

February 25, 2014

Example 1 A reversible engine converts one-sixth of the heat input into work. When the temperature of the sink is reduced by 70°C , its efficiency is doubled. Determine the temperature of the source and the sink.

Let's first establish that T_1 and T_2 [K] are the source and sink temperatures, respectively. For a **reversible engine**, converting $1/6$ of the heat into work means

$$\frac{T_1 - T_2}{T_1} = \frac{1}{6} \implies T_1 = 1.2T_2$$

Now if the sink temperature is reduced to $70^{\circ}\text{C} = 343.15\text{ K}$, ie, $T'_2 = T_2 - 343.15$ then the efficiency of the cycle is doubled

$$\begin{aligned} \frac{T_1 - T'_2}{T_1} &= 2 \times \frac{1}{6} \\ 2T_1 &= 3T_2 - 1029.45 \implies T_2 = 1715\text{ K} \text{ and } T_1 = 2058.90\text{ K} \end{aligned}$$

Example 2 The minimum pressure and temperature in an Otto cycle are 100 kPa and 27°C . The amount of heat added to the air per cycle is 1500 kJ/kg . Calculate:

- (a) Pressures and temperatures at all stages of the air standard Otto cycle;
- (b) Specific work and thermal efficiency of the cycle for a compression ratio of $8 : 1$.

Assuming an isentropic compression stage 1–2 and an isentropic expansion stage 3–4 with a compression ratio $V_1/V_2 = 8$. Initial temperature of 300.15 K and pressure of 100 kPa .

- *adiabatic compression (1–2):*

$$\frac{T_2}{T_1} = \left(\frac{V_1}{V_2}\right)^{\gamma-1} = r^{\gamma-1} = 8^{1.4-1} = 2.297 \implies T_2 = 689.10\text{ K}$$

and

$$\frac{P_1}{P_2} = \left(\frac{V_1}{V_2}\right)^{\gamma} = 18.379 \implies P_2 = 18.379\text{ bar}$$

- *constant volume (2–3): heat added: $C_v (T_3 - T_2) = 1500 \implies T_3 = 2772.4\text{ K}$ and*

$$\frac{P_2}{T_2} = \frac{P_3}{T_3} \implies P_3 = 73.94\text{ bar}$$

- *adiabatic expansion (3–4):*

$$\frac{T_3}{T_4} = \left(\frac{V_4}{V_3}\right)^{\gamma-1} = r^{\gamma-1} \implies T_4 = 1206.9\text{ K}$$

and

$$P_3 V_3^{\gamma} = P_4 V_4^{\gamma} \implies P_4 = 4.023\text{ bar}$$

- *Specific work* = heat added - heat rejected

$$W = C_v (T_3 - T_2) - C_v (T_4 - T_1) = 847 \text{ kJ/kg}$$

- *Thermal efficiency:*

$$\eta_{\text{Otto}} = 1 - \frac{1}{r^{\gamma-1}} = 0.5647$$

Problem 1 An ideal engine operates on the Carnot cycle using a perfect gas as the working fluid. The ratio of the greatest to the least volume is fixed as $x : 1$, the lower temperature of the cycle is also fixed, but the volume compression ratio r of the reversible adiabatic compression is variable. The ratio of the specific heats is γ . Show that if the work done in the cycle is a maximum then,

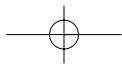
$$(\gamma - 1) \ln \frac{x}{r} + \frac{1}{r^{\gamma-1}} - 1 = 0$$

Problem 2 An ideal Otto cycle has a volumetric compression ratio of 6, the lowest cycle pressure of 0.1 MPa and operates between temperature limits of 300.15 and 1842.15 K.

- (a) Calculate the temperature and pressure after the isentropic expansion;
- (b) Since the values in (Problem 2a) are well above the lowest cycle operating conditions, the expansion process was allowed to continue down to a pressure of 0.1 MPa. Which process is required to complete the cycle ?
- (c) Determine the percentage in which the cycle efficiency has improved.

Problem 3 The volume ratios of compression and expansion for a diesel engine are 15.3 and 7.5, respectively. The pressure and temperature at the beginning of the compression are 1 bar and 27 °C. Assuming an ideal engine, determine the (a) MEP, (b) ratio of maximum pressure to MEP and (c) cycle efficiency. Also find the fuel consumption per kWh if the indicated thermal efficiency is 0.5 of ideal efficiency, mechanical efficiency is 0.8 and the calorific value of oil 42000 kJ/kg.

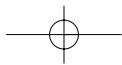
Problem 4 An air-standard dual cycle has a compression ratio of 9. At the beginning of compression, $P_1 = 100 \text{ kPa}$, $T_1 = 300 \text{ K}$ and $V_1 = 14 \text{ liters}$. The heat addition is 22.7 KJ with one-half added at constant volume and one-half added at constant pressure. Determine: (a) the temperatures at the end of each heat addition process; (b) the net work of the cycle per unit mass of air; (c) the thermal efficiency and; (d) MEP.



Appendix 1

PROPERTY TABLES AND CHARTS (SI UNITS)

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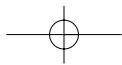


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TABLE A-17

Ideal-gas properties of air

| <i>T</i> K | <i>h</i> kJ/kg | <i>P_r</i> | <i>u</i> kJ/kg | <i>v_r</i> | <i>s^o</i> kJ/kg · K | <i>T</i> K | <i>h</i> kJ/kg | <i>P_r</i> | <i>u</i> kJ/kg | <i>v_r</i> | <i>s^o</i> kJ/kg · K |
|---------------|-------------------|----------------------|-------------------|----------------------|-----------------------------------|---------------|-------------------|----------------------|-------------------|----------------------|-----------------------------------|
| 200 | 199.97 | 0.3363 | 142.56 | 1707.0 | 1.29559 | 580 | 586.04 | 14.38 | 419.55 | 115.7 | 2.37348 |
| 210 | 209.97 | 0.3987 | 149.69 | 1512.0 | 1.34444 | 590 | 596.52 | 15.31 | 427.15 | 110.6 | 2.39140 |
| 220 | 219.97 | 0.4690 | 156.82 | 1346.0 | 1.39105 | 600 | 607.02 | 16.28 | 434.78 | 105.8 | 2.40902 |
| 230 | 230.02 | 0.5477 | 164.00 | 1205.0 | 1.43557 | 610 | 617.53 | 17.30 | 442.42 | 101.2 | 2.42644 |
| 240 | 240.02 | 0.6355 | 171.13 | 1084.0 | 1.47824 | 620 | 628.07 | 18.36 | 450.09 | 96.92 | 2.44356 |
| 250 | 250.05 | 0.7329 | 178.28 | 979.0 | 1.51917 | 630 | 638.63 | 19.84 | 457.78 | 92.84 | 2.46048 |
| 260 | 260.09 | 0.8405 | 185.45 | 887.8 | 1.55848 | 640 | 649.22 | 20.64 | 465.50 | 88.99 | 2.47716 |
| 270 | 270.11 | 0.9590 | 192.60 | 808.0 | 1.59634 | 650 | 659.84 | 21.86 | 473.25 | 85.34 | 2.49364 |
| 280 | 280.13 | 1.0889 | 199.75 | 738.0 | 1.63279 | 660 | 670.47 | 23.13 | 481.01 | 81.89 | 2.50985 |
| 285 | 285.14 | 1.1584 | 203.33 | 706.1 | 1.65055 | 670 | 681.14 | 24.46 | 488.81 | 78.61 | 2.52589 |
| 290 | 290.16 | 1.2311 | 206.91 | 676.1 | 1.66802 | 680 | 691.82 | 25.85 | 496.62 | 75.50 | 2.54175 |
| 295 | 295.17 | 1.3068 | 210.49 | 647.9 | 1.68515 | 690 | 702.52 | 27.29 | 504.45 | 72.56 | 2.55731 |
| 298 | 298.18 | 1.3543 | 212.64 | 631.9 | 1.69528 | 700 | 713.27 | 28.80 | 512.33 | 69.76 | 2.57277 |
| 300 | 300.19 | 1.3860 | 214.07 | 621.2 | 1.70203 | 710 | 724.04 | 30.38 | 520.23 | 67.07 | 2.58810 |
| 305 | 305.22 | 1.4686 | 217.67 | 596.0 | 1.71865 | 720 | 734.82 | 32.02 | 528.14 | 64.53 | 2.60319 |
| 310 | 310.24 | 1.5546 | 221.25 | 572.3 | 1.73498 | 730 | 745.62 | 33.72 | 536.07 | 62.13 | 2.61803 |
| 315 | 315.27 | 1.6442 | 224.85 | 549.8 | 1.75106 | 740 | 756.44 | 35.50 | 544.02 | 59.82 | 2.63280 |
| 320 | 320.29 | 1.7375 | 228.42 | 528.6 | 1.76690 | 750 | 767.29 | 37.35 | 551.99 | 57.63 | 2.64737 |
| 325 | 325.31 | 1.8345 | 232.02 | 508.4 | 1.78249 | 760 | 778.18 | 39.27 | 560.01 | 55.54 | 2.66176 |
| 330 | 330.34 | 1.9352 | 235.61 | 489.4 | 1.79783 | 780 | 800.03 | 43.35 | 576.12 | 51.64 | 2.69013 |
| 340 | 340.42 | 2.149 | 242.82 | 454.1 | 1.82790 | 800 | 821.95 | 47.75 | 592.30 | 48.08 | 2.71787 |
| 350 | 350.49 | 2.379 | 250.02 | 422.2 | 1.85708 | 820 | 843.98 | 52.59 | 608.59 | 44.84 | 2.74504 |
| 360 | 360.58 | 2.626 | 257.24 | 393.4 | 1.88543 | 840 | 866.08 | 57.60 | 624.95 | 41.85 | 2.77170 |
| 370 | 370.67 | 2.892 | 264.46 | 367.2 | 1.91313 | 860 | 888.27 | 63.09 | 641.40 | 39.12 | 2.79783 |
| 380 | 380.77 | 3.176 | 271.69 | 343.4 | 1.94001 | 880 | 910.56 | 68.98 | 657.95 | 36.61 | 2.82344 |
| 390 | 390.88 | 3.481 | 278.93 | 321.5 | 1.96633 | 900 | 932.93 | 75.29 | 674.58 | 34.31 | 2.84856 |
| 400 | 400.98 | 3.806 | 286.16 | 301.6 | 1.99194 | 920 | 955.38 | 82.05 | 691.28 | 32.18 | 2.87324 |
| 410 | 411.12 | 4.153 | 293.43 | 283.3 | 2.01699 | 940 | 977.92 | 89.28 | 708.08 | 30.22 | 2.89748 |
| 420 | 421.26 | 4.522 | 300.69 | 266.6 | 2.04142 | 960 | 1000.55 | 97.00 | 725.02 | 28.40 | 2.92128 |
| 430 | 431.43 | 4.915 | 307.99 | 251.1 | 2.06533 | 980 | 1023.25 | 105.2 | 741.98 | 26.73 | 2.94468 |
| 440 | 441.61 | 5.332 | 315.30 | 236.8 | 2.08870 | 1000 | 1046.04 | 114.0 | 758.94 | 25.17 | 2.96770 |
| 450 | 451.80 | 5.775 | 322.62 | 223.6 | 2.11161 | 1020 | 1068.89 | 123.4 | 776.10 | 23.72 | 2.99034 |
| 460 | 462.02 | 6.245 | 329.97 | 211.4 | 2.13407 | 1040 | 1091.85 | 133.3 | 793.36 | 23.29 | 3.01260 |
| 470 | 472.24 | 6.742 | 337.32 | 200.1 | 2.15604 | 1060 | 1114.86 | 143.9 | 810.62 | 21.14 | 3.03449 |
| 480 | 482.49 | 7.268 | 344.70 | 189.5 | 2.17760 | 1080 | 1137.89 | 155.2 | 827.88 | 19.98 | 3.05608 |
| 490 | 492.74 | 7.824 | 352.08 | 179.7 | 2.19876 | 1100 | 1161.07 | 167.1 | 845.33 | 18.896 | 3.07732 |
| 500 | 503.02 | 8.411 | 359.49 | 170.6 | 2.21952 | 1120 | 1184.28 | 179.7 | 862.79 | 17.886 | 3.09825 |
| 510 | 513.32 | 9.031 | 366.92 | 162.1 | 2.23993 | 1140 | 1207.57 | 193.1 | 880.35 | 16.946 | 3.11883 |
| 520 | 523.63 | 9.684 | 374.36 | 154.1 | 2.25997 | 1160 | 1230.92 | 207.2 | 897.91 | 16.064 | 3.13916 |
| 530 | 533.98 | 10.37 | 381.84 | 146.7 | 2.27967 | 1180 | 1254.34 | 222.2 | 915.57 | 15.241 | 3.15916 |
| 540 | 544.35 | 11.10 | 389.34 | 139.7 | 2.29906 | 1200 | 1277.79 | 238.0 | 933.33 | 14.470 | 3.17888 |
| 550 | 555.74 | 11.86 | 396.86 | 133.1 | 2.31809 | 1220 | 1301.31 | 254.7 | 951.09 | 13.747 | 3.19834 |
| 560 | 565.17 | 12.66 | 404.42 | 127.0 | 2.33685 | 1240 | 1324.93 | 272.3 | 968.95 | 13.069 | 3.21751 |
| 570 | 575.59 | 13.50 | 411.97 | 121.2 | 2.35531 | | | | | | |

**TABLE A-17**Ideal-gas properties of air (*Concluded*)

| <i>T</i> K | <i>h</i> kJ/kg | <i>P_r</i> | <i>u</i> kJ/kg | <i>v_r</i> | <i>s°</i> kJ/kg · K | <i>T</i> K | <i>h</i> kJ/kg | <i>P_r</i> | <i>u</i> kJ/kg | <i>v_r</i> | <i>s°</i> kJ/kg · K |
|---------------|-------------------|----------------------|-------------------|----------------------|------------------------|---------------|-------------------|----------------------|-------------------|----------------------|------------------------|
| 1260 | 1348.55 | 290.8 | 986.90 | 12.435 | 3.23638 | 1600 | 1757.57 | 791.2 | 1298.30 | 5.804 | 3.52364 |
| 1280 | 1372.24 | 310.4 | 1004.76 | 11.835 | 3.25510 | 1620 | 1782.00 | 834.1 | 1316.96 | 5.574 | 3.53879 |
| 1300 | 1395.97 | 330.9 | 1022.82 | 11.275 | 3.27345 | 1640 | 1806.46 | 878.9 | 1335.72 | 5.355 | 3.55381 |
| 1320 | 1419.76 | 352.5 | 1040.88 | 10.747 | 3.29160 | 1660 | 1830.96 | 925.6 | 1354.48 | 5.147 | 3.56867 |
| 1340 | 1443.60 | 375.3 | 1058.94 | 10.247 | 3.30959 | 1680 | 1855.50 | 974.2 | 1373.24 | 4.949 | 3.58335 |
| 1360 | 1467.49 | 399.1 | 1077.10 | 9.780 | 3.32724 | 1700 | 1880.1 | 1025 | 1392.7 | 4.761 | 3.5979 |
| 1380 | 1491.44 | 424.2 | 1095.26 | 9.337 | 3.34474 | 1750 | 1941.6 | 1161 | 1439.8 | 4.328 | 3.6336 |
| 1400 | 1515.42 | 450.5 | 1113.52 | 8.919 | 3.36200 | 1800 | 2003.3 | 1310 | 1487.2 | 3.994 | 3.6684 |
| 1420 | 1539.44 | 478.0 | 1131.77 | 8.526 | 3.37901 | 1850 | 2065.3 | 1475 | 1534.9 | 3.601 | 3.7023 |
| 1440 | 1563.51 | 506.9 | 1150.13 | 8.153 | 3.39586 | 1900 | 2127.4 | 1655 | 1582.6 | 3.295 | 3.7354 |
| 1460 | 1587.63 | 537.1 | 1168.49 | 7.801 | 3.41247 | 1950 | 2189.7 | 1852 | 1630.6 | 3.022 | 3.7677 |
| 1480 | 1611.79 | 568.8 | 1186.95 | 7.468 | 3.42892 | 2000 | 2252.1 | 2068 | 1678.7 | 2.776 | 3.7994 |
| 1500 | 1635.97 | 601.9 | 1205.41 | 7.152 | 3.44516 | 2050 | 2314.6 | 2303 | 1726.8 | 2.555 | 3.8303 |
| 1520 | 1660.23 | 636.5 | 1223.87 | 6.854 | 3.46120 | 2100 | 2377.7 | 2559 | 1775.3 | 2.356 | 3.8605 |
| 1540 | 1684.51 | 672.8 | 1242.43 | 6.569 | 3.47712 | 2150 | 2440.3 | 2837 | 1823.8 | 2.175 | 3.8901 |
| 1560 | 1708.82 | 710.5 | 1260.99 | 6.301 | 3.49276 | 2200 | 2503.2 | 3138 | 1872.4 | 2.012 | 3.9191 |
| 1580 | 1733.17 | 750.0 | 1279.65 | 6.046 | 3.50829 | 2250 | 2566.4 | 3464 | 1921.3 | 1.864 | 3.9474 |

Note: The properties P_r (relative pressure) and v_r (relative specific volume) are dimensionless quantities used in the analysis of isentropic processes, and should not be confused with the properties pressure and specific volume.

Source: Kenneth Wark, *Thermodynamics*, 4th ed. (New York: McGraw-Hill, 1983), pp. 785–86, table A-5. Originally published in J. H. Keenan and J. Kaye, *Gas Tables* (New York: John Wiley & Sons, 1948).

Gas Power CyclesProblem 01

$$\frac{\sqrt{3}}{\sqrt{4}} = \gamma \quad \text{and}$$

$$\text{compression ratio } l = \sqrt{4}/\sqrt{3}$$

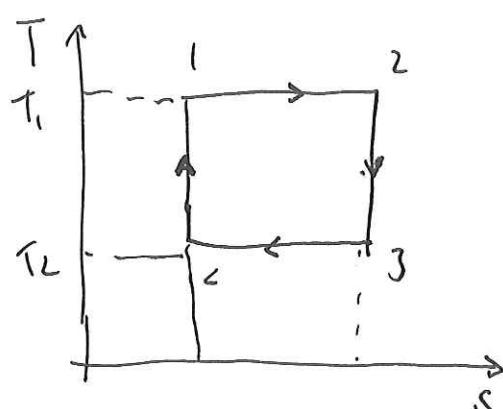
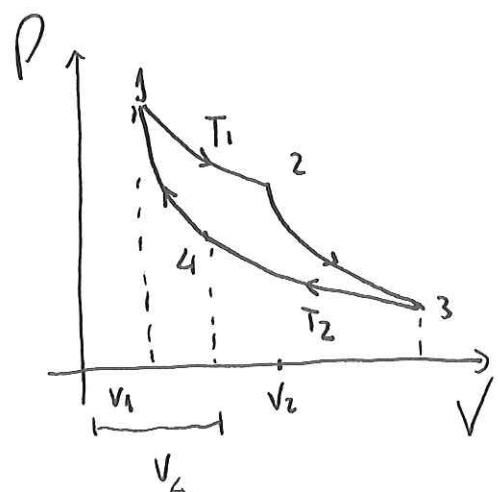
In the isotherms, since

$$\text{compression ratio} = \text{expansion ratio}$$

$$\frac{\sqrt{3}}{\sqrt{4}} = \frac{\sqrt{2}}{\sqrt{3}}$$

But we can redefine $\sqrt{3}/\sqrt{4}$ with the known variables:

$$\frac{\sqrt{3}}{\sqrt{4}} = \frac{\sqrt{3}}{\sqrt{1}} \times \frac{\sqrt{1}}{\sqrt{4}} = \gamma \cdot \frac{1}{l} = \gamma/n$$



The work done per unit mass of gas is

$\omega = \text{heat supplied} - \text{heat rejected}$

$$= R T_1 \ln(x/r) - R T_2 \ln(x/r)$$

$$= R (T_1 - T_2) \ln(x/r)$$

However we know that the isentropic change

$$\frac{T_1}{T_2} = \left(\frac{V_1}{V_2} \right)^{\gamma-1} = r^{\gamma-1}$$

Thus

$$\omega = R T_2 (r^{\gamma-1}) \ln(x/r)$$

The maximum work is obtained by differentiating

W wrt r :

$$\frac{dW}{dr} = R T_2 \left[(r^{\gamma-1}) \frac{r}{x} (-x^{-2}) + \ln \frac{x}{r} (\gamma-1) r^{\gamma-2} \right] = 0$$

This equation becomes

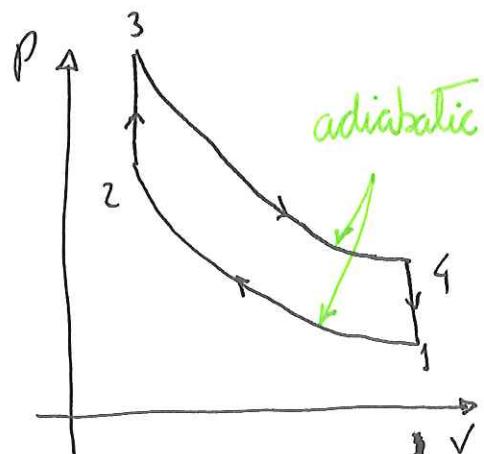
$$(\gamma-1) \ln(x/r) + \frac{1}{r^{\gamma-1}} - 1 = 0$$

Problem 02:

Given: $\frac{V_1}{V_2} = \frac{V_4}{V_3} = \gamma = 6$

$$P_1 = 0.1 \text{ MPa}$$

$$T_1 = 300.15 \text{ K} ; T_3 = 1842.15 \text{ K}$$



(a) $T_4, P_4 : ?$

For the compression 1-2

$$P_1 V_1^\gamma = P_2 V_2^\gamma \therefore P_2 = P_1 \left(\frac{V_1}{V_2} \right)^\gamma = 12.28 \text{ bar}$$

and

$$\frac{T_2}{T_1} = \left(\frac{V_1}{V_2} \right)^{\gamma-1} = 6^{\gamma-1} = 20477$$

$$T_2 = 614.61 \text{ K}$$

For the constant-volume process 2-3:

$$\frac{P_2}{T_2} = \frac{P_3}{T_3} \therefore P_3 = \frac{P_2 T_3}{T_2} = 36.82 \text{ bar}$$

And for the expansion 3-4:

$$\frac{T_3}{T_4} = \left(\frac{\sqrt{4}}{\sqrt{3}} \right)^{\delta-1} = 2.0477$$

$$T_4 = T_3 \left(\frac{\sqrt{3}}{\sqrt{4}} \right)^{\delta-1} = \boxed{899.63 \text{ K} = T_4}$$

and

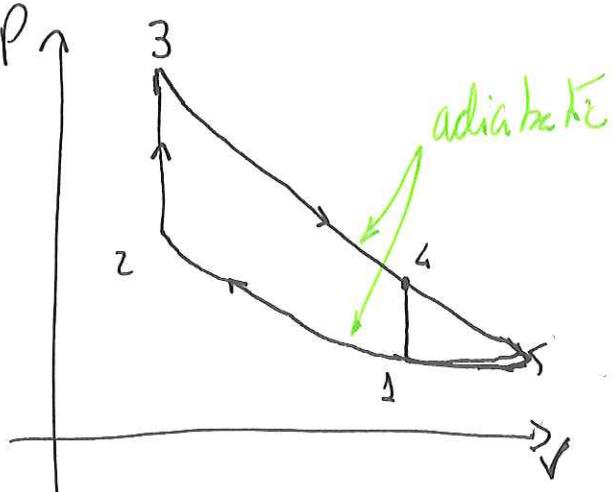
$$P_3 V_3^\delta = P_4 V_4^\delta \therefore P_4 = P_3 \left(\frac{\sqrt{3}}{\sqrt{4}} \right)^\delta$$

$$\boxed{P_4 = 2.997 \text{ bar}}$$

(b) We can use a heat rejection at constant pressure after further expansion -

We saw this as the

Atkinson cycle



(c)

$$\eta = 1 - \frac{1}{n^{\delta-1}} = 0.5116 \therefore 51.16\%$$

$$\eta_{Atkinson} = \frac{\text{Heat Supplied} - \text{Heat Rejected}}{\text{Heat Supplied}}$$

$$\eta_{Atkinson} = \frac{C_v(T_3 - T_2) - C_p(T_5 - T_1)}{C_v(T_3 - T_2)}$$

$$= 1 - \frac{C_p}{C_v} \frac{(T_5 - T_1)}{(T_3 - T_2)} = 1 - \gamma \frac{T_5 - T_1}{T_3 - T_2}$$

The isentropic process 3-5

$$\frac{T_5}{T_3} = \left(\frac{P_5}{P_3} \right)^{\frac{\gamma-1}{\gamma}} \therefore T_5 = 657.47 \text{ K}$$

$$\eta_{Atkinson} = 1 - 1.4 \frac{657.47 - 300.15}{1842.15 - 614.61}$$

$$\eta_{Atkinson} = 0.5925 \therefore 59.25\%$$

Improvement of 8.08%

Problem 05

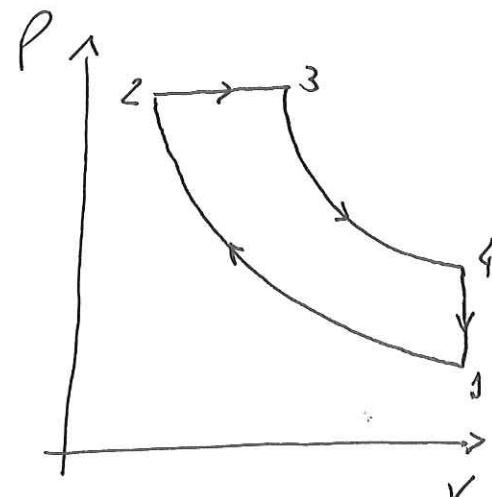
05/03
06

Given: $\sqrt{3}/\sqrt{2} = 15.3$

$$\sqrt{4}/\sqrt{3} = 7.5$$

$$P_3 = 1 \text{ bar}$$

$$T_3 = 300.15 \text{ K}$$



(a) MEP : ?

$$\begin{aligned} \text{MEP} &= \frac{\text{Work done by the cycle}}{\text{Swept Volume}} = \frac{\omega}{\sqrt{1}-\sqrt{2}} \\ &= \frac{m C_p (T_3 - T_2) - m C_v (T_4 - T_1)}{\sqrt{1} - \sqrt{2}} \end{aligned}$$

Before we are able to calculate MEP, we need to determine all properties of the cycle, i.e., T_i , ~~V_i~~ , and P_i ($i=1,4$). Thus

(i) ~~is~~ adiabatic compression 1-2 :

$$\frac{T_2}{T_1} = \left(\frac{V_1}{V_2} \right)^{\gamma-1} \therefore T_2 = 893.75 \text{ K}$$

$$P_1 V_1^\gamma = P_2 V_2^\gamma \therefore P_2 = 45.56 \text{ bar}$$

(ii) Constant pressure (2-3):

$$\frac{V_2}{T_2} = \frac{V_3}{T_3} \therefore T_3 = \frac{V_3}{V_2} T_2 = \frac{\sqrt{3}}{\sqrt{2}} \frac{\sqrt{1}}{\sqrt{2}} \frac{\sqrt{3}}{\sqrt{4}} T_2$$

$$T_3 = \frac{15.3}{7.5} T_2 \therefore T_3 = 1823.25 K$$

Let's assume that $V_2 = 1m^3$, thus the mass of air is (assuming ideal gas behaviour)

$$m = \frac{P_2 V_2}{R T_2} = \frac{45.56 \text{ bar} \times 1 \text{ m}^3}{8.314 \times 10^5 \frac{\text{m}^3 \text{ bar}}{\text{K mol}} \times 893.75 \text{ K} \times \frac{1 \text{ g/mol}}{29 \text{ g}}}$$

$$m = 17792.532 \text{ g} = 17.79 \text{ kg}$$

(iii). For adiabatic expansion (3-4):

$$\frac{T_4}{T_3} = \left(\frac{\sqrt{3}}{\sqrt{4}} \right)^{\gamma-1} = \left(\frac{1}{7.5} \right)^{\gamma-1} = 0.4467 \quad \text{---}$$

$$T_4 = 814.37 K$$

Thus

$$MEP = \frac{m [C_p(T_3 - T_2) - C_v(T_4 - T_1)]}{V_1 - V_2}$$

1.005 $\frac{KJ}{kg.K}$

0718 $\frac{KJ}{kg.K}$

01/03
08

$$MEP = \frac{10051.64 \text{ KJ}}{14.3 V_2} = 702.91 \frac{\text{KJ}}{\text{m}^3} = 7.029 \text{ bar}$$

↳ dm^3

(b) $\frac{P_2}{MEP}$: ?
MEP

$$\frac{P_2}{MEP} = 6.48$$

(c)

$$\eta = \frac{\text{Work done}}{\text{heat supplied}} = \frac{10051.64}{m C_p (T_3 - T_2)}$$

$$\eta = 0.6049 \therefore 60.49\%$$

(d) Fuel consumption : ? $\eta_T = 0.5 \eta_{\text{ideal}}$

$$\eta_T = 0.5 \times 0.6049 = 0.3024 \therefore 30.24\%$$

However if the mechanical efficiency is 0.8,

the "real" thermal efficiency is

$$\eta'_T = \eta_T \times \eta_M = 0.3024 \times 0.8 = 0.2419$$

$$\eta'_T = \frac{\text{Brake Power}}{\dot{m}_f \times C}$$

\hookrightarrow calorific heat of the oil

$$0.2419 = \frac{1}{\dot{m}_f \frac{42000}{3600}}$$

$$\therefore \dot{m}_f = 0.354 \text{ kg/kWh}$$

Problem 01: Air-Standard Dual-Cycle

Given:

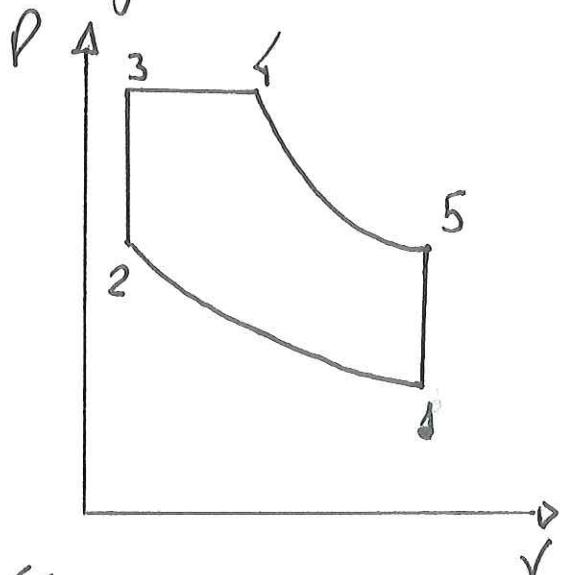
- compression ratio : $\sqrt{3}/\sqrt{2} = 9$

- $P_1 = 100 \text{ kPa}$

- $T_1 = 300 \text{ K}$

- $V_1 = 14 \text{ l} = 0.014 \text{ m}^3$

- $Q_{34} = Q_{23} = \frac{22.7}{2} \text{ kJ} = 11.35 \text{ kJ}$



Compute: T_i ? MEP_{Dual} ?

Wcycle/m?

M_{Dual}

(a) Thermal analysis of all stages:

- Stage 01: from "ideal-gas properties of air"
- (Table in attachment)

$$\begin{aligned} T_1 &= 300 \text{ K} & u_1 &= 214.97 \text{ kJ/kg} \\ P_1 &= 100 \text{ kPa} & v_{n_1} &= 625.2 \quad (\text{dimensionless}) \end{aligned}$$

↳ relative specific volume
(see note)

• Stage 02 : isentropic compression

$$\frac{\sqrt{r_2}}{\sqrt{r_1}} = \frac{\sqrt{2}}{\sqrt{1}} \therefore \sqrt{r_2} = \sqrt{r_1} \left(\frac{\sqrt{2}}{\sqrt{1}} \right) = 621.2 \times \frac{1}{9}$$
$$\sqrt{r_2} = 69.022$$

From the table :

| T(K) | U(KJ/kg) | $\sqrt{r_2}$ |
|------|----------|--------------|
| 700 | 512.33 | 69.76 |
| 710 | 520.23 | 67.07 |

With linear interpolation :

$$T_2 = 702.74 \text{ K}$$

$$U_2 = 514.50 \text{ KJ/kg}$$

• Stage 03 : For 2-3, heat addition at

constant volume : $J_3 = \frac{Q_{23}}{m} + J_2$

$$Q_{23} = m(U_3 - U_2) \therefore \cancel{m Q_{23}}$$

As we know all relevant variables at Stage 01, we can calculate the mass, m, of the fluid :

$$PV = mRT = \frac{m}{MW} RT \quad \xrightarrow{\text{molecular weight}} \quad \text{molecular weight} \\ \hookrightarrow 28.97 \text{ kg/Kmole}$$

$$m = \frac{P_1 V_1}{R T_1} \quad \overline{MW}$$

$$m = (100 \text{ kPa}) (0.014 \text{ m}^3) \left(\frac{1}{300 \text{ K}} \right) \left(\frac{\text{Kmole} \cdot \text{K}}{8.314 \text{ kJ}} \right) \times \left(28.97 \frac{\text{kg}}{\text{Kmole}} \right) \\ \times \left(\frac{10^3 \text{ N/m}^3}{1 \text{ kPa}} \right) \left(\frac{1 \text{ kJ}}{10^3 \text{ Nm}} \right)$$

$m = 0.01626 \text{ kg}$

Thus in

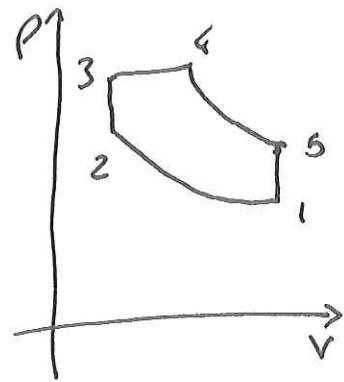
$$U_3 = \underline{Q_{23}} + J_2 = \frac{11.35 \text{ kJ}}{0.01626 \text{ kg}} + 514.50 \text{ kJ/kg} = 1212. \cancel{53} \text{ kJ}$$

From the Table :

| $T(\text{K})$ | $U(\text{kJ/kg})$ | $U(\text{kJ/kg})$ | Through linear interpolation: |
|---------------|-------------------|-------------------|----------------------------------|
| 1500 | 1635.97 | 1205.41 | $T_3 = 1507.71 \text{ K}$ |
| 1520 | 1660.23 | 1223.87 | $U_3 = 1645.33 \text{ kJ/kg}$ |

- State 04: For constant pressure heat addition

$$\text{Q}_{34} = m (h_4 - h_3)$$



$$h_4 = \frac{\text{Q}_{34}}{m} + h_3$$

$$h_4 = \frac{11.35 \text{ kJ}}{0.01626 \text{ kg}} + 1645.33 \frac{\text{kJ}}{\text{kg}} = 2343.36 \text{ kJ/kg}$$

~~Interpolate~~ From the table:

| T(K) | U(kJ/kg) | V ₂ | Through linear interpolation: |
|------|----------|----------------|-------------------------------|
| 2050 | 2314.6 | 2.555 | $T_4 = 2072.79 \text{ K}$ |
| 2100 | 2377.7 | 2.356 | $V_{R4} = 2.41643$ |

- State 05: isentropic expansion

$$\frac{\sqrt{v_{R5}}}{\sqrt{v_{R4}}} = \frac{\sqrt{5}}{\sqrt{4}} \therefore \sqrt{v_{R5}} = \sqrt{v_{R4}} \frac{\sqrt{5}}{\sqrt{4}} = \sqrt{v_{R4}} \frac{\sqrt{1}}{\sqrt{4}} = \sqrt{v_{R4}} \left(\frac{\sqrt{1}}{\sqrt{2}} \frac{\sqrt{3}}{\sqrt{4}} \right)$$

$$\Delta s \quad P_3 = P_4 \Rightarrow \frac{V_3}{V_4} = \frac{T_3}{T_4}$$

$$\sqrt{r_5} = \sqrt{r_4} \left(\frac{\sqrt{1}}{\sqrt{2}} \frac{T_3}{T_4} \right) = 24643 \left(9 \times \frac{1507.71}{\cancel{2072.79}} \right)$$

19

$$\sqrt{r_5} = 16.13$$

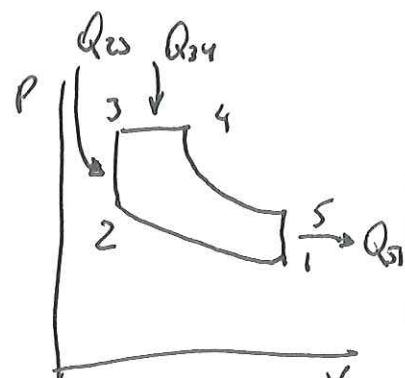
From the table :

| $T(K)$ | $U(kJ/kg)$ | \sqrt{r} |
|--------|------------|------------|
| 1140 | 880.35 | 16.946 |
| 1160 | 897.91 | 16.064 |

Through linear interpolation
 $T_5 = 1158.50K$
 $U_5 = 896.60 \text{ kJ/kg}$

Now all T_i are calculated, and W_{cycle}/m :

$$\frac{W_{cycle}}{m} = \frac{Q_{cycle}}{m} = \frac{Q_{23} + Q_{34} - Q_{51}}{m}$$



$$\frac{W_{cycle}}{m} = \frac{22.70 - m(U_5 - U_1)}{m}$$

$$= \frac{22.70 - 0.01626 (896.60 - 214.07)}{0.01626} = 713.53 \text{ kJ/kg}$$

And the efficiency:

$$\eta = \frac{W_{cycle}}{Q_{in}} = \frac{11.60}{22.70} = 0.511 \rightarrow 51.1\%$$

\downarrow

$Q_{out} + Q_{23}$

And the Mean Effective Pressure (MEP):

$$MEP = \frac{W_{cycle}}{\sqrt{(1 - V_2/V_1)}} = \frac{11.60 \text{ kJ}}{0.014 \text{ m}^3 (1 - 1/9)} \times \frac{10^3 \text{ Nm}}{1 \text{ kJ}} \frac{1 \text{ kPa}}{10^3 \text{ N/m}^2}$$

$$\underline{\underline{MEP = 932.14 \text{ kPa}}}$$

Back from the Past - A defining $\underline{\underline{A}}$

From the second law we know that for ideal gases:

$$\left. \begin{array}{l} dU = TdS - PdV \\ dH = TdS + VdP \end{array} \right\} \quad (1)$$

$$\left. \begin{array}{l} dU = C_V dT \\ dH = C_P dT \end{array} \right\} \quad (2)$$

Replacing $dU = C_V dT$, $dH = C_P dT$ and $PV = RT$

in (1) and (2):

$$\left. \begin{array}{l} dS = C_V \frac{dT}{T} + R \frac{dV}{V} \\ dS = C_P \frac{dT}{T} - R \frac{dP}{P} \end{array} \right\} \quad (3)$$

$$\left. \begin{array}{l} dS = C_V \frac{dT}{T} + R \ln \left(\frac{V_2}{V_1} \right) \\ dS = C_P \frac{dT}{T} - R \ln \left(\frac{P_2}{P_1} \right) \end{array} \right\} \quad (4)$$

with: $C_P + C_V = R$. Integrating (3-4) from (T_1, P_1) to (T_2, P_2) conditions

$$\left. \begin{array}{l} S(T_2, V_2) - S(T_1, V_1) = \int_{T_1}^{T_2} C_V \frac{dT}{T} + R \ln \left(\frac{V_2}{V_1} \right) \\ S(T_2, P_2) - S(T_1, P_1) = \int_{T_1}^{T_2} C_P \frac{dT}{T} - R \ln \left(\frac{P_2}{P_1} \right) \end{array} \right\} \quad (5)$$

$$\left. \begin{array}{l} S(T_2, V_2) - S(T_1, V_1) = \int_{V_1}^{V_2} C_V dV + R \ln \left(\frac{T_2}{T_1} \right) \\ S(T_2, P_2) - S(T_1, P_1) = \int_{P_1}^{P_2} C_P dP - R \ln \left(\frac{T_2}{T_1} \right) \end{array} \right\} \quad (6)$$

B
Now, if we define

$$S^\circ(T) = \int_0^T C_p \frac{dT}{T}$$

as the entropy at 1 atm and temperature T .
Also assuming that Temperature is continuous and
differentiable through all domain:

$$\int_{T_1}^{T_2} C_p \frac{dT}{T} = \int_0^{T_2} C_p \frac{dT}{T} - \int_0^{T_1} C_p \frac{dT}{T} = S^\circ(T_2) - S^\circ(T_1)$$

Thus, replacing in (6):

$$S(T_2, P_2) - S(T_1, P_1) = S^\circ(T_2) - S^\circ(T_1) - R \ln\left(\frac{P_2}{P_1}\right) \quad (7)$$

~~Replacing in (5):~~

~~$S(T, P)$~~

$E_{Jm}(7)$ is valid if we assume that C_p is a function of T , ie, $C_p = C_p(T)$ and we only need to integrate the $C_p(T) \frac{dT}{T}$ term. However if C_p is constant, the (5) and (6) can be easily integrated resulting in:

$$\left\{ S(T_2, V_2) - S(T_1, V_1) = C_v \ln\left(\frac{T_2}{T_1}\right) + R \ln\left(\frac{V_2}{V_1}\right) \quad (8) \right.$$

$$\left. (S(T_2, P_2) - S(T_1, P_1)) = C_p \ln\left(\frac{T_2}{T_1}\right) + R \ln\left(\frac{P_2}{P_1}\right) \quad (9) \right.$$

Now for isentropic processes — ~~$S_2 = S_1$~~ , E_{Jm} .

(8) becomes

$$C_v \ln\left(\frac{T_2}{T_1}\right) + R \ln\left(\frac{V_2}{V_1}\right) = 0 \quad (\cdot 1/C_v)$$

$$\ln\left(\frac{T_2}{T_1}\right) = - \frac{R}{C_v} \ln\left(\frac{V_2}{V_1}\right) = \ln\left(\frac{V_1}{V_2}\right)^{\frac{R}{C_v}}$$

2

Since : $R = C_p - C_v \quad \left\{ \frac{R}{C_v} = \gamma - 1 \right.$
 $\gamma = C_p/C_v \quad \left\{ \frac{R}{C_v} = \gamma - 1 \right.$

$$\left(\frac{T_2}{T_1} \right)_{S=\text{const}} = \left(\frac{\sqrt{1}}{\sqrt{2}} \right)^{\gamma-1} \quad (10)$$

Applying the same procedure to Eqn. (9):

$$\left(\frac{T_2}{T_1} \right)_{S=\text{const}} = \left(\frac{P_2}{P_1} \right)^{\frac{\gamma-1}{\gamma}} \quad (11)$$

Now (10)=(11):

$$\left(\frac{P_2}{P_1} \right)_S = \left(\frac{\sqrt{1}}{\sqrt{2}} \right)^{\gamma} \quad (12)$$

Eqns (10-12) - isentropic relations for ideal gas
 may also be re-written as:

$$\left| \begin{array}{l} T \sqrt{\gamma^{\gamma-1}} = \text{constant} \\ T P^{\frac{1-\gamma}{\gamma}} = \text{constant} \\ P V^\gamma = \text{constant} \end{array} \right| \quad (13)$$

In most cases specific heats are not constant, thus we need to use ~~$\Delta H_m(T)$~~ , $S(T_2, P_2) - S(T_1, P_1) = S^\circ(T_2) - S^\circ(T_1) - R \ln\left(\frac{P_2}{P_1}\right)$ (7)

$$S(T_2, P_2) - S(T_1, P_1) = S^\circ(T_2) - S^\circ(T_1) - R \ln\left(\frac{P_2}{P_1}\right) \quad (7)$$

and for ^{any} isentropic process 1-2:

$$\delta = S^\circ(T_2) - S^\circ(T_1) - R \ln(P_2/P_1)$$

$$S^\circ(T_2) = S^\circ(T_1) + R \ln(P_2/P_1) \quad (14)$$

Now taking ~~δ~~ , making P_2/P_1 explicit:

$$\exp\left[\frac{[S^\circ(T_2) - S^\circ(T_1)]}{R}\right] = \frac{P_2}{P_1}$$

$$\frac{P_2}{P_1} = \frac{\exp[S^\circ(T_2)/R]}{\exp[S^\circ(T_1)/R]} \quad (15)$$

The term $\exp[S^\circ/R]$ is the relative pressure $\underline{\underline{P_r}}$; therefore

$$\left[\left(\frac{P_2}{P_1} \right)_s = \frac{P_{2s}}{P_{1s}} \right] \quad (16)$$

As S° is a function of the temperature only,

P_2 is also a function of temperature only. P_2 can be obtained from tabulated tables.

Now using the ideal-gas equation:

$$\frac{PV}{T} = \text{constant}$$

$$\frac{P_1 \sqrt{V_1}}{T_1} = \frac{P_2 \sqrt{V_2}}{T_2} \therefore \frac{\sqrt{V_2}}{\sqrt{V_1}} = \frac{T_2}{T_1} \frac{P_1}{P_2} = \frac{T_2}{T_1} \frac{P_{21}}{P_{12}}$$

$$\frac{\sqrt{V_2}}{\sqrt{V_1}} = \frac{T_2 / P_{21}}{T_1 / P_{12}}$$

The term T/P_2 is called relative specific volume $\frac{V_2}{V_1}$, and just as the relative pressure is a function of the temperature only:

$$\boxed{\left(\frac{\sqrt{V_2}}{\sqrt{V_1}} \right)_S = \frac{\sqrt{V_2}}{\sqrt{V_1}}} \quad (17)$$