

# Dynamics at an exceptional point in an interacting quantum dot system

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# 1 Introduction

The Hermiticity of the Hamiltonian in quantum mechanics is thought of being a necessity for physical behavior of a quantum system. However, in recent decades, the field of non-Hermitian quantum physics has shown major progress, both theoretically and experimentally. Systems of an "open nature", i.e systems being connected to an environment, in nuclear, atomic and optical physics have been accurately described by non-Hermitian Hamiltonians [1]. One of the main properties of non-Hermitian operators is the possibility of exceptional points (EPs). These correspond to points in parameter space where two or more eigenvalues and their corresponding eigenvectors of the operator simultaneously coalesce. Exceptional points have been proposed to have several useful technical applications, and along with the optical microring experiments in 2017, exceptional point sensors successfully increased the sensitivity of current and nano-particle detection [2, 3].

Another operator in quantum mechanics which more recently has been studied for its native non-Hermiticity is the Liouvillian superoperator [4–6]. The Liouvillian and its corresponding quantum master equation describe the dynamics in open quantum systems, where the system is coupled to a larger environment. Due to dissipation to the environment in such systems, the Liouvillian is not necessarily Hermitian, which again brings the possibility of exceptional points. EPs in Liouvillian physics affect the dynamics of the density operator, and therefore the evolution of all relevant observables in the system. Recently, exceptional points in the Liouvillian have been theoretically studied in quantum thermal machines, where critical decay towards the steady state was found at the EP [5]; and in quantum control, where an EP corresponded to optimal steering toward a target quantum state [6].

Another application of interest for quantum master equations and Liouvillian physics is electron transport in systems of quantum dots connected to metallic leads [7]. A quantum dot is a fabricated semiconductor structure containing a small number of electrons and is typically in the order of 100s nanometres in size (?). The size of the dot needs to be small in comparison with the thermal wavelength of the electrons, which is why experiments are generally realized at temperatures close to absolute zero [8]. One method of creating this isolation of electrons is to apply voltages via nanoscale electrodes, called gates, which depletes the number of electrons in a small region. If the voltages are tuned successfully, the quantum dot can be tunnel coupled to the surrounding conducting regions of the semiconductor, known as the source and drain. The electrons can then tunnel from the source into the dot and then exiting it by tunneling into the drain, producing a current through the system [9]. (figure of SEM)

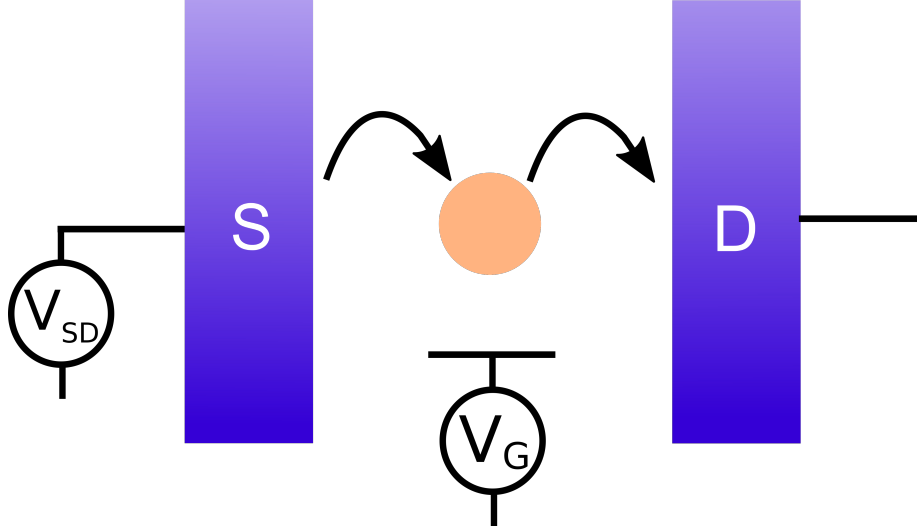


Figure 1: A schematic figure of a quantum dot system, where the quantum dot is in the center, tunnel coupled to the source (S) and drain (D). The gate voltage  $V_G$  and the source-drain voltage  $V_{SD}$ , are tuned to define the quantum dot.

The dynamics of the current through a quantum dot system can be derived from the microscopic properties of the system, resulting in a quantum master equation and a corresponding Liouvillian. The matrix representation of the Liouvillian will depend on the tunnel coupling strength between the dot and the source and drain, the gate and drain voltages, and all other parameters which describe the system. Since the Liouvillian is not necessarily Hermitian, it is possible to find points in this parameter space which correspond to an exceptional point.

The system considered in this thesis consists of two quantum dots coupled to two metallic leads, corresponding to the source and drain. The dots are assumed to be coupled in parallel, i.e not directly coupled to each other but only indirectly through the two leads, see figure 2.

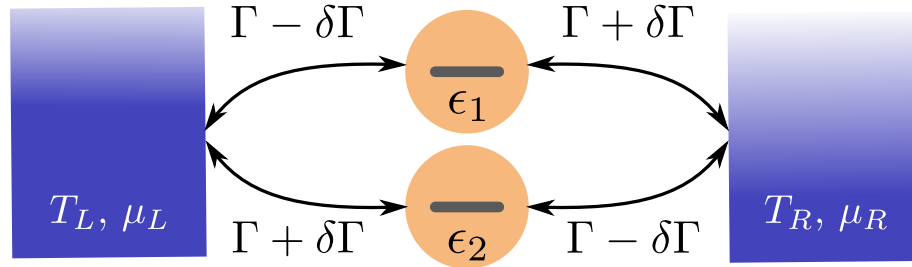


Figure 2: A model of the system containing two quantum dots (yellow) coupled in parallel to metallic leads (blue).

## 2 Theory

The theory part of the thesis is divided into six sections. Firstly, the model of the double dot system will be introduced. Then, a brief overview of the theory behind transport through quantum dots is given. The focus is then put on some of the tools used in open quantum systems, ultimately ending in a section about quantum master equations. The final three sections will explain the notion of exceptional points, the Jordan normal form and lastly present a useful analytic form of the solution of a system of ordinary differential equations.

### 2.1 The system

It is also assumed that each dot is restricted to be either empty or contain one electron, which reduces the dimension of the problem to a (?) degree. Furthermore, the coupling between the dots and the leads are assumed to have two properties: that it can be described by a Markovian interaction, and that it is weak. (reasonable assumptions?) These last assumptions enables the use of the Lindblad master equation to calculate the Liouvillian operator of the system.

### 2.2 Transport through quantum dots

#### 2.2.1 Single electron transistor

To begin understanding the transport of electrons through a system of quantum dots, it is insightful to first study a simple system, known as a single electron transistor. The following explanation is inspired by the article in reference [8].

Consider a single quantum dot with  $N$  electrons capacitively coupled to a gate and coupled to source and drain reservoirs through tunnel junctions, see figure 1. The main transport properties of the system can then be understood in terms of the chemical potentials of the quantum dot and the source and drain reservoirs. These are often depicted in electro-chemical potential diagrams, see figure X. There,  $\mu(N)$  is the energy required to add the  $N$ th electron to the quantum dot, and  $\mu_S$  and  $\mu_D$  are the Fermi levels of the source and drain. The shaded areas represent the Fermi-Dirac distributions of the source and drain.

Using this picture, the electron transport through the quantum dot can be easily visualized. If the chemical potential levels are located as in figure Y,  $\mu(N+1)$  is below  $\mu_S$ , and an electron will likely tunnel from the source onto the dot, increasing the number of electrons on the dot from  $N$  to  $N+1$ . After the tunneling event, there is an even lower chemical potential available for the electron,  $\mu_D$ , and the electron will with a high probability leave the quantum dot and enter the drain. In this fashion, the system will cycle through having  $N$  and  $N+1$  electrons on the dot, producing a current.

It turns out that  $\mu(N)$  depends linearly on the gate voltage  $V_G$ , so by changing it, the "ladder" of states in figure X can be lowered or raised. The drain voltage  $V_D$  on the other hand, changes the spacing between  $\mu_S$  and  $\mu_D$ . Through

these two processes, it is possible to change the electro-chemical potential landscape and therefore control the current through the system.

### 2.2.2 Master equations

An important tool for simulating the current through a quantum dot system is the master equation approach. A classical master equation is defined by

$$\frac{dp_i}{dt} = \text{Rate of entering state } i - \text{Rate of leaving state } i = \sum_j R_{ji}p_j - \sum_j R_{ij}p_i \quad (1)$$

with probabilities  $p_i$  of being in state  $i$ , and transition rates  $R_{ij}$  describing the rate of transitions from state  $i$  to state  $j$ . For the single electron transistor, the states  $i$  correspond to the number of electrons on the dot. The transition rates are typically related to the Fermi-Dirac distributions of the reservoirs and the tunneling rates of the barriers.

Collecting the rate equations for each  $i$  and adding the condition  $\sum_i p_i = 1$ , a system of ordinary differential equations is obtained. This can be solved numerically using standard methods, such as the Runge-Kutta method. Once the probabilities are obtained, the evolution of the current can be calculated by considering the amount of charge tunneling through the barriers at each time point.

While a classical master equation may give accurate results for a quantum system, it leaves out some important physics. A *quantum master equation* generalizes the notion of a classical master equation in the following sense: The unknown in a quantum master equation does not only contain the probabilities  $p_i$  of being in a quantum state  $|\psi_i\rangle$ , but also *coherences*, which are necessary to describe a quantum system fully. This topic will further be delved upon in section 2.3.3.

## 2.3 The theory of open quantum systems

Open quantum systems are systems which are non-isolated and connected to some sort of environment. Often, it considers a total system consisting of the (sub)system of interest, and an environment. The total system is closed, and therefore it obeys the regular quantum mechanical equations of motion. The goal of the theory of open quantum systems is to infer the equation of motion of the smaller system from the equations of the total system.

### 2.3.1 The von Neumann equation and the reduced density operator

An important tool used in open quantum theory is the *density operator*  $\hat{\rho}$ . In open quantum systems, it might only be known that the system is in states  $|\psi_i\rangle$  with probabilities  $p_i$ . Then the system is said to be in an *ensemble*  $\{|\psi_i\rangle, p_i\}$ . The density operator describes this information in a compact way:

$$\hat{\rho} = \sum_i p_i |\psi_i\rangle \langle \psi_i|. \quad (2)$$

Fixing a basis  $\{|\phi_i\rangle\}$ , the density operator can be represented by a matrix with elements  $\rho_{ij} = \langle \phi_i | \hat{\rho} | \phi_j \rangle$ . The diagonal elements  $\rho_{ii}$  represents classical probabilities of being in a state  $|\phi_i\rangle$  and the off-diagonal elements  $\rho_{ij}$  are the coherences of the system.

The average measured value of an observable  $\hat{O}$  of an ensemble represented by  $\hat{\rho}$  can be shown to be given by equation (3). From the density operator one can therefore extract all significant information from the ensemble, and it can be described as the "state" of the system.

$$\langle \hat{O} \rangle = \text{Tr}(\hat{\rho} \hat{O}). \quad (3)$$

It is also possible to calculate how the density operator evolves over time. Under a Hamiltonian  $\hat{H}$ , the evolution is given by the von Neumann equation

$$i\hbar \frac{d\hat{\rho}}{dt} = [\hat{H}, \hat{\rho}] \equiv \mathcal{L}\hat{\rho}, \quad (4)$$

where  $\mathcal{L}$  is the *Liouvillian superoperator*, or just the Liouvillian.

Connecting back to the mentioned goal of open quantum theory - How does one calculate the evolution of the density operator describing a subsystem connected to an environment? This question is solved by the *reduced density operator*. If the total system  $T$  consists of a subsystem  $S$  and an environment  $E$ , then the reduced density matrix for the subsystem is given by

$$\hat{\rho}_S = \text{Tr}_E(\hat{\rho}_T), \quad (5)$$

where  $\text{Tr}_E$  is the partial trace over the environment. It is therefore possible to "trace out" the environment and obtain a density operator for the subsystem of interest. The reduced density operator contain all measurement statistics of interest in the subsystem, placing it at the core of the theory of open quantum systems.

### 2.3.2 The Fock-Liouville space

The Liouvillian  $\mathcal{L}$  in equation (4) is an operator acting on an operator, which is why it is called a superoperator. To be able to more easily handle the Liouvillian, it is possible to treat it as a normal operator represented by a matrix, acting on a Hilbert space spanned by the density matrices. The density matrix is therefore effectively transformed into a vector, usually written as  $|\rho\rangle\rangle$ , with a scalar product  $\langle\langle \rho_1 | \rho_2 \rangle\rangle$ . The Hilbert space spanned by these vectors is called the Fock-Liouville space.

### 2.3.3 The Lindblad Master equation

The starting point for quantum master equations is the von Neumann equation (eq. (4)) for the density matrix of the total system. Using the partial trace, a differential equation for the reduced density matrix  $\hat{\rho}$  of the system of interest can be obtained. However, the resulting equation generally still includes the total density operator and is impossible to solve. The approximations required to find a solution can be done in many different ways and is essentially what differs the quantum master equations apart.

One of the main examples of a quantum master equation is the *Lindblad equation*. The scheme used to derive the Lindblad equation is based on two major assumptions. The first is to assume that the strength of the interaction between the system and the environment is weak. It is therefore possible to treat the interaction perturbatively, which in the Lindblad approach is used to drop terms of  $\mathcal{O}(\alpha^3)$  or higher. (what is alpha?) The second main approximation is to assume that the interaction between the system and the environment is Markovian. A Markovian process is one in which the next time step only depends on the current state of the system. In a quantum dot system, this can be interpreted to mean that the leads "forget" that they have lost or gained an electron, right after the tunneling process has happened. Since the number of electrons in the leads typically is large, this can be a reasonable approximation for such systems.

Applying these approximations to the von Neumann equation, the Lindblad equation can be derived:

$$\frac{d\hat{\rho}}{dt} = i[\hat{\rho}, \hat{H}_{\text{eff}}] + \sum_i \hat{L}_i \hat{\rho} \hat{L}_i^\dagger - \frac{1}{2} \hat{\rho} \hat{L}_i^\dagger \hat{L}_i - \frac{1}{2} \hat{L}_i^\dagger \hat{L}_i \hat{\rho} \equiv \mathcal{L}\hat{\rho}, \quad (6)$$

where  $\hat{H}_{\text{eff}}$  consists of two terms: the subsystem Hamiltonian  $\hat{H}_S$  and the Lamb shift Hamiltonian  $\hat{H}_{LS}$ , a renormalization of energy levels due to the interaction with the environment. The  $\hat{L}_i$ 's are the so called *jump operators*, which capture the coupling processes to the environment. ( $\hbar = 1$ ) The first term of the Lindblad equation is Hermitian and the second is non-Hermitian. The matrix representation of the Liouvillian is therefore not necessarily Hermitian, which will enable the possibility of exceptional points. More about this in section 2.4.

One of the attractive properties of the Lindblad equation is that it ensures the positivity and unity trace of the density operator throughout the evolution (explain in prev section). Other quantum master equations do not have this property, which may cause negative probabilities and other unphysical results.

One method for calculating the jump operators is described in ref X. There, a phenomenological approach is proposed, where the jump operators "carries information of both the spatial and energetic properties of the jump process". Using this PERLind approach, the jump operators are calculated using well defined steps, and can be applied in a wide range of systems. Once the jump operators are constructed and the first term of equation (6) is calculated, an

expression for the Liouvillian  $\mathcal{L}$  can be obtained. Finally, equation (6) can be solved, resulting in the evolution of the density matrix to be known. Using equation (3), the dynamics of any relevant observable of the system can be simulated.

## 2.4 Exceptional points

In quantum mechanics, a common way to perform calculations is to diagonalize the matrix representation of the Hamiltonian. The diagonalization process can be understood as a change of basis to linearly independent eigenvectors. In this basis, the linear transformation of the matrix is very simple: it scales each eigenvector by the corresponding eigenvalue. The matrix in the new basis is therefore diagonal, explaining the name of the process. For a matrix  $A$  and its diagonal form  $D$ , this can be written as

$$A = SDS^{-1}, \quad (7)$$

where  $S = (v_1, \dots, v_n)$  consists of the eigenvectors. The diagonalization process can always be done for Hamiltonians in quantum mechanics. The reason is that the Hamiltonian is an observable, which by the postulates in quantum mechanics always are Hermitian. This in turn means that the matrix representation is normal, and therefore diagonalizable.

In Liouvillian physics, there are non-Hermitian parts in the Liouvillian superoperator. The matrix representation of the Liouvillian is hence not necessarily diagonalizable, but can be defective. If a matrix is defective, two or more eigenvalues and their corresponding eigenvectors coalesce.

The Liouvillian matrix depends on a set of parameters describing the system. If a point in this parameter space produces a defective matrix, the point is called an *exceptional point* (EP). Exceptional points can be of different *orders*, reflecting how many eigenvectors coalesce at the point. For example, if three eigenvectors coalesce, the EP is said to be of order three. The notion of exceptional points is tightly connected to a type of matrix decomposition, the Jordan normal form, which will be explained in the following section.

## 2.5 Jordan normal form

For a defective matrix, there does not exist a basis of eigenvectors, and the diagonalization process is not possible. Fortunately, there is a notion of an "almost diagonal" form, called the Jordan normal form. Recall that in the diagonalizable case, the basis was changed to the linearly independent eigenvectors. To construct the Jordan form for a defective matrix, this basis has to be completed in some way to span the full space. This can be done using *Jordan chains*, which for each defective eigenvector  $r_i$  with eigenvalue  $\lambda_i$ , consists of vectors  $r'_i, r''_i \dots$  defined by equation 8. The length of chain,  $n_i$ , is the same as the order of the corresponding EP.



$$\begin{aligned}
(A - \lambda_i I)r_i &= 0 \\
(A - \lambda_i I)r'_i &= r_i \\
(A - \lambda_i I)r''_i &= r'_i \\
&\vdots
\end{aligned} \tag{8}$$

Using this new basis and creating the change-of-basis matrix  $M$  by forming

$$M = (\mathbf{r}_1 \dots \mathbf{r}_q), \quad \text{where } \mathbf{r}_i = (r_i, r'_i, r''_i \dots), \tag{9}$$

the Jordan normal form  $J$  of the matrix  $A$  is obtained:

$$A = MJM^{-1},$$

$$\text{where } J = \begin{bmatrix} J_{n_1}(\lambda_1) & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & J_{n_q}(\lambda_q) \end{bmatrix} \quad \text{and } J_{n_i}(\lambda_i) = \begin{bmatrix} \lambda_i & 1 & \dots & 0 \\ \vdots & \ddots & \ddots & \vdots \\ \vdots & & \ddots & 1 \\ 0 & \dots & \dots & \lambda_i \end{bmatrix}. \tag{10}$$

The Jordan form hence consist of  $q$  blocks on the diagonal, each block of size  $n_i$  consisting of its eigenvalue on the diagonal and ones on the super diagonal. Note that if all blocks are of size one, i.e there are no exceptional points, the Jordan form is diagonal.

## 2.6 Solution of ODEs

An ordinary differential equation (ODE) is a linear differential equation of the form  $x' = Ax$ . Often, the unknown  $x$  is a vector, and  $A$  a matrix. The solution can be written as a matrix exponential in the following way:

$$x(t) = e^{At}x(0), \quad \text{where } e^{At} = \sum_{k=0}^{\infty} \frac{(At)^k}{k!}. \tag{11}$$

The matrix exponential can be simplified using Jordan decomposition. It can be shown that  $e^{At} = Me^{Jt}M^{-1}$  where

$$e^{Jt} = \begin{bmatrix} e^{J_{n_1}(\lambda_1)t} & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & e^{J_{n_q}(\lambda_q)t} \end{bmatrix}, \quad \text{and}$$

$$e^{J_{n_i}(\lambda_i)t} = e^{\lambda_i t} \begin{bmatrix} 1 & t & \dots & t^{n_i-1}/(n_i-1)! \\ \vdots & \ddots & \ddots & \vdots \\ \vdots & & \ddots & t \\ 0 & \dots & \dots & 1 \end{bmatrix}. \tag{12}$$

In general, the matrix exponential therefore consists of entries with terms of the form  $t^k e^{\lambda_i t}$ . Note that if  $A$  is diagonalizable, all blocks are of size one and the entries consist of pure exponentials on the diagonal.

Using this result, the solution can be written as

$$x(t) = M e^{Jt} M^{-1} x(0). \quad (13)$$

This can further be decomposed if one considers the generalized modes of the system. Suppose first that the initial state is in a linear combination of vectors in one of the Jordan chains,  $x(0) = a_1 r_i + a_2 r'_i + a_3 r''_i + \dots = \mathbf{r}_i a$ , where  $a = (a_1, \dots, a_{n_i})^T$  is a constant vector and  $\mathbf{r}_i$  is defined in equation (9). By rewriting  $M^{-1}$  as (EXPLAIN IN PREV SECTION?)

$$M^{-1} = \begin{bmatrix} \mathbf{l}_1 \\ \vdots \\ \mathbf{l}_{n_i} \end{bmatrix} \quad (14)$$

and using equation (13), the solution can be written as (MAYBE WRITE MORE STEPS)

$$x(t) = M e^{Jt} M^{-1} x(0) = M e^{Jt} \begin{bmatrix} \mathbf{l}_1 \\ \vdots \\ \mathbf{l}_{n_i} \end{bmatrix} \mathbf{r}_i a = \mathbf{r}_i e^{J_{n_i}(\lambda_i)t} a, \quad (15)$$

since  $\mathbf{l}_j \mathbf{r}_i = \delta_{ij} I$ , where  $I$  is the  $n_j \times n_i$  identity matrix. Hence, the solution stays in the space spanned by the initial condition throughout the evolution. For an arbitrary initial condition, the solution can be written as a sum over the generalized modes:

$$x(t) = \sum_{i=1}^q \mathbf{r}_i e^{J_{n_i}(\lambda_i)t} \mathbf{l}_i x(0). \quad (16)$$

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