First principle defect calculation

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1 Defect formation energy

Defect formation energy of an isolated defect ${\cal H}_q^f$ is expressed by the reaction equation:

$$\sum_{i} N_{i}\mu_{i} + E_{f}dn_{e} + H_{q}^{f} \to E_{iso,q}^{def} \tag{1}$$

where the summation is over all the components i. N and μ is the counts and chemical potential of each components. dn_e is the number of electrons that we additionally added or removed to the system, E_f is their energy. The chemical potential is calculated by considering the equilibrium between coexisting phase, which is obtained from the phase diagram[1]. $E_{iso,q}^{def}$ is the energy of a supercell that contain the isolated defect with total charge charge q, q > 0 for removing electrons and q < 0 for adding electrons ($dn_e = -q$ from definition).

In actual calculation, we have:

$$E_{iso,q}^{def} = E_{periodic,q}^{def} - E_{q}^{corr} \tag{2}$$

 E_q^{corr} is the correction necessary to remove the contribution of calculated total energy in the defected supercell. Rearranging, we have the expression for the defect formation energy:

$$H_q^f = E_{periodic,q}^{def} - \sum_i N_i \mu_i + q E_f - E_q^{corr}$$
(3)

2 Correction of the image charge

In this section, we provide a summary of the popular methods for the finite size correction in charged defect calculation. We review the method by Freysoldt[2] and Kumagai[3]

Freysoldt's method

In Freysoldt's article, the creation of a charged point defect is split into three steps, In the first step, we consider a charge neutral cell which contain the defect. Take the As defect in Silicon as an example, compare to the perfect crystal, the additional ionic charge on the As atoms and an additional electron are introduced. This extra electron will form an orbital $\phi_d(r)$. Then we attend to remove this additional electron from this orbital (to the conduction band bottom in some far away unitcell), leading to ionization. In the defect supercell, we have a charge:

$$q_d(r) = q|\phi_d(r)|2\tag{4}$$

Next, we relax the system, allowing the electrons (and ions) to screen the introduced charge. The overall resulting electron distribution give raise to a change in electrostatic potential, compared to the case of neutral defect:

$$V_{a/0}(r) = V_{defect,q}(r) - V_{defect,0}(r)$$

$$\tag{5}$$

 $V_{q/0}(r)$ is the potential introduced from ionization. Finally, we introduce the artificial periodicity and add a compensating homogeneous background charge density $n = -q/\Omega$, where Ω is the volume of the supercell. The compensating background removes the divergence of the electrostatic potential:

$$V_{q/0}(G) = \int d^3r V_{q/0}(r) e^{-iGr}$$
 (6)

$$\tilde{V}_{q/0}(r) = \frac{1}{\Omega} \sum_{G \neq 0} V_{q/0}(G) e^{iGr}$$
 (7)

where $q \neq 0$ removes the divergence with compensating background charge and $\tilde{V}_{q/0}(r)$ is now the additional periodic potential (compared to charge neutral case) from charged periodic defects.

We keep track of the artificial elements we introduced in our calculation: the periodic image charge and their potentials as well as the neutral compensating background. Thus the energy correction consist of two parts:

$$E^{inter} = \frac{1}{2} \int_{\Omega} [q_d(r) + n] [\tilde{V}_{q/0}(r) - V_{q/0}(r)] d^3r$$
 (8)

$$E^{intra} = \int_{\Omega} nV_{q/0}(r)d^3r = -\frac{q}{\Omega} \left(\int_{\Omega} V_{q/0}(r)d^3r \right)$$
(9)

The first part is the interaction between the defect charge and compensating background in the home supercell with the potential from imaginary charge and background from image cells. The second part is the interaction of the defect potential in the supercell introduced by the actual defect and the compensating background. The integral is confined to the supercell, whose total energy is used.

We now separate the potential $V_{q/0}(r)$ into two parts: 1) a long ranged part $V_q^{lr}(r)$, that we can approximate using simple form. 2) a short ranged part $V_q^{sr}(r)$ that is zero outside of the supercell:

$$V_{q/0}(r) = V_q^{lr}(r) + V_{q/0}^{sr}(r)$$
(10)

$$V_q^{lr}(r) = \frac{1}{\epsilon} \int \frac{q_d(r')}{|r - r'|} d^3r' \tag{11}$$

For their periodic form, we have:

$$\tilde{V}_{q/0}^{sr}(r) = \sum_{R} V_{q/0}^{sr}(r+R) + C \tag{12}$$

$$\approx V_a^{sr}(r) + C$$

$$\tilde{V}_{q/0}(r) = \tilde{V}_q^{lr}(r) + \tilde{V}_{q/0}^{sr}(r) = \tilde{V}_q^{lr}(r) + V_q^{sr}(r) + C$$
(13)

the periodic part, $\tilde{V}_q^{lr}(r)$ is obtained from Eq.11 and Eq.7. The constant C to reproduce the absolute value of $\tilde{V}_{q/0}(r)$ (the value of $\tilde{V}_{q/0}^{lr}(r)$ is fixed once we choose the form the $q_d(r)$).

The correction energy can be simplified to be:

$$E^{corr} = E^{inter} + E^{intra} = E_q^{lat} + q\Delta_{q/0}$$
(14)

with each part

$$E_q^{lat} = \int_{\Omega} \left(\frac{1}{2} [q_d(r) + n] [\tilde{V}_q^{lr}(r) - V_q^{lr}(r)] + n V_q^{lr}(r) \right) d^3r$$
 (15)

$$\Delta_{q/0} = \frac{1}{\Omega} \int_{\Omega} V_{q/0}^{sr}(r) d^3r \tag{16}$$

$$V_{q/0}^{sr}(r) = \tilde{V}_{q/0}(r) - \tilde{V}_{q}^{lr}(r) - C \tag{17}$$

The correction energy thus depends only on values of ϵ , $q_d(r)$, $\tilde{V}_{q/0}$ and C. the dielectric constant ϵ can be calculated by first principle methods such as DFPT, $\tilde{V}_{q/0}$ are calculated by Eq.5 using the potential calculated for the charged supercell and neutral defect (or perfect bulk potential). The constant C is calculated by Eq.13. For the charge density $q_d(r)$, reasonable approximation suffices (point charge $q_d(r) = \delta(r)q$ or a Gaussian charge).

Finally, the form of defect formation energy is:

$$H_q^f = E^{def} - \sum_{i} N_i \mu_i - E_q^{lat} + q(E_f + \Delta_{q/0})$$
(18)

q is positive if we remove the electron from the supercell (The charge of the whole supercell is positive).

Modified Freysoldt's method (Kumagai)

In Kumagai's paper, some method as Freysoldt is used. However, some modification is introduced to consider the anisotropy and the difficult in potential alignment when the ions are polarizable.

References

- [1] First Principles Calculation of Defect Formation Energies in Sr and Mg-Doped LaGaO₃, Akihide Kuwabara and Isao Tanaka. J. Phys. Chem. B 108, 9168-9172 (2004)
- [2] Fully Ab-initio Finite-size Corrections for Charged-Defect Supercell Calculations, Christoph Freysoldt, et. al. PRL 102, 016402 (2019)
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