# Semi-Quantum Key Distribution with Limited Measurement Capabilities

Walter O. Krawec

Computer Science & Engineering Department
University of Connecticut
Storrs, CT USA

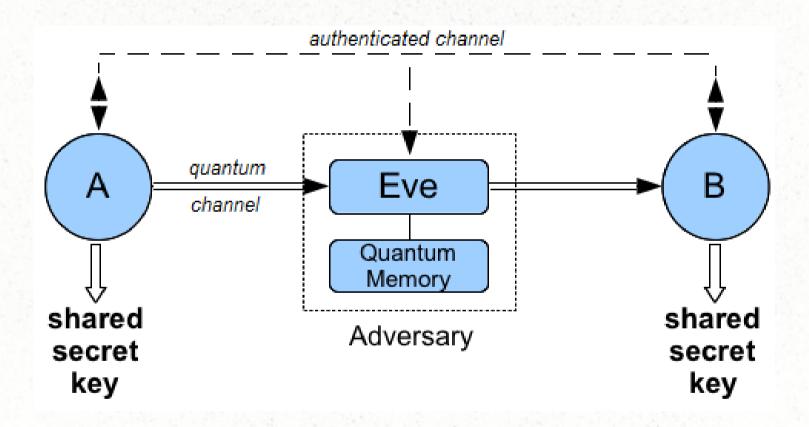
Email: walter.krawec@gmail.com

**ISITA 2018** 

# Quantum Key Distribution (QKD)

- Allows two users Alice (A) and Bob (B) to establish a shared secret key
- Secure against an all powerful adversary
  - Does not require any computational assumptions
  - Attacker bounded only by the laws of physics
  - Something that is not possible using classical means only
- Accomplished using a quantum communication channel

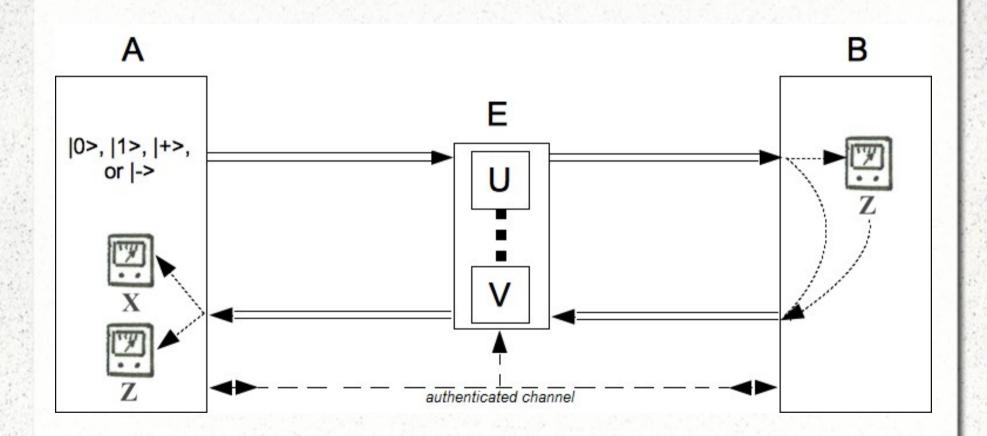
# Quantum Key Distribution



# Semi-Quantum Key Distribution

- In 2007, Boyer et al., introduced semi-quantum key distribution (SQKD)
- Now Alice (A) is quantum, but Bob (B) is limited or "classical"
  - He can only directly work with the  $Z = \{|0\rangle, |1\rangle\}$  basis.
- Theoretically interesting:
  - "How quantum does a protocol need to be in order to gain an advantage over a classical one?"
- Practically interesting:
  - What if equipment breaks down or is never installed?
- Requires a two-way quantum communication channel

# Semi-Quantum Key Distribution



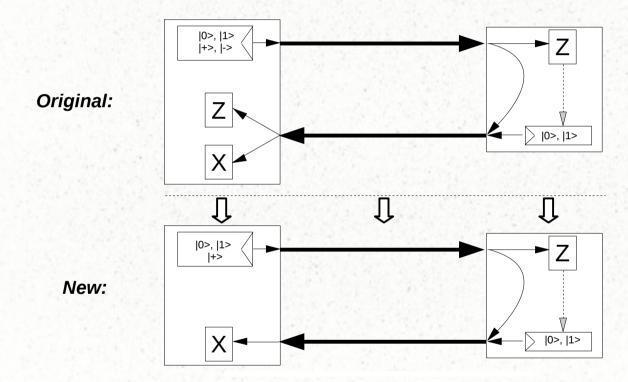
# **SQKD Security**

- Model introduced in 2007, with many protocols developed
  - But security proofs were in terms of "robustness"
- Not until 2015 that rigorous security proofs became available for some protocols along with noise tolerances and key-rate bounds
  - Noise tolerance shown to be 6.1% if using only error-statistics
  - Tolerance is 11% if using mismatched measurements [5,9,10]
    - Requires 18 different measurement statistics

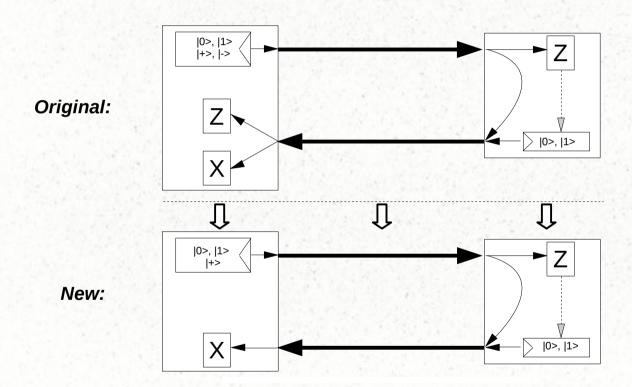
#### New Protocol

- All SQKD protocols require a two-way quantum channel
- All SQKD protocols so far have required the quantum user to measure in two (or more) bases
- We show this is not necessary
- Furthermore, the noise tolerance of our new protocol is just as high as BB84 assuming symmetric attacks!

#### New Protocol



#### New Protocol



Interestingly, protocol is **insecure** if we only look at error rates – looking at mismatched measurements is **necessary** for security of this protocol!

#### **Our Contributions**

- We propose a new SQKD protocols where **both** users have severe restrictions placed on their measurement capabilities
- We show how the technique of **mismatched measurements** [9,10] can be applied to this two-way protocol to produce very optimistic key-rate bounds
  - We also show that it is necessary to look at these mismatched statistics!
- We show our protocol has the same noise tolerance as other SQKD and fully-quantum QKD protocols

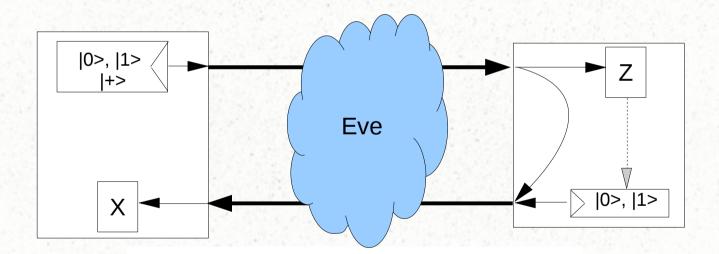
[9] S. M. Barnett, B. Huttner, and S. J. Phoenix, "Eavesdropping strategies and rejected-data protocols in quantum cryptography," Journal of Modern Optics, vol. 40, no. 12, pp. 2501–2513, 1993.

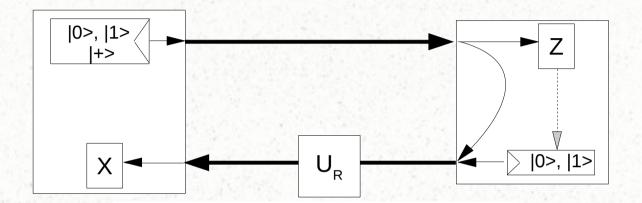
# The Protocol

#### The Protocol

- Alice's Restrictions:
  - Can only send |0>, |1>, or |+>
  - Can only measure in the X basis {|+>, |->}
- Bob's Restrictions:
  - Measure-and-Resend: Measure in the Z basis and resend the observed result
  - Reflect: Disconnect from the quantum channel and ignore the incoming state

# The Protocol (in a nutshell)

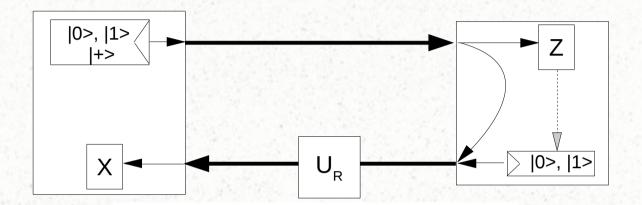




Forward Channel: Ignore (no noise)

Reverse Channel, apply  $U_R$ :

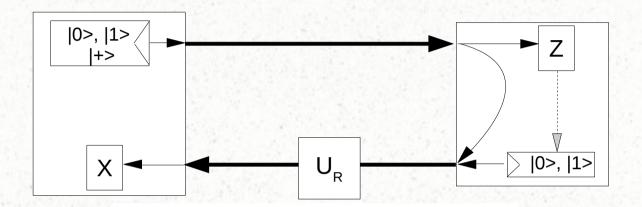
$$U_R|+>=|+,0>$$
  
 $U_R|->=|+,1>$ 



Forward Channel: Ignore (no noise)

Reverse Channel, apply  $U_R$ :

$$U_R|+>=|+,0>$$
 $U_R|->=|+,1>$ 
No detectable noise!

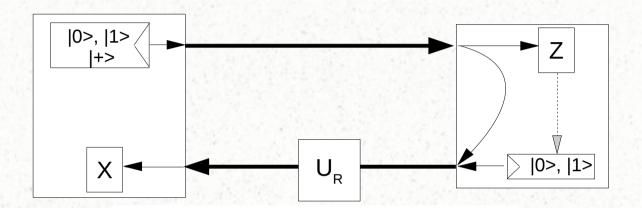


Forward Channel: Ignore (no noise)

Reverse Channel, apply  $U_R$ :

$$U_R|+>=|+,0>$$
  
 $U_R|->=|+,1>$ 

$$U_R | 0> = |+,+>$$
  
 $U_R | 1> = |+,->$ 



$$U_R|+>=|+,0>$$
 $U_R|0>=|+,+>$ 
 $U_R|1>=|+,->$ 

Two Fixes:

- •Increase complexity of protocol by having A send |->
- •Use mismatched measurements [5,9,10]

# **Security Proof**

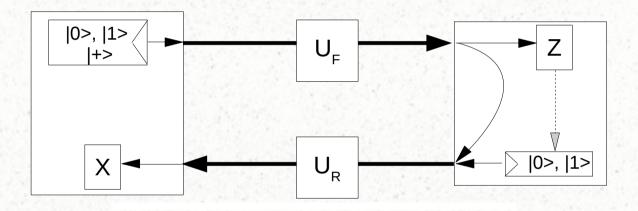
# General QKD Security

- We consider collective attacks (and comment on general attacks later)
- After the quantum communication stage and parameter estimation stage, A and B hold an N bit raw key; E has a quantum system
- They then run an error correcting protocol and privacy amplification protocol
- Result is an l(n)-bit secret key of interest is Devetak-Winter key-rate:

$$r = \lim_{N \to \infty} \frac{l(N)}{N} = \inf \left( S(A|E) - H(A|B) \right)$$

#### Two Attacks

Eve is allowed to opportunities to probe the qubit:

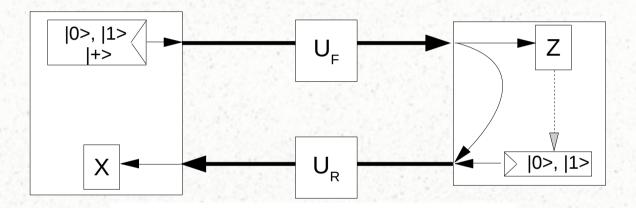


Forward: 
$$\begin{array}{c} U_{F}|0,0>_{TE}=|0,e_{0}>+|1,e_{1}>\\ U_{F}|1,0>_{TE}=|1,e_{2}>+|1,e_{3}> \end{array}$$

Reverse: 
$$U_R | i, e_j >_{TE} = | 0, e_{i,j}^0 > + | 1, e_{i,j}^1 >$$

#### Two Attacks

Eve is allowed to opportunities to probe the qubit:



Forward: 
$$U_F |0,0>_{TE} = |0,e_0> + |1,e_1>$$

$$U_F |1,0>_{TE} = |1,e_2> + |1,e_3>$$

Not necessarily normalized or orthogonal

Reverse: 
$$U_R|i,e_j>_{TE}=|0,e_{i,j}^0>+|1,e_{i,j}^1>$$

#### Quantum State ABE

 With this notation, simple algebra allows us to derive the following density operator describing one iteration (conditioning on a keybit being distilled):

$$\begin{split} \rho_{ABE} &= \frac{1}{2} [0,\!0]_{AB} \otimes ([e^0_{0,0}] \!+\! [e^1_{0,0}]) \!+\! \frac{1}{2} [0,\!1]_{AB} \otimes ([e^0_{1,1}] \!+\! [e^1_{1,1}]) \\ &+ \frac{1}{2} [1,\!0]_{AB} \otimes ([e^0_{0,2}] \!+\! [e^1_{0,2}]) \!+\! \frac{1}{2} [1,\!1]_{AB} \otimes ([e^0_{1,3}] \!+\! [e^1_{1,3}]) \end{split}$$

Note: [x]=|x>< x|

$$\begin{split} \rho_{ABE} &= \frac{1}{2} [0,\!0]_{AB} \otimes ([e^0_{0,0}] \!+\! [e^1_{0,0}]) \!+\! \frac{1}{2} [0,\!1]_{AB} \otimes ([e^0_{1,1}] \!+\! [e^1_{1,1}]) \\ &+ \frac{1}{2} [1,\!0]_{AB} \otimes ([e^0_{0,2}] \!+\! [e^1_{0,2}]) \!+\! \frac{1}{2} [1,\!1]_{AB} \otimes ([e^0_{1,3}] \!+\! [e^1_{1,3}]) \end{split}$$

Using a result in [5] allows us to bound:

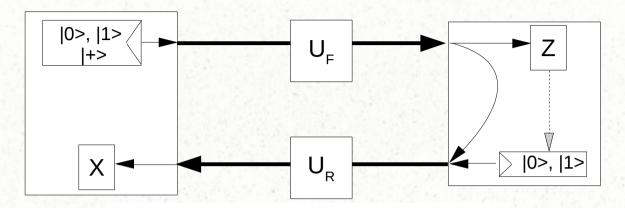
$$\begin{split} S(A|E) &\geq \frac{\langle e^0_{0,0}|e^0_{0,0}\rangle + \langle e^1_{1,3}|e^1_{1,3}\rangle}{2} \big(h\big(\frac{\langle e^0_{0,0}|e^0_{0,0}\rangle}{\langle e^0_{0,0}|e^0_{0,0}\rangle + \langle e^1_{1,3}|e^1_{1,3}\rangle}\big) - h\big(\lambda_1\big)\big) \\ &+ \langle e^1_{0,0}|e^1_{0,0}\rangle + \langle e^0_{1,3}|e^0_{1,3}\rangle}{2} \big(h\big(\frac{\langle e^1_{0,0}|e^1_{0,0}\rangle}{\langle e^1_{0,0}|e^1_{0,0}\rangle + \langle e^0_{1,3}|e^0_{1,3}\rangle}\big) - h\big(\lambda_2\big)\big) \\ &+ \langle e^1_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}{2} \big(h\big(\frac{\langle e^1_{1,1}|e^1_{1,1}\rangle}{\langle e^1_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}\big) - h\big(\lambda_3\big)\big) \\ &+ \langle e^0_{1,1}|e^0_{1,1}\rangle + \langle e^1_{0,2}|e^0_{0,2}\rangle} \big(h\big(\frac{\langle e^0_{1,1}|e^0_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}{\langle e^1_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}\big) - h\big(\lambda_4\big)\big) \end{split}$$

Unlike past SQKD protocols, we can only bound these (based on the noise in the **forward channel**)

$$\begin{split} S(A|E) &\geq \frac{\langle e^0_{0,0}|e^0_{0,0}\rangle + \langle e^1_{1,3}|e^1_{1,3}\rangle}{2} \big(h\big(\frac{\langle e^0_{0,0}|e^0_{0,0}\rangle}{\langle e^0_{0,0}|e^0_{0,0}\rangle + \langle e^1_{1,3}|e^1_{1,3}\rangle}\big) - h(\lambda_1)\big) \\ &\frac{+\langle e^1_{0,0}|e^1_{0,0}\rangle + \langle e^0_{1,3}|e^0_{1,3}\rangle}{2} \big(h\big(\frac{\langle e^1_{0,0}|e^0_{0,0}\rangle}{\langle e^1_{0,0}|e^1_{0,0}\rangle}\big) - h(\lambda_2)\big) \\ &\frac{+\langle e^1_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}{2} \big(h\big(\frac{\langle e^1_{1,1}|e^1_{1,1}\rangle}{\langle e^1_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}\big) - h(\lambda_3)\big) \\ &\frac{+\langle e^0_{1,1}|e^0_{1,1}\rangle + \langle e^1_{0,2}|e^0_{0,2}\rangle}{2} \big(h\big(\frac{\langle e^0_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}{\langle e^0_{1,1}|e^1_{1,1}\rangle + \langle e^0_{0,2}|e^0_{0,2}\rangle}\big) - h(\lambda_4)\big) \end{split}$$

Function of 
$$\Re < e_{0,0}^0 | e_{1,3}^1 >$$

$$\begin{split} S(A|E) \ge & \frac{< e^0_{0,0} | e^0_{0,0} > + < e^1_{1,3} | e^1_{1,3} >}{2} (h(\frac{< e^0_{0,0} | e^0_{0,0} >}{< e^0_{0,0} | e^0_{0,0} > + < e^1_{1,3} | e^1_{1,3} >}) - h(\lambda_1)) \\ & \frac{+ < e^1_{0,0} | e^1_{0,0} > + < e^0_{1,3} | e^0_{1,3} >}{2} (h(\frac{< e^1_{0,0} | e^1_{0,0} >}{< e^1_{0,0} | e^1_{0,0} >} + < e^0_{1,3} | e^0_{1,3} >}) - h(\lambda_2)) \\ & \frac{+ < e^1_{1,1} | e^1_{1,1} > + < e^0_{0,2} | e^0_{0,2} >}{2} (h(\frac{< e^1_{1,1} | e^1_{1,1} >}{< e^1_{1,1} | e^1_{1,1} >} + < e^0_{0,2} | e^0_{0,2} >}) - h(\lambda_3)) \\ & \frac{+ < e^0_{1,1} | e^0_{1,1} > + < e^1_{0,2} | e^0_{0,2} >}{2} (h(\frac{< e^0_{1,1} | e^1_{1,1} >}{< e^1_{1,1} | e^1_{1,1} >} + < e^0_{0,2} | e^0_{0,2} >}) - h(\lambda_4)) \end{split}$$



Forward: 
$$U_F |0,0>_{TE} = |0,e_0>+|1,e_1>$$
  
 $U_F |1,0>_{TE} = |1,e_2>+|1,e_3>$ 

Reverse: 
$$U_R|i,e_j>_{TE}=|0,e_{i,j}^0>+|1,e_{i,j}^1>$$

$$p_{0,0}^{A \to B} = \langle e_0 | e_0 \rangle$$
 $p_{0,0}^{A \to B} = \langle e_{0,0}^0 | e_{0,0}^0 \rangle + \langle e_{0,0}^1 | e_{0,0}^1 \rangle$ 

$$\begin{split} \text{Bound based on } p_{0,0}^{A \to B} &= < e_{0,0}^0 | e_{0,0}^0 > + < e_{1,0}^1 | e_{0,0}^1 > \\ S(A|E) &\geq \frac{< e_{0,0}^0 | e_{0,0}^0 > + < e_{1,3}^1 | e_{1,3}^1 >}{2} (h(\frac{< e_{0,0}^0 | e_{0,0}^0 >}{< e_{0,0}^0 | e_{0,0}^0 > + < e_{1,3}^1 | e_{1,3}^1 >}) - h(\lambda_1)) \\ &\frac{+ < e_{0,0}^1 | e_{0,0}^1 > + < e_{1,3}^0 | e_{1,3}^0 >}{2} (h(\frac{< e_{0,0}^1 | e_{0,0}^1 >}{< e_{0,0}^1 | e_{0,0}^1 > + < e_{1,3}^0 | e_{1,3}^0 >}) - h(\lambda_2)) \\ &\frac{+ < e_{1,1}^1 | e_{1,1}^1 > + < e_{0,2}^0 | e_{0,2}^0 >}{2} (h(\frac{< e_{1,1}^1 | e_{1,1}^1 >}{< e_{1,1}^1 | e_{1,1}^1 >} + < e_{0,2}^0 | e_{0,2}^0 >}) - h(\lambda_3)) \\ &\frac{+ < e_{1,1}^0 | e_{1,1}^0 > + < e_{0,2}^1 | e_{0,2}^1 >}{2} (h(\frac{< e_{1,1}^0 | e_{1,1}^0 >}{< e_{1,1}^0 | e_{1,1}^0 >} + < e_{0,2}^0 | e_{0,2}^0 >}) - h(\lambda_4)) \end{split}$$

Similarly, we can look at:  $p_{i,j}^{A \to B}$ 

$$\begin{split} S(A|E) \ge & \frac{< e^0_{0,0}|e^0_{0,0}> + < e^1_{1,3}|e^1_{1,3}>}{2} (h(\frac{< e^0_{0,0}|e^0_{0,0}>}{< e^0_{0,0}|e^0_{0,0}> + < e^1_{1,3}|e^1_{1,3}>}) - h(\lambda_1)) \\ & \frac{+ < e^1_{0,0}|e^1_{0,0}> + < e^0_{1,3}|e^0_{1,3}>}{2} (h(\frac{< e^1_{0,0}|e^1_{0,0}>}{< e^1_{0,0}|e^1_{0,0}> + < e^0_{1,3}|e^0_{1,3}>}) - h(\lambda_2)) \\ & \frac{+ < e^1_{1,1}|e^1_{1,1}> + < e^0_{0,2}|e^0_{0,2}>}{2} (h(\frac{< e^1_{1,1}|e^1_{1,1}>}{< e^1_{1,1}|e^1_{1,1}> + < e^0_{0,2}|e^0_{0,2}>}) - h(\lambda_3)) \\ & \frac{+ < e^0_{1,1}|e^0_{1,1}> + < e^1_{0,2}|e^0_{0,2}>}{2} (h(\frac{< e^0_{1,1}|e^0_{1,1}>}{< e^1_{1,1}|e^0_{1,1}> + < e^0_{0,2}|e^0_{0,2}>}) - h(\lambda_4)) \\ & \frac{+ < e^0_{1,1}|e^0_{1,1}> + < e^1_{0,2}|e^0_{0,2}>}{2} (h(\frac{< e^0_{1,1}|e^0_{1,1}>}{< e^0_{1,1}|e^0_{1,1}> + < e^1_{0,2}|e^0_{0,2}>}) - h(\lambda_4)) \end{split}$$

Just leaves: 
$$\Re < e_{0,0}^0 | e_{1,3}^1 >$$

$$\begin{split} S(A|E) \ge & \frac{< e^0_{0,0} | e^0_{0,0} > + < e^1_{1,3} | e^1_{1,3} >}{2} (h(\frac{< e^0_{0,0} | e^0_{0,0} >}{< e^0_{0,0} | e^0_{0,0} > + < e^1_{1,3} | e^1_{1,3} >}) - h(\lambda_1)) \\ & \frac{+ < e^1_{0,0} | e^1_{0,0} > + < e^0_{1,3} | e^0_{1,3} >}{2} (h(\frac{< e^1_{0,0} | e^1_{0,0} >}{< e^1_{0,0} | e^1_{0,0} >} + < e^0_{1,3} | e^0_{1,3} >}) - h(\lambda_2)) \\ & \frac{+ < e^1_{1,1} | e^1_{1,1} > + < e^0_{0,2} | e^0_{0,2} >}{2} (h(\frac{< e^1_{1,1} | e^1_{1,1} >}{< e^1_{1,1} | e^1_{1,1} >} + < e^0_{0,2} | e^0_{0,2} >}) - h(\lambda_3)) \\ & \frac{+ < e^0_{1,1} | e^0_{1,1} > + < e^1_{0,2} | e^0_{0,2} >}{2} (h(\frac{< e^0_{1,1} | e^1_{1,1} >}{< e^1_{1,1} | e^1_{1,1} >} + < e^0_{0,2} | e^0_{0,2} >}) - h(\lambda_4)) \end{split}$$

However, we show that techniques applying mismatched measurements for two-way semi-quantum protocols derived in [5] can be applied to this scenario.

By looking at the error-rate in the "reflection" case, we find:

$$p_{+,R,-}^{A \to A} = 1 - \frac{1}{2} (L_1 + L_2 + L_3 + L_4 + \eta_1 + \eta_2) - \frac{1}{2} (p_{0,R,+}^{A \to A} + p_{1,R,+}^{A \to A})$$

However, we show that techniques applying mismatched measurements for two-way semi-quantum protocols derived in [5] can be applied to this scenario.

By looking at the error-rate in the "reflection" case, we find:

$$p_{+,R,-}^{A \to A} = 1 - \frac{1}{2} \left( L_1 + L_2 + L_3 + L_4 + \eta_1 + \eta_2 \right) - \frac{1}{2} \left( p_{0,R,+}^{A \to A} + p_{1,R,+}^{A \to A} \right)$$

Needed to compute  $\lambda_i$ 

e.g., 
$$L_1 = \Re \langle e_{0,0}^0 | e_{1,3}^1 \rangle$$

However, we show that techniques applying mismatched measurements for two-way semi-quantum protocols derived in [5] can be applied to this scenario.

By looking at the error-rate in the "reflection" case, we find:

$$p_{+,R,-}^{A \to A} = 1 - \frac{1}{2} (L_1 + L_2 + L_3 + L_4 + \eta_1 + \eta_2) - \frac{1}{2} (p_{0,R,+}^{A \to A} + p_{1,R,+}^{A \to A})$$

Mismatched Measurements – in a symmetric attack, these are ½ each

However, we show that techniques applying mismatched measurements for two-way semi-quantum protocols derived in [5] can be applied to this scenario.

By looking at the error-rate in the "reflection" case, we find:

$$p_{+,R,-}^{A \to A} = 1 - \frac{1}{2} (L_1 + L_2 + L_3 + L_4 + \eta_1 + \eta_2) - \frac{1}{2} (p_{0,R,+}^{A \to A} + p_{1,R,+}^{A \to A})$$

Functions of five different mismatched statistics (each).

If symmetric attack, it holds that:  $\eta_1 = \eta_2 = 0$ 

# **Entropy Computation**

• Our entropy bound on S(A|E) is a function of eight variables:

$$< e_{0,0}^1 | e_{0,0}^1 >$$
 ,  $< e_{1,3}^1 | e_{1,3}^1 >$  ,  $< e_{0,2}^1 | e_{0,2}^1 >$  ,  $< e_{1,1}^1 | e_{1,1}^1 >$  ,  $L_{1,} L_{2,} L_{3,} L_{4}$ 

With restrictions:

Restriction	Reason
$\langle e_{i,j}^k   e_{i,j}^k \rangle \geq 0$	Property of inner-product
$ \begin{aligned} & < e_{0,0}^1   e_{0,0}^1 > \le p_{0,0}^{A \to B} \\ & < e_{1,3}^1   e_{1,3}^1 > \le p_{1,1}^{A \to B} \\ & < e_{0,2}^1   e_{0,2}^1 > \le p_{1,0}^{A \to B} \\ & < e_{1,1}^1   e_{1,1}^1 > \le p_{0,1}^{A \to B} \end{aligned} $	Unitarity of U <sub>R</sub>
$\begin{split}  L_1  \leq & \sqrt{<} e_{0,0}^0  e_{0,0}^0 > <} e_{1,3}^1  e_{1,3}^1 >} \\  L_2  \leq & \sqrt{<} e_{0,0}^1  e_{0,0}^1 > <} e_{1,3}^0  e_{1,3}^0 >} \\  L_3  \leq & \sqrt{<} e_{1,1}^1  e_{1,1}^1 > <} e_{0,2}^0  e_{0,2}^0 >} \\  L_4  \leq & \sqrt{<} e_{1,1}^0  e_{1,1}^0 > <} e_{0,2}^1  e_{0,2}^1 >} \end{split}$	Cauchy-Schwarz
$\begin{aligned} p_{+,R,-}^{A \to A} &= 1 - \frac{1}{2} (L_1 + L_2 + L_3 + L_4 + \eta_1 + \eta_2) \\ &- \frac{1}{2} (p_{0,R,+}^{A \to A} + p_{1,R,+}^{A \to A}) \end{aligned}$	Mismatched Measurements

# **Evaluation + Summary**

#### Results

- We numerically minimize S(A|E) based on the above constraints
  - Need to minimize as we must assume the worst case
- Computing H(A|B) is trivial given observable data
- Thus, we can compute the key-rate r = S(A|E) H(A|B)

	Independent: $Q_x = 2Q(1-Q)$	Dependent: $Q_x = Q$
Max. Q:	Q < 7.9%	Q < 11%

# Required Measurement Statistics

#### **Error Rates**

$A \rightarrow B$
$p_{0,0}^{A \rightarrow B}$
$p_{0,1}^{A \rightarrow B}$
$p_{0,1}$
$p_{1,0}^{A \rightarrow B}$
$p_{1,0}$
$p_{1,1}^{A \rightarrow B}$
$P_{1,1}$

$$p_{+,R,-}^{A \rightarrow A}$$

#### **Mismatched Events**

$$p_{+,0}^{A \rightarrow B}$$

$$p_{+,1}^{A \rightarrow B}$$

$$p_{0,R,+}^{A \to A} \\ p_{1,R,+}^{A \to A} \\ p_{1,R,+}^{A \to A} \\ p_{0,0,+}^{A \to A} \\ p_{0,0,+}^{A \to A} \\ p_{1,0,+}^{A \to A} \\ p_{0,1,+}^{A \to A} \\ p_{0,1,+}^{A \to A} \\ p_{1,1,+}^{A \to A}$$

# Required Measurement Statistics

#### **Error Rates**

$$p_{0,0}^{A op B}$$
 $p_{0,1}^{A op B}$ 
 $p_{1,0}^{A op B}$ 
 $p_{1,1}^{A op B}$ 

$$p_{+,R,-}^{A \rightarrow A}$$

While we only evaluated on a symmetric channel, our equations apply to arbitrary channels.

#### **Mismatched Events**

$$p_{+,0}^{A \rightarrow B} \\ p_{+,1}^{A \rightarrow B}$$

$$p_{0,R,+}^{A \to A}$$

$$p_{1,R,+}^{A \to A}$$

$$p_{1,R,+}^{A \to A}$$

$$p_{0,0,+}^{A \to A}$$

$$p_{1,0,+}^{A \to A}$$

$$p_{1,0,+}^{A \to A}$$

$$p_{0,1,+}^{A \to A}$$

$$p_{0,1,+}^{A \to A}$$

$$p_{1,1,+}^{A \to A}$$

#### **Future Work**

- How does the protocol compare to others over nonsymmetric attacks?
- We only considered collective attacks does the usual techniques of applying de Finetti work here?
  - Or some other way to extend to general attacks
- What about a finite-key analysis?
  - Especially comparing with other SQKD or fully quantum protocols.

Thank you! Questions?

# References

- [2] M. Boyer, D. Kenigsberg, T. Mor. Quantum key distribution with classical Bob. PRL 99:140510, 2007
- [5] W. O. Krawec, "Quantum key distribution with mismatched measurements over arbitrary channels," Quantum Information and Computation, vol. 17, pp. 209–241, 2017.
- [9] S. M. Barnett, B. Huttner, and S. J. Phoenix, "Eavesdropping strategies and rejected-data protocols in quantum cryptography," Journal of Modern Optics, vol. 40, no. 12, pp. 2501–2513, 1993.
- [10] S. Watanabe, R. Matsumoto, and T. Uyematsu, "Tomography increases key rates of quantum-key distribution protocols," Physical Review A, vol. 78, no. 4, p. 042316, 2008.
- [14] W. O. Krawec. Security proof of a semi-quantum key distribution protocol. In IEEE ISIT 2015, 686-690.
- [17] W. O. Krawec. Quantum key distribution with mismatched measurements over arbitrary channels. Quantum Information and Computation. 17 (3&4) 209-241. 2017.
- [21] N. Beaudry, M. Lucamarini, S. Mancini, and R. Renner. Security of two-way quantum key distribution. PRA 88(6)062302, 2013
- [23] I. Devetak and A. Winter. Distillation of secret key and entanglement from quantum states. Proc. Royal Society A 461(2053) 207-235, 2005.
- [24] M. Berta, M. Christandl, R. Colbeck, J. Renes, R. Renner. The uncertainty principle in the presence of quantum memory. Nature Physics 6(9):659-662, 2010.
- [25] A. Winter. Tight uniform continuity bounds for quantum entropies: conditional entropy, relative entropy distance and energy constraints. Communications in Mathematical Physics. 347(1):291-313,2016.

# References (cont.)

- C.H. Bennett and G. Brassard, 1984, Quantum cryptography: Public key distribution and coin tossing. in Proc. IEEE Int. Conf. on Computers, Systems, and Signal Processing. Vol 175, NY.
- C.H. Bennett, 1992, Quantum cryptography using any two nonorthogonal states. Phys. Rev. Lett., 68:3121-3124.
- M. Boyer, D. Kenigsberg, and T. Mor, 2007, Quantum Key Distribution with classical bob, in ICQNM.
- M. Christandl, R. Renner, and A. Ekert, A generic security proof for quantum key distribution.
- I. Devetak and A. Winter, Distillation of secret key and entanglement from quantum states. Proc. R. Soc. A 2005 461.
- W.O. Krawec, 2014, Restricted attacks on semi-quantum key distribution protocols. Quantum Information Processing, 13(11):2417-2436.

# References (cont.)

- H. Lu and Q.-Y. Cai, 2008, Quantum key distribution with classical Alice, Int. J. Quantum Information 6, 1195.
- R. Renner, N. Gisin, and B. Kraus, 2005, Information-theoretic security proof for QKD protocols. Phys. Rev. A, 72:012332.
- R. Renner, 2007, Symmetry of large physical systems implies independence of subsystems, Nat. Phys. 3, 645.
- V. Scarani, A. Acin, G. Ribordy, and N. Gisin, 2004, Phys. Rev. Lett. 92, 057901.
- Z. Xian-Zhou, G. Wei-Gui, T. Yong-Gang, R. Zhen-Zhong, and G. Xiao-Tian, 2009, Quantum key distribution series network protocol with m-classical bobs, Chin. Phys. B 18, 2143.
- Xiangfu Zou, Daowen Qiu, Lvzhou Li, Lihua Wu, and Lvjun Li, 2009, Semiquantum key distribution using less than four quantum states. Phys. Rev. A, 79:052312.