

Chapter 31-5: Elliptic PDEs Using the Monte Carlo Method.

Monte Carlo Methods are a class of computational solutions that can be applied to various problems which rely on repeated random sampling to provide generally approximate solutions. Monte Carlo is a stochastic approach, in which a series of simulations (trials), representing the analyzed problem, with randomly selected input values, are performed. Among these trials, a specified number of properly defined successes is achieved. The ratio between the number of success trials to the number of all trials, scaled by dimensional quantity (e.g., area or function value) allows for the estimation of the unknown solution, providing the number of trials is large enough.

A capacitor is constructed out of two infinite metal plates spaced 10 cm apart from one another. One plate is placed at 5V potential with respect to the other and the space between the plates is free of charges. The number of random walks is taken to be 400 and the number of lattice points is taken to be 30.

(The description of the problem and the code solution was taken from the repository of s-ankur.)

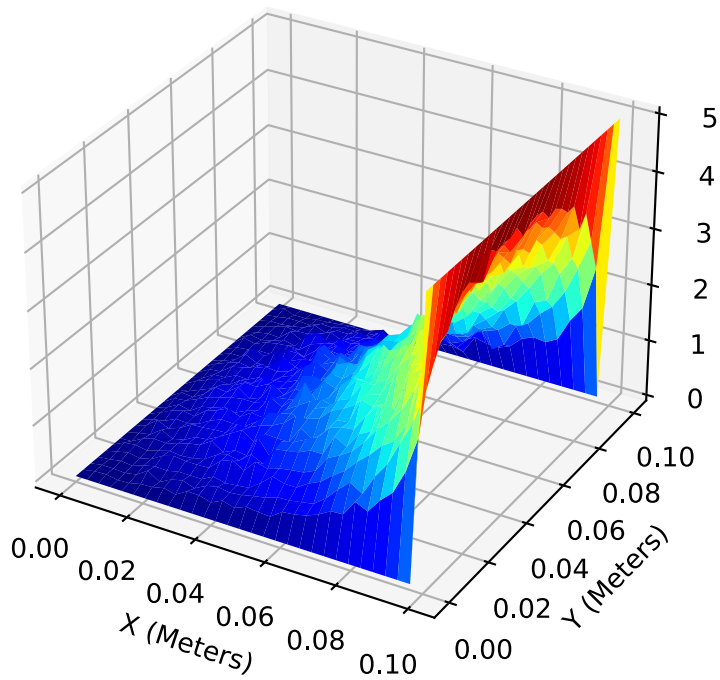
```
In [3]: 1 """
2 All calculations in SI units
3 """
4 import random
5 from mpl_toolkits.mplot3d import axes3d
6 import matplotlib.pyplot as plt
7 from matplotlib import cm
8 import numpy as np
9 %config InlineBackend.figure_formats = ['svg']
10
11 h = 10e-2 # Distance between plates = 10cm
12 lattice_points = 30 # Number of Points in lattice
13 d = h / lattice_points # Lattice size
14 boundary_voltage_high = 5.0 # 5 Volts at Positive Plate
15 boundary_voltage_low = 0.0 # 0 Volts at Negative Plate
16 epsilon_naught = 8.854e-12 # Permittivity of Vacuum
17 charge_density = 6e-16 # Coulomb per meter cube
18 N = 400 # Number of Random Walks
19
20
21 def f(x):
22     # The Function \nabla^2(phi) = f
23     # For Laplace f = 0
24     return 0
25
26
27 def g(x):
28     # Two Dimensional Boundary Conditions: two parallel metal plates at x=0,x=h
29     # the plate at x=h is at high potential and x=0 is low potential
30     # Assume that there are metal plates along y=0 and y=h (uncharged)
31     # this is because I dont know how to simulate open boundry conditions
32     if x[0] <= 0:
33         return boundary_voltage_low
34     if x[0] >= h:
35         return boundary_voltage_high
36     if x[1] <= 0 or x[1] >= h:
37         return boundary_voltage_low
38
39 """
40 def f_2(x):
41     # Alternative charge distribution: A charged Sphere in the centre of metal box
42     if (h / 2 - x[0])** 2 + (h / 2 - x[1])** 2 <= (h / 5) ** 2:
43         return -charge_density * 5 / epsilon_naught
44     else:
45         return 0
46
47
48 def g_2(x):
49     # Two Dimentional Alternative Boundary Conditions: uncharged metal box
50     return 0
51 """
52
53
54 def f_3(x):
55     # Alternative charge distribution: TWO charged Sphere in the centre of metal box
56     if (h / 3 - x[0])** 2 + (h / 2 - x[1])** 2 <= (h / 5) ** 2:
57         return -charge_density * 5 / epsilon_naught
58     if (2 * h / 3 - x[0])** 2 + (h / 2 - x[1])** 2 <= (h / 5) ** 2:
59         return charge_density * 5 / epsilon_naught
60     else:
61         return 0
62
63
```

```

64 @np.vectorize
65 def poisson_approximation_fixed_step(*A):
66     # Returns the Value of Potential Feild at a given point A with N random walks
67     result = 0
68     F = 0
69     for i in range(N):
70         x = list(A)
71         while True:
72             if x[0] <= 0 or x[0] >= h or x[1] <= 0 or x[1] >= h:
73                 break
74             random_number = random.randint(0, 3)
75             if random_number == 0:
76                 x[0] += d
77             elif random_number == 1:
78                 x[0] -= d
79             elif random_number == 2:
80                 x[1] += d
81             elif random_number == 3:
82                 x[1] -= d
83             F += f(x) * h ** 2
84         result += g(x) / N
85     result = result - F
86     return result
87
88
89 def plot(x, y, z):
90     # Function for plotting the potential
91     fig = plt.figure()
92     ax = fig.add_subplot(111, projection="3d")
93
94     ax.plot_surface(x, y, np.array(z), cmap=cm.jet, linewidth=0.1)
95     plt.xlabel("X (Meters)")
96     plt.ylabel("Y (Meters)")
97     ax.set_zlabel("Potential (Volts)")
98     plt.show()
99
100
101 if __name__ == "__main__":
102     # Experiment E: 2D Capacitor
103     print(
104         f"Calculating Monte Carlo with {lattice_points}x{lattice_points} lattice points and {N} r
105     )
106     lattice_x, lattice_y = np.mgrid[
107         0 : h : lattice_points * 1j, 0 : h : lattice_points * 1j
108     ]
109     z = poisson_approximation_fixed_step(lattice_x.ravel(), lattice_y.ravel()).reshape(
110         lattice_x.shape
111     )
112     plot(lattice_x, lattice_y, z)
113
114     # Experiment F: Metal box with positively charged metal ball inside
115     # f = f2
116     # g = g2
117     print(
118         f"Calculating Monte Carlo with {lattice_points}x{lattice_points} lattice points and {N} r
119     )
120     lattice_x, lattice_y = np.mgrid[
121         0 : h : lattice_points * 1j, 0 : h : lattice_points * 1j
122     ]
123     z = poisson_approximation_fixed_step(lattice_x.ravel(), lattice_y.ravel()).reshape(
124         lattice_x.shape
125     )
126     plot(lattice_x, lattice_y, z)
127
128     # Experiment G: Metal Box with two spheres (positive and negative)
129
130     f = f_3
131     g = g_2
132     print(
133         f"Calculating Monte Carlo with {lattice_points}x{lattice_points} lattice points and {N} r
134     )
135     lattice_x, lattice_y = np.mgrid[
136         0 : h : lattice_points * 1j, 0 : h : lattice_points * 1j
137     ]
138     z = poisson_approximation_fixed_step(lattice_x.ravel(), lattice_y.ravel()).reshape(
139         lattice_x.shape
140     )
141     plot(lattice_x, lattice_y, z)
142
143
144
145
146
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148
149
150
151
152

```

Calculating Monte Carlo with 30x30 lattice points and 400 random walks



```

-----
NameError                                Traceback (most recent call last)
Cell In[3], line 118
    112 plot(lattice_x, lattice_y, z)
    114 # Experiment F: Metal box with positively charged metal ball inside
    115 #     f = f2
    116 #     g = g2
    117 print(
--> 118     f"Calculating Monte Carlo with {lattice_points}x{lattice_points} lattice points and
{N} random walks for {'Laplace' if laplace else 'Poisson'}"
    119 )
    120 lattice_x, lattice_y = np.mgrid[
    121     0 : h : lattice_points * 1j, 0 : h : lattice_points * 1j
    122 ]
    123 z = poisson_approximation_fixed_step(lattice_x.ravel(), lattice_y.ravel()).reshape(
    124     lattice_x.shape
    125 )

```

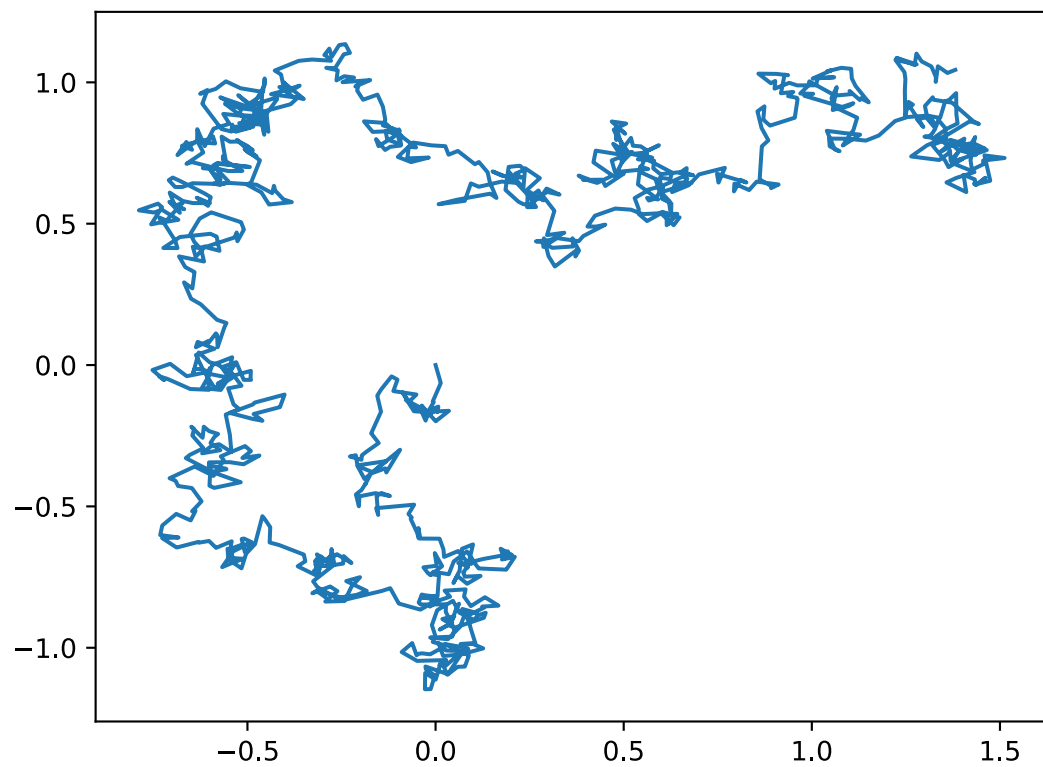
NameError: name 'laplace' is not defined

At the address: <https://srome.github.io/On-Solving-Partial-Differential-Equations-with-Brownian-Motion-in-Python/> (<https://srome.github.io/On-Solving-Partial-Differential-Equations-with-Brownian-Motion-in-Python/>) there is an interesting blog about doing pde solving with Brownian motion approximation. However, BM is not exactly Monte Carlo, is it? Actually, the two stochastic approaches are sometimes covered in the same place, for example: http://www.columbia.edu/~mh2078/MonteCarlo/MCS_Generate_RVars.pdf (http://www.columbia.edu/~mh2078/MonteCarlo/MCS_Generate_RVars.pdf). So it will be assumed here that using BM is essentially the same as using MC. And besides, it took 7.5 hrs for a 13700K processor to spit out the graphics for it, so they are not about to be wasted.

Perhaps a place to reaffirm that the Laplace equation is in the parabolic category.

```
In [7]: 1  ##matplotlib inline
2  import numpy as np
3  import matplotlib.pyplot as plt
4
5  t=1
6  n = 1000
7  delta_t = np.sqrt(t/n)
8
9  # Create each dW step
10 dw1 = delta_t * np.random.normal(0,1,size=n)
11 dw2 = delta_t * np.random.normal(0,1,size=n)
12 W = np.zeros((n+1,2))
13
14 # Add W_{j-1} + dW_{j-1}
15 W[1:,0] = np.cumsum(dw1)
16 W[1:,1] = np.cumsum(dw2)
17
18 plt.plot(W[:,0],W[:,1], '-')
19
20
```

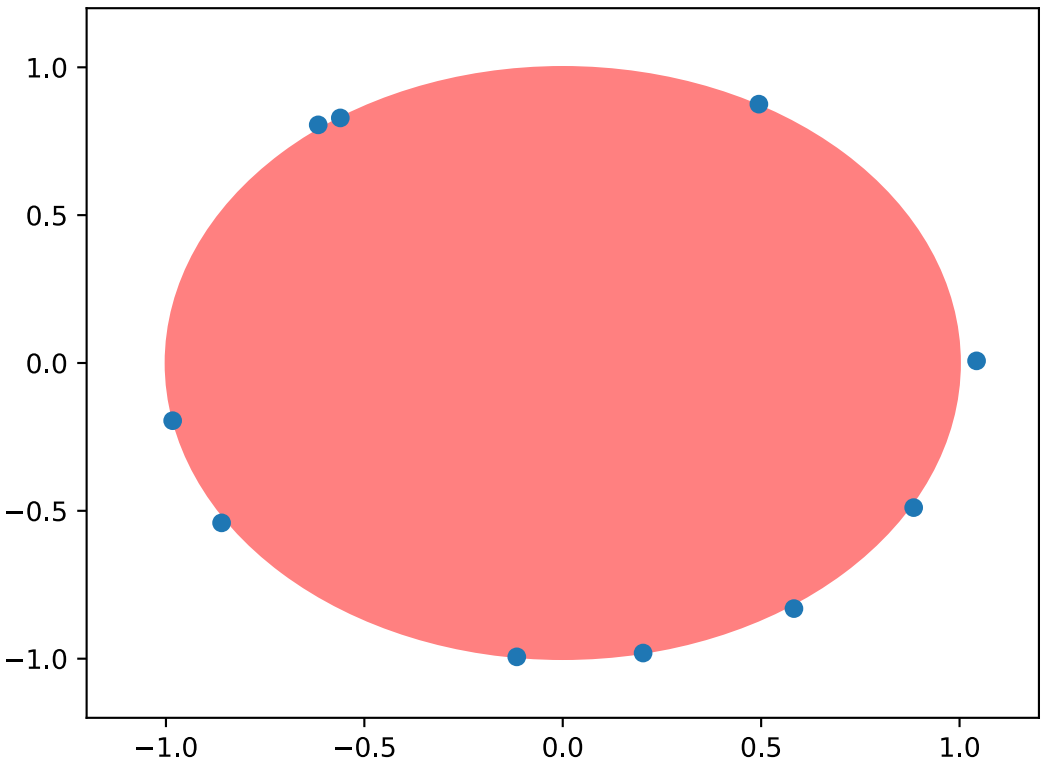
Out[7]: [



Above: a random Brownian walk, which has the same appearance as a Monte Carlo generated walk.

```
In [8]: 1 def check_if_exit(v):
2         # Takes a vector v=(x,y)
3         # Checks if v has intersected with the boundary of D
4         if np.linalg.norm(v,2) >=1:
5             return True
6         return False
7
8 def simulate_exit_time(v):
9     # Simulates exit time starting at v=(x,y), returns exit position
10    delta_t = np.sqrt(.001)
11    exit = False
12    x = v.copy()
13    while not exit:
14        x = x + delta_t * np.random.normal(0,1,size=2)
15        exit = check_if_exit(x)
16    return x
17
18 v=np.array((0,0)) # The origin
19 exit_times = np.array([simulate_exit_time(v) for k in range(0,10) ])
20
21
22 circle1=plt.Circle((0,0),1,color='r', alpha=.5)
23 plt.gcf().gca().add_artist(circle1)
24 plt.axis([-1.2, 1.2, -1.2, 1.2])
25 plt.scatter(exit_times[:,0],exit_times[:,1])
26
27
```

Out[8]: <matplotlib.collections.PathCollection at 0x2a2a8a943a0>



Above: simulated exit points on an ellipse, fairly similar to the epsilon criterion for deciding when a random step reaches the domain boundary in the WOS (walk on spheres) version of popular pde solvers.

In []:

```
In [4]: 1 np.random.seed(8) #Side Infinity
2
3 def check_if_exit(v):
4     # Takes a vector v=(x,y)
5     # Checks if v has intersected with the boundary of D
6     if np.linalg.norm(v,2) >=1:
7         return True
8     return False
9
10 def simulate_exit_time(v):
11     # Simulates exit time starting at v=(x,y), returns exit position
12     delta_t = np.sqrt(.001)
13     exit = False
14
15     # Copy because simulation modifies in place
16     if hasattr(v,'copy'): # For NumPy arrays
17         x = v.copy()
18     else:
19         x = np.array(v) # We input a non-NumPy array
20     while not exit:
21         x += delta_t * np.random.normal(0,1,size=2) # += modifies in place
22         exit = check_if_exit(x)
23     return x
24
25 v=np.array((.5,.5)) # The origin
26 u = lambda x : np.linalg.norm(x,2)*np.cos(np.arctan2(x[1],x[0]))
27 f = lambda x : np.cos(np.arctan2(x[1],x[0]))
```

```
28
29     def get_exp_f_exit(starting_point, n_trials):
30         return np.mean([f(simulate_exit_time(starting_point)) for k in range(0,n_trials)])
31
32     exp_f_exit = get_exp_f_exit(v,2000) # Expected value of f(Exit(x,d))
33     print('The value u(v) = %s\nThe value of Exp(f(Exit))=%s' %(u(v), exp_f_exit))
34
35
```

The value u(v) = 0.500000000000000001

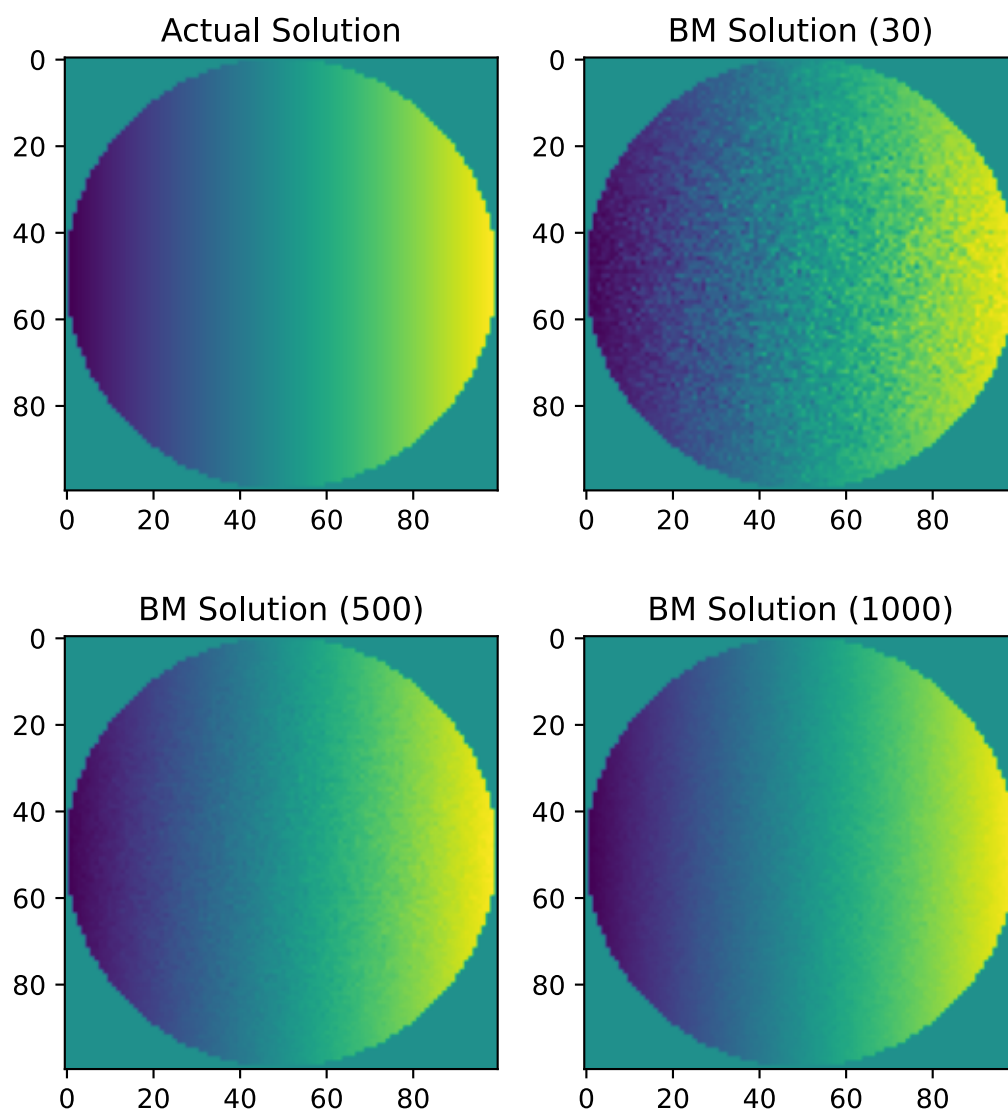
The value of Exp(f(Exit))=0.4975601145395467

```

In [5]: 1  lin = np.linspace(-1, 1, 100)
2  x, y = np.meshgrid(lin, lin)
3  print(x.shape)
4  u_vec = np.zeros(x.shape)
5  bm_vec_30 = np.zeros(x.shape)
6  bm_vec_500 = np.zeros(x.shape)
7  bm_vec_1000 = np.zeros(x.shape)
8
9  # Convert u to a solution in x,y coordinates
10 u_x = lambda x,y : np.linalg.norm(np.array([x,y]),2)*np.cos(np.arctan2(y,x))
11
12 # Calculate actual and approximate solution for (x,y) in D
13 for k in range(0,x.shape[0]):
14     for j in range(0,x.shape[1]):
15         x_t = x[k,j]
16         y_t = y[k,j]
17
18         # If the point is outside the circle, the solution is undefined
19         if np.sqrt((x_t)**2 + (y_t)**2) > 1:
20             continue
21
22         # Calculate function value at this point for each image
23         u_vec[k,j] = u_x(x_t,y_t)
24         bm_vec_30[k,j] = get_exp_f_exit((x_t,y_t),30)
25         bm_vec_500[k,j] = get_exp_f_exit((x_t,y_t),500)
26         bm_vec_1000[k,j] = get_exp_f_exit((x_t,y_t),1000)
27
28 fig = plt.figure()
29 ax = fig.add_subplot(121)
30 plt.imshow(u_vec)
31 plt.title('Actual Solution')
32
33 ax = fig.add_subplot(122)
34 plt.title('BM Solution (30)')
35 plt.imshow(bm_vec_30)
36
37 fig = plt.figure()
38
39 ax = fig.add_subplot(121)
40 plt.title('BM Solution (500)')
41 plt.imshow(bm_vec_500)
42
43 ax = fig.add_subplot(122)
44 plt.title('BM Solution (1000)')
45 plt.imshow(bm_vec_1000)
46
47 (100, 100)

```

Out[5]: <matplotlib.image.AxesImage at 0x2a2a852cc10>



Above: Brownian motion solutions, rather similar to the walk on spheres output for a parabolic pde performed with Monte Carlo.