Kepler's laws: a project for Calculus

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Abstract:

This article is in order to derive Kepler's three laws for planetary motion with the method which we have learnt in university's Calculus class. What we do is the same to do what Newton did several hundreds of years. To derive these laws, we may used the law of universal gravitation, Newton's laws, vector differential and integral polar coordinates, cross and dot product in vectors and so on. At last, we used Kepler's laws to compute the height of synchronous satellites from the earth.

Keywords:

Kepler's laws, gravitation, Newton laws, ellipse, orbits, differential, velocity, radial vector, geosynchronous satellite

1. About Kepler's laws

In celestial mechanics, Kepler laws describe the motion of an orbiting body as an ellipse, parabola, or hyperbola, which forms a two-dimensional orbital plane in three-dimensional space. It considers only the gravitational attraction of two bodies, neglecting interactions with other objects, atmospheric drag, and so on. It is said to be a solution of a special case of the two-body problem, known as the Kepler problem. As a theory in classical mechanics, it also does not take into account the effects of general relativity.

These three laws are following:

Kepler's elliptical orbit law (first law)

The planets orbit the sun in elliptical orbits with the sun at one focus.

Kepler's equal-area law (second law)

The line connecting a planet to the sun sweeps out equal areas in equal amounts of time.

Kepler's law of periods (third law)

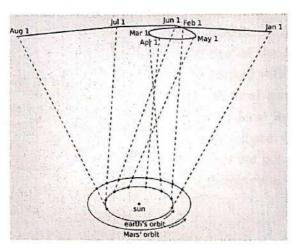
The time required for a planet to orbit the sun, called its period, is proportional to the long axis of the ellipse raised to the 3/2 power. The constant of proportionality is the same for all the planets.

The second law was discovered before the first one, for it can only make conclusion by data, while the first laws change a period as well as people's mind about universe; and the third one many years (for ten years) later than the first two, because it's a rule for not only planets but also other celestial bodies.

The discussion part is informatival, o the regard contains most nother matrices is correct, except for tixt.

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2. How Kepler found these laws



The astronomer Tycho Brahe, the last of the great naked-eye astronomers and his protege Johannes Kepler, who together came up with the first simple and accurate description of how the planets

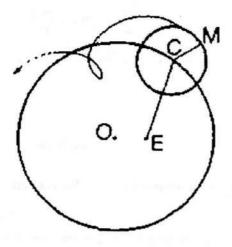
actually do move. The difficulty of

their task is suggested by the figure left, which shows how the relatively

simple orbital motions of the earth and Mars combine so that as seen from earth Mars appears to be staggering in loops like a drunken sailor.

However, After Brahe's death, it fell to his former assistant Kepler to try to make some sense out of the volumes of data. Kepler, in contradiction to his late boss, had formed a prejudice, a correct one as it turned out, in favor of the theory that the earth and planets revolved around the sun, rather than the earth staying fixed and everything rotating about it. Although such direct experiments were not carried out until the 19th century, what convinced everyone of the sun-centered system in the 17th century was that Kepler was able to come up with a surprisingly simple set of mathematical and geometrical rules for describing the planets' motion using the sun-centered assumption. After 900 pages of calculations and many false starts and dead-end ideas, Kepler finally synthesized the data into his three laws. [1]

Before Kepler's declaring his ellipse orbits of planets, everyone believed the wonderful theory by Ptolemy, which orbit just like below.

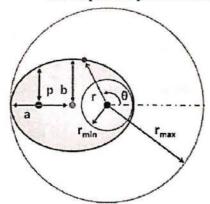


An epicycle combines with a bisector

Conjecture

3. Deriving Kepler's first law

An ellipse is a particular class of mathematical shapes that is most similar to a circle.



The Sun is not at the center of the ellipse but is at one of the focal points. The other focal point is marked with a lighter dot but is a point that has no physical significance for the orbit. Ellipses have two focal points neither of which is in the center of the ellipse. Circles are a special case of an ellipse that are not stretched out and in which both focal points coincide at the center.

By Newton's second laws

$$\overline{F_{net}} = m \frac{d^2 \vec{r}}{dt^2}$$

While by the law of universal gravitation

$$\overrightarrow{F_{net}} = \overrightarrow{F_g} = -\frac{GMm}{|\overrightarrow{r}|^2} \hat{r} = -\frac{GMm}{|\overrightarrow{r}|^2} \frac{\overrightarrow{r}}{|\overrightarrow{r}|} = -\frac{GMm\overrightarrow{r}}{|\overrightarrow{r}|^3}$$

Thus

$$m\frac{d^2\vec{r}}{dt^2} = -\frac{GMm\vec{r}}{|\vec{r}|^3}$$
Take the cross product of \vec{r} in both sides of (1)

We can get

$$\vec{r} \times m \frac{d^2 \vec{r}}{dt^2} = \vec{r} \times -\frac{GMm\vec{r}}{|\vec{r}|^3} = \vec{0} \Rightarrow \vec{r} \times \frac{d\vec{v}}{dt} = \vec{0}$$

$$\vec{r} \times \frac{d\vec{v}}{dt} + \frac{d\vec{r}}{dt} \times \vec{v} = \frac{d(\vec{r} \times \vec{v})}{dt} = \vec{0} + \frac{d\vec{r}}{dt} \times \vec{v} = \vec{v} \times \vec{v} = \vec{0}$$

Hence

$$\vec{r} \times \vec{v} = constant\vec{c}$$
 (2)

For constant \vec{c} has a fixed direction vertical to \vec{r} and \vec{v} , that's to say \vec{r} and \vec{v} are in the same plane.

So the sun and the plane stay on the same plane.

In addition,

$$\frac{d}{dt} \left(\frac{\vec{r}}{|\vec{r}|} \right) = \frac{\frac{d\vec{r}}{dt} |\vec{r}| - \frac{d|\vec{r}|}{dt} \vec{r}}{|\vec{r}|^2} = \frac{d\vec{r}}{dt} \frac{1}{|\vec{r}|} - \frac{d|\vec{r}|}{dt} |\vec{r}| = \frac{\vec{v}}{|\vec{r}|^2} = \frac{\vec{v}}{|\vec{r}|} - \frac{\vec{v}_r |\vec{r}|}{|\vec{r}|^2} = \frac{\vec{v}_r - \vec{v}_r}{|\vec{r}|} = \frac{\vec{v}_r - \vec{v}_r}{|\vec{r}|} = \frac{\vec{r}|}{|\vec{r}|} = \frac{\vec{r}|}{|\vec{r}|} = \frac{d\theta}{dt} =$$

 $\overrightarrow{v_r}$ is radial velocity, $\overrightarrow{v_\theta}$ is the velocity which is perpendicular to $\overrightarrow{v_r}$, and $\overrightarrow{v_r} + \overrightarrow{v_\theta} = \overrightarrow{v} \cdot \overrightarrow{e_r}$, $\overrightarrow{e_\theta}$ are the unit vector of $\overrightarrow{v_r}$, $\overrightarrow{v_\theta}$.

Or

$$\frac{d}{dt} \left(\frac{\overrightarrow{r}}{|\overrightarrow{r}|} \right) = \frac{\frac{d\overrightarrow{r}}{dt} |\overrightarrow{r}| - \frac{d}{dt} |\overrightarrow{r}|}{|\overrightarrow{r}|^2} = \frac{d\overrightarrow{r}}{dt} \frac{1}{|\overrightarrow{r}|} - \frac{d}{dt} |\overrightarrow{r}| \frac{d}{|\overrightarrow{r}|} |\overrightarrow{r}| e_r}{|\overrightarrow{r}|^2} = \frac{\overrightarrow{v}}{|\overrightarrow{r}|} - \frac{\overrightarrow{v_r} |\overrightarrow{r}|}{|\overrightarrow{r}|^2} = \frac{\overrightarrow{v}}{|\overrightarrow{r}|} - \frac{(\overrightarrow{r} \cdot \overrightarrow{v}) e_r}{|\overrightarrow{r}|^2}$$

From (1),
$$m\frac{d^2\vec{r}}{dt^2} = -\frac{GMm\vec{r}}{|\vec{r}|^3}, \frac{d\vec{v}}{dt} = -\frac{GM\vec{r}}{|\vec{r}|^3}$$

Take the cross product of \bar{c} in both sides of it.

$$\frac{d\vec{v}}{dt} \times \vec{c} = -\frac{GM\vec{r}}{|\vec{r}|^3} \times \vec{c} = -\frac{GM\vec{r}}{|\vec{r}|^3} \times (\vec{r} \times \vec{v}) = -\frac{GM}{|\vec{r}|^3} \left[(\vec{r} \cdot \vec{v})\vec{r} - (\vec{r} \cdot \vec{r})\vec{v} \right]$$

On the other hand, by question 3,

$$\frac{d}{dt} \left(\frac{GM\vec{r}}{|\vec{r}|} \right) = GM \left(\frac{\vec{v}}{|\vec{r}|} - \frac{(\vec{r} \cdot \vec{v})\vec{e_r}}{|\vec{r}|^3} \right) = \frac{GM}{|\vec{r}|^3} \left(|\vec{r}|^2 \vec{v} - (\vec{r} \cdot \vec{v})\vec{e_r}|\vec{r}| \right) = \frac{GM}{|\vec{r}|^3} \left((\vec{r} \cdot \vec{r})\vec{v} - (\vec{r} \cdot \vec{v})\vec{r} \right) = -\frac{GM}{|\vec{r}|^3} \left((\vec{r} \cdot \vec{v})\vec{r} - (\vec{r} \cdot \vec{r})\vec{v} \right)$$

$$\frac{d\vec{v}}{dt} \times \vec{c} = \frac{d}{dt} \left(\frac{GM\vec{r}}{|\vec{r}|} \right)$$

Thus

And we can compute: $\frac{d(\vec{v} \times \vec{c})}{dt} = \frac{d\vec{v}}{dt} \times \vec{c} + \vec{v} \times \frac{d\vec{c}}{dt}$

For \vec{c} is a constant vector, $\frac{d\vec{c}}{dt} = \vec{0}$, so $\frac{d(\vec{v} \times \vec{c})}{dt} = \frac{d\vec{v}}{dt} \times \vec{c}$

Hence,

$$\int \left(\frac{d\vec{v}}{dt} \times \vec{c}\right) dt = \int \left(\frac{d(\vec{v} \times \vec{c})}{dt}\right) dt = \int d(\vec{v} \times \vec{c}) = \vec{v} \times \vec{c} + constant\vec{A}$$

$$\int \frac{d}{dt} \left(\frac{GM\vec{r}}{|\vec{r}|}\right) dt = \frac{GM\vec{r}}{|\vec{r}|} + constant\vec{B}$$

$$thus_{-} \int \left(\frac{d\vec{v}}{dt} \times \vec{c}\right) dt = \int \frac{d}{dt} \left(\frac{GM\vec{r}}{|\vec{r}|}\right) dt + constant\vec{I}$$

$$\vec{v} \times \vec{c} = \frac{GM\vec{r}}{|\vec{r}|} + constant\vec{I} \quad (3)$$

Take the dot product of \vec{c} in the both sides of equality (3)

$$(\vec{v} \times \vec{c}) \cdot \vec{c} = \frac{GM\vec{r}}{|\vec{r}|} \cdot \vec{c} + \vec{l} \cdot \vec{c}$$

$$0 = \frac{GM\vec{r}}{|\vec{r}|} \cdot (\vec{r} \times \vec{v}) + \vec{l} \cdot \vec{c} = 0 + \vec{l} \cdot \vec{c}$$
thus $\vec{l} \cdot \vec{c} = 0, \vec{l}$ is perpendicular to \vec{c}

Take the dot product of \vec{r} in the both sides of equality (3)

$$(\vec{v} \times \vec{c}) \cdot \vec{r} = \frac{GM\vec{r}}{|\vec{r}|} \cdot \vec{r} + \vec{l} \cdot \vec{r} = GM |\vec{r}| + l(\vec{i} \cdot \vec{r}) = GM |\vec{r}| \left(1 + \frac{l}{GM} \cos \theta\right) = GM |\vec{r}| \left(1 + e \cos \theta\right)$$

$$thus |\vec{r}| = \frac{\vec{r} \cdot (\vec{v} \times \vec{c})}{GM \left(1 + e \cos \theta\right)}$$

$$\vec{r} \cdot (\vec{v} \times \vec{c}) = \vec{r} \cdot (\vec{v} \times (\vec{r} \times \vec{v})) = \vec{r} \cdot \left((\vec{v} \cdot \vec{v})\vec{r} - (\vec{v} \cdot \vec{r})\vec{v}\right) = |\vec{v}|^2 |\vec{r}|^2 - (\vec{v} \cdot \vec{r})^2$$

$$= |\vec{v}|^2 |\vec{r}|^2 \left(1 - \cos^2 \theta\right) = |\vec{v}|^2 |\vec{r}|^2 \sin^2 \theta = (\vec{v} \times \vec{r})^2 = (-\vec{c})^2 = c^2$$

$$thus |\vec{r}| = \frac{\vec{r} \cdot (\vec{v} \times \vec{c})}{GM \left(1 + e \cos \theta\right)} = \frac{c^2}{GM \left(1 + e \cos \theta\right)}$$

Assume
$$m = \frac{c^2}{GM}$$
, $|\vec{r}| = \frac{m}{1 + e \cos \theta}$

$$m = |\vec{r}| (1 + e \cos \theta) = \sqrt{x^2 + y^2} + ex$$

$$\frac{\left(x + \frac{em}{1 - e^2}\right)^2}{\left(\frac{m}{1 - e^2}\right)^2} + \frac{y^2}{\frac{m^2}{1 - e^2}} = 1$$

$$thus_the_eccentricity_is_\frac{\sqrt{\left(\frac{m}{1 - e^2}\right)^2 - \frac{m^2}{1 - e^2}}}{\frac{m}{1 - e^2}} = e$$

From above, we have already proof that the orbits of planets was an ellipse, and our sun at one focus.

4. Deriving Kepler's second law

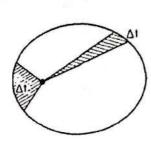


Fig. 7–2. Kepler's law of areas.

Kepler's second law is equivalent to the fact that the force perpendicular to the radius vector is zero. The "areal velocity" is proportional to angular momentum, and so for the same reasons, Kepler's second law is also in effect a statement of the conservation of angular momentum.

It would be nice if you could elaborate this is your prost on the next page

For a small time, dA could be thought was a triangle.

By the formula of triangle, we can get $dA = \frac{1}{2}r(rd\theta) = \frac{1}{2}r^2d\theta$ $\frac{dA}{dt} = \frac{1}{2}r^2\frac{d\theta}{dt}$ $\frac{c}{2} = \frac{dA}{dt} = \frac{1}{2}r^2 \frac{d\theta}{dt} \Rightarrow c = r^2 \frac{d\theta}{dt} \Rightarrow \vec{c} = r^2 \frac{d\theta}{dt} \vec{k}$ (for z-axis_to_be_in_the_direction_c) = Ir r do sing o k

 $A = \int \frac{dA}{dt} dt = \int \frac{c}{2} dt = \frac{c}{2} t + C,$

Because $for_{when_{t}=0,A=0}$, $for_{when_{t}=0,A=0}$

= 12 do 1/2

So when $t = t_0$, A=constant= $\frac{c}{2}t_0$ Hence, we have proof the second laws.

5. Deriving Kepler's third law

The third law used to be known as the harmonic law, because Kepler enunciated it in a laborious attempt to determine what he viewed as the "music of the spheres" according to precise laws, and express it in terms of musical notation. [2] This third law currently receives additional attention as it can be used to estimate the distance from a planet to its central star, and help to decide if this distance is inside the habitable zone of that star.

Compute the whole area of the ellipse,

$$S_{ellipse} = \frac{dA}{dt} \cdot T = \frac{c}{2}T = \pi ab$$

$$thus_{T} = \frac{2\pi ab}{c}$$

By problem 8, $\frac{c^2}{GM} = r(1 + e\cos\theta) = const$

Let $\theta = \frac{\pi}{2}$, then we get the semi-latus rectum

$$\frac{c^2}{GM} = r(1 + e\cos\theta) = r = semi - latus_rectum = \sqrt{b^2 \left(1 - \frac{c^2}{a^2}\right)} = b\sqrt{1 - e^2}$$

$$\frac{ac^2}{GM} = ab\sqrt{1 - e^2} = b\sqrt{a^2 - a^2 e^2} = b^2$$

By question 13,
$$b = \frac{cT}{2\pi a}$$

Then

$$b^2 = \left(\frac{cT}{2\pi a}\right)^2 = \frac{ac^2}{GM}$$
$$T^2 = \frac{ac^2}{GM} \frac{4\pi^2 a^2}{c^2} = \frac{4\pi^2}{GM} a^3$$

QED

6. General situation

These laws approximately describe the motion of any two bodies in orbit around each other. The masses of the two bodies can be nearly equal, e.g. Charon—Pluto, in a small proportion, or in a great proportion, e.g. Mercury—Sun. [3] In all cases of two-body motion, rotation is about the barycenter of the two bodies, with neither one having its center of mass exactly at one focus of an ellipse. However, both orbits are ellipses with one focus at the barycenter. When the ratio of masses is large, the barycenter may be deep within the larger object, close to its center of mass.

7. An application using Kepler's three laws

Compute the altitude of a geosynchronous satellite.

some_needed_constants $G = 6.67 \times 10^{-11} kg^{-1} \cdot m^3 \cdot s^{-2}$ $M = 5.976 \times 10^{24} kg$ T = 86164s R = 6378137m $\pi = 3.14159265358$

A geosynchronous satellite stays at a circle which e=0 instead of a ellipse

so _ a = b,
$$\overline{r}$$
 _ is _ perpendicular _ to _ \overline{v}
 $T = \frac{2\pi a^2}{c}$, $a = \frac{c^2}{GM} \Rightarrow Ta^2 = \frac{2\pi a^2}{c} \frac{c^4}{G^2M^2} \Rightarrow T = \frac{2\pi c^3}{G^2M^2}$
 $\Rightarrow \frac{TG^2M^2}{2\pi} = (R+h)^3 v^3 = (R+h)^3 (R+h)^3 \omega^3 = (R+h)^6 \frac{(2\pi)}{T^3}$
 $\Rightarrow h = \sqrt[3]{\frac{T^2GM}{4\pi^2}} - R$
 $take_those_constants_into_this_equality$
 $we_can_get_the_answer$
 $h = 3.5786 \times 10^7 m$

Actually, a geosynchronous satellite's orbit is incompletely a circle, but with a little like ellipse. In one day, the radius is bigger in a time and smaller in another time. When the radius comes to be bigger, the speed of the satellite will decrease, what's more, its orbit would drift to west a little, otherwise, it would drift to east.

In addition, the obliquity of a geosynchronous satellite is actually not zero. At this time, the satellite may drift in north-south way. If the orbit is just like a ellipse as well as with some obliquity, the track of it may be the combination of them, which makes the stars move like a '8'.

8. Sense in doing this project

-----No pain, no gain

The members of our group completed the project task with cooperation. In the project research, we got some viewpoints and feelings:

1) Concerning the physics:

The reason why Newton is greater than Kepler is that he does something more normal, which not only in cosmetology but also the whole motion system. However, Kepler is also greater than Tycho, for he can find mathematical rules by the data offered by Tycho. In the modern age, Einstein's theory extended the scope of the object motion which can describe the object with a high-speed. So the trend of physics' development is to develop a theory to adapt To make true breakthroughs, concrete cases to more general circumstance. still need to be studied; general theories without special & important examples are empty non-sense.

2) Concerning calculus

We know that the calculus is an important subject for our study, but we knew less than we should know-how Newton can found Calculus such a great building etc. Through this project study, we realize calculus can be used in such broad scope, and we are also amazed at the new form of calculus-vector calculus. Therefore, we may learn more things than we used to do, and we are looking forward to other beautiful applications of Calculus.

3) Concerning group

The project is not only homework, which like of some of students think, however good opportunity for us to exercise ourselves, how to write an article, how to solve a big and real problem, more importantly, how to work with our partners.

Actually, my team was founded for a long time, when we didn't know each other well, so we faced something which had broken our process. However, this headachy problem was solved by us teammates and we felt that a problem from person themselves is more complex than a project problem. Therefore, a good project report comes from every teammate efforts, no pain, no gain; working is more interesting than reaping where one has not sown.

4) Concerning doing problem

We have given all the proofs by ourselves; however, we only solve these questions step by step according to the materials. Although it's our first time to do such a project, next time, we think that we can try our best to solve a rather big problem without a lot of materials. This may let us think like a science, for example, how Kepler found these rules and proved them, and we may get more experience from that.

In a word, we harvested a lot during this wonderful project.

Reference:

2006

[1][Pictures]Benjamin Crowell, Newtonian Physics, www.lightandmatter.com, edition 2.3,

[2] Murray and Dermott, Solar System Dynamics, Cambridge University Press 1999,

[3] V.I. Arnold, Mathematical Methods of Classical Mechanics, Chapter 2. Springer 1989