

PHYS 251: Honours Classical Mechanics

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Abstract

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1 Introduction

2 Prerequisite knowledge

3 Mathematical Tools for Mechanics

3.1 Levi-Civita Symbol and Cross Product

Definition 1 (Levi-Civita Symbol). *The Levi-Civita tensor is a tensor that is used to calculate the cross product of two vectors.*

$$(A \times B)_i = \epsilon_{ijk} A_j B_k,$$

where ϵ_{ijk} is the Levi-Civita symbol, and is defined to be

$$\epsilon_{ijk} = \begin{cases} 1, & \text{for even permutations of } i, j, k \\ -1, & \text{for odd permutations of } i, j, k \\ 0, & \text{otherwise} \end{cases}.$$

Problem 1 (Levi-Civita Symbol). *Consider the cross product $\vec{A} \times \vec{B} = \vec{C}$. Show that the components of \vec{C} obtained with the determinant-based definition of the cross product are the same as those obtained with $C_i = \epsilon_{ijk} A_j B_k$, where ϵ_{ijk} is the Levi-Civita tensor and the Einstein convention is implied.*

3.2 Motion in Polar Coordinates

In planar motion it is often convenient to describe position using polar coordinates (r, θ) instead of (x, y) . The relations between Cartesian and polar coordinates are

$$r^2 = x^2 + y^2, \quad \tan \theta = \frac{y}{x}.$$

The associated unit vectors are

$$\hat{r} = \cos \theta \hat{i} + \sin \theta \hat{j}, \quad \hat{\theta} = -\sin \theta \hat{i} + \cos \theta \hat{j}.$$

Unlike \hat{i} and \hat{j} , these unit vectors depend on time through $\theta(t)$. The position vector is

$$\vec{r} = r \hat{r}.$$

Differentiating with respect to time gives the velocity:

$$\vec{v} = \frac{d\vec{r}}{dt} = \dot{r} \hat{r} + r \dot{\theta} \hat{\theta}.$$

Differentiating again yields the acceleration:

$$\vec{a} = (\ddot{r} - r \dot{\theta}^2) \hat{r} + (2\dot{r} \dot{\theta} + r \ddot{\theta}) \hat{\theta}.$$

The radial component $(\ddot{r} - r \dot{\theta}^2)$ contains the centripetal term $-r \dot{\theta}^2$, while the angular component $(2\dot{r} \dot{\theta} + r \ddot{\theta})$ contains the Coriolis term $2\dot{r} \dot{\theta}$ and the tangential term $r \ddot{\theta}$.

3.3 Uniform Circular Motion as a Special Case

Uniform circular motion corresponds to motion with constant radius $r = r_0$ and constant angular velocity $\dot{\theta} = \omega$. In Cartesian coordinates the position is given by

$$x = r_0 \cos(\omega t), \quad y = r_0 \sin(\omega t),$$

so that

$$\vec{r} = x\hat{i} + y\hat{j}.$$

Differentiating with respect to time gives the velocity,

$$\vec{v} = -r_0\omega \sin(\omega t)\hat{i} + r_0\omega \cos(\omega t)\hat{j},$$

which is always tangent to the circle. Differentiating once more yields the acceleration,

$$\vec{a} = -r_0\omega^2 \cos(\omega t)\hat{i} - r_0\omega^2 \sin(\omega t)\hat{j} = -r_0\omega^2 \hat{r}.$$

Thus the acceleration has constant magnitude $r_0\omega^2$ and is directed radially inward. This inward acceleration is the centripetal acceleration required to maintain circular motion.

4 Newtonian Mechanics

4.1 Newton's Laws of Motion

Definition 2 (Law of Inertia). *An object remains at rest or continues to move at a constant velocity unless acted upon by a net external force.*

Definition 3 (Law of Acceleration). *The net force acting on an object is directly proportional to its acceleration and directly proportional to its mass*

$$F_{\text{net}} = ma.$$

Definition 4 (Law of Action-Reaction). *For every action, there is an equal and opposite reaction; if object A exerts a force on object B, object B simultaneously exerts a force of equal magnitude and opposite direction on object A.*

Remark 1 (Momentum Conservation). *Newton's third law implies conservation of momentum for an isolated system. When summing the forces over all objects in the system, internal forces occur in equal and opposite pairs and therefore cancel. It is the net internal force on the system that vanishes, not the force acting on each individual object.*

Remark 2 (Free-Body Diagram and Action-Reaction Pairs). *Consider a block A resting on block B, with block B on a surface. For block A, the forces are its weight w_A downward and the normal force N_{BA} exerted by block B upward. For block B, the forces are its weight w_B downward, the normal force N_{AB} exerted by block A downward, and the normal force N from the ground upward. The forces N_{BA} and N_{AB} are an action-reaction pair. They are equal in magnitude and opposite in direction, but they act on different bodies. Action-reaction forces never appear in the same free-body diagram.*

4.2 Atwood's Machines

An Atwood's machine consists of two or more masses connected by a light, inextensible string passing over an ideal pulley. Because the string has fixed length, the motion of the two masses is constrained.

If one mass moves downward by a displacement x , the other must move upward by the same amount. Writing the total string length as a constant and differentiating with respect to time gives a velocity constraint, and differentiating again gives an acceleration constraint. This shows that the magnitudes of the accelerations of the two masses are equal.

The essential steps in these problems are:

- use the string-length constraint to relate accelerations,
- draw separate free-body diagrams for each mass,
- apply $F_{\text{net}} = ma$ to each object.

4.3 Gravity

5 Gravitational Potential Inside a Uniform Solid Sphere

Consider a solid sphere of radius R , total mass M , and uniform density ρ . We compute the gravitational potential at a point located a distance r from the center, where $r < R$.

Step 1: Split the mass into two regions. We divide the sphere into:

- An inner solid sphere of radius r ,
- An outer spherical shell extending from r to R .

By symmetry (shell theorem):

- The inner sphere acts like a point mass located at the center.
- The outer shell produces zero *force*, but still contributes to the potential.

Step 2: Contribution from the inner sphere. The density is

$$\rho = \frac{M}{\frac{4}{3}\pi R^3}.$$

The mass enclosed within radius r is

$$M_{\text{in}} = \frac{4}{3}\pi r^3 \rho = M \frac{r^3}{R^3}.$$

Since we are outside this inner sphere, its contribution to the potential is

$$U_1 = -\frac{GmM_{\text{in}}}{r} = -GmM \frac{r^2}{R^3}.$$

Step 3: Contribution from the outer shell. A thin spherical shell of radius r' and thickness dr' has mass

$$dM = 4\pi r'^2 \rho dr'.$$

The potential contribution from this shell is

$$dU = -\frac{Gm dM}{r'}.$$

Integrating from r to R gives

$$U_2 = -Gm \int_r^R \frac{4\pi r'^2 \rho}{r'} dr' = -4\pi Gm \rho \int_r^R r' dr'.$$

Evaluating the integral,

$$U_2 = -2\pi Gm \rho (R^2 - r^2).$$

Substituting $\rho = \frac{3M}{4\pi R^3}$,

$$U_2 = -\frac{3GmM}{2R^3} (R^2 - r^2).$$

Step 4: Total potential inside the sphere. The total potential is

$$U(r < R) = U_1 + U_2 = -\frac{GmM}{2R^3} (3R^2 - r^2).$$

Remarks.

- The potential is finite at $r = 0$.
- At $r = R$, this expression matches the exterior result $U = -\frac{GmM}{R}$.
- Inside the sphere, the potential varies quadratically with r .

6 Appendix

7 Solutions

Solution 1 (Levi-Cevita Tensor). Let $\vec{A} = (a_1, a_2, a_3)^T$, $\vec{B} = (b_1, b_2, b_3)^T$ and $\vec{C} = (c_1, c_2, c_3)^T$. Using the determinant-based definition of the cross product, we have

$$\vec{A} \times \vec{B} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix} = \begin{pmatrix} a_2 b_3 - a_3 b_2 \\ a_3 b_1 - a_1 b_3 \\ a_1 b_2 - a_2 b_1 \end{pmatrix} = \begin{pmatrix} c_1 \\ c_2 \\ c_3 \end{pmatrix}.$$

Using the Einstein convention, we have

$$c_1 = \epsilon_{123} a_2 b_3 + \epsilon_{132} a_3 b_2 = a_2 b_3 - a_3 b_2,$$

$$c_2 = \epsilon_{231}a_3b_1 + \epsilon_{213}a_1b_3 = a_3b_1 - a_1b_3,$$

$$c_3 = \epsilon_{312}a_1b_2 + \epsilon_{321}a_2b_1 = a_1b_2 - a_2b_1.$$

Therefore, the components of \vec{C} obtained with the determinant-based definition of the cross product are the same as those obtained with $C_i = \epsilon_{ijk}A_jB_k$.

8 Useful Links