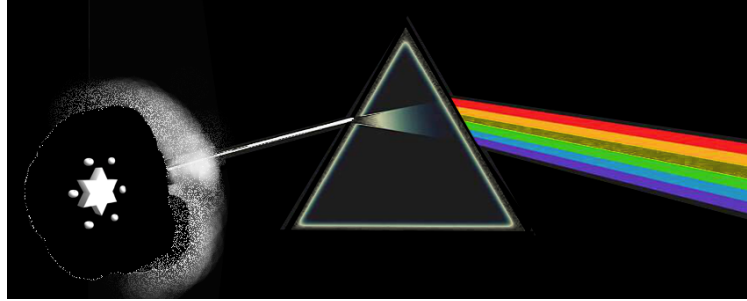


The Magic Star of Exceptional Periodicity



Piero Truini¹, Michael Rios², Alessio Marrani^{3,4}

¹ *Dipartimento di Fisica, Università di Genova,
and INFN, sez. di Genova,
via Dodecaneso 33, I-16146 Genova, Italy*
truini@ge.infn.it

² *Dyonica, ICMQG, USA*
mrios@dyonicatech.com

³ *Museo Storico della Fisica e Centro Studi e Ricerche “Enrico Fermi”,
Via Panisperna 89A, I-00184, Roma, Italy*

⁴ *Dipartimento di Fisica e Astronomia “Galileo Galilei”, Università di Padova,
and INFN, Sez. di Padova,
Via Marzolo 8, I-35131 Padova, Italy*
jazzphyzz@gmail.com

ABSTRACT

We present a periodic infinite chain of finite generalisations of the exceptional structures, including \mathbf{e}_8 , the exceptional Jordan algebra (and pair), and the octonions. We demonstrate that the exceptional Jordan algebra is part of an infinite family of finite-dimensional matrix algebras (corresponding to a particular class of cubic Vinberg’s T-algebras). Correspondingly, we prove that \mathbf{e}_8 is part of an infinite family of algebras (dubbed “Magic Star” algebras) that resemble lattice vertex algebras.

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1 Introduction

In the mind of every mathematician, there is a tension between the general rule and exceptional cases. Our conscience tells us we should strive for general theorems, yet we are fascinated and seduced by beautiful exceptions.

J. Stillwell [1]

Since the inception of quantum mechanics operator algebras and their symmetries have proved to be of central importance. Jordan, Wigner and von Neumann [2], in a rigorous exploration of extended forms of quantum mechanics, classified finite dimensional self-adjoint operator algebras, now known as formally real Jordan algebras. The exceptional case they found contained the largest division algebra, the octonions \mathbb{O} [3], first discovered by J. T. Graves in 1843.

With Gell-Mann's Eightfold way [4], and subsequent study of the su_3 quark and gluon structures [5, 6], Lie algebras solidified a firm role in the study of elementary particles. After the classifications by Killing and Cartan, all finite-dimensional complex simple Lie algebras, along with all their non-compact real forms, are known. Besides the infinite classical series \mathfrak{a}_n , \mathfrak{b}_n , \mathfrak{c}_n , \mathfrak{d}_n , five *exceptional* Lie algebras exist : \mathfrak{g}_2 , \mathfrak{f}_4 , \mathfrak{e}_6 , \mathfrak{e}_7 and \mathfrak{e}_8 . Gursay, Ramond and others argued exceptional Lie algebras played a role in unification of fundamental particles [7]. With such studies the octonions made a return, as the algebra of their derivations is the smallest exceptional Lie algebra \mathfrak{g}_2 . Moreover the maximal self-adjoint operator algebra over the octonions, the exceptional Jordan algebra $\mathbf{J}_3^{\mathbb{O}}$, has the next largest exceptional Lie algebra \mathfrak{f}_4 describing its derivations [8].

Various non-compact real forms of exceptional Lie algebras play the role of electric-magnetic duality (U -duality¹) algebras in locally supersymmetric theories of gravity² (see *e.g.* [12] for theories with 8 supersymmetries).

The largest exceptional Lie algebra \mathfrak{e}_8 became central to the heterotic string [13] construction, that assigned 16 of the 26 dimensions of the bosonic string to the $\mathfrak{e}_8 \oplus \mathfrak{e}_8$ even self-dual lattice. Later, as $D = 11$ supergravity merged with $D = 10$ string theory in Witten's [14] mysterious M -theory, Ramond *et al.* noticed the $D = 11$ supermultiplet had a hidden \mathfrak{f}_4 symmetry [15], an observation which was further investigated by Sati [16, 17]. Ramond even

¹Here U -duality is referred to as the “continuous” symmetries of [9]. Their discrete versions are the U -duality non-perturbative string theory symmetries introduced in [10].

²Some non-compact real forms of exceptional algebras also occur in absence of local supersymmetry (*cfr.* [11], and Refs. therein).

speculated whether the exceptional Jordan algebra $\mathbf{J}_3^\mathbb{O}$ might be a special charge space related to the $D = 11$ lightcone [18], as it has natural $\mathfrak{so}(9)$ and \mathfrak{f}_4 symmetry³. Schwarz and Kim then recast the BFSS matrix model [20] for M -theory in terms of octonion variables [21], and Smolin invoked the full exceptional Jordan algebra in a Chern-Simons string matrix model [22] for Horowitz and Susskind’s conjectured “bosonic M -theory” in $D = 27$ [23].

With breakthroughs in algebraic geometry, coming from Connes and others in the investigation of noncommutative geometry [24, 25, 26], a program for emergent spacetime is underway and goes beyond Riemannian geometry in favor of recovering manifolds from operator algebras. This move to an algebraic derived spacetime proved to be novel, as it circumvented the usual problems with Lorentz symmetry via discretization by making geometry inherently fuzzy [27]. It was even found that extended objects in string theory, D -branes, had natural coordinates that were noncommutative [20, 28]. This made the role of C^* -algebras and K -homology of central importance to the study of D -branes [29, 30], while also providing the spectral triples for traditional noncommutative geometry and its applications in deriving the Standard Model of particle physics [31, 32].

In 2007, Lisi proposed a unified model [33] using \mathfrak{e}_8 , which was later discovered to be troubled [34]. However, Lisi’s theory inspired Truini to study a special star-like projection - named “*Magic Star*” (MS) - of \mathfrak{e}_8 under \mathfrak{a}_2 [35]; also motivated by the search of a unified way to characterize the fourth row of the Freudenthal-Rozen-Tits Magic Square [36], the Magic Star made the structure of Jordan algebras of degree three manifest in each exceptional Lie algebra [35] (see also [37, 38]). An interaction based model using this \mathfrak{e}_8 projection has been speculated in [35], and proposed in [39].

Our aim, in this contribution to the Proceedings of 4th Mile High Conference on Nonassociative Mathematics, as well as in subsequent forthcoming works, is to define and study a consistent *generalization* of exceptional Lie algebras, by crucially exploiting some remarkable properties of the MS.

We will show how it is possible to go beyond the largest finite-dimensional exceptional Lie algebra \mathfrak{e}_8 by relying on the basic structures underlying the MS. This will result in the formulation of the so-called “*Exceptional Periodicity*” (EP), namely a generalization of exceptional Lie algebras which is parametrized by a natural number $n \in \mathbb{N}$, and which also enjoys a periodicity (ultimately traced back to the well known Bott periodicity). At each “level” of EP, namely for each fixed value of n , the dimension of the resulting algebra will be *finite*, raising however the intriguing question (left for future studies) of investigating its $n \rightarrow \infty$ limit. The impetus for the EP generalization came from certain 3- and 5- gradings of the exceptional Lie algebras, along with spinor structures⁴ beyond \mathfrak{e}_8 .

We will prove the existence of a *periodic infinite chain of finite generalisations of the exceptional structures*, including \mathfrak{e}_8 , the Exceptional Jordan (or Albert) Algebra $\mathbf{J}_3^\mathbb{O}$, and the octonions \mathbb{O} themselves. Remarkably, for $n = 1$, the MS-shaped structure of known finite-dimensional exceptional Lie algebras is recovered.

As it will be evident from the subsequent treatment, the price to be paid for such an elegant and periodic, n -parametrized and MS-shaped, generalization of exceptional algebras is that *the resulting algebras will no longer be of Lie type*, namely they will not satisfy Jacobi identities anymore. Indeed, within EP we will not be dealing with root lattices, but rather with “*extended*” root lattices, which will be thoroughly defined further below.

It is here worth remarking that EP provides a way to go beyond \mathfrak{e}_8 which is radically

³Recently, the maximally non-compact (*i.e.*, split) real form $\mathfrak{f}_{4(4)}$ has been conjectured as the global symmetry of an exotic ten-dimensional theory in the context of the study of “Magic Pyramids” [19]. See Mike Duff’s talk at this Conference.

⁴Discussion with Eric Weinstein, during the “Advances in Quantum Gravity” symposium, San Francisco, July 2016.

different from the way provided by affine and (extended) Kac-Moody generalizations, such as⁵ $\mathbf{e}_8^+ =: \mathbf{e}_9$, $\mathbf{e}_8^{++} =: \mathbf{e}_{10}$, $\mathbf{e}_8^{+++} =: \mathbf{e}_{11}$, which also appeared as symmetries for (super)gravity models reduced to $D = 2, 1, 0$ dimensions [41, 42], respectively, as well as near spacelike singularities in supergravity [43]. In fact, while such extensions of \mathbf{e}_8 are still of Lie type but *infinite-dimensional*, the generalization of exceptional Lie algebras provided by EP is *not of Lie type*, but nevertheless is *finite-dimensional*, for each level of the EP itself.

The paper is organized as follows.

In section 2 we recall the notions of Jordan Algebras and Pairs with particular emphasis on their relationship with all the exceptional Lie algebras, except \mathbf{e}_8 . In section 3 we introduce the Magic Star and show its main features, among which the appearance of a triple of Jordan Pairs at the core of \mathbf{e}_8 . This is, in our opinion, the best way of expressing the link between \mathbf{e}_8 and Jordan algebras. In sections 4, 5 and 6 we extend the features of the Magic Star, in particular its relationship with the exceptional Lie algebras, to an infinite chain of finite dimensional algebras called *Exceptional Periodicity* (EP). In the case of the extension $\mathbf{e}_8^{(n)}$ of \mathbf{e}_8 , the algebra has rank $N = 4(n + 1)$, $n = 1, 2, \dots$. In section 4 we introduce a system of *extended roots* for EP, depending on n , which generalizes the root system of the exceptional Lie algebras. In section 5 we associate an algebra, for each n , to the extended root system. It appears as a finite dimensional generalization of the exceptional Lie algebras, in which the Jacobi identity is only partially true. In section 6 we show that in the Magic Star for EP the role of the Jordan algebras is played by matrix algebras, first introduced by Vinberg [44], called T-algebras. All in all, EP not only generalizes exceptional Lie algebras, but also the exceptional Jordan algebra via its T-algebra [44] structure. And, moreover, this structure is made manifest in the MS projections of the extended root systems. In section 7 we list some future developments on the study of EP and its applications to a model for Quantum Gravity. Most of these developments are at an advanced stage at the time we are completing the present paper.

2 Jordan Structures and related Lie algebras

There are no Jordan algebras, there are only Lie algebras.

I. L. Kantor

The modern formulation of Jordan algebras, [45], involves a quadratic map U_{xy} (like xyx for associative algebras) instead of the original symmetric product $x \cdot y = \frac{1}{2}(xy + yx)$, [2]. The quadratic map and its linearization $V_{x,y}z = (U_{x+z} - U_x - U_z)y$ (like $xyz + zyx$ in the associative case) reveal the mathematical structure of Jordan Algebras much more clearly, through the notion of inverse, inner ideal, generic norm, *etc.* The quadratic formulation is also particularly suited for the connection with Lie algebras, as we will explain in the next section. The axioms for (quadratic) Jordan algebras are:

$$U_1 = Id \quad , \quad U_x V_{y,x} = V_{x,y} U_x \quad , \quad U_{U_x y} = U_x U_y U_x \quad (2.1)$$

The quadratic formulation led to the concept of Jordan Triple systems [46], an example of which is a pair of modules represented by rectangular matrices. There is no way of multiplying two matrices x and y , say $n \times m$ and $m \times n$ respectively, by means of a bilinear product. But one can do it using a product like xyx , quadratic in x and linear in y . Notice that, like in the case of rectangular matrices, there needs not be a unity in these structures. The axioms are in this case:

$$U_x V_{y,x} = V_{x,y} U_x \quad , \quad V_{U_x y, y} = V_{x, U_y x} \quad , \quad U_{U_x y} = U_x U_y U_x \quad (2.2)$$

⁵For recent development on \mathbf{e}_{11} and beyond, see [40].

Finally, a Jordan Pair is defined just as a pair of modules (V^+, V^-) acting on each other (but not on themselves) like a Jordan Triple:

$$\begin{aligned} U_{x^\sigma} V_{y^{-\sigma}, x^\sigma} &= V_{x^\sigma, y^{-\sigma}} U_{x^\sigma} \\ V_{U_{x^\sigma} y^{-\sigma}, y^{-\sigma}} &= V_{x^\sigma, U_{y^{-\sigma}} x^\sigma} \\ U_{U_{x^\sigma} y^{-\sigma}} &= U_{x^\sigma} U_{y^{-\sigma}} U_{x^\sigma} \end{aligned} \quad (2.3)$$

where $\sigma = \pm$ and $x^\sigma \in V^{+\sigma}$, $y^{-\sigma} \in V^{-\sigma}$.

Jordan Pairs are strongly related to the Tits-Kantor-Koecher construction of Lie Algebras \mathfrak{L} [47]-[49] (see also the interesting relation to Hopf algebras, [50]):

$$\mathfrak{L} = J \oplus \text{str}(J) \oplus \bar{J} \quad (2.4)$$

where J is a Jordan algebra and $\text{str}(J) = L(J) \oplus \text{Der}(J)$ is the structure algebra of J [51]; $L(x)$ is the left multiplication in J : $L(x)y = x \cdot y$ and $\text{Der}(J) = [L(J), L(J)]$ is the algebra of derivations of J (the algebra of the automorphism group of J) [52][53].

In the case of complex exceptional Lie algebras, this construction applies to \mathfrak{e}_7 , with $J = \mathbf{J}_3^8 \equiv \mathbf{J}_3^\mathbb{O}$, the 27-dimensional exceptional Jordan algebra of 3×3 Hermitian matrices over the complex octonions, and $\text{str}(J) = \mathfrak{e}_6 \otimes \mathbb{C} - \mathbb{C}$ denoting the complex field. The algebra \mathfrak{e}_6 is called the *reduced structure algebra* of J , $\text{str}_0(J)$, namely the structure algebra with the generator corresponding to the multiplication by a complex number taken away: $\mathfrak{e}_6 = L(J_0) \oplus \text{Der}(J)$, with J_0 denoting the traceless elements of J .

Let us focus on exceptional Jordan and Lie algebras. We consider the Jordan algebras and Pairs based on the 3×3 matrices $\mathbf{J}_3^\nu, \nu = 1, 2, 4, 8$:

$$\mathbf{J}_3^\nu := \begin{pmatrix} \alpha & a & \bar{b} \\ \bar{a} & \beta & c \\ b & \bar{c} & \gamma \end{pmatrix} \quad \begin{array}{l} \alpha, \beta, \gamma \in \mathbb{C}; a, b, c \in \mathbb{A} \\ \nu = 1, 2, 4, 8 \text{ for } \mathbb{A} = \mathbb{R}, \mathbb{C}, \mathbb{H}, \mathbb{O} \end{array} \quad (2.5)$$

where $\mathbb{R}, \mathbb{C}, \mathbb{H}, \mathbb{O}$ denote the Hurwitz algebras of reals, complex, quaternions and octonions, respectively.

If $x, y \in \mathbf{J}_3^\nu$ and xy denotes their standard matrix product, we denote by $x \cdot y := \frac{1}{2}(xy + yx)$ the Jordan product of x and y . The Jordan identity is the power associativity with respect to this product:

$$x^2 \cdot (x \cdot z) - x \cdot (x^2 \cdot z) = 0, \quad (2.6)$$

Another fundamental product is the *sharp* product $\#$, [51]. It is the linearization of $x^\# := x^2 - t(x)x - \frac{1}{2}(t(x^2) - t(x)^2)I$, with $t(x)$ denoting the trace of $x \in \mathbf{J}_3^\nu$, in terms of which we may write the fundamental cubic identity for $\mathbf{J}_3^\nu, \nu = 1, 2, 4, 8$:

$$x^\# \cdot x = \frac{1}{3}t(x^\#, x)I \quad \text{or} \quad x^3 - t(x)x^2 + t(x^\#)x - \frac{1}{3}t(x^\#, x)I = 0 \quad (2.7)$$

where we use the notation $t(x, y) := t(x \cdot y)$ and $x^3 = x^2 \cdot x$ (notice that for \mathbf{J}_3^8 , because of non-associativity, $x^2 x \neq x x^2$ in general).

The triple product is defined as, [51]:

$$\begin{aligned} \{x, y, z\} := V_{x,y}z &= t(x, y)z + t(z, y)x - (x \# z) \# y \\ &= 2[(x \cdot y) \cdot z + (y \cdot z) \cdot x - (z \cdot x) \cdot y] \end{aligned} \quad (2.8)$$

Notice that the last equality of (2.8) is not trivial at all. $V_{x,y}z$ is the linearization of the quadratic map U_{xy} . The equation (2.3.15) at page 484 of [51] shows that:

$$U_{xy} = t(x, y)x - x^\# \# y = 2(x \cdot y) \cdot x - x^2 \cdot y \quad (2.9)$$

The following identities can be derived from the Jordan Pair axioms, [54]:

$$[V_{x,y}, V_{z,w}] = V_{V_{x,y}z, w} - V_{z, V_{x,y}w} \quad (2.10)$$

and, for $D = (D_+, D_-)$ a derivation of the Jordan Pair V and $\beta(x, y) = (V_{x,y}, -V_{y,x})$,

$$[D, \beta(x, y)] = \beta(D_+(x), y) + \beta(x, D_-(y)) \quad (2.11)$$

We have: For $J_P := (J, \bar{J})$: $\text{str}(J) = \text{Der}(J_P)$

$$U_{x^+}y^- = [[x^+, y^-], x^+] \quad ; \quad U_{x^-}y^+ = [[x^-, y^+], x^-]$$

JP axioms follow from the Jacobi identity of the Lie algebra $\text{str}(J)$.

To summarize the relationship between the exceptional Jordan algebra and the exceptional Lie algebras is:

$$\mathfrak{g}_2 = \text{Der}(\mathbb{O})$$

$$\mathfrak{f}_4 = \text{Der}(\mathbf{J}_3^8) ; \text{Der}(J) = [L(J), L(J)]$$

$$\mathfrak{e}_6 = \text{str}_o(\mathbf{J}_3^8) = L(J_0) \oplus \text{Der}(J) = \text{Der}_o(\mathbf{J}_3^8, \bar{\mathbf{J}}_3^8)$$

$$\mathfrak{e}_7 = \mathbf{J}_3^8 \oplus \text{Der}(\mathbf{J}_3^8, \bar{\mathbf{J}}_3^8) \oplus \bar{\mathbf{J}}_3^8 \quad \text{superstructure algebra}$$

$$\mathfrak{e}_8 = ?$$

The link between \mathfrak{e}_8 and Jordan Algebras and Pairs is expressed, in our opinion, by the Magic Star, [35].

3 The Magic Star: Exceptional Lie algebras and Jordan structures

The star that can bring peace and order back to the chaotic world.

citation from the “Magic Star” TV series, China, 2017.

The Magic Star, [35] [37], shown in Figure 1, is the projection of the roots of the exceptional Lie algebras on a complex $\mathfrak{su}(3) = \mathfrak{a}_2$ plane, recognizable by the dots forming the external hexagon, and it exhibits the *Jordan Pair* content of each exceptional Lie algebra. There are three Jordan Pairs $(\mathbf{J}_3^\nu, \bar{\mathbf{J}}_3^\nu)$, each of which lies on an axis symmetrically with respect to the center of the diagram. Each pair doubles a simple Jordan algebra of rank 3, \mathbf{J}_3^ν , with involution - the conjugate representation $\bar{\mathbf{J}}_3^\nu$, which is the algebra of 3×3 Hermitian matrices over \mathbb{A} , where $\mathbb{A} = \mathbb{R}, \mathbb{C}, \mathbb{H}, \mathbb{O}$ for $\nu = 1, 2, 4, 8$ respectively. The exceptional Lie algebras $\mathfrak{f}_4, \mathfrak{e}_6, \mathfrak{e}_7, \mathfrak{e}_8$ are obtained for $\nu = 1, 2, 4, 8$, respectively. \mathfrak{g}_2 can be also represented in the same way, with the Jordan algebra reduced to a single element. The Jordan algebras \mathbf{J}_3^ν (and their conjugate $\bar{\mathbf{J}}_3^\nu$) globally behave like a $\mathbf{3}$ (and a $\bar{\mathbf{3}}$) dimensional representation of the outer \mathfrak{a}_2 . The algebra denoted by \mathfrak{g}_0^ν in the center (plus the Cartan generator associated with the axis along which the pair lies) is the algebra of the automorphism group of the Jordan Pair; namely, \mathfrak{g}_0^ν is the *reduced* structure group of the corresponding Jordan algebra \mathbf{J}_3^ν : $\mathfrak{g}_0^\nu = \text{str}_0(\mathbf{J}_3^\nu)$.

The *base field* considered throughout the present paper is \mathbb{C} . The various real compact and non-compact forms of the exceptional Lie algebras follow as a consequence, using some more or less laborious tricks, whose treatment we leave to a future study, and they do not affect the

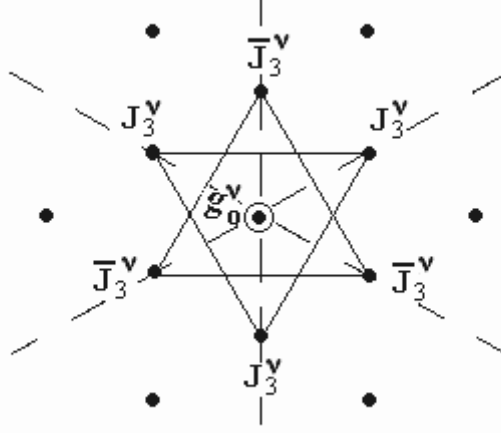


Figure 1: Root diagram of \mathbf{f}_4 , \mathbf{e}_6 , \mathbf{e}_7 , \mathbf{e}_8 ($\nu = 1, 2, 4, 8$) projected on the plane of one \mathbf{a}_2

essential structure.

What is magic in the Magic Star? Let us list a number of features which are clearly shown by the picture.

- 1) The Magic Star represents a unifying view of all the exceptional algebras.
- 2) Since $Der_o(\mathbf{J}_3^2, \bar{\mathbf{J}}_3^2) = \mathbf{a}_2 \oplus \mathbf{a}_2$ and $Der_o(\mathbf{J}_3^8, \bar{\mathbf{J}}_3^8) = \mathbf{e}_6$, Figure 1 depicts the following decomposition, [35] :

$$\begin{aligned}
 \mathbf{e}_8 &= \mathbf{a}_2^c + (3, \mathbf{J}_3^8) + (\bar{3}, \bar{\mathbf{J}}_3^8) + Der_o(\mathbf{J}_3^8, \bar{\mathbf{J}}_3^8) \\
 &= \mathbf{a}_2^c + (3, \mathbf{J}_3^8) + (\bar{3}, \bar{\mathbf{J}}_3^8) + \mathbf{a}_2^f + (3, \mathbf{J}_3^2) + (\bar{3}, \bar{\mathbf{J}}_3^2) + Der_o(\mathbf{J}_3^2, \bar{\mathbf{J}}_3^2) \\
 &= \mathbf{a}_2^c + (3, \mathbf{J}_3^8) + (\bar{3}, \bar{\mathbf{J}}_3^8) + \mathbf{a}_2^f + (3, \mathbf{J}_3^2) + (\bar{3}, \bar{\mathbf{J}}_3^2) + \mathbf{a}_2^{(1)} + \mathbf{a}_2^{(2)}
 \end{aligned} \tag{3.1}$$

which shows that \mathbf{e}_8 is made of 4 orthogonal \mathbf{a}_2 subalgebras, plus 3 Jordan Pairs $(\mathbf{J}_3^8, \bar{\mathbf{J}}_3^8)$ over the octonions, plus 3 Jordan Pairs $(\mathbf{J}_3^2, \bar{\mathbf{J}}_3^2)$ over the complex. The role of the \mathbf{a}_2 's is interchangeable: any \mathbf{a}_2 would produce a Magic Star. Select one and call it \mathbf{a}_2^c for *color* \mathbf{a}_2 and another one, \mathbf{a}_2^f , inside \mathbf{e}_6 as *flavor* \mathbf{a}_2 , then the Jordan Pairs in 3.1 represent three families of colored and three families of flavor degrees of freedom (notice that the flavor part is a singlet of \mathbf{a}_2^c). The two remaining \mathbf{a}_2 's may be related to gravity. We are currently working on this aspect with the hope to produce an \mathbf{e}_8 model of quantum gravity with an emergent discrete spacetime.

- 3) By considering a Jordan Pair $(\mathbf{J}_3^8, \bar{\mathbf{J}}_3^8)$ plus $\mathbf{e}_6 \oplus \mathbb{C}$ we get the 3-graded Lie algebra \mathbf{e}_7 , clearly shown in Figure 2. This shows why quadratic maps and Jordan Pairs are the natural Jordan product and structure inside \mathbf{e}_8 : $U_{x^\sigma} y^{-\sigma} = [[x^\sigma, y^{-\sigma}], x^\sigma]$, $\sigma = \pm$ whereas $[x^\sigma, y^\sigma] = 0$ as clearly shown by the 3-grading.
- 4) Other important algebraic structures easily visible in the Magic Star are related to the 5-grading of \mathbf{e}_8 : the Freudenthal Triple Systems (FTS), [55], in the case of \mathbf{e}_8 an irreducible 56-dimensional irreducible representation of \mathbf{e}_7 , and the Kantor Pair, [56], formed by two FTS'. They are shown in Figure 3. Denoting $\mathfrak{g} = \mathfrak{g}_{-2} + \mathfrak{g}_{-1} + \mathfrak{g}_0 + \mathfrak{g}_1 + \mathfrak{g}_2$ the 5-grading, we have \mathbf{e}_7 in grade 0 = $Der(FTS)$, an FTS in grade 1 (a 56-dim irrep of \mathbf{e}_7) and 1-dimensional $\mathfrak{g}_{\pm 2} := \mathbb{C}x_{\pm\rho}$. The non-deg skew bilinear $<, >$ and a fully symmetric quartic

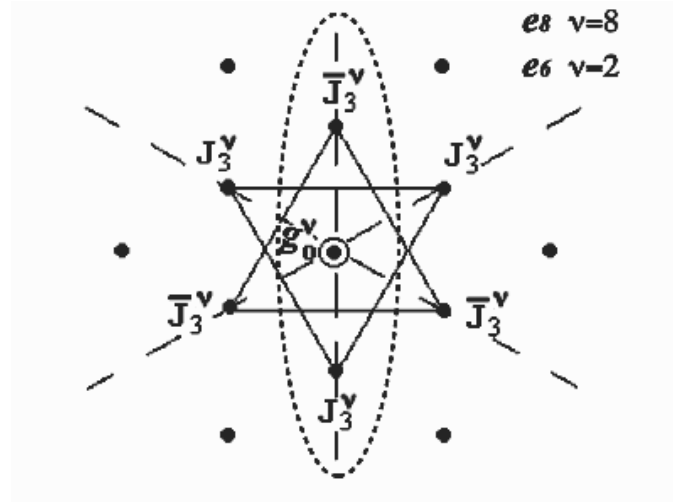


Figure 2: \mathbf{e}_7 (resp. \mathbf{a}_5) as 3-graded inside \mathbf{e}_8 (resp. \mathbf{e}_6)

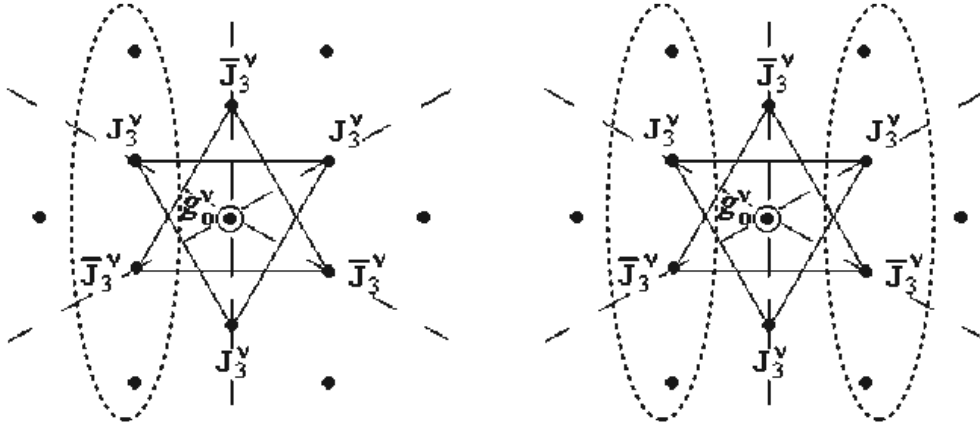


Figure 3: FTS (left) and KP (right) within \mathbf{e}_8

form q on \mathfrak{g}_1 , characteristic of the FTS, are obtained by:

$$\begin{aligned} [x, y] &= \langle x, y \rangle x_\rho & x, y \in \mathfrak{g}_1 \\ q(x, x, x, x) &= [x, [x, [x, [x, x_{-\rho}]]]] & x \in \mathfrak{g}_1 \end{aligned} \quad (3.2)$$

from which $\Rightarrow q(x_{\beta_1}, x_{\beta_2}, x_{\beta_3}, x_{\beta_4})$ follows by linearization.

The Kantor Pair is a pair of modules and within \mathbf{e}_8 is

$$KP = \mathbf{e}_8 \ominus (\mathbf{e}_7 \oplus \mathbf{a}_1)$$

4 Exceptional Periodicity and Extended Roots

Divinity drew Earth out of Emptiness as he drew One from Nothing to create many.
Pythagoras

We force the definition of root system to include what we call *extended* roots not obeying the symmetry by Weyl reflection, nor the fact that $2 \frac{(\alpha, \beta)}{(\alpha, \alpha)}$ be integer for all roots α, β .

Let $\{k_1, \dots, k_N\}$ be an orthonormal basis of an Euclidean space V of dimension N . For any $n = 1, 2, \dots$ we introduce $N = 4(n + 1)$ and define the extended roots of $\mathbf{e}_8^{(n)}$ as:

$$\mathbf{e}_8^{(n)}: \quad \begin{array}{lll} \pm k_i \pm k_j & 1 \leq i < j \leq N & 2N(N-1) \text{ roots} \\ \frac{1}{2}(\pm k_1 \pm k_2 \pm \dots \pm k_N) & \text{even } \# \text{ of } + & 2^{N-1} \text{ roots} \end{array} \quad (4.1)$$

This is a root system only in the case $n = 1$ and $\mathbf{e}_8^{(1)} = \mathbf{e}_8$.

The *extended* roots of $\mathbf{g}_2^{(n)}$ are obtained by projecting on the space spanned by $k_1 - k_2$ and $k_1 + k_2 - 2k_3$, hence $\mathbf{g}_2^{(n)} = \mathbf{g}_2$; those of $\mathbf{f}_4^{(n)}$ are the projection on the space spanned by k_1, k_2, \dots, k_{N-4} :

$$\mathbf{f}_4^{(n)}: \quad \begin{array}{ll} \pm k_i, \pm k_i \pm k_j & 1 \leq i < j \leq N-4 \\ \frac{1}{2}(\pm k_1 \pm k_2 \pm \dots \pm k_{N-4}) & \end{array} \quad (4.2)$$

and finally:

$$\mathbf{e}_6^{(n)}: \quad \begin{array}{ll} \pm k_i \pm k_j & 1 \leq i < j \leq N-3 \\ \frac{1}{2}(\pm k_1 \pm k_2 \pm \dots \pm (k_{N-2} + k_{N-1} + k_N)) & \text{even } \# \text{ of } + \end{array} \quad (4.3)$$

$$\mathbf{e}_7^{(n)}: \quad \begin{array}{ll} \pm(k_{N-1} + k_N) & \\ \pm k_i \pm k_j & 1 \leq i < j \leq N-2 \\ \frac{1}{2}(\pm k_1 \pm k_2 \pm \dots \pm (k_{N-1} + k_N)) & \text{even } \# \text{ of } + \end{array} \quad (4.4)$$

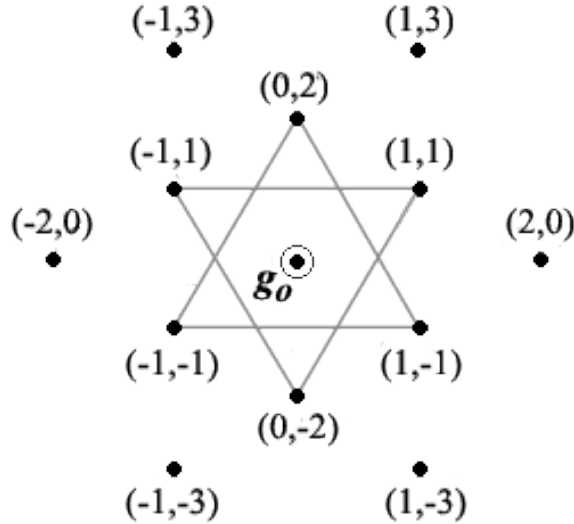


Figure 4: The Magic Star

All these sets of *extended* roots form a Magic star of figure 4, once projected on the plane spanned by $k_1 - k_2$ and $k_1 + k_2 - 2k_3$. In the figure the pair of integers (a, b) are the (Euclidean) scalar products of each root with $k_1 - k_2$ and $k_1 + k_2 - 2k_3$.

One can readily check that, upon a relabeling of the k 's, $\mathbf{e}_6^{(n)}$ is the center of the Magic Star of $\mathbf{e}_8^{(n)}$ and that $\mathbf{e}_7^{(n)} = \mathbf{e}_6^{(n)} \oplus T_{(a,b)} \oplus T_{(-a,-b)}$, for a fixed pair $(a,b) \in \{(1,1), (-1,1), (0,-2)\}$, where $\mathbf{e}_6^{(n)}$ is the center of the Magic Star and $T_{(a,b)}$ is the set of roots of $\mathbf{e}_8^{(n)}$ whose scalar products with $k_1 - k_2$ and $k_1 + k_2 - 2k_3$ are (a,b) .

The *rank* of $\mathbf{g}_2^{(n)}$, $\mathbf{f}_4^{(n)}$, $\mathbf{e}_6^{(n)}$, $\mathbf{e}_7^{(n)}$, $\mathbf{e}_8^{(n)}$ is the dimension of the vector space V spanned by their roots, namely 2, $N - 4$, $N - 2$, $N - 1$, N respectively.

By abuse of definition we shall often say *root*, for short, instead of *extended root*.

From now on we restrict to $\mathbf{e}_6^{(n)}$, $\mathbf{e}_7^{(n)}$, $\mathbf{e}_8^{(n)}$ and denote by $\mathcal{L}_{\mathbf{MS}}$ anyone of them, by Φ the set of extended roots of $\mathcal{L}_{\mathbf{MS}}$ and by R its rank. We recall that $N = 4(n+1)$, $n = 1, 2, \dots$, hence $R = N - 2 = 4n + 2$ for $\mathbf{e}_6^{(n)}$, $R = N - 1 = 4n + 3$ for $\mathbf{e}_7^{(n)}$, $R = N = 4n + 4$ for $\mathbf{e}_8^{(n)}$.

We denote by Φ_O and Φ_S the following subsets of Φ :

$$\begin{aligned}\Phi_O &= \{(\pm k_i \pm k_j) \in \Phi\} \\ \Phi_S &= \{\tfrac{1}{2}(\pm k_1 \pm k_2 \pm \dots \pm k_N) \in \Phi\}\end{aligned}\tag{4.5}$$

Remark 4.1. Notice that Φ_O is the root system of \mathbf{d}_{N-3} in the case of $\mathbf{e}_6^{(n)}$, of $\mathbf{d}_{N-2} \oplus a_1$ in the case of $\mathbf{e}_7^{(n)}$ and of \mathbf{d}_N in the case of $\mathbf{e}_8^{(n)}$.

We now prove the following results.

Proposition 4.2. For all $\rho \in \Phi_O$ and $x \in \Phi$: $2 \frac{(x, \rho)}{(\rho, \rho)} \in \mathbb{Z}$ and $w_\rho(x) = x - 2 \frac{(x, \rho)}{(\rho, \rho)} \rho \in \Phi$ (the set of extended roots is closed under the Weyl reflections by all $\rho \in \Phi_O$). The set of extended roots is closed under the Weyl reflections by all $\rho \in \Phi$ if and only if $n = 1$.

Proof. If $\rho \in \Phi_O$ then $(\rho, \rho) = 2$ and $(x, \rho) \in \{0, \pm 1, \pm 2\}$, hence $2 \frac{(x, \rho)}{(\rho, \rho)} \in \mathbb{Z}$ and $w_\rho(x) = x - (x, \rho)\rho$. If $(x, \rho) = 0$ then $w_\rho(x) = x \in \Phi$. If $(x, \rho) = \pm 1$ then $w_\rho(x) = x \mp \rho \in \Phi$ as we shall prove in proposition 4.4. If $(x, \rho) = \pm 2$ then $\rho = \pm x$ and $w_\rho(x) = -x \in \Phi$. Suppose now that both $x, \rho \in \Phi_S$ and write $x = \frac{1}{2} \sum \lambda_i k_i$, $\rho = \frac{1}{2} \sum \mu_i k_i$ where $\lambda_i, \mu_i \in \{-1, 1\}$. We can certainly pick an x such that $(x, \rho) = -n$ which occurs whenever for two indices j, ℓ $\lambda_j = \mu_j$ and $\lambda_\ell = \mu_\ell$ while $\lambda_i = -\mu_i$ for $i \neq j, \ell$. Hence $w_\rho(x) = x + 2 \frac{n}{n+1} \rho = \frac{1}{2} \sum \nu_i k_i$. We have that $|\nu_j| = \left| \lambda_j + 2 \frac{n}{n+1} \mu_j \right| = \frac{3n+1}{n+1} \geq 2$ and $|\nu_j| = 2$ if and only if $n = 1$, in which case Φ is the root system of a simple Lie algebra. For $n > 1$ there is no root with such a coefficient ν_j . \square

We introduce the basis $\Delta = \{\alpha_1, \dots, \alpha_R\}$ of Φ , with $\alpha_i = k_i - k_{i+1}$, $1 \leq i \leq R - 2$, $\alpha_{R-1} = k_{R-2} + k_{R-1}$ and $\alpha_R = -\frac{1}{2}(k_1 + k_2 + \dots + k_N)$; we order them by setting $\alpha_i > \alpha_{i+1}$:

$$\Delta = \{k_1 - k_2, k_2 - k_3, \dots, k_{R-2} - k_{R-1}, k_{R-2} + k_{R-1}, -\tfrac{1}{2}(k_1 + k_2 + \dots + k_N)\}\tag{4.6}$$

Proposition 4.3. The set Δ in (4.6) is a set of **simple extended roots**, by which we mean:

- i) Δ is a basis of the Euclidean space V of finite dimension R ;
- ii) every root β can be written as a linear combination of roots of Δ with all positive or all negative integer coefficients: $\beta = \sum \ell_i \alpha_i$ with $\ell_i \geq 0$ or $\ell_i \leq 0$ for all i .

For all roots $\beta = \sum_{i=1}^R m_i \alpha_i$ the coefficient m_R is such that:

$$\begin{aligned} m_R &\in \{0, \pm 2\} & \text{if } \beta \in \Phi_O \\ m_R &= \pm 1 & \text{if } \beta \in \Phi_S \end{aligned} \quad (4.7)$$

Proof. The set $\Delta = \{\alpha_1, \dots, \alpha_R\}$ is obviously a basis in V . Let $\mathbf{u} = k_{N-2} + k_{N-1} + k_N, k_{N-1} + k_N, k_N$ for $\mathbf{e}_6^{(n)}, \mathbf{e}_7^{(n)}, \mathbf{e}_8^{(n)}$ respectively. We have:

$$\begin{aligned} k_{R-1} &= \frac{1}{2}(\alpha_{R-1} - \alpha_{R-2}) \\ k_i &= \alpha_i + k_{i+1} = \sum_{\ell=i}^{R-2} \alpha_\ell + \frac{1}{2}(\alpha_{R-1} - \alpha_{R-2}), \quad 1 \leq i \leq R-2 \\ \mathbf{u} &= -2\alpha_R - \sum_{\ell=1}^{R-2} \ell \alpha_\ell - \frac{R-1}{2}(\alpha_{R-1} - \alpha_{R-2}) \end{aligned} \quad (4.8)$$

from which we obtain, for $1 \leq i < j \leq R-1$ and forcing $\sum_{\ell=r}^s \alpha_\ell = 0$ if $r > s$:

$$\left. \begin{aligned} k_i + k_j &= \sum_{\ell=i}^{R-3} \alpha_\ell + \sum_{\ell=j}^{R-2} \alpha_\ell + \alpha_{R-1} \\ k_i - k_j &= \sum_{\ell=i}^{j-1} \alpha_\ell \end{aligned} \right\} 1 \leq i < j \leq R-1 \quad (4.9)$$

for $\mathbf{e}_7^{(n)}$, namely for $R = N-1 = 4n+3$ and $\mathbf{u} = k_{N-1} + k_N$:

$$-\mathbf{u} = 2\alpha_R + \sum_{\ell=1}^{R-3} \ell \alpha_\ell + 2n \alpha_{R-2} + (2n+1) \alpha_{R-1}, \quad \text{for } \mathbf{e}_7^{(n)} \quad (4.10)$$

for $\mathbf{e}_8^{(n)}$, namely for $R = N$:

$$\begin{aligned} \pm k_i - k_N &= 2\alpha_N + \sum_{\ell=1}^{i-1} \ell \alpha_\ell + \sum_{\ell=i}^{N-3} (\ell \pm 1) \alpha_\ell + (2n + \frac{1 \pm 1}{2}) \alpha_{N-2} \\ &\quad + (2n + 1 + \frac{1 \pm 1}{2}) \alpha_{N-1}, \quad i \leq N-2 \end{aligned} \quad (4.11)$$

$$\pm k_{N-1} - k_N = 2\alpha_N + \sum_{\ell=1}^{N-3} \ell \alpha_\ell + (2n + \frac{1 \mp 1}{2}) \alpha_{N-2} + (2n + 1 + \frac{1 \mp 1}{2}) \alpha_{N-1}$$

We see that all the roots in (4.9), (4.10), (4.11) are the sum of simple roots with all positive integer coefficients. These are half of the roots in Φ_O and they are all positive roots. The rest of the roots in Φ_O are negative and are obviously the sum of simple roots with integer coefficients that are all negative.

Finally all the roots in Φ_S that contain $-\frac{1}{2}\mathbf{u}$ can be obtained from α_R by flipping an even number of signs and this is done by adding to α_R a certain number of terms of the type $k_i + k_j$, $1 \leq i < j \leq R-1$. These are all positive roots, a half of the roots in Φ_S and are linear combination of simple roots with integer coefficients that are all positive.

The negative roots are similarly obtained by adding to $-\alpha_N$ a certain number of terms of the type $-(k_i + k_j)$, $1 \leq i < j \leq R-1$, and are linear combination of simple roots with integer coefficients that are all negative.

As a consequence we easily obtain 4.7. \square

Proposition 4.4. For each $\alpha \in \Phi_O, \beta \in \Phi$ the scalar product $(\alpha, \beta) \in \{\pm 2, \pm 1, 0\}$; $\alpha + \beta$ (respectively $\alpha - \beta$) is a root if and only if $(\alpha, \beta) = -1$ (respectively $+1$); if both $\alpha + \beta$ and $\alpha - \beta$ are not in $\Phi \cup \{0\}$ then $(\alpha, \beta) = 0$.

For each $\alpha, \beta \in \Phi_S$ the scalar product $(\alpha, \beta) \in \{\pm(n+1), \pm n, \pm(n-1), \dots, 0\}$; $\alpha + \beta$ (respectively $\alpha - \beta$) is a root if and only if $(\alpha, \beta) = -n$ (respectively $+n$).

For $\alpha, \beta \in \Phi$ if $\alpha + \beta$ is a root then $\alpha - \beta$ is not a root.

Proof. If both $\alpha, \beta \in \Phi_O$ the proof follows from the fact that Φ_O is the root system of a simply laced Lie algebra. If $\alpha \in \Phi_O, \beta \in \Phi_S$ then obviously $(\alpha, \beta) \in \{\pm 1, 0\}$. Moreover, let us write $\alpha = \sigma_i k_i + \sigma_j k_j, i < j, \sigma_{i,j} \in \{-1, 1\}$. Then $(\alpha, \beta) = -1$ if and only if $\beta = \frac{1}{2}(\pm k_1 \pm \dots - \sigma_i k_i \pm \dots - \sigma_j k_j \pm \dots \pm k_N)$, which is true if and only if $\alpha + \beta \in \Phi$ (in particular $\alpha + \beta \in \Phi_S$). Similarly $(\alpha, \beta) = 1$ if and only if $\alpha - \beta \in \Phi$. As a consequence, if both $\alpha \pm \beta \notin \Phi \cup \{0\}$ then $(\alpha, \beta) \neq \pm 1$ and also $(\alpha, \beta) \neq \pm 2$ because $(\alpha, \beta) = \pm 2$ if and only if $\alpha = \pm \beta$; therefore $(\alpha, \beta) = 0$.

If both $\alpha, \beta \in \Phi_S$, then all their signs but an even number $2m$ must be equal, $m = 0, \dots, N/2 = 2(n+1)$ and we get $(\alpha, \beta) = \frac{1}{4}(N - 2m - 2m) = n + 1 - m = n + 1, n, n - 1, \dots, -(n + 1)$ for $m = 0, \dots, 2(n + 1)$. Moreover, since $\pm k_i \pm k_i \in \{0, \pm 2k_i\}, i = 1, \dots, N$ then $\alpha + \beta \in \Phi$ if and only if all signs are opposite but 2 (in which case $\alpha + \beta$ is actually in Φ_O) and this is true if and only if $(\alpha, \beta) = -\frac{1}{4}(N - 4) = -n$. Similarly $(\alpha, \beta) = n$ if and only if $\alpha - \beta \in \Phi$. The last statement of the Proposition follows trivially. \square

The preceding propositions extend to EP analogous results for the root system of simply laced Lie algebras. They are essential in determining the structure of the algebra $\mathcal{L}_{\mathbf{MS}}$ that we associate to the extended roots in the next section.

5 The $\mathcal{L}_{\mathbf{MS}}$ algebra

Mathematics is the art of giving the same name to different things.

H. Poincaré

We define the $\mathcal{L}_{\mathbf{MS}}$ algebra (as before $\mathcal{L}_{\mathbf{MS}}$ is either $\mathbf{e}_6^{(n)}$ or $\mathbf{e}_7^{(n)}$ or $\mathbf{e}_8^{(n)}$) by extending the construction of a Lie algebra from a root system, [57] [58] [59]. In particular we generalize the algorithm in [59] for simply laced Lee algebras, since also in our set of extended roots the β chain through α , namely the set of roots $\alpha + c\beta, c \in \mathbb{Z}$, has length one.

We give $\mathcal{L}_{\mathbf{MS}}$ an algebra structure of rank R over a field extension \mathfrak{F} of the rational integers \mathbb{Z} in the following way ⁶ :

- a) we select the set of simple extended roots $\Delta = \{\alpha_1, \dots, \alpha_R\}$ of Φ
- b) we select a basis $\{h_1, \dots, h_R\}$ of the R -dimensional vector space H over \mathfrak{F} and set $h_\alpha = \sum_{i=1}^R c_i h_i$ for each $\alpha \in \Phi$ such that $\alpha = \sum_{i=1}^R c_i \alpha_i$
- c) we associate to each $\alpha \in \Phi$ a one-dimensional vector space L_α over \mathfrak{F} spanned by x_α
- d) we define $\mathcal{L}_{\mathbf{MS}} = H \oplus_{\alpha \in \Phi} L_\alpha$ as a vector space over \mathfrak{F}
- e) we give $\mathcal{L}_{\mathbf{MS}}$ an algebraic structure by defining the following multiplication on the basis $\mathfrak{B}_{LS} = \{h_1, \dots, h_R\} \cup \{x_\alpha \mid \alpha \in \Phi\}$, extended by linearity to a bilinear multiplication $\mathcal{L}_{\mathbf{MS}} \times \mathcal{L}_{\mathbf{MS}} \rightarrow \mathcal{L}_{\mathbf{MS}}$:

$$\begin{aligned}
[h_i, h_j] &= 0, \quad 1 \leq i, j \leq R \\
[h_i, x_\alpha] &= -[x_\alpha, h_i] = (\alpha, \alpha_i) x_\alpha, \quad 1 \leq i \leq R, \quad \alpha \in \Phi \\
[x_\alpha, x_{-\alpha}] &= -h_\alpha \\
[x_\alpha, x_\beta] &= 0 \text{ for } \alpha, \beta \in \Phi \text{ such that } \alpha + \beta \notin \Phi \text{ and } \alpha \neq -\beta \\
[x_\alpha, x_\beta] &= \varepsilon(\alpha, \beta) x_{\alpha+\beta} \text{ for } \alpha, \beta \in \Phi \text{ such that } \alpha + \beta \in \Phi
\end{aligned} \tag{5.1}$$

⁶Specifically, we will take \mathfrak{F} to be the complex field \mathbb{C} .

where $\varepsilon(\alpha, \beta)$ is the *asymmetry function*, introduced in [60], see also [59], defined as follows:

Definition 5.1. Let \mathbb{L} denote the lattice of all linear combinations of the simple extended roots with integer coefficients

$$\mathbb{L} = \left\{ \sum_{i=1}^R c_i \alpha_i \mid c_i \in \mathbb{Z}, \alpha_i \in \Delta \right\} \quad (5.2)$$

the asymmetry function $\varepsilon(\alpha, \beta) : \mathbb{L} \times \mathbb{L} \rightarrow \{-1, 1\}$ is defined by:

$$\varepsilon(\alpha, \beta) = \prod_{i,j=1}^R \varepsilon(\alpha_i, \alpha_j)^{\ell_i m_j} \quad \text{for } \alpha = \sum_{i=1}^R \ell_i \alpha_i, \beta = \sum_{j=1}^R m_j \alpha_j \quad (5.3)$$

where $\alpha_i, \alpha_j \in \Delta$ and

$$\varepsilon(\alpha_i, \alpha_j) = \begin{cases} -1 & \text{if } i = j \\ -1 & \text{if } \alpha_i + \alpha_j \text{ is a root and } \alpha_i < \alpha_j \\ +1 & \text{otherwise} \end{cases} \quad (5.4)$$

We now show some properties of the asymmetry function. In particular we show that, for $\alpha, \beta, \alpha + \beta \in \Phi$, $\varepsilon(\alpha, \beta) = -\varepsilon(\beta, \alpha)$ which implies that the bilinear product (5.1) is antisymmetric.

Proposition 5.2. The asymmetry function ε satisfies, for $\alpha, \beta, \gamma, \delta \in \mathbb{L}$, $\alpha = \sum m_i \alpha_i$ and $\beta = \sum n_i \alpha_i$:

$$\begin{aligned} i) \quad & \varepsilon(\alpha + \beta, \gamma) = \varepsilon(\alpha, \gamma) \varepsilon(\beta, \gamma) \\ ii) \quad & \varepsilon(\alpha, \gamma + \delta) = \varepsilon(\alpha, \gamma) \varepsilon(\alpha, \delta) \\ iii) \quad & \varepsilon(\alpha, \alpha) = (-1)^{\frac{1}{2}(\alpha, \alpha) - m_R^2 \frac{n-1}{2}} \\ iv) \quad & \varepsilon(\alpha, \beta) \varepsilon(\beta, \alpha) = (-1)^{(\alpha, \beta) - m_R n_R (n-1)} \\ v) \quad & \varepsilon(0, \beta) = \varepsilon(\alpha, 0) = 1 \\ vi) \quad & \varepsilon(-\alpha, \beta) = \varepsilon(\alpha, \beta)^{-1} = \varepsilon(\alpha, \beta) \\ vii) \quad & \varepsilon(\alpha, -\beta) = \varepsilon(\alpha, \beta)^{-1} = \varepsilon(\alpha, \beta) \end{aligned}$$

Proof. The first two properties follow directly from the definition. In order to prove *iii)* we first notice that for $\alpha_i, \alpha_j \in \Delta$, $i \neq j$, $(\alpha_i, \alpha_j) \in \{0, -1\}$. Therefore:

$$\begin{aligned} \varepsilon(\alpha, \alpha) &= \prod_{1 \leq i, j \leq R} \varepsilon(\alpha_i, \alpha_j)^{m_i m_j} = \prod_{1 \leq i < j \leq R} (-1)^{m_i m_j (\alpha_i, \alpha_j)} \prod_{1 \leq i \leq R} (-1)^{m_i^2} \\ &= (-1)^{\sum_{1 \leq i < j \leq R} m_i m_j (\alpha_i, \alpha_j) + \frac{1}{2} \sum_{1 \leq i \leq R} m_i^2 (\alpha_i, \alpha_i) - \frac{1}{2} \sum_{1 \leq i \leq R} m_i^2 ((\alpha_i, \alpha_i) - 2)} \\ &= (-1)^{\frac{1}{2}(\alpha, \alpha) - \frac{1}{2} m_R^2 ((\alpha_R, \alpha_R) - 2)} = (-1)^{\frac{1}{2}(\alpha, \alpha) - m_R^2 \frac{n-1}{2}} \end{aligned} \quad (5.5)$$

Property *iv)* follows from *iii)* by replacing α with $\alpha + \beta$ and using the first two. If $\alpha = \sum m_i \alpha_i$, $\beta = \sum n_i \alpha_i$ and $\alpha + \beta = \sum \ell_i \alpha_i = \sum (m_i + n_i) \alpha_i$ we get:

$$\begin{aligned} \varepsilon(\alpha + \beta, \alpha + \beta) &= (-1)^{\frac{1}{2}(\alpha + \beta, \alpha + \beta) - \ell_R^2 \frac{n-1}{2}} = (-1)^{\frac{1}{2}(\alpha, \alpha) + \frac{1}{2}(\beta, \beta) + (\alpha, \beta) - \ell_R^2 \frac{n-1}{2}} \\ &= (-1)^{\frac{1}{2}(\alpha, \alpha) - m_R^2 \frac{n-1}{2}} (-1)^{\frac{1}{2}(\beta, \beta) - n_R^2 \frac{n-1}{2}} \varepsilon(\alpha, \beta) \varepsilon(\beta, \alpha) \end{aligned} \quad (5.6)$$

from which the property follows. Property *v)* is a trivial consequence of the definition, whereas properties *vi)* and *vii)* follow from property *v)* together with *i)* and *ii)*. \square

Proposition 5.3. *If $\alpha, \beta, \alpha + \beta \in \Phi$ then:*

- i) $\varepsilon(\alpha, \alpha) = -1$ $\alpha \in \Phi$
- ii) $\varepsilon(\alpha, \beta) = -\varepsilon(\beta, \alpha)$ $\alpha, \beta, (\alpha + \beta) \in \Phi$ *antisymmetry*
- iii) $\varepsilon(\alpha, \beta) = \varepsilon(\beta, \alpha + \beta)$ *if $\alpha, \alpha + \beta \in \Phi$, $\beta \in \mathbb{L}$*
- iv) $\varepsilon(\alpha, \beta) = \varepsilon(\beta, \alpha - \beta)$ *if $\alpha, \alpha - \beta \in \Phi$, $\beta \in \mathbb{L}$*

Proof. By (4.7) if $\alpha \in \Phi_O$ then $(\alpha, \alpha) = 2$ and $m_R^2/2$ is even, hence $\varepsilon(\alpha, \alpha) = -1$. If $\alpha \in \Phi_S$ then $(\alpha, \alpha) = n + 1$ and $m_R^2 = 1$. Therefore if $(-1)^{\frac{1}{2}(\alpha, \alpha) - m_R^2 \frac{n-1}{2}} = (-1)^{\frac{1}{2}(n+1-n+1)} = -1$. As a consequence, if $\alpha, \beta, \alpha + \beta \in \Phi$, then $-1 = \varepsilon(\alpha + \beta, \alpha + \beta) = \varepsilon(\alpha, \beta)\varepsilon(\beta, \alpha)$, hence $\varepsilon(\alpha, \beta) = -\varepsilon(\beta, \alpha)$

Finally we prove iii) - similar proof for iv):

$$\varepsilon(\alpha, \beta) = \varepsilon(\alpha, \alpha - \alpha + \beta) = \varepsilon(\alpha, \alpha)\varepsilon(\alpha, \alpha + \beta) = -\varepsilon(\alpha + \beta - \beta, \alpha + \beta) = \varepsilon(\beta, \alpha + \beta)$$

□

6 T-algebras

The essence of Mathematics lies in its freedom.

G. Cantor

We now concentrate on the set of roots $T_{(a,b)}$ on a tip of the Magic Star. In order to fix one, let us consider $T_{(1,1)}$ and denote it simply by T . The reader will forgive us if we shall be concise whenever there is no risk of misunderstanding and use T to denote both the set of roots and the set of elements in $\mathcal{L}_{\mathbf{MS}}$ associated to those roots. An element of T is an \mathfrak{F} -linear combination of x_α, x_β, \dots for α, β, \dots in $T_{(1,1)}$. We give T an algebraic structure with a symmetric product, thus mimicking the case $n = 1$ when T is a Jordan algebra.

Let us first show what happens in the case $n = 1$, in particular for \mathbf{e}_8 . It is proven in [37] that $T = J_{(1,1)}$ is the Jordan Algebra \mathbf{J}_3^8 of 3×3 Hermitean matrices over the octonions \mathbb{O} . Let us denote by $P_1 := E_{11}$, $P_2 := E_{22}$, $P_3 := E_{33}$ the three trace-one idempotents, whose sum is the identity in \mathbf{J}_3^8 , E_{ii} representing the matrix with a 1 in the (ii) position and zero elsewhere. Let us identify the algebra $\mathbf{d}_4 \subset \mathbf{f}_4 \subset \mathbf{e}_6 \subset \mathbf{e}_6 \oplus \mathbf{a}_2^c \subset \mathbf{e}_8$ with the roots $\pm k_i \pm k_j$, $4 \leq i < j \leq 7$ and P_1, P_2, P_3 with the elements of $J_{(1,1)}$ which are left invariant by \mathbf{d}_4 , [61]. This uniquely identifies P_1, P_2, P_3 with the roots $k_1 + k_8$, $k_1 - k_8$ and $-k_2 - k_3$. Since $\mathbf{e}_7 \subset \mathbf{e}_8$ is 3-graded, one can define a 3-linear product on $J_{(1,1)}$, [51], and a Jordan product through the correspondence with the quadratic formulation of Jordan algebras.

Still in the case of \mathbf{e}_8 ($\mathbf{e}_8^{(n)}$ for $n = 1$), if we identify the \mathbf{e}_8 generators corresponding to the roots in Φ_O with the bosons and those corresponding to the roots in Φ_S with the fermions, then the 11 bosons in T are $k_1 \pm k_j$, $j = 4, \dots, 8$ and $-k_2 - k_3$. The three bosons $k_1 \pm k_8$ and $-k_2 - k_3$, correspond to the three diagonal idempotents $x_{P_1}, x_{P_2}, x_{P_3}$, which are left invariant by \mathbf{d}_4 , whose roots are $\pm k_i \pm k_j$, $4 \leq i < j \leq 7$.

The remaining 8 vectors (bosons) and the 2x8 spinors (fermions) are linked by triality. Notice that one spinor has $\frac{1}{2}(k_1 - k_2 - k_3 + k_8)$ fixed and even number of + signs among $\pm k_4 \pm k_5 \pm k_6 \pm k_7$, while the other one has $\frac{1}{2}(k_1 - k_2 - k_3 - k_8)$ fixed and odd number of + signs among $\pm k_4 \pm k_5 \pm k_6 \pm k_7$. So they are both 8 dimensional representations of \mathbf{d}_4 .

Therefore a generic element of \mathbf{J}_3^8 is a linear superposition of 3 diagonal elements plus one vector and two spinors (or a bispinor). The vector can be viewed as the (12) octonionic entry of the

matrix (plus its octonionic conjugate in the (21) position) and the two 8-dimensional spinors as the (31) and (23) entry (plus their respective octonionic conjugate in the (31) and (23) position).

We do the same for $n \geq 1$ and consider the most general case $\mathbf{e}_8^{(n)}$ (the other two cases $\mathbf{e}_6^{(n)}$ and $\mathbf{e}_7^{(n)}$ being a restriction of it).

Let us define x_{P_1} , x_{P_2} and x_{P_3} as the elements of $\mathcal{L}_{\mathbf{MS}}$ in T associated to the roots $\rho_1 := k_1 + k_N$, $\rho_2 := k_1 - k_N$ and $\rho_3 := -k_2 - k_3$:

$$x_{P_1} \leftrightarrow \rho_1 := k_1 + k_N ; x_{P_2} \leftrightarrow \rho_2 := k_1 - k_N ; x_{P_3} \leftrightarrow \rho_3 := -k_2 - k_3 \quad (6.1)$$

They are left invariant by the Lie subalgebra $\mathbf{d}_{N-4} = \mathbf{d}_{4n}$, whose roots are $\pm k_i \pm k_j$, $4 \leq i < j \leq N-1$.

We denote by T_O the set of roots in $T \cap \Phi_O$, by T'_O the set of roots in T_O that are not ρ_1, ρ_2, ρ_3 and by T_S the set of roots in $T \cap \Phi_S$. In the case we are considering, where $T = T_{(1,1)}$ we have $T'_O = \{k_1 \pm k_j, j = 4, \dots, N-1\}$ and $T_S = \{\frac{1}{2}(k_1 - k_2 - k_3 \pm k_4 \pm \dots \pm k_N)\}$, even # of $+$. We further split T_S into $T_S^+ = \{\frac{1}{2}(k_1 - k_2 - k_3 \pm k_4 \pm \dots + k_N)\}$ and $T_S^- = \{\frac{1}{2}(k_1 - k_2 - k_3 \pm k_4 \pm \dots - k_N)\}$. Then $v \in T'_O$ is an $8n$ -dimensional vector and $s^\pm \in T_S^\pm$ are 2^{4n-1} -dimensional spinors of \mathbf{d}_{4n} .

We write a generic element x of T as $x = \sum \lambda_i x_{P_i} + x_v + x_{s^+} + x_{s^-}$ where

$$x_v = \sum_{\alpha \in T'_O} \lambda_\alpha^v x_\alpha \quad (6.2)$$

$$x_{s^\pm} = \sum_{\alpha \in T_S^\pm} \lambda_\alpha^{s^\pm} x_\alpha \quad (6.3)$$

We view λ_α^v as a coordinate of the vector λ^v and $\lambda_\alpha^{s^\pm}$ as a coordinate of the spinor λ^{s^\pm} ; we denote by $\bar{\lambda}^v$ ($\bar{\lambda}^{s^\pm}$) the vectors (spinor) in the dual space with respect to the appropriate bilinear form and view x as a 3×3 Hermitean matrix:

$$\begin{pmatrix} \lambda_1 & \lambda_v & \bar{\lambda}_{s^+} \\ \bar{\lambda}_v & \lambda_2 & \lambda_{s^-} \\ \lambda_{s^+} & \bar{\lambda}_{s^-} & \lambda_3 \end{pmatrix} \quad (6.4)$$

whose entries have the following \mathfrak{F} -dimensions:

- 1 for the *scalar* diagonal elements $\lambda_1, \lambda_2, \lambda_3$;
- $8n$ for the *vector* λ_v ;
- 2^{4n-1} for the *spinors* λ_{s^\pm} .

We see that only for $n = 1$ the dimension of the vector and the spinors is the same, whereas for $n > 1$ the entries in the (12), (21) position have different dimension than those in the (31), (13), (23), (32) position. Nevertheless we now define a symmetric product of the elements in T , which then becomes a generalization of the Jordan algebra \mathbf{J}_3^8 in a very precise sense. This type of generalization of the Jordan Algebra is known in the literature, [44], where it is called *T-algebra*.

We denote by I the element $I := x_{P_1} + x_{P_2} + x_{P_3}$ and by I^- the element $I^- := -\bar{x}_{P_1} - \bar{x}_{P_2} - \bar{x}_{P_3}$ of $\bar{T} := T_{(-1,-1)}$, where \bar{x}_{P_1} , \bar{x}_{P_2} and \bar{x}_{P_3} are associated to the roots $-k_1 - k_N$, $-k_1 + k_N$ and $k_2 + k_3$.

We give T an algebraic structure by introducing the symmetric product, see Proposition 6.2:

$$x \circ y := \frac{1}{2}[[x, I^-], y] \quad , \quad x, y \in T \quad (6.5)$$

We introduce the trace (see Proposition 6.2) $tr(x) \in \mathfrak{F}$ for $x \in T$ in the following way:

$$\text{for } x = \ell_1 x_{P1} + \ell_2 x_{P2} + \ell_3 x_{P3} + \sum_{\alpha \neq \rho_1, \rho_2, \rho_3} \ell_\alpha x_\alpha \quad tr(x) = \ell_1 + \ell_2 + \ell_3 \quad (6.6)$$

We denote by $tr(x, y) := tr(x \circ y)$ and by $x^2 := x \circ x$. For each $x \in T$ we define

$$x^\# = x^2 - tr(x)x - \frac{1}{2}(tr(x^2) - tr(x)^2)I \quad (6.7)$$

and we say that x is rank-1 if $x^\# = 0$. Notice that $tr(x^\#) = -\frac{1}{2}(tr(x^2) - tr(x)^2)$, therefore $x^\# = 0$ implies $x^2 = tr(x)x$: a rank-1 element of T is either a nilpotent or a multiple of a *primitive* idempotent $u \in T$: $u^2 = u$ and $tr(u) = 1$.

Let us also introduce $N(x) \in \mathfrak{F}$ for $x \in T$ in the following way:

$$N(x) = \frac{1}{6} \{tr(x)^3 - 3tr(x)tr(x^2) + 2tr(x^3)\} = \frac{1}{3}tr(x^\#, x) \quad (6.8)$$

where $x^2 = x \circ x$ and $x^3 = x \circ x^2$.

Since by (6.8) a rank-1 element x has $N(x) = 0$ any element of T falls into the following classification:

$$\begin{aligned} \text{rank-1: } & x^\# = 0 \\ \text{rank-2: } & x^\# \neq 0, N(x) = 0 \\ \text{rank-3: } & N(x) \neq 0 \end{aligned} \quad (6.9)$$

Examples of rank-2 and rank-3 elements are $x_{P1} + x_{P2}$ and I respectively. We finally now show the main properties of the algebra T .

Proposition 6.1. *Let $x, y \in T$ then $[[x, I^-], y] + [[I^-, y], x] = 0$.*

Proof. By the linearity it is sufficient to prove the proposition for basis elements $X_0 \in \{\bar{x}_{P1}, \bar{x}_{P2}, \bar{x}_{P3}\}$ and basis elements X_1, X_2 of T and the statement of the Proposition amounts to $J_0 + J_2 = 0$, with $J_0 := [[X_0, X_1], X_2]$ and $J_2 := [[X_2, X_0], X_1]$. Set $X_0 = x_\alpha$, $\alpha \in \{-k_1 \pm k_N, k_2 + k_3\}$, $X_1 = x_\beta$, $X_2 = x_\gamma$. Since $X_1, X_2 \in T$ then $[X_1, X_2] = 0$ and $\beta + \gamma \notin \Phi \cup \{0\}$; we also have that $\alpha \neq \beta, \alpha \neq \gamma$.

If none of the sums $\alpha + \beta$, $\alpha + \gamma$ is in Φ nor is 0, then $J_0 + J_2 = 0$. Also if $\beta = \gamma$ then $J_0 + J_2 = 0$ is trivially satisfied.

From now on at least one of $\alpha + \beta$, $\alpha + \gamma$, is in $\Phi \cup \{0\}$. Suppose first that $\alpha + \beta = 0$. Then $J_0 = -(\gamma, \alpha)x_\gamma$ and we have the following possibilities:

- a1) if $\alpha + \gamma = 0$ then $\beta = \gamma$ and $J_0 + J_2 = 0$ becomes trivial;
- a2) if $\alpha + \gamma \in \Phi$ then $(\gamma, \alpha) = -1$ by Proposition 4.4 and $J_0 = x_\gamma$; moreover $J_2 = \varepsilon(\gamma, \alpha)\varepsilon(\alpha + \gamma, -\alpha)x_\gamma = \varepsilon(\alpha, -\alpha)x_\gamma = -x_\gamma$ (because of Propositions 5.2 and 5.3) and $J_0 + J_2 = 0$ is verified;
- a3) if both $\alpha \pm \gamma \notin \Phi \cup \{0\}$ then $(\alpha, \gamma) = 0$, by Proposition 4.4 being $\alpha \in \Phi_O$, hence $J_0 = J_2 = 0$.

Similarly if we suppose $\alpha + \gamma = 0$.

From now on $\alpha, \beta, \gamma \in \Phi$, $\alpha \neq \pm\beta \neq \pm\gamma \neq \pm\alpha$ and $\alpha + \beta + \gamma \in \Phi$ (notice that $\alpha + \beta + \gamma$ cannot be zero for $\alpha \in T_{(-1,-1)}$ and $\beta, \gamma \in T_{(1,1)}$ and if $\alpha + \beta + \gamma \notin \Phi$ then obviously $J_0 = J_2 = 0$).

If none of $\alpha + \beta, \alpha + \gamma$ is in Φ then $J_0 = J_2 = 0$. If $\alpha + \beta \in \Phi$ then also $\alpha + \gamma \in \Phi$ (and viceversa): if $\gamma \in \Phi_O$ then, by Proposition 4.4, $(\alpha + \beta, \gamma) = -1 = (\alpha, \gamma) + (\beta, \gamma) = (\alpha, \gamma)$, implying $\alpha + \gamma \in \Phi$; if $\beta \in \Phi_O$ and $\gamma \in \Phi_S$ then $(\beta, \gamma) \in \{0, 1\}$ and $(\alpha + \beta, \gamma) = -1 = (\alpha, \gamma) + (\beta, \gamma)$ implies $(\alpha, \gamma) = -1$ and $(\beta, \gamma) = 0$; if $\beta, \gamma \in \Phi_S$ then $(\alpha + \beta) \in \Phi_S$ and $(\alpha + \beta, \gamma) = -n = (\alpha, \gamma) + (\beta, \gamma)$, with $(\alpha, \gamma) \in \{0, \pm 1\}$ since $\alpha \in \Phi_0$. But $(\alpha, \gamma) = 0$ implies $\beta + \gamma \in \Phi$ and $(\alpha, \gamma) = 1$ implies $(\beta, \gamma) = -(n + 1)$ hence $\beta = -\gamma$ which both contraddict the hypothesis. Hence $(\alpha, \gamma) = -1$ and $\alpha + \gamma \in \Phi$.

Finally, if $\alpha + \beta, \alpha + \gamma \in \Phi$ then

$$[[x_\alpha, x_\beta], x_\gamma] + [[x_\gamma, x_\alpha], x_\beta] = -\varepsilon(\alpha, \beta)\varepsilon(\gamma, \alpha)\varepsilon(\beta, \gamma) + \varepsilon(\gamma, \alpha)\varepsilon(\gamma, \beta)\varepsilon(\alpha, \beta) = \varepsilon(\alpha, \beta)\varepsilon(\gamma, \alpha)(\varepsilon(\gamma, \beta) - \varepsilon(\beta, \gamma)) \quad (6.10)$$

By *iv*) of Proposition 5.2 plus Proposition 4.4 and (4.7) we have $\varepsilon(\beta, \gamma)\varepsilon(\gamma, \beta) = (-1)^{(\beta, \gamma) - m_N n_N (n-1)}$. If $\beta \in \Phi_O$ or $\gamma \in \Phi_O$ then $\varepsilon(\beta, \gamma)\varepsilon(\gamma, \beta) = (-1)^{(\beta, \gamma)}$ and $(\beta, \gamma) \in \{0, 1\}$. Suppose $(\beta, \gamma) = 1$. Then either $\alpha + \gamma \in \Phi_O$ or $\beta \in \Phi_O$ and $(\alpha + \gamma, \beta) = -1$; but also $(\alpha + \gamma, \beta) = (\alpha, \beta) + (\gamma, \beta) = -1 + 1 = 0$, a contraddiction. If both $\beta, \gamma \in \Phi_S$ then $(\alpha + \gamma, \beta) = -n$ and $(\alpha, \beta) = -1$. So $(\beta, \gamma) = -(n - 1)$ and $\varepsilon(\beta, \gamma)\varepsilon(\gamma, \beta) = (-1)^{(\beta, \gamma) - m_N n_N (n-1)} = (-1)^{-(n-1) \pm (n-1)} = 1$, therefore $\varepsilon(\beta, \gamma) = \varepsilon(\gamma, \beta)$ and (6.10) is zero.

This ends the proof of Proposition 6.1. \square

Proposition 6.2. *The product $x, y \rightarrow x \circ y$ is symmetric. With respect to this product the element $I \in T$ is the identity and the elements $x_{P1}, x_{P2}, x_{P3} \in T$ are trace-one idempotents, hence are rank-1 elements of T . All $x_\alpha \in T$ but x_{P1}, x_{P2}, x_{P3} are nilpotent and trace-0, hence they are also rank-1 elements of T . The form $x \rightarrow \text{tr}(x)$ is a trace form, namely $\text{tr}(x, y)$ is bilinear and symmetric in x, y and the associative property $\text{tr}(x \circ y, z) = \text{tr}(x, y \circ z)$ holds for every $x, y, z \in T$.*

Proof. By Proposition 6.1:

$$[[x, I^-], y] + [[I^-, y], x] = 0 \Rightarrow [[x, I^-], y] = [[y, I^-], x] \Rightarrow x \circ y = y \circ x \quad (6.11)$$

We also have

$$[x_{Pi}, -\bar{x}_{Pj}] = \delta_{ij} h_{\rho_i}, \quad i, j = 1, 2, 3 \quad (6.12)$$

and for any $\alpha \in T_{(r,s)}$, being $\rho_1 + \rho_2 + \rho_3 = 2k_1 - k_2 - k_3$:

$$I \circ x_\alpha = \frac{1}{2}[[I, I^-], x_\alpha] = \frac{1}{2}[h_{\rho_1} + h_{\rho_2} + h_{\rho_3}, x_\alpha] = \frac{1}{2}(\rho_1 + \rho_2 + \rho_3, \alpha)x_\alpha = x_\alpha \quad (6.13)$$

By linearity this extends to any $x \in T$.

Moreover:

$$x_{Pi} \circ x_{Pi} = \frac{1}{2}[[x_{Pi}, I^-], x_{Pi}] = \frac{1}{2}[h_{\rho_i}, x_{Pi}] = \frac{1}{2}(\rho_i, \rho_i)x_{Pi} = x_{Pi} \quad (6.14)$$

Obviously $\text{tr}(x_{P1}) = \text{tr}(x_{P2}) = \text{tr}(x_{P3}) = 1$.

We now show that $x_\alpha \in T$ is nilpotent if $\alpha \neq k_1 \pm k_N, -k_2 - k_3$.

Suppose $x_\alpha \circ x_\alpha \neq 0$ then either $2\alpha - k_1 - k_N$ or $2\alpha - k_1 + k_N$ or $2\alpha + k_2 + k_3$ is a root. If $\alpha \in \Phi_O$ then $2\alpha - k_1 \mp k_N = k_1 \pm 2k_i \mp k_N$ is a root if and only if $k_i = k_N$, whereas $2\alpha + k_2 + k_3$ is not a root. If $\alpha \in \Phi_S$ then nor $2\alpha - k_1 \mp k_N$ nor $2\alpha + k_2 + k_3$ are roots. Therefore $\alpha \neq k_1 \pm k_N, -k_2 - k_3$ implies $x_\alpha^2 = 0$ and, obviously, $\text{tr}(x_\alpha) = 0$ and x_α be rank-1.

Finally, from the definition of the product $x, y \rightarrow x \circ y$ and the definition of the trace it easily follows that $\text{tr}(x, y)$ is bilinear. The associativity of the trace is proven by direct calculation. \square

7 Future Developments

There are more things in heaven and earth, Horatio, than are dreamt of in your philosophy.

“Hamlet” 1.5.167-8, W. Shakespeare

We have presented a periodic infinite chain of finite generalisations of the exceptional structures, including \mathfrak{e}_8 , the exceptional Jordan algebra \mathbf{J}_3^\oplus (and pair), and the octonions. We have demonstrated that the exceptional Jordan algebra \mathbf{J}_3^\oplus is part of an infinite family of finite-dimensional matrix algebras (corresponding to a particular class of Vinberg’s cubic T-algebras [44]). Correspondingly, we have proved that \mathfrak{e}_8 is part of an infinite family of MS algebras, that resemble lattice vertex algebras.

We are currently working on several topics concerning the mathematical structure presented in this paper. Many of them are at an advanced stage, some are completed. In particular we will expand in future papers on the algebra of derivations of the algebras \mathcal{L}_{MS} and T , the lattice, the modular forms and the theta functions related to EP.

There are also several topics that we are planning to develop in the future, strictly related to Quantum Gravity. In particular, a model for interactions based on EP which includes gravity and the expansion of space-time, starting from a singularity. We aim at a new perspective of elementary particle physics at the early stages of the Universe based on the idea that interactions, defined in a purely algebraic way, are the fundamental objects of the theory, whereas space-time, hence gravity, are derived structures.

An interesting venue of research stemming from EP is to study the higher dimensional weight vectors of algebras akin to lattice vertex algebras (the original motivation for Borcherds’ definition of vertex algebras [62]), that project to a “Magic Star” structure. Such extended “Magic Star” algebras make use of an *asymmetry function* [60, 59], that acts like the cocycle of a lattice vertex algebra which gives a twisted group ring $\mathbb{C}_\epsilon[\Lambda]$ over an even lattice Λ . This gives algebras that span higher dimensional lattices, beyond that of the self-dual $D = 8$ lattice of \mathfrak{e}_8 , and allows one to potentially probe the symmetries of the heterotic string and moonshine [63, 64], as well as to formulate a T -algebra generalization of noncommutative geometry. Specifically, EP gives a novel algebraic method to study even self-dual lattices, such as the $\mathfrak{e}_8 \oplus \mathfrak{e}_8$ and Leech lattices, which already have a well known connection to the Monster vertex algebra and $D = 24$ bosonic string compactifications [65, 66, 67].

With an infinite family of new algebras that extend the exceptional Lie algebras, EP gives a fresh new toolkit for studying emergent spacetime and quantum gravity, in dimensions beyond those previously explored, using spectral techniques applied to an infinite class of cubic T -algebras.

An interaction based approach might also have a Yang-Mills interpretation, as EP hints at higher-dimensional (possibly super) Yang-Mills theories (see *e.g.* [68]), that may exhibit some kind of triality.

As the exceptional Jordan algebra \mathbf{J}_3^\oplus is to \mathfrak{e}_8 , so is the particular infinite class of cubic T-algebras to the higher EP algebras, when projected along an \mathbf{a}_2 . Therefore, physically, the class of cubic T-algebras under consideration allows one to consistently generalize exceptional quantum mechanics of Jordan, Wigner and von Neumann [2], because they go beyond the formally real Jordan algebraic classification of quantum-mechanical self-adjoint operators; in a sense, the classification of quantum-mechanical observables can only probe the T-algebras up to the exceptional Jordan algebra.

It is the authors' hope that a synthesis of spectral algebraic geometry and EP can unify many approaches to unification and quantum gravity, and provide a new lens for searching a non-perturbative theory of all matter, forces and spacetime.

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