

# MEK1100 - Mandatory assignment 1

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## 1 Scaling

A ball thrown with initial velocity  $v_0$  and at an angle  $\theta$  to the horizontal at time  $t = 0$  will follow the curve defined by

$$x(t) = v_0 t \cos \theta \tag{1}$$

$$y(t) = v_0 t \sin \theta - \frac{1}{2} g t^2. \tag{2}$$

**a.**

The time taken for the ball to reach  $y = 0$  is

$$\begin{aligned} y(t) &= 0 \\ v_0 t \sin \theta &= \frac{1}{2} g t^2 \\ t &= \frac{2 v_0 \sin \theta}{g} = t_m \end{aligned}$$

The x-component at this time is

$$\begin{aligned} x(t_m) &= \frac{2 v_0^2 \sin \theta \cos \theta}{g} \\ &= \frac{v_0^2 \sin(2\theta)}{g} = x_m \end{aligned}$$

**b.**

We introduce the following dimensionless variables:

$$\begin{aligned} x^* &= \frac{x}{x_m} \implies x = x^* x_m \\ y^* &= \frac{y}{x_m} \implies y = y^* x_m \end{aligned}$$

We insert these into (1) and get

$$\begin{aligned}
v_0 t \cos \theta &= x^* \frac{v_0^2 \sin(2\theta)}{g} \\
x^* &= \frac{gt}{2v_0 \sin \theta} \\
v_0 t \sin \theta - \frac{1}{2}gt^2 &= y^* \frac{v_0^2 \sin(2\theta)}{g} \\
y^* &= \frac{g(v_0 t \sin \theta - \frac{1}{2}gt^2)}{v_0^2 \sin \theta \cos \theta} \\
y^* &= \frac{gt}{2v_0 \cos \theta} - \frac{g^2 t^2}{4v_0^2 \sin \theta \cos \theta} \\
y^* &= \frac{gt}{2v_0 \sin \theta} \cdot \frac{\sin \theta}{\cos \theta} \cdot \left(1 - \frac{gt}{2v_0 \sin \theta}\right) \\
y^* &= x^* \tan \theta (1 - x^*)
\end{aligned}$$

Scaling time is done by  $t^* = t/t_m = x^*$ .  $\theta$  is dimensionless, hence we do not need to scale it.

**C.**

```

1 import numpy as np
2 import matplotlib.pyplot as plt
3
4 v0 = 5
5 g = 9.81
6
7 for theta in (np.pi/6, np.pi/4, np.pi/3):
8     t = np.linspace(0, 2*v0*np.sin(theta)/g, 100)
9     x = g*t/(2*v0*np.sin(theta))
10    y = x*np.tan(theta)*(1-x)
11    plt.plot(x, y)
12
13 plt.legend(['pi/6', 'pi/4', 'pi/3'])
14 plt.xlabel('x*')
15 plt.ylabel('y*')
16 plt.show()

```

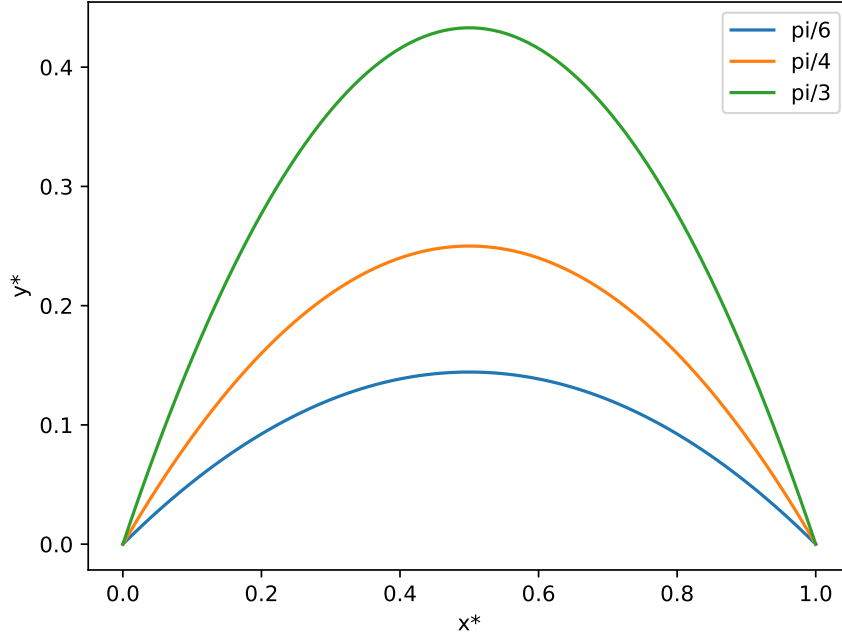


Figure 1: Plot of  $x^*$ ,  $y^*$  for  $\theta = \pi/6, \pi/4, \pi/3$ .

Since we have used dimensionless variables, the trajectory of the ball will only depend on the angle  $\theta$ .

## 2 Stream lines for a two-dimensional field

We are given the velocity field

$$\mathbf{v} = v_x \mathbf{i} + v_y \mathbf{j} = xy \mathbf{i} + y \mathbf{j}. \quad (3)$$

**a.**

To find the stream lines we integrate both sides as described in chapter 2.4 in Gjevik & Fagerland (2021).

$$\begin{aligned} \int xy dy &= \int y dx \\ \int dy &= \int \frac{1}{x} dx \\ y &= \ln |x| + C \end{aligned}$$

It is clear that if  $y = 0$  the whole  $x$ -axis is a solution to the initial differential equation.

b.

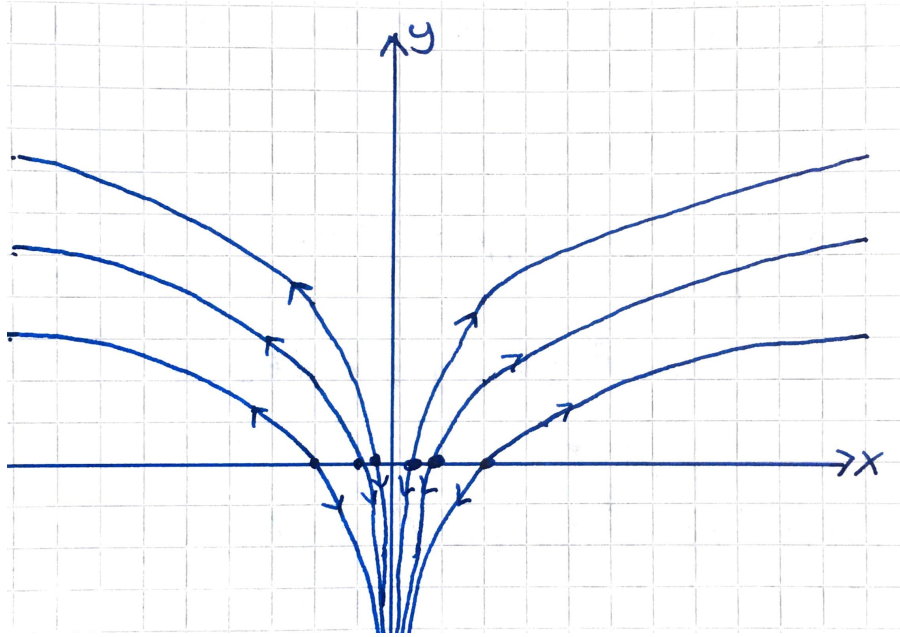


Figure 2: Hand-drawn stream lines for velocity field in (3)

The contour lines are defined when  $y = \ln|x| + C$  is constant, meaning  $z = \ln|x| - y$ . All stagnation points will lie along the  $x$ -axis.

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3
4 t = np.linspace(-2, 2, 1001)
5 x, y = np.meshgrid(t, t)
6 z = np.log(np.abs(x)) - y
7 plt.contour(x, y, z)
8 plt.show()
```

See also Figure 3 on next page.

c.

According to 4.6 in Gjevik & Fagerland (2021), there is a stream function  $\psi = \psi(x, y)$  if

$$\nabla \cdot \mathbf{v} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} = 0 \quad (4)$$

We can calculate the divergence of  $\mathbf{v}$  in (3).

$$\frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} = y + 1 \neq 0$$

Hence, there is no  $\psi$  for the given velocity field.

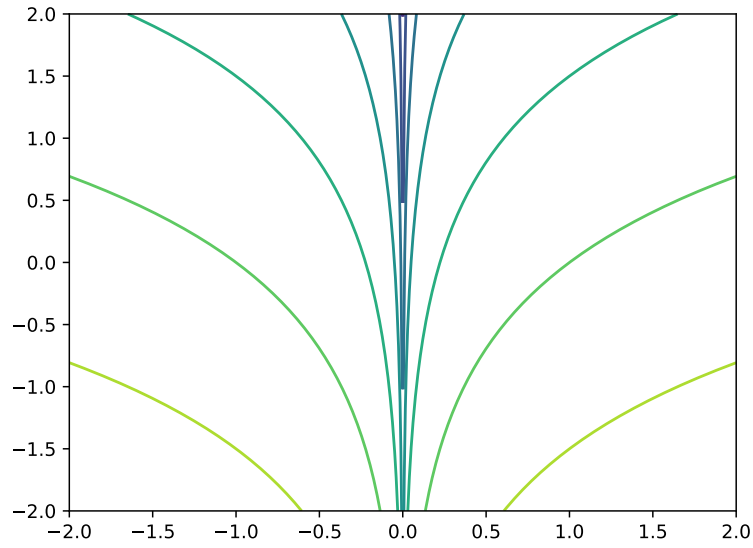


Figure 3: Stream lines drawn in python for velocity field in (3)

### 3 Another two dimensional stream field

A velocity field in the  $xy$ -plane is given by  $\mathbf{v} = v_x \mathbf{i} + v_y \mathbf{j}$  where

$$v_x = \cos(x) \sin(y), \quad v_y = -\sin(x) \cos(y). \quad (5)$$

**a.**

Divergence:

$$\nabla \cdot \mathbf{v} = -\sin(x) \sin(y) + \sin(x) \sin(y) = 0$$

Curl:

$$\nabla \times \mathbf{v} = \left( \frac{\partial v_x}{\partial x} - \frac{\partial v_y}{\partial y} \right) \mathbf{k} = (-2 \cos(x) \cos(y)) \mathbf{k}$$

**b.**

The given task was quite vague, so I simply plotted the field in python.

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3
4 t = np.linspace(-1, 1, 6)
5 x, y = np.meshgrid(t, t)
6 vx, vy = np.cos(x)*np.sin(y), -np.sin(x)*np.cos(y)
7
8 plt.quiver(x, y, vx, vy)
9 plt.axis('equal')
10 plt.show()
```

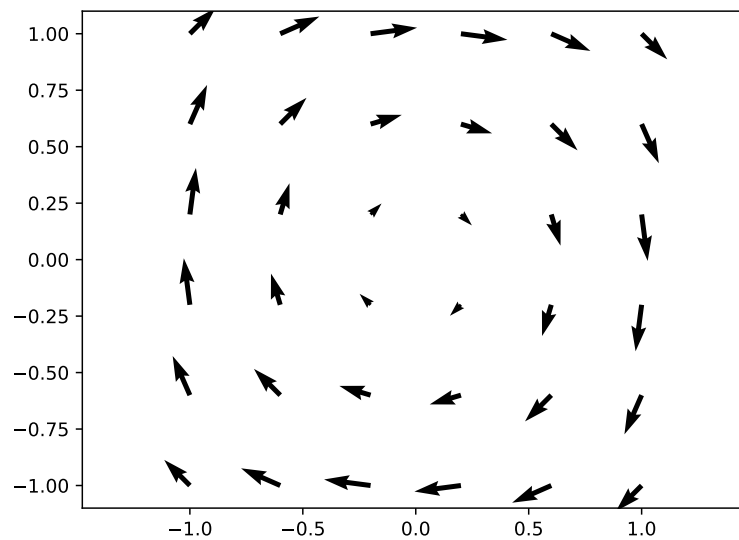


Figure 4: Plot of  $\mathbf{v}$  for  $x, y \in [-1, 1]$

**c.**

In this section, we will use the formula for a line integral of a vector field over a curve

$$\int_C \mathbf{F} \cdot d\mathbf{r} \quad (6)$$

We separate the curve into four separate pieces. Consider the paramerisation  $\mathbf{r}(t) = (t, -\pi/2)$

for  $t \in [-\pi/2, \pi/2]$ . Differentiating this we get  $d\mathbf{r} = (1, 0)$ . Hence, our integral is

$$\begin{aligned} & \int_{-\pi/2}^{\pi/2} (v_x, v_y) \cdot (1, 0) dt \\ &= \int_{-\pi/2}^{\pi/2} \cos(t) \sin(-\pi/2) dt \\ &= [-\sin(t)]_{-\pi/2}^{\pi/2} \\ &= -2 \end{aligned}$$

We perform similar calculations for the three other sides. This results in the circulation around the square defined as  $-\pi/2 \leq x, y \leq \pi/2$  to be  $-2 * 4 = -8$ .

**d.**

To check if

$$\psi = \cos(x) \cos(y) \tag{7}$$

is a valid stream function for  $\mathbf{v}$ , we turn to 4.6 in Gjevik & Fagerland (2021) once again. We have that

$$v_x = -\frac{\partial \psi}{\partial y}, \quad v_y = \frac{\partial \psi}{\partial x} \tag{8}$$

if  $\psi(x, y)$  is a stream function for  $\mathbf{v} = (v_x, v_y)$ . We test for our given  $\psi$  and  $\mathbf{v}$ :

$$\begin{aligned} \frac{\partial \psi}{\partial y} &= -\cos(x) \sin(y) = -v_x \\ \frac{\partial \psi}{\partial x} &= \sin(x) \cos(y) = v_y \end{aligned}$$

which shows that (7) is a stream function for (5).

**e.**

To find the Taylor approximation of  $\psi(0, 0)$  we use the formula stated in chapter 2.2 in Gjevik & Fagerland (2021).

$$\begin{aligned} f(x, y) &\cong f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0)(x - x_0) + \frac{\partial f}{\partial y}(x_0, y_0)(y - y_0) \\ &+ \frac{1}{2} \frac{\partial^2 f}{\partial x^2}(x_0, y_0)(x - x_0)^2 + \frac{1}{2} \frac{\partial^2 f}{\partial y^2}(x_0, y_0)(y - y_0)^2 \\ &+ \frac{\partial^2 f}{\partial x \partial y}(x - x_0)(y - y_0) \end{aligned}$$

To save on some typing, I will only include the parts that do not become zero. We use  $(x_0, y_0) = (0, 0)$ .

$$\begin{aligned}\psi(0, 0) &= 1 \\ \frac{\partial^2 \psi}{\partial x^2} &= -\cos(x) \cos(y) \implies \frac{\partial^2 \psi}{\partial x^2}(0, 0) = -1 \\ \frac{\partial^2 \psi}{\partial y^2} &= -\cos(x) \cos(y) \implies \frac{\partial^2 \psi}{\partial y^2}(0, 0) = -1\end{aligned}$$

If we insert these values into our formula we get

$$\psi(x, y) \cong 1 - \frac{x^2}{2} - \frac{y^2}{2} \quad (9)$$

## 4 Stream lines and velocity field in Python

a.

strlin.py:

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3 from streamfun import streamfun
4
5 for n in (5, 30):
6     x, y, psi = streamfun(n)
7     plt.contour(x, y, psi)
8     plt.axis('equal')
9     plt.show()
```

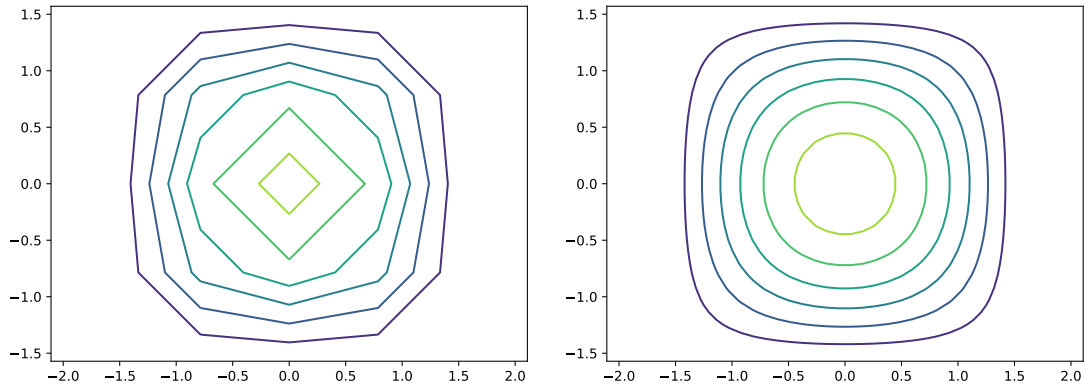


Figure 5: Showing streamlines for  $n = 5, 30$ .

We observe that if we use a higher  $n$ , such as 30, the shape is more close to the circles described by (9).



b.

velfield.py:

```
1 import numpy as np
2
3 def velfield(n):
4     t = np.linspace(-np.pi/2, np.pi/2, n)
5     x, y = np.meshgrid(t, t)
6     u, v = np.cos(x)*np.sin(y), np.sin(x)*np.cos(y)
7
8     return x, y, u, -v
```

vec.py:

```
1 import numpy as np
2 import matplotlib.pyplot as plt
3 from velfield import velfield
4
5 x, y, u, v = velfield(21)
6 plt.quiver(x, y, u, v)
7 plt.axis('equal')
8 plt.show()
```

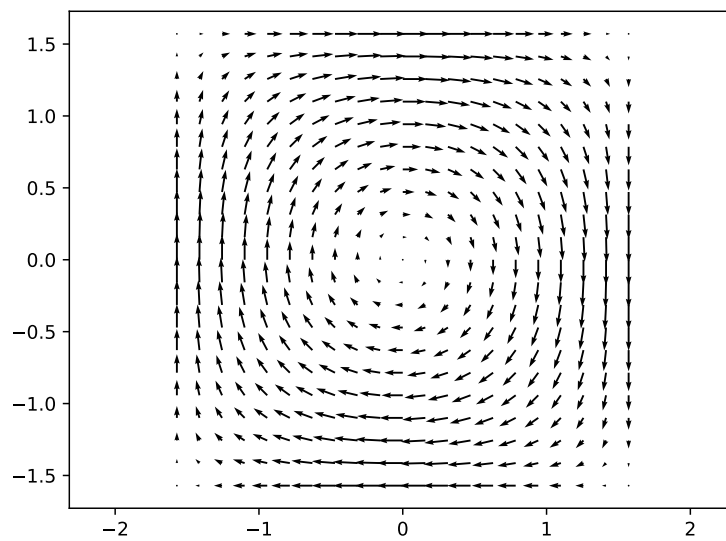


Figure 6: Plot of velocity field with  $n = 21$ .