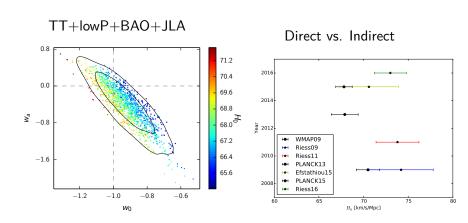


Theoretical attempts?



Issues in the statistical analysis?

Determining H_0 with Bayesian hyper-parameters: preliminary results

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April 15, 2016

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Distance modulus

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$$m_{X,i,j} = \mu_{0,i} + M_X + b_X(\log P_{i,j} - 1) + Z_X \Delta \log[O/H]_{i,j},$$

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► SN Ia magnitudes

$$5\log_{10}H_0 = m_V - \mu_0 + 5a_V + 25$$
.

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- Assume Gaussian likelihood for the data D_i

$$P_G(D_i|\vec{w}) = \tilde{N}_i \frac{\exp(-\chi_i^2(\vec{w})/2)}{\sqrt{2\pi}},$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- where

$$\chi_i^2 \equiv rac{\left(x_{{
m obs},i} - x_{{
m pred},i}(ec{w})
ight)^2}{\sigma_i^2} \quad {
m and} \quad ilde{N}_i = rac{1}{\sqrt{\sigma_i^2}}.$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- ► Control relative weight of data points in the likelihood with HP α_i , likelihood becomes

$$P(D_i|\vec{w},\alpha_i) = \tilde{N}_i \alpha_i^{1/2} \frac{\exp(-\alpha_i \chi_i^2(\vec{w})/2)}{\sqrt{2\pi}}.$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- ▶ For a given model with parameters \vec{w} and a set of N data points $\{D_i\}$

$$P(\vec{w}|\{D_i\}) = \int \ldots \int P(\vec{w}, \{\alpha_i\}|\{D_i\}) d\alpha_1 \ldots d\alpha_N,$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- From Bayes' theorem

$$P(\vec{w}, \{\alpha_i\} | \{D_i\}) = \frac{P(\{D_i\} | \vec{w}, \{\alpha_i\}) P(\vec{w}, \{\alpha_i\})}{P(D_1, \dots, D_N)},$$

$$P(\vec{w}, \{\alpha_i\}) = P(\vec{w} | \{\alpha_i\}) P(\{\alpha_i\}),$$

$$P(\{D_i\} | \vec{w}, \{\alpha_i\}) = P(D_1 | \vec{w}, \alpha_1) \dots P(D_N | \vec{w}, \alpha_N),$$

$$P(\vec{w} | \{\alpha_i\}) = \text{constant}$$

$$P(\{\alpha_i\}) = P(\alpha_1) \dots P(\alpha_N)$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- ▶ Further assuming uniform prior for HP $(P(\alpha_i) = 1)$ and that errors are never smaller than their reported value $(\alpha_i \in [0, 1])$

$$P(\vec{w}, \{D_i\}) = \frac{P(D_1|\vec{w}) \dots P(D_N|\vec{w})}{P(D_1, \dots, D_N)},$$

$$P(D_i|\vec{w}) \equiv \int_0^1 P(D_i|\vec{w}, \alpha_i) d\alpha_i.$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- In the case of a Gaussian HP likelihood

$$P(D_i|\vec{w}) = \tilde{N}_i \left(\frac{\operatorname{Erf}\left(\frac{\chi_i(\vec{w})}{\sqrt{2}}\right) - \sqrt{\frac{2}{\pi}}\chi_i(\vec{w}) \exp(-\chi_i^2(\vec{w})/2)}{\chi_i^3(\vec{w})} \right)$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- Rewriting

In
$$P(\vec{w}, \{D_i\}) = \sum_i \ln \tilde{N}_i + \ln \tilde{\chi}_i^2$$

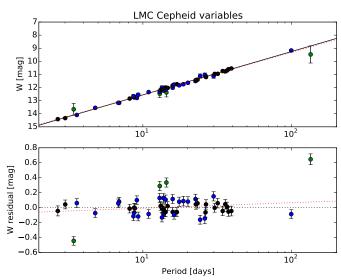
$$\tilde{\chi}_i^2(\chi_i^2(\vec{w})) \equiv \frac{\operatorname{Erf}\left(\frac{\chi_i(\vec{w})}{\sqrt{2}}\right) - \sqrt{\frac{2}{\pi}}\chi_i(\vec{w}) \exp(-\chi_i^2(\vec{w})/2)}{\chi_i^3(\vec{w})}$$

- ▶ Trust no one. Some error bars may have been underestimated (e.g., unrecognised systematic effects), then allow for a rescaling $\sigma_i \to \sigma_i/\sqrt{\alpha_i}$.
- Effective HP for each data point are given by

$$\begin{array}{lll} \alpha_i^{\text{eff}} &= 1, & \text{if} & \chi_i^2 \leq 1 \\ \alpha_i^{\text{eff}} &= \frac{1}{\chi_i^2}, & \text{if} & \chi_i^2 > 1. \end{array}$$

Applying HP

Large Magellanic Cloud (LMC)



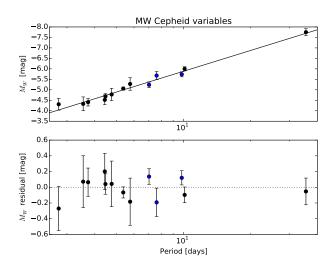
Applying HP

Large Magellanic Cloud (LMC)

LMC Cepheid variables								
Fit	Α	b_W	$\sigma_{ m int}^{ m LMC}$	Period cut				
а	12.570 (0.035)	-3.32(0.10)	0.06	10 < P < 60				
b	12.562 (0.016)	-3.30(0.05)	0.06	P < 60				
С	12.562 (0.016)	-3.31(0.05)	0.06	<i>P</i> < 205				
d	12.555 (0.019)	-3.24(0.06)	0.12	P < 205				

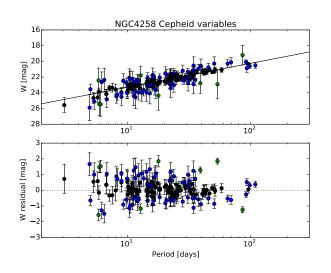
$$\mu_{0, {
m LMC}}^{
m M} = 18.49 \pm 0.05, \quad \textit{M}_{W} = -5.93 \pm 0.07$$

Applying HP Milky Way (MW)



$$M_W = -5.88 \pm 0.07,$$

$$b_W = -3.30 \pm 0.26, \qquad \sigma_{\rm int}^{\rm MW} = 0.02,$$



$$\mu_{0,4258} = 29.40 \pm 0.07, \qquad \textit{M}_{\textit{W}} = -6.12 \pm 0.15, \quad \textit{b}_{\textit{W}} = -3.02 \pm 0.17,$$

Applying HP

Megamaser system NGC 4258

I do not see any reason to discard those data sets

▶ 8 SN Ia hosts: apparent magnitudes

- ▶ 8 SN la hosts: apparent magnitudes
- MW Cepheid variables (parallax measurements): period, metallicity, magnitudes

Data

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- MW Cepheid variables (parallax measurements): period, metallicity, magnitudes
- ► LMC Cepheid variables: period, metallicity, magnitudes
- Distance modulus to LMC derived from eclipsing binaries
- Distance modulus to NGC 4258 derived from geometric maser distance and from standardised candle method for type IIP SNe

The likelihood

Full likelihood

$$\ln P(\vec{w}, \{D_i\}) = \ln P^{\text{Cepheid}} + \ln P^{\text{SNe Ia}} + \ln P^{\text{Anchors}}.$$

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$$\ln P^{ ext{Cepheid}} = \sum_{ii} \ln \tilde{\chi}^2(\chi_{ij}^{2, ext{Cepheid}}) + \ln \tilde{N}_{ij}^{ ext{Cepheid}},$$

The likelihood

► Full likelihood

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► For instance

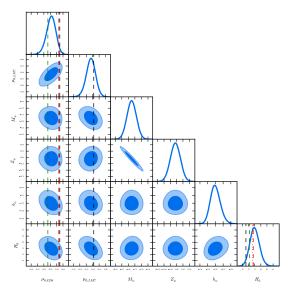
$$\ln P^{ ext{Cepheid}} = \sum_{ii} \ln \tilde{\chi}^2(\chi_{ij}^{2, ext{Cepheid}}) + \ln \tilde{N}_{ij}^{ ext{Cepheid}},$$

► SNe Ia, Anchors,...

The Hubble constant H_0

		NGC 4258 $+$ LMC $+$ MW anchors			
Fit	<i>H</i> ₀	M_W	b_W	Z_W	
$M1^a$	75.0 (3.9)	-5.87(0.18)	-3.20(0.05)[N]	-0.005(0.020)[S]	
$M1^{af}$	74.9 (3.9)	-5.88(0.18)	-3.20(0.05)[N]	-0.005(0.020)[S]	
$M1^{ag}$	73.3 (2.4)	-5.88(0.18)	-3.19(0.05)[N]	-0.004(0.020)[S]	
$M1^{ah}$	74.6 (1.9)	-5.89(0.17)	-3.21(0.05)[N]	-0.005(0.020)[S]	
$M1^{aj}$	72.4 (2.2)	-5.90(0.18)	-3.20(0.05)[N]	-0.004(0.020)[S]	
$M1^b$	76.1 (3.1)	-5.86(0.17)	-3.27(0.04)[N]	-0.007(0.019)[S]	
$M1^{be}$	75.1 (3.4)	-5.89(0.18)	-3.25(0.07)[N]	-0.004(0.020)[S]	
$M2^a$	74.7 (4.0)	-4.68(0.97)	-3.20(0.05)[N]	-0.141(0.110)[W	
$M2^{b}$	75.1 (3.8)	-3.84(1.05)	-3.28(0.04)[N]	-0.236(0.119)[W	
$M3^a$	75.0 (3.9)	-5.92(0.04)	-3.20(0.05)[N]	0	
М3 ^ь	75.4 (3.7)	-5.92(0.04)	-3.27(0.04)[N]	0	
M4 ^a	74.7 (4.0)	-4.35(1.11)	-3.20(0.05)[N]	-0.178(0.125)[N]	
$M4^b$	75.2 (3.9)	-3.18(1.25)	-3.28(0.04)[N]	-0.311(0.141)[N]	

Results: fit M1^a



Conclusions

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- ▶ Results agree with R16: 18 SN Ia hosts, better distance modulus to NGC4258, consistent values of H₀, no outlier rejection needed any longer (they claim), no outliers among SN Ia hosts

