

Adiabatic photoevaporation flow

Bernoulli equation: $\frac{1}{2} v^2 + w = \text{constant}$
 "What Landau & Lifshitz use. Whereas Shu uses h "

Enthalpy, $w = \begin{cases} \text{ISOTHERMAL} & c_s^2 \ln p/p_0 \equiv w_\infty \\ \text{ADIABATIC} & \left(\frac{\gamma}{\gamma-1}\right) c_0^2 \left[\left(\frac{p}{p_0}\right)^{\gamma-1} - 1\right] \equiv w_\gamma \end{cases}$

$$h = \int_{p_0}^p \frac{dp}{\rho}$$

So

p/p_0	$\frac{w_\infty}{c_0^2}$	$\frac{v}{c_0} \frac{w_{5/3}}{c_0^2} \frac{T}{T_0}$	$\frac{w_{4/3}}{c_0^2}$	$\frac{w_{11}}{c_0^2}$
1	0	0	0	0
0.5	1.55	1.68	-0.82	-0.73
0.1	2.36	2.22	-2.14	-2.26
0.01	3.19	2.41	-3.14	-4.06
0	$-\infty$	2.45	-4	-11

So in the adiabatic case, the velocity increase is limited, even as $p \rightarrow 0$.

$$\frac{1}{2} v^2 - 2.5 c_0^2 = \frac{1}{2} c_0^2 \Rightarrow v^2 = 6 c_0^2 \Rightarrow v = 2.45 c_0$$

$\frac{T}{T_0} = \left(\frac{p}{p_0}\right)^{\gamma-1}$ So T falls to 500K when p falls by 100 in case of $\gamma=5/3$.

This will also affect calculation of h_0 since α_B goes up as T goes down. Assuming $\alpha_B \propto T^{-1}$ then we will have $\int n^{4/3} dr$ instead of $\int n^2 dr$

So, for instance at constant velocity we would have $\int_1^\infty x^{-8/3} dx$ instead of $\int_1^\infty x^{-4} dx = -\frac{1}{3} [x^{-3}]_1^\infty = \frac{1}{3}$
 $= \frac{3}{5}$, so it nearly doubles the value of h_0

$$h_0 = \frac{1}{2} \frac{v^2}{c_0^2} \frac{1}{1-2\gamma} \Rightarrow h_0 = \frac{1}{2} \frac{v^2}{c_0^2} \frac{1}{1-2\gamma}$$

