

FYS3150

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## Analytical solution

The diffusion equations is given as:

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2}$$

We want the boundary conditions equal to 0, so we define a new function

$$v(x, t) = u(x, t) - u_s(x, t)$$

and solve the diffusion equation for this function. This equation can be solved by separation of variables. Setting  $v(x, t) = X(x)T(t)$  gives:

$$\frac{dT}{dt} \frac{1}{T(t)} \frac{1}{D^2} = \frac{1}{X(x)} \frac{d^2 X}{dx^2}$$

Both LHS and RHS must be equal to the same constant for this equation to hold at all times. We get:

$$\frac{dT}{dt} \frac{1}{T(t)} \frac{1}{D^2} = -\lambda^2 \text{ and } \frac{1}{X(x)} \frac{d^2 X}{dx^2} = -\lambda^2$$

The solution for the time equation is

$$T(t) = e^{-D^2 \lambda^2 t}$$

while the spacial solution is

$$X(x) = A \sin(\lambda x) + B \cos(\lambda x)$$

Setting  $D = 1$ . By the boundary conditions  $v(0, t) = v(1, t) = 0$  we see that  $B = 0$  and that  $\lambda = n\pi$ . So we get

$$v(x, t) = \sum_{n=1}^{\infty} C_n \sin(n\pi x) e^{-n^2 \pi^2 t}$$

The initial condition is given on the form  $Ax + b$ , so  $u_s(x) = 1 - x$  fulfills the boundary conditions. We have defined  $v(x, t) = u(x, t) - u_s(x)$  so  $u(x, t) = v(x, t) + u_s(x)$ ,

$$u(x, t) = 1 - x + \sum_{n=1}^{\infty} C_n \sin(n\pi x) e^{-n^2 \pi^2 t}$$

To find the  $C_n$ s, I solve the equation for time equal to 0. I know that  $u(x, 0) = 0$  for  $0 < x < 1$ , so

$$v(x, 0) = \sum_{n=1}^{\infty} C_n \sin(n\pi x) = x - 1$$

Using the fourier expansion we see that

$$C_n = \frac{-2}{\pi n}$$

and we get the full solution

$$v(x, t) = \sum_{n=1}^{\infty} \frac{-2}{\pi n} \sin(n\pi x) e^{-n^2 \pi^2 t}$$

Since this should be solved for  $u(x, t)$ , we get

$$u(x, t) = \sum_{n=1}^{\infty} \frac{-2}{\pi n} \sin(n\pi x) e^{-n^2 \pi^2 t} + x - 1$$

# Algorithms

I will use three different schemes to solve the diffusion equation. First I have implemented the explicit scheme, then the implicit scheme, and at last I will use Crank-Nicolson.

## The explicit scheme

The Explicit scheme is a very straight-forward coding of the following equation.

$$v_{i,j+1} = \alpha v_{i-1,j} + (1 - 2\alpha)v_{i,j} + \alpha v_{i+1,j}$$

where

$$\alpha = \frac{\Delta t}{\Delta x^2}$$

and the error going as

$$O(\Delta t, \Delta x^2)$$

One can see that without boundary and initial conditions, this equation is unsolvable, but since we have the conditions given, we do not need to compute for  $j = 0$  and  $i = 0$  or  $i = j$ .

$$\text{boundary conditions: } v(0) = v(L) = 0$$

Since we are solving for  $v(x, t)$ , we need to use the initial state given from the analytical section above:

$$v(x, 0) = x - 1, \text{ for } 0 < x < 1 \text{ and } v(0, 0) = v(1, 0) = 0$$

## The Implicit scheme

The implicit scheme is obtained by using the backward formula for the first derivative. This will result in the matrix-vector multiplication:

$$Av_j = v_{j-1}$$

where A is a tridiagonal matrix with off-diagonal elements  $-\alpha$  and diagonal elements  $1 + 2\alpha$

I solve this equation by using the solver from project 1 by setting  $v_0$  by the initial conditions and the boundary conditions. Then I solve the equation  $N_t$  times to .

**The Crank-Nicolson scheme** The Crank-Nicolson scheme is a combination of the implicit and the explicit schemes. This scheme is shortly written as

$$(2I + \alpha B)V_j = (2I - \alpha B)V_{j-1}$$

## Numerical Notes

When using  $\Delta x = \frac{1}{10}$  and  $\Delta t = \frac{1}{100}$ , the Explicit scheme will go berzerk and give horrible results. This is because this scheme is unstable for  $\Delta t \geq \frac{1}{2}\Delta x^2$ . So when increasing  $\Delta t$  to  $\Delta t = \frac{1}{200}$  I get stable results.

None of the numerical approximations manage to overlap with the analytical solution, even though Crank-Nicolson is pretty close. The error for Crank-Nicolson is smaller because the error goes as  $\Delta t^2$  and  $\Delta x^2$ , but for explicit and implicit scheme, the error goes as  $\Delta x^2$  and  $\Delta t^2$ .

## Results and plots



