$\Lambda_b \to p$ form factors

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Relativistic form factors 1

The authors of Ref. [1] define in their Eq. (73)

$$\langle N^{+}(p')|\overline{u}\gamma^{\mu}b|\Lambda_{b}(p)\rangle = \overline{u}_{N}(p')\left[\tilde{F}_{1}^{V}\gamma^{\mu} + \tilde{F}_{2}^{V}v^{\mu} + \tilde{F}_{3}^{V}v'^{\mu}\right]u_{\Lambda_{b}}(p), \tag{1}$$

$$\langle N^{+}(p')|\overline{u}\gamma^{\mu}\gamma_{5}b|\Lambda_{b}(p)\rangle = -\overline{u}_{N}(p')\left[\tilde{F}_{1}^{A}\gamma^{\mu} + \tilde{F}_{2}^{A}v^{\mu} + \tilde{F}_{3}^{A}v'^{\mu}\right]\gamma_{5}u_{\Lambda_{b}}(p), \tag{2}$$

where $v^{\mu} = p^{\mu}/m_{\Lambda_b}$, $v'^{\mu} = p'^{\mu}/m_N$.

2 **HQET** form factors

In Ref. [2] we use HQET for the b quark, and calculate the form factors F_1 and F_2 , which are defined as

$$\langle N^{+}(p')| \, \bar{u}\Gamma b \, |\Lambda_b(p)\rangle = \overline{u}_N(p') \left[F_1 + \psi \, F_2 \right] \Gamma \, u_{\Lambda_b}(p). \tag{3}$$

3 Relativistic form factors in terms of HQET form factors

We insert the vector current into Eq. (3):

$$\langle N^{+}(p')| \, \bar{u}\gamma^{\mu}b \, | \Lambda_{b}(p) \rangle = \bar{u}_{N}(p') \left[F_{1} + \psi \, F_{2} \right] \gamma^{\mu} \, u_{\Lambda_{b}}(p)$$

$$= \bar{u}_{N}(p') \left[F_{1}\gamma^{\mu} + F_{2} \left(-\gamma^{\mu}\psi + \gamma^{\mu}\psi + \psi\gamma^{\mu} \right) \right] \, u_{\Lambda_{b}}(p)$$

$$= \bar{u}_{N}(p') \left[F_{1}\gamma^{\mu} + F_{2} \left(-\gamma^{\mu}\psi + v_{\nu}(\gamma^{\mu}\gamma^{\nu} + \gamma^{\nu}\gamma^{\mu}) \right) \right] \, u_{\Lambda_{b}}(p)$$

$$= \bar{u}_{N}(p') \left[F_{1}\gamma^{\mu} + F_{2} \left(-\gamma^{\mu}\psi + 2v^{\mu} \right) \right] \, u_{\Lambda_{b}}(p)$$

$$= \bar{u}_{N}(p') \left[F_{1}\gamma^{\mu} + F_{2} \left(-\gamma^{\mu} + 2v^{\mu} \right) \right] \, u_{\Lambda_{b}}(p)$$

$$= \bar{u}_{N}(p') \left[(F_{1} - F_{2})\gamma^{\mu} + 2 \, F_{2} \, v^{\mu} \right] \, u_{\Lambda_{b}}(p).$$

$$(4)$$

Here we have used the Dirac equation $\psi u_{\Lambda_b} = u_{\Lambda_b}$. By comparing Eqs. (4) and (1), we see that, in HQET,

$$\tilde{F}_{1}^{V} = F_{1} - F_{2},$$

$$\tilde{F}_{2}^{V} = 2 F_{2},$$

$$\tilde{F}_{3}^{V} = 0.$$
(5)
(6)

$$\tilde{F}_2^V = 2F_2, \tag{6}$$

$$F_3^V = 0. (7)$$

Similarly, for the the axial vector current, one finds

$$\langle N^{+}(p')| \, \bar{u}\gamma^{\mu}\gamma_{5}b \, |\Lambda_{b}(p)\rangle = \overline{u}_{N}(p') \left[(F_{1} + F_{2})\gamma^{\mu} + 2 \, F_{2} \, v^{\mu} \right] \gamma_{5} \, u_{\Lambda_{b}}(p). \tag{8}$$

By comparing Eqs. (8) and (2), we see that, in HQET,

$$\hat{F}_1^A = -(F_1 + F_2), \tag{9}$$

$$F_2^A = -2F_2, (10)$$

$$\tilde{F}_{1}^{A} = -(F_{1} + F_{2}),$$

$$\tilde{F}_{2}^{A} = -2F_{2},$$

$$\tilde{F}_{3}^{A} = 0.$$
(9)
(10)

The form factors F_1 and F_2 in Ref. [2] are written as functions of $E_N - m_N$, where E_N is the energy of the proton in the Λ_b rest frame. In terms of q^2 , we have

$$E_N - m_N = \frac{m_{\Lambda_b}^2 + m_N^2 - q^2}{2m_{\Lambda_b}} - m_N.$$
 (12)

References

- $[1]\ {\rm F.\ Hussain,\ D.-S.\ Liu,\ M.\ Kramer,\ J.\ G.\ K\"{\rm orner,\ and\ S.\ Tawfiq,\ Nucl.\ Phys.\ B}$ ${\bf 370},\ 259\ (1992).$
- $[2]\,$ W. Detmold, C.-J. D. Lin, S. Meinel, and M. Wingate, arXiv:1306.0446 [hep-lat].