

$\Lambda_b \rightarrow p\mu\nu$ Form Factors



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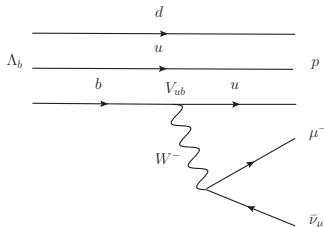
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The tree-level matrix element for this decay may be written as:

$$\mathcal{M} = -i \frac{G_F}{\sqrt{2}} V_{ub} H_\nu \bar{u}_\mu \gamma^\nu (1 - \gamma_5) \nu_{\nu_\mu}$$

where $H_\nu = \langle N^+(p', s') | \bar{u} \gamma_\nu (1 - \gamma_5) b | \Lambda_b^0(p, s) \rangle$



- ▶ In general for baryon to baryon V-A transitions: $H_\nu = \bar{u}_N(p') [F_1^V \gamma_\nu + F_2^V v_\nu + F_3^V v'_\nu - (F_1^A \gamma_\nu + F_2^A v_\nu + F_3^A v'_\nu) \gamma_5] u_{\Lambda_b}(p)$
- ▶ This convention is used by EvtGen. Current papers presenting predictions for $\Lambda_b \rightarrow p$ form factors use different conventions.

For instance A. Khodjamirian et al. in arXiv:1108.2971 (LCSR predictions) use:

dimensionless. The hadronic matrix elements with the vector and axial-vector transition currents contain three form factors each:

$$\langle \Lambda_c(P-q) | \bar{c} \gamma_\mu u | N(P) \rangle = \bar{u}_{\Lambda_c}(P-q) \left\{ f_1(q^2) \gamma_\mu + i \frac{f_2(q^2)}{m_{\Lambda_c}} \sigma_{\mu\nu} q^\nu + \frac{f_3(q^2)}{m_{\Lambda_c}} q_\mu \right\} u_N(P), \quad (11)$$

$$\langle \Lambda_c(P-q) | \bar{c} \gamma_\mu \gamma_5 u | N(P) \rangle = \bar{u}_{\Lambda_c}(P-q) \left\{ g_1(q^2) \gamma_\mu + i \frac{g_2(q^2)}{m_{\Lambda_c}} \sigma_{\mu\nu} q^\nu + \frac{g_3(q^2)}{m_{\Lambda_c}} q_\mu \right\} \gamma_5 u_N(P). \quad (12)$$

Meanwhile W. Detmold, S. Meinel et al. in arXiv:1306.0446 (LQCD predictions) use the form:

The $\Lambda_b \rightarrow p$ matrix elements of the vector and axial vector $b \rightarrow u$ currents are parametrized in terms of six independent form factors (see, e.g., Ref. [8]). In leading-order heavy-quark effective theory (HQET), which becomes exact in the limit $m_b \rightarrow \infty$ and is a good approximation at the physical value of m_b , only two independent form factors remain, and the matrix element with arbitrary Dirac matrix Γ in the current can be written as [8-10]

$$\langle N^+(p', s') | \bar{u} \Gamma Q | \Lambda_Q(v, s) \rangle = \bar{u}_N(p', s') [F_1 + \not{v} F_2] \Gamma u_{\Lambda_Q}(v, s). \quad (3)$$

- Use HQET under which only 2 form factors are required.

Also include next to first order loop corrections:

where A_μ is the hadronic matrix element

$$A_\mu = \langle N^+(p', s') | \bar{u} \gamma_\mu (1 - \gamma_5) b | \Lambda_b(p, s) \rangle. \quad (31)$$

Because we have computed the form factors in HQET, we need to match the QCD current $\bar{u} \gamma_\mu (1 - \gamma_5) b$ in Eq. (31) to the effective theory. This gives (at leading order in $1/m_b$)

$$A_\mu = \sqrt{m_{\Lambda_b}} \langle N^+(p', s') | (c_\gamma \bar{u} \gamma_\mu Q + c_v \bar{u} v_\mu Q - c_\gamma \bar{u} \gamma_\mu \gamma_5 Q + c_v \bar{u} v_\mu \gamma_5 Q) | \Lambda_Q(v, s) \rangle, \quad (32)$$

where Q is the static heavy-quark field, and to one loop, the matching coefficients are given by (40)

$$c_\gamma = 1 - \frac{\alpha_s(\mu)}{\pi} \left[\frac{4}{3} + \ln \left(\frac{\mu}{m_b} \right) \right], \quad (33)$$

$$c_v = \frac{2}{3} \frac{\alpha_s(\mu)}{\pi}. \quad (34)$$

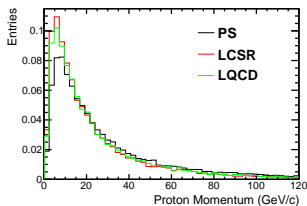
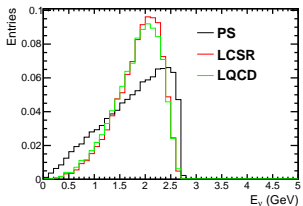
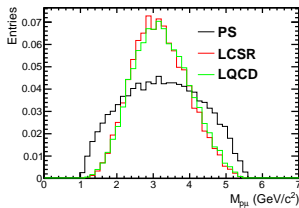
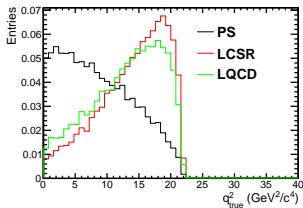
Here we set $\mu = m_b$. We can now use Eq. (3) to express the matrix element A_μ in terms of the form factors F_1 and F_2 :

$$A_\mu = \bar{u}_N(p', s') \left(F_1 + \not{v} F_2 \right) \left(c_\gamma \gamma_\mu + c_v v_\mu - c_\gamma \gamma_\mu \gamma_5 + c_v v_\mu \gamma_5 \right) \sqrt{m_{\Lambda_b}} u_{\Lambda_Q}(v, s). \quad (35)$$

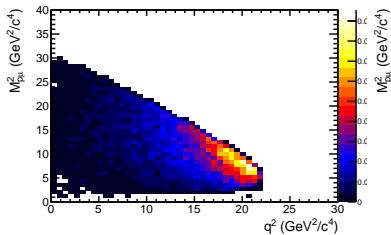
- ▶ Need to match the form factors calculated in arXiv:1108.2971 (LCSR predictions) and arXiv:1306.0446 (LQCD predictions) to convention used by EvtGen.
- ▶ Contacted Stefan Meinel (theorist at MIT) and he kindly provided me with a number of derivations relating the various form factors.

Generator level MC

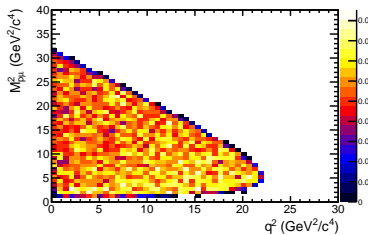
Now compare generator level distributions for LQCD, LCSR and Phase Space predictions:



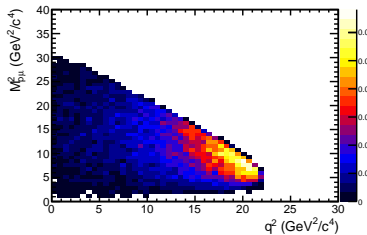
LCSR



Phase Space

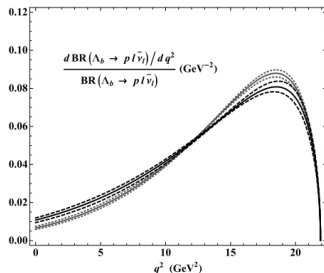


LQCD

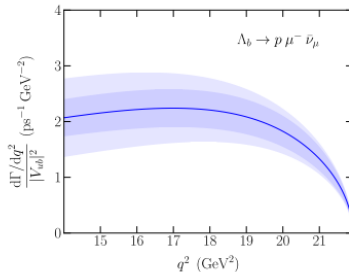


- Consistency check for these distributions?
- Can at least compare the q^2 distributions with the predicted differential rate, which is derived in terms of the various form factors in both papers.

LCSR, arXiv:1108.2971
A. Khodjamirian et al.



LQCD, arXiv:1306.0446
W. Detmold et al.



For LQCD:

the spinor $u_{\Lambda_b}(p, s)$ has the standard relativistic normalization. A straightforward calculation then gives the following differential decay rate,

$$\frac{d\Gamma}{dq^2} = \frac{|V_{ub}|^2 G_F^2}{768\pi^3 q^6 m_{\Lambda_b}^5} (q^2 - m_\ell^2)^2 \sqrt{((m_{\Lambda_b} + m_N)^2 - q^2)((m_{\Lambda_b} - m_N)^2 - q^2)} \\ \times \left[(4c_\gamma^2 + 4c_\gamma c_v + 2c_v^2)m_\ell^2 \mathcal{I} + (2c_\gamma^2(\mathcal{I} + 3q^2 m_{\Lambda_b}^2) + c_v(2c_\gamma + c_v)(\mathcal{I} - 3q^2 m_{\Lambda_b}^2))q^2 \mathcal{F} + 4c_\gamma(c_\gamma + c_v)\mathcal{K} \right], \quad (36)$$

where we have defined the combinations

$$\mathcal{F} = ((m_{\Lambda_b} + m_N)^2 - q^2)F_+^2 + ((m_{\Lambda_b} - m_N)^2 - q^2)F_-^2, \quad (37)$$

$$\mathcal{I} = m_{\Lambda_b}^4 - 2m_N^2(m_{\Lambda_b}^2 + q^2) + q^2 m_{\Lambda_b}^2 + m_N^4 + q^4, \quad (38)$$

$$\mathcal{K} = (2m_\ell^2 + q^2)((m_{\Lambda_b} + m_N)^2 - q^2)((m_{\Lambda_b} - m_N)^2 - q^2)(m_{\Lambda_b}^2 - m_N^2 + q^2)F_+F_-, \quad (39)$$

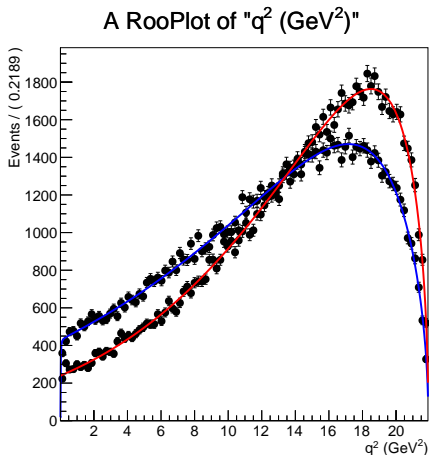
For LCSR:

$$\frac{d\Gamma}{dq^2}(\Lambda_b \rightarrow p \ell \nu) = \frac{G_F^2 m_{\Lambda_b}^3}{192\pi^3} |V_{ub}|^2 \lambda^{1/2}(1, r^2, t) \left\{ [(1-r)^2 - t][(1+r)^2 + 2t]|f_1(q^2)|^2 \right. \\ \left. + [(1+r)^2 - t][(1-r)^2 + 2t]|g_1(q^2)|^2 - 6t[(1-r)^2 - t](1+r)f_1(q^2)f_2(q^2) \right. \\ \left. - 6t[(1+r)^2 - t](1-r)g_1(q^2)g_2(q^2) + t[(1-r)^2 - t][2(1+r)^2 + t]|f_2(q^2)|^2 \right. \\ \left. + t[(1+r)^2 - t][2(1-r)^2 + t]|g_2(q^2)|^2 \right\}, \quad (63)$$

where $r = m_N/m_{\Lambda_b}$, $t = q^2/m_{\Lambda_b}^2$ and $\lambda(a, b, c) = a^2 + b^2 + c^2 - 2ab - 2ac - 2bc$. Substituting the form factors (62) and integrating over q^2 we obtain the total branching fraction

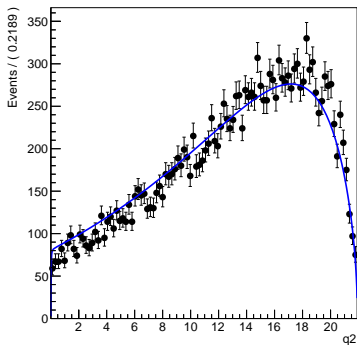
- ▶ Make 2 new RooPdf classes for LCSR and LQCD differential BF predictions.
- ▶ Can then fit these to the generator level q^2 distribution.
- ▶ Only leave normalisation free.

- ▶ Generate 100,000 MC toys using the theoretically predicted differential BF.
- ▶ Here LCSR (red) and LQCD (blue):

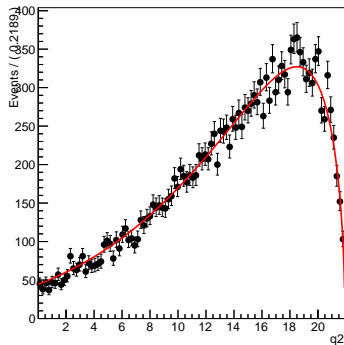


Fit generator level data:

LQCD
A RooPlot of "q2"



LCSR
A RooPlot of "q2"



Conclusion

- ▶ Thanks to Stefan Meinel have relations between the form factors defined by EvtGen and those defined in papers.
- ▶ Plots seem to make sense.
- ▶ However, LQCD fit is not perfect. Further checks + greater statistics required.
- ▶ Need to get changes into EvtGen soon so I can get new and improved MC