

Notes for PHYS 233: Interstellar Medium

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1 Introduction

Monday, January 6, 2013

This course is called “Interstellar Medium” (ISM), but really it is about the more general topic of the physics of diffuse gas in the universe. By diffuse, we mean *really* diffuse, down to number densities of 1 cm^{-3} , much more diffuse than the best vacuums in terrestrial laboratories. This includes the interstellar gas in galaxies as well as the intergalactic gas in between the galaxies.

1.1 Organization of the ISM in the Milky Way

Interstellar gas is what forms the stars in galaxies that are the dominant sources of energy and light in the universe. Thus, understanding the physics of the ISM helps us understand (and predict) the visible appearance of galaxies.

The space between the stars is occupied by gas, dust, cosmic rays, electromagnetic radiation (from stars, the CMB, and radiation from interstellar matter), neutrinos, dark matter particles (whatever they are), and magnetic and gravitational fields. Do note that dust is distinct from gas, as dust grains have a lengthscale (micron) four orders of magnitude greater than gas particles (angstrom).

In the Milky Way, there is roughly $10^{11} M_{\odot}$ in various components:

- $5 \times 10^{10} M_{\odot}$ of stars
- $5 \times 10^{10} M_{\odot}$ of dark matter
- $7 \times 10^9 M_{\odot}$ of interstellar gas

So the Milky Way is a rather gas-poor galaxy. Other galaxies like the Magellanic clouds have most of their baryons in the gas phase. The obscuration in the central disk of the Milky Way is produced by the dust along the midplane (see color plate in 1.2, 1.65, 2.2 μm images). While the stars do form in a very thin disk (thin compared to its width), the starlight heats up dust and expands the dust disk to higher altitudes from the midplane.

The magnetic fields in the Milky Way can be detected by looking for synchrotron emission (caused by electrons spiraling around galactic field lines). We do observe this emission (see additional plate) and believe the magnetic fields are caused by a dynamo effect due to the circulating galactic material.

1.1.1 Gas Content

There is about $5 \times 10^9 M_{\odot}$ of hydrogen in the Milky Way, in the following states

- $2.9 \times 10^9 M_{\odot}$ of HI gas (neutral)
- $1.12 \times 10^9 M_{\odot}$ of HII gas (ionized)
- $0.84 \times 10^9 M_{\odot}$ of H₂ gas (molecular)

We can track the presence of interstellar gas by looking at C^+ emission in the Milky Way, which tracks hydrogen presence very well, though the weighting is different since the emission strength is proportional to the density squared.

Finally, we track galactic gas via CO emission, which is very concentrated to the disk, but also has strong filaments outside the central disk of the galaxy.

The interstellar medium exhibits turbulence that is the result of supernovae and other violent hydrodynamic events.

Phases of Interstellar Gas We categorize the ISM into various phases according to the ionization/molecular state of its hydrogen (descriptions taken from Draine 1.1):

- **Warm H I** : Neutral hydrogen gas at temperatures around $T \approx 10^{3.7}$ K. Typically at densities of $n_H \approx 0.6 \text{ cm}^{-3}$. Occupies about 40% of the galactic disk. Often referred to as the warm neutral medium (WNM)
- **Cool H I** : Neutral hydrogen gas at temperatures around $T \approx 100$ K with higher densities of $n_H \approx 30 \text{ cm}^{-3}$. Occupies around 1% of the local ISM, sometimes referred to as the cold neutral medium (CNM)
- **Diffuse H₂**: Like the CNM, but with larger densities and column densities, allowing molecular hydrogen (H_2) to be more abundant.
- **Dense H₂**: Gravitationally bound clouds with number densities exceeding 10^3 cm^{-3} . Distinguished by their dark appearance (strong optical extinction) in their central regions. Hosts of star formation. Not “dense” by terrestrial standards.
- **Ionized H II** at 10^4 K: Gas consisting largely of photoionized hydrogen, likely from ultraviolet radiation from a nearby massive star. May be dense material from a cloud (H II region) or diffuse intercloud material (diffuse H II). Lifetimes typically on the order of that of a massive star (Myrs). Extended diffuse photoionized regions, sometimes called the warm ionized medium (WIM), contain *much* more total mass than the more visible H II regions. Planetary nebulae (PNe) are also in this class.
- **Coronal Gas** (Ionized H II at $\log T > 5.5$): Very hot, ionized gas that has been shock-heated by blastwaves from supernova explosions. It is mostly collisionally ionized with some exotic ionization states present. Typically very low density, filling about half the volume of the galactic disk. Cools on Myr timescales. Often referred to as the hot ionized medium (HIM).

Elemental Composition The gas is largely made of hydrogen and helium from the early universe, but an additional 1-2% (by mass) is heavy elements ($Z > 2$), or to astronomers, “metals”. The metallicity is actually a declining function from the center of the galaxy outwards since it is the result of stellar processing, and the center of the galaxy is more evolved than its relatively young outer regions. The metallicity near the solar system is about half of that at the galactic center. We know the local composition of the ISM by measuring solar photospheric abundances and the composition of meteorites (See table 1.4 in Draine).

Energy Type	Density (eV cm^{-3})
Thermal	0.49
Bulk Kinetic	0.22
Cosmic Ray	1.39
Magnetic	0.89
CMB	0.265
Far infrared from dust	0.31
Starlight	0.54

Table 1: Energy densities in the ISM.

1.1.2 Energy Content

Energy in the CMB has a strange equipartition where nearly all components are near 1 eV cm^{-3} . Some of this is self-sustaining, but since the CMB energy density increases as $(1+z)^4$, its equality is coincidental. The rough energy densities by categories are given in Table 1. Since the galactic magnetic fields are driven by bulk motions, it is unsurprising that the turbulent (bulk) and magnetic energies are nearly equal. If the cosmic ray energy density were significantly greater, the magnetic fields could not confine them, so they would simply escape the galaxy. Similarly, if the energy from starlight were significantly higher, the ISM would expand and likely dampen star formation, causing the near equipartition that we observe. If any of these components grows much larger than the gravitational binding energy, then hydrostatic equilibrium is disrupted, driving a wind. All galaxies drive winds at some time.

1.1.3 Galactic Endgame

All of the ISM constituents are present between galaxies, and the same physical processes apply to studying the intergalactic medium (IGM). In our present state of the field, we can't account for the relative overabundance of dark matter in the vicinity of a galaxy. It is thought that the “missing baryons” are in the circum-galactic material (CGM). The only place where the mass is “right” is in massive clusters. What is unknown is what sort of lengthscale determines how far out this CGM might extend. Interestingly, the material observed in the CGM contains metals, so either galactic winds are responsible for pushing material processed by stellar evolution out of the galaxy, or very massive population III stars formed at $z \sim 20 - 40$ to form the first metals.

1.2 Basics of Diffuse Gas

Perhaps the prettiest regions of diffuse gas are the H II regions, like the Orion Nebula, which is illuminated by four central stars (the trapezium). The typical H II region has a number density around 1 cm^{-3} and a temperature of 10^4 K . A big difference between these H II regions (and in fact most instances of interstellar gas) and stellar interiors is that they are *not* in local thermodynamic equilibrium (LTE).

The NLTE-ness of diffuse gas is due to the much, much lower collision rate than in planetary atmospheres or stellar interiors. In order to achieve LTE, a gas must satisfy the following conditions:

Environment	Number Density (cm^{-3})	Collision Timescale (s)
Earth Atmosphere	2×10^{13}	$\sim 10^{-9}$
Molecular Cloud	10^6	10^8 (several days)
ISM	1	10^{11} (10^4 yr)

Table 2: Some characteristic collision timescales.

- Population of excited states given by Boltzmann equilibrium
- Particle energies distributed according to the Maxwell distribution
- Ionization balance given by Saha Equation
- Photon energies described by Planck Function

In diffuse gas, we may indeed be able to identify a kinetic temperature that gives the average kinetic energy of the particles, but this temperature need not be the same as the Boltzmann temperature, which gives the occupation of energy states in the system. Likewise, the ionization temperature, Planck temperature (of a blackbody), and other such derived temperatures need not match such a diffuse system.

2 Collisionally Excited Lines

Wednesday, January 8, 2014

A typical timescale for collisions between particles can be estimated via

$$\tau = \frac{1}{\sigma v n} \quad (2.1)$$

That is, it's essentially the inverse of the expected collision rate. Here σ is the cross section for interaction, v is the relative velocity of particles, and n is the number density of particles. Some typical timescale for various environments are given in Table 2.

Also of importance is the volumetric rate of emission of photons due to recombination, which will be proportional to the product of electron and ion number densities, or $r_{\text{emis}} \propto n_{\text{ion}} n_e$. We define the **Emission Measure** as

$$\text{EM} = \int n_e^2 ds \quad (2.2)$$

where the integral is over some path in space. The “natural” units for this measure is then pc cm^{-6} .

Example: Emission Measure of a Nova Shell For a nova shell with mass $\Delta M = 10^{-4} M_{\odot}$ expanding with a velocity $v = 10^3 \text{ km s}^{-1}$, we get an approximate emission measure of

$$\text{EM} = (3 \times 10^{-4} \text{ pc})(10^7 \text{ cm}^{-3}) \sim 3 \times 10^{10} \text{ pc cm}^{-6} \quad (2.3)$$

So at early times nova shells are very bright. We could play this game with other objects like planetary nebulae (PNe), H II Regions, Diffuse warm-ionized gas, or circumgalactic gas to get an idea of how bright these objects would appear.

2.1 Two-Level Atom

If we consider an atom that has only two states, 1 and 2, with degeneracies g_1 and g_2 with an energy gap ϵ_{12} , the collisional cross section is given by

$$\sigma_{12}(E) = \frac{h^2}{8\pi m_e E} \frac{\Omega_{12}(E)}{g_1} \quad (2.4)$$

where Ω is the **collision strength**, which is a dimensionless quantity with values around order unity. The **Principle of Detailed Balance**, which we have not yet proven, tells us that $\Omega_{12} = \Omega_{21}$.

Indeed, the principle of detailed balance tells us that the rate of upwards and downward collisional transitions are equal, or

$$R_{12} = R_{21} \quad (2.5)$$

$$n_1 n_e \alpha_{12} = n_2 n_e \alpha_{21} \quad (2.6)$$

where α is the collisional excitation coefficient for that reaction. To derive this principle, we turn to statistical mechanics, and the concept of a **partition function**, which is the normalization to a boltzmann distribution:

$$Z(T) \equiv \sum_s e^{-E(s)/(kT)} \quad (2.7)$$

where the sum is over all microstates. In a dilute gas, the partition function is actually a product of the partition functions from the internal system and that of the translational system:

$$Z(T) = Z_{\text{trans}}(T) \times Z_{\text{int}}(T) \quad (2.8)$$

Then we can define the partition function per unit volume in the usual way

$$f(T) = \frac{Z(T)}{V} \quad (2.9)$$

$$= \left(\frac{2\pi M_x kT}{h^3} \right)^{3/2} \times Z_{\text{int}}(T) \quad (2.10)$$

Where we've glossed over some details present in Draine (Chapter 3). Using statistical mechanics, we are more or less assuming we are in LTE, so the number densities of various states for the reaction



is given in the usual way:

$$\frac{n_{\text{LTE}}(C)}{n_{\text{LTE}}(A)n_{\text{LTE}}(B)} = \frac{f(C)}{f(A)f(B)} \quad (2.12)$$

which constrains the rate coefficients in a specific way. Armed with this tool, we can actually try to compute collisional excitation rates:

$$R_{12} = n_1 n_e \int_{E_{12}}^{\infty} \sigma(v) v f(v) dv \quad (2.13)$$

$$R_{21} = n_2 n_e \int_0^{\infty} \sigma(v) v f(v) dv \quad (2.14)$$

Now using the principle of detailed balance, we equate these rates and get

$$n_1 e^{-E_{12}/(kT)} \frac{\Omega_{12}}{g_1} = \frac{n_2 \Omega_{21}}{g_2} \quad (2.15)$$

Solving for the ratio of the number densities, we get

$$\frac{n_2}{n_1} = \frac{\alpha_{12}}{\alpha_{21}} = \frac{\Omega_{21}}{g_1} \frac{g_2}{\Omega_{12}} e^{-E_{12}/(kT)} \quad (2.16)$$

Or, more simply

$$\frac{n_2}{n_1} = \frac{\Omega_{21}}{\Omega_{12}} \frac{g_2}{g_1} e^{-E_{12}/(kT)} \quad (2.17)$$

Now if we are in LTE, we immediately recover $\Omega_{21} = \Omega_{12}$ since the Boltzmann ratio is already present. However, the collision strength is an intrinsic property of the ion, so this must be true in all cases. That is, we don't require LTE to recover $\Omega_{12} = \Omega_{21}$ simply because we proved that it *does* apply in LTE, and thus at all times.

In the low density limit, spontaneous decay dominates emission since collisionally excited transitions are few and far between. Then the transition rate can be written as

$$n_e n_1 \alpha_{12} = n_2 A_{21} \quad (2.18)$$

where A is the **Einstein Coefficient** for spontaneous decay from state 2 to state 1.

$$F_{12} = E_{12} A_{21} n_2 \quad (2.19)$$

Subbing in some Boltzmann algebra, n_2 can be expressed as

$$n_2 = \frac{n_e n_1}{A_{21}} \left(\frac{2\pi\hbar^4}{km_e^3} \right)^{1/2} \frac{\Omega_{12}}{g_1} \frac{e^{-E_{12}/(kT)}}{\sqrt{T}} \quad (2.20)$$

Plugging (2.20) into (2.19) and evaluating some of the constants, we get

$$F_{12} = E_{12} n_e n_1 (8.62942 \times 10^{-6}) \frac{1}{\sqrt{T}} \frac{\Omega_{12}}{g_1} e^{-E_{12}/(kT)} \quad (2.21)$$

Now in the high density limit, we do something similar, getting

$$n_e n_1 \alpha_{12} = n_2 A_{21} + n_e n_2 \alpha_{21} \quad (2.22)$$

Now we get

$$F_{12} = E_{21} A_{21} n_1 \frac{g_2}{g_1} e^{-E_{12}/(kT)} \quad (2.23)$$

These two regimes beg for us to define a critical density that separates these two limits where $n_2 A_{21} = n_2 n_e \alpha_{21}$, which is, after some algebra,

$$n_{e,\text{crit}} = \frac{A_{21} g_2 T^{1/2}}{\beta \Omega_{12}} \quad (2.24)$$

where

$$\beta \equiv \left[\frac{2\pi \hbar^4}{k m_e^3} \right]^{1/2} \quad (2.25)$$

For forbidden lines, this critical density is in the range of $10^2 - 10^7 \text{ cm}^{-3}$. For intercombination lines, the critical density is around 10^{10} cm^{-3} . For resonance lines (aka permitted lines), the critical density is 10^{15} cm^{-3}

2.2 Three-Level Atom

Friday, January 10, 2014

Suggested References

- Dopita & Sutherland Chapter 3

Now we consider the slightly more complicated three-level atom. Considering that the collisional excitation rate is, in general, given by

$$C_{ij} = n_e \alpha_{ij} \quad (2.26)$$

for the appropriate rate factors α_{ij} for excitations from state i to state j . Now we'll employ statistical equilibrium to figure out the equilibrium rates and concentrations. For level 3, we get the following

$$n_1 C_{13} + n_2 C_{23} = n_3 (C_{31} + C_{32} + A_{32} + A_{31}) \quad (2.27)$$

now for level 2, we get

$$n_1 C_{12} + n_3 (C_{32} + A_{32}) = n_2 (C_{21} + C_{23} + A_{21}) \quad (2.28)$$

and for normalization purposes, let's assume that the total number density is given by n :

$$n_1 + n_2 + n_3 = n \quad (2.29)$$

(2.27) and (2.28) essentially just equate the rates of reactions (collisional or spontaneous) for the "formation" of that energy state (left side) to the rates of "destruction" rates (right side).

Now we'll make some assumptions, namely that the energy levels are roughly equally spaced, or $E_{12} \approx E_{23}$. This tells us that the rate of excitations from state 1 to state 3 will be much lower than

that from state 1 to state 2, or $C_{13} \gg C_{12}$. Now in a low density limit, we may ignore collisional excitation, which reduces our equations to a much cleaner form:

$$n_1 C_{13} = n_3 (A_{32} + A_{31}) \quad (2.30)$$

$$n_1 C_{12} + n_3 A_{32} = n_2 A_{21} \quad (2.31)$$

$$n_1 + n_2 + n_3 = n \quad (2.32)$$

Solving these equations for n_3 and n_2 gives

$$n_3 = \frac{n_1 C_{13}}{A_{32} + A_{31}} \quad (2.33)$$

$$n_2 = \left[n_1 C_{12} + \frac{A_{32}}{A_{32} + A_{31}} n_1 C_{13} \right] \frac{1}{A_{21}} \quad (2.34)$$

$$\approx \frac{n_1 C_{12}}{A_{21}} \quad (2.35)$$

Then the ratio of fluxes is

$$\frac{F_{32}}{F_{21}} = \frac{n_3 A_{32} E_{32}}{n_2 A_{21} E_{21}} = \dots = \frac{C_{13}}{C_{12}} \frac{A_{32}}{A_{32} + A_{31}} \frac{E_{32}}{E_{21}} \quad (2.36)$$

Now recalling that the C 's are related to our α 's, via, for example,

$$C_{13} = n_3 \left(\frac{2\pi\hbar^4}{km_e^3} \right)^{1/2} \frac{1}{\sqrt{T}} \frac{\Omega_{13}}{g_1} e^{-E_{13}/kT}, \quad (2.37)$$

(2.36) reduces to a function of a bunch of physical constants and the temperature

$$\boxed{\frac{F_{32}}{F_{21}} = \frac{E_{32}}{E_{21}} \frac{A_{32}}{A_{32} + A_{21}} \frac{\Omega_{13}}{\Omega_{12}} e^{-E_{23}/kT}} \quad (2.38)$$

So now given a line strength ratio, we can actually compute a temperature for that diffuse gas. Some examples of ions where these assumptions are valid are N^+ , O^{+2} , Ne^{+4} , and S^{+2} .

Now instead if we assume that $E_{23} \ll E_{12}$ (i.e., the second and third states are very close in energy). This means, though it's not obvious, that $A_{32} \ll A_{31}$. Still in the low density limit (ignoring collisional de-excitation), we'll simplify (2.27) and (2.28) again:

$$n_1 C_{13} = n_3 (A_{32} + A_{31}) \quad (2.39)$$

$$\approx n_3 A_{31} \quad (2.40)$$

$$n_1 C_{12} = n_2 A_{21} \quad (2.41)$$

Then the flux ratios are simply

$$\frac{F_{31}}{F_{21}} = \frac{n_1 C_{13}}{A_{31}} \frac{A_{21}}{n_1 C_{12}} \frac{E_{31}}{E_{21}} \frac{A_{31}}{A_{21}} \approx \frac{C_{13}}{C_{12}} \quad (2.42)$$

So the flux rates depend *only* on the collisional excitation rates in this limit. Re-expressing the flux ratio in terms of more fundamental quantities:

$$\frac{F_{31}}{F_{21}} \approx \frac{C_{13}}{C_{12}} \approx \frac{\Omega_{13} e^{-E_{13}/kT}}{\Omega_{12} e^{-E_{12}/kT}} \approx \frac{\Omega_{13}}{\Omega_{12}} \approx \frac{g_3}{g_2} \quad (2.43)$$

for somewhat unobvious reasons. So if we measure line strengths and get this magic number, all we really know is that we're in the low density limit. The magic comes in when we are in the high density limit:

$$\frac{F_{31}}{F_{21}} = \frac{n_3}{n_2} \frac{A_{31}}{A_{21}} \frac{E_{31}}{E_{21}} \quad (2.44)$$

Since in the high density limit there are many collisions, the number densities should be Boltzmann-like:

$$\frac{n_3}{n_2} = \frac{g_3}{g_2} e^{-E_{23}/kT} \approx \frac{g_3}{g_2} \quad (2.45)$$

So the flux ratio is now

$$\frac{F_{31}}{F_{21}} \approx \frac{A_{31}}{A_{21}} \frac{g_3}{g_2} \quad (2.46)$$

So far, this isn't really a density diagnostic, though we've identified the values at both limits. The transition between these limits should be at the critical density discussed earlier, defined by:

$$A_{21}n_2 = n_e n_2 \alpha_{21} \quad (2.47)$$

Earlier, we found that this critical density was found to be

$$n_{e,\text{crit}} = \frac{A_{21}g_2 T^{1/2}}{\beta \Omega_{12}} \quad (2.48)$$

Different transitions will have different critical densities that depend on their intrinsic physics, for electron densities below the smaller critical density, the flux ratio is constant, and the same holds for electron densities greater than the larger critical density. In between, though, there is a transition zone wherein we can approximate a density. Examples of ions that allow for such a diagnostic are O^+ , S^+ , Ne^{+3} , and Ar^{+3} .

2.3 Energy Level Diagrams

Suggested References

- Draine Chapter 4

Often we will refer to ions by their atomic name followed by a roman numeral that is one greater than its oxidation state. For instance, O^{+2} is referred to as O III. There will be some energy level diagram figures after this demonstrating where some of these energy level assumptions work.

Let's briefly review the electronic configurations of some of these ions. O I has 8 electrons, so in its ground state it has an electronic configuration of $1s^2 2s^2 2p^4$. O III however, has a configuration of $1s^2 2s^2 2p^2$, where the outermost electrons are actually in different orbitals (remember Hund's rules!). This p^2 configuration disallows the two outermost electrons from having the same quantum numbers (n, ℓ, m, s). In this case, we know that $n = 2$ (the "2" in $2p^2$), and $\ell = \pm 1$ (the p). For electrons, $s = \pm 1/2$. So the total spin can take on values of 0, or 1 and the total orbital angular momentum can be 0, 1, or 2. We'll often see this configurations labeled via $^{2S+1}L_J$ where S is the total spin angular momentum, L is the orbital name (s, p, d, f , etc.), and J is the total angular momentum. The available states for O II are shown in Table 3. Not all of these states are allowed transitions though, for reasons I still need to record here.

$^1D_{J=2}$	$L = 2$	$S = 0$
$^3D_{J=3,2,1}$	$L = 2$	$S = 1$
$^1P_{J=1}$	$L = 1$	$S = 0$
$^3P_{J=2,1,0}$	$L = 1$	$S = 1$
$^1S_{J=0}$	$L = 0$	$S = 0$
$^3S_{J=1}$	$L = 0$	$S = 1$

Table 3: Possible electron configurations for O II.

2.4 Infrared Line Diagnostics

Transitions between fine-structure levels of p^2 and p^4 ions are the dominant cooling processes in gas at $T = 100 - 3000$ K. Examples of such lines are [CII] $158 \mu\text{m}$ and [OIII] $88.36 \mu\text{m}$. Atmospheric water vapor blocks these lines in the 25 to $300 \mu\text{m}$ range, but airborne observatories can observe them, and ALMA is making great headway in observing mm-scale radiation.

However, note that in the real world (like in H II regions), the ISM is clumpy, so these diagnostics actually only measure the number density of the clumps, $n_{e,c}$. The emission measure, however, cares about the average electron density: $\text{EM} \sim \langle n_e^2 \rangle \times (\text{length})$. We parameterize this clumpiness with a **volume filling factor** for the clumps, f , defined by

$$\langle n_e^2 \rangle \approx f n_{e,c}^2 \quad (2.49)$$

Typically we find $f \sim 0.1 - 0.1$ in nearby H II regions.

2.5 Resonance Lines and Selection Rules

Remember that energy transitions are surrounded in square brackets when they are the so-called “forbidden transitions”. The more run-of-the-mill resonance transitions, or allowed transitions, are those that obey the following rules:

1. Only 1 electron involved in the transition
2. Initial and final states have different parity
3. Emitted photon carries 1 unit of angular momentum, so $\Delta L = \pm 1$
4. Electron spin does not change
5. Change in the total angular momentum of the active electron is $\Delta J = \pm 1$ or 0, with $J = 0 \rightarrow 0$ being forbidden.

The statistical weight of any level is $g = 2J + 1$.

In addition to the forbidden lines, which break rule 4, there are the intercombination, or **semi-forbidden** lines. These are transitions that break any of the other rules. These transitions are about a million times lower than the resonance transitions ($A \sim 10^2 - 10^3 \text{ s}^{-1}$ compared to $\sim 10^8 \text{ s}^{-1}$ for resonance lines). Forbidden lines are several orders of magnitude less prevalent still, as they are

magnetic dipole transitions.

For forbidden or semi-forbidden transitions to be important, we must have very low densities, since higher densities will make collisional processes dominate over these weak transitions. As such, the emitted photons from forbidden transitions are unlikely to be absorbed by another ion, so forbidden line photons usually escape from a nebula and are thus very important coolants. Examples of such transitions are [OIII] 4959, 5007 and [OII] 3726,29

2.5.1 Hydrogen Line Series

The line transitions in a de-exciting electron (for instance, a recombination electron that came in at high n , ℓ to form an excited hydrogen atom) often produce a series of lines. Transitions that end at $n = 1$ are called the **Lyman Series**. Transitions from $n = 2$ to $n = 1$ are Lyman- α transitions, and those from $n = 3$ to $n = 1$ are Lyman- β transitions, etc. Similarly, transitions to the $n = 2$ state form the Balmer series (denoted $H\alpha$, $H\beta$, etc.), and the transitions to $n = 3$ form the Paschen series.

3 Molecular Spectra

Suggested References

- Draine Chapter 5

Molecules present additional complication since rotational and vibrational energies become important. We'll first explore rotating molecules.

3.1 Rotating Molecules

Like nearly all forms of atomic energy, the energies of rotating molecules are quantized to energy levels. These rotational energy levels are related to the moments of inertia of the molecules along the various axes of symmetry.

Example: Linear Diatomic Molecules Suppose we have a linear diatomic molecule with characteristic lengthscale r and reduced mass μ . Then the moment of inertia of this molecule is

$$I = \mu r^2 \quad (3.1)$$

And the quantized rotational energy levels are

$$E = J(J+1) \frac{h^2}{8\pi^2 I} \quad (3.2)$$

3.2 Vibrating Molecules

3.3 Molecular Hydrogen

H_2 does not radiate strongly because rotational transitions require a heterogeneous linear molecule like CO. In the H_2 molecule, transitions occur via electric quadrupole interaction. The least energetic transition is $J = 0$ to $J = 2$. Thus, lifetimes of excited states are much, much longer than for

the ions (e.g. about 1000 years for the $J = 2$ level). Hence, the rotational elvels are populated by collisions.

4 Transition Rates

We return now to our two-level atom with two states, 1, and 2, separated by an energy gap E_{12} . We can characterize the rate of decays from state 2 through some characteristic e -folding time via

$$N_2 = N_2(0)e^{-A_{21}t} \quad (4.1)$$

Here, A_{21} is proportional to the electric dipole matrix element, or

$$A_{21} \propto e^2 |\langle \psi_1 | \psi_2 \rangle|^2 \quad (4.2)$$

An atom with an electron in state 1 has a probabliity of absorbing a photon given by

$$p_{12} = B_{12}u(\nu_{12}) \quad (4.3)$$

where B_{12} is the Einstein absorption coefficient and $u(\nu)$ is the energy density in radiation at a given frequency. We usually derive this energy density from the **specific intensity**, I_ν (I_λ), which has units of energy per area per time per sold angle per frequency (wavelength). The specific intensity is a conserved quantity along a ray of light. To convert this intensity to a flux through some surface, we must integrate this quantity (really, the portion that is normal to the surface) over all incident angles, or

$$F_\nu = \int I_\nu \cos \theta d\Omega \quad (4.4)$$

To get an energy density from the specific intensity (a process we'll discuss later), we use

$$u_\nu = \frac{4\pi}{c} I_\nu \quad (4.5)$$

Similarly, the probability for stimulated emission to occur is

$$p_{21} = B_{21}u(\nu_{12}) \quad (4.6)$$

where the two Einstein coefficients are related via

$$g_1 B_{12} = g_2 B_{21} \quad (4.7)$$

Using LTE, we can derive the spontaneous emission Einstein coefficient to be