sources of energy and spin which are compact and diffractionless

- We would like to extend the picture to general null sources of energy and spin which are compact and diffractionless
- In the limit of fast timelike sources it can be meaningful to consider the 'rest frame of light (or gravity)'

- We would like to extend the picture to general null sources of energy and spin which are compact and diffractionless
- In the limit of fast timelike sources it can be meaningful to consider the 'rest frame of light (or gravity)'
- In the rest frame, the source energy gives gravitoelectric fields and the spin gravitomagnetic fields

$$h_{\mu\nu} = e_{\mu\nu} + h_{\mu\nu} \tag{5}$$

- We would like to extend the picture to general null sources of energy and spin which are compact and diffractionless
- In the limit of fast timelike sources it can be meaningful to consider the 'rest frame of light (or gravity)'
- In the rest frame, the source energy gives gravitoelectric fields and the spin gravitomagnetic fields

$$h_{\mu\nu} = e_{\mu\nu} + h_{\mu\nu} \tag{5}$$

 An ultraboost takes us back to the laboratory frame and the GEM multipoles become flattened abreast of the source trajectory



Final result for the electrical fields

$$e^{\mu\nu} = -8\frac{k^{\mu}k^{\nu}}{k^{2}} \left[ a_{0}^{(\Phi)}(z-t) \ln \frac{1}{\rho} + \sum_{l=1}^{\infty} \frac{a_{l}^{(\Phi)}(z-t) \cos l\varphi + b_{l}^{(\Phi)}(z-t) \sin l\varphi}{l\rho^{l}} \right],$$

$$(6)$$

Final result for the electrical fields

$$e^{\mu\nu} = -8\frac{k^{\mu}k^{\nu}}{k^{2}} \left[ a_{0}^{(\Phi)}(z-t) \ln \frac{1}{\rho} + \sum_{l=1}^{\infty} \frac{a_{l}^{(\Phi)}(z-t) \cos l\varphi + b_{l}^{(\Phi)}(z-t) \sin l\varphi}{l\rho^{l}} \right],$$

$$(6)$$

• i.e. the gravitational fields are a 2D transverse multipole expansion for each slice of the null source

Final result for the electrical fields

$$e^{\mu\nu} = -8\frac{k^{\mu}k^{\nu}}{k^{2}} \left[ a_{0}^{(\Phi)}(z-t) \ln \frac{1}{\rho} + \sum_{l=1}^{\infty} \frac{a_{l}^{(\Phi)}(z-t) \cos l\varphi + b_{l}^{(\Phi)}(z-t) \sin l\varphi}{l\rho^{l}} \right],$$

$$(6)$$

- i.e. the gravitational fields are a 2D transverse multipole expansion for each slice of the null source
- Magnetic fields are more complicated, but have the same 2D structure



■ The fields depend in an obvious way on the sources, but the role of retarded time is still unclear

- The fields depend in an obvious way on the sources, but the role of retarded time is still unclear
- Applying the retarded potential to the fast limit of eternally falling timelike sources, the physical fields and gauge waves appear together correctly, the same fields as given by the ultraboost

- The fields depend in an obvious way on the sources, but the role of retarded time is still unclear
- Applying the retarded potential to the fast limit of eternally falling timelike sources, the physical fields and gauge waves appear together correctly, the same fields as given by the ultraboost

 Everything becomes clear if the source is produced by a definite emission process

- Everything becomes clear if the source is produced by a definite emission process
- For an abruptly emitted delta-pulse source, the metric perturbation is confined to the spherical shell that expands at v=1 from the emission point...

- Everything becomes clear if the source is produced by a definite emission process
- For an abruptly emitted delta-pulse source, the metric perturbation is confined to the spherical shell that expands at v=1 from the emission point...
- ... this shell also contains the pulse

- Everything becomes clear if the source is produced by a definite emission process
- For an abruptly emitted delta-pulse source, the metric perturbation is confined to the spherical shell that expands at v=1 from the emission point...
- ...this shell also contains the pulse
- According to the retarded potential, the emission signal embedded in this shell is **not radiative**

- Everything becomes clear if the source is produced by a definite emission process
- For an abruptly emitted delta-pulse source, the metric perturbation is confined to the spherical shell that expands at v=1 from the emission point...
- ...this shell also contains the pulse
- According to the retarded potential, the emission signal embedded in this shell is **not radiative**
- At large times and close to the pulse trajectory, the known gravitational fields are resurrected



# Causal disconnect and the emission mirage

 Hence the gravitational fields of lightlike sources are just an afterglow of their emission process

# Causal disconnect and the emission mirage

- Hence the gravitational fields of lightlike sources are just an afterglow of their emission process
- The fields accompany the sources as if produced by them, but can be thought of as inevitable 'decoration'

# Causal disconnect and the emission mirage

- Hence the gravitational fields of lightlike sources are just an afterglow of their emission process
- The fields accompany the sources as if produced by them, but can be thought of as inevitable 'decoration'
- Observable field strength:

Equivalent gravity of a half-infinite beam of null waves with constant spin/energy profile equal to that of the part of the source undergoing emission at retarded time



## Test geodesics

Magnetic fields perturb the Lagrangian of a test particle, e.g. for a thin circularly polarized wavepulse of intensity profile f(z)

$$\mathcal{L} = \frac{1}{2} \left( \dot{w}_{+} \dot{w}_{-} - \dot{\rho}^{2} - \rho^{2} \dot{\varphi}^{2} \right) - f(w_{-}) \dot{w}_{-} \dot{\varphi}$$
 (7)

## Test geodesics

Magnetic fields perturb the Lagrangian of a test particle, e.g. for a thin circularly polarized wavepulse of intensity profile f(z)

$$\mathcal{L} = \frac{1}{2} \left( \dot{w}_{+} \dot{w}_{-} - \dot{\rho}^{2} - \rho^{2} \dot{\varphi}^{2} \right) - f \left( w_{-} \right) \dot{w}_{-} \dot{\varphi} \tag{7}$$

In turn the radial equation of motion is perturbed

$$\ddot{\rho} = \frac{J^2}{\rho^3} - 2\frac{J}{\rho^3} f(w_-) \tag{8}$$

### Test geodesics

Magnetic fields perturb the Lagrangian of a test particle, e.g. for a thin circularly polarized wavepulse of intensity profile f(z)

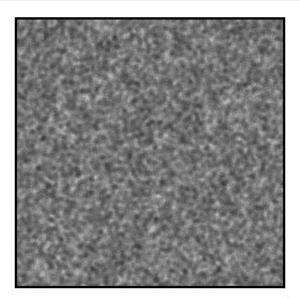
$$\mathcal{L} = \frac{1}{2} \left( \dot{w}_{+} \dot{w}_{-} - \dot{\rho}^{2} - \rho^{2} \dot{\varphi}^{2} \right) - f(w_{-}) \dot{w}_{-} \dot{\varphi}$$
 (7)

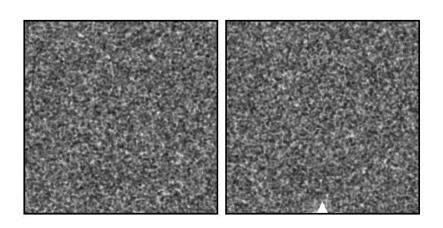
In turn the radial equation of motion is perturbed

$$\ddot{\rho} = \frac{J^2}{\rho^3} - 2\frac{J}{\rho^3} f(w_-) \tag{8}$$

■ Hence spinning waves imply velocity memory







#### Outlook?

Maximum wake diameter independent of gas temperature

$$\rho_0 \approx \sqrt{GE\lambda}/c^2 \tag{9}$$

#### Outlook?

Maximum wake diameter independent of gas temperature

$$\rho_0 \approx \sqrt{GE\lambda}/c^2 \tag{9}$$

■ But this is far too small, most powerful optical laser (LFEX in Japan) fires  $E=10^9 {\rm erg}$  pulses and would carve a  $10^-24 {\rm m}$  hole in a cold gas!

#### Outlook?

Maximum wake diameter independent of gas temperature

$$\rho_0 \approx \sqrt{GE\lambda}/c^2 \tag{9}$$

- But this is far too small, most powerful optical laser (LFEX in Japan) fires  $E=10^9{\rm erg}$  pulses and would carve a  $10^-24{\rm m}$  hole in a cold gas!
- But the process might play a role in the depolarization and absorbtion of gravitational waves passing through the ISM, in this context the calculation is an exercise in quadratic gravity