

# Partially coherent ptychography

Tong Wu

September 28, 2021



$$f_j = |\mathcal{F}(\mathcal{S}_j u \circ \omega)|^2 \quad (1)$$

In a discrete setting,  $u \in \mathbb{C}^n$  is a 2D image with  $\sqrt{n} \times \sqrt{n}$  pixels,  $\omega \in \mathbb{C}^{\bar{m}}$  is a localized 2D probe with  $\sqrt{\bar{m}} \times \sqrt{\bar{m}}$  pixels.

$f_j \in \mathbb{R}_+^{\bar{m}} (\forall 0 \leq j \leq J-1)$  is a stack of phaseless measurements.

Here  $|\cdot|$  represents the element-wise absolute value of a vector,  $\circ$  denotes the elementwise multiplication, and  $\mathcal{F}$  denotes the normalized 2 dimensional discrete Fourier transform. Each  $\mathcal{S}_j \in \mathbb{R}^{\bar{m} \times n}$  is a binary matrix that crops a region  $j$  of size  $\bar{m}$  from the image  $u$ .

# Partially coherent model: a general one

Thibault, Pierre, and Andreas Menzel. "Reconstructing state mixtures from diffraction measurements." Nature 494.7435 (2013): 68-71.

Blind ptychography model + quantum state tomography<sup>1</sup>.

Probe  $w$  is assumed to be in mixed state to represent partially coherent effect.

Find  $u, r$  orthogonal  $w_k$  s.t.

$$f_{pc,j} = \sum_{k=1}^r |\mathcal{F}(\mathcal{S}_j u \circ (\omega_k))|^2 \quad (2)$$

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<sup>1</sup>[https://homepage.univie.ac.at/reinhold.bertlmann/pdfs/T2\\_Skript\\_Ch\\_9corr.pdf](https://homepage.univie.ac.at/reinhold.bertlmann/pdfs/T2_Skript_Ch_9corr.pdf) Many symbols in quantum mechanics are included here. 🔍🔍🔍

## Step1: Fourier magnitude projection

We have used the shorthand notation  $\tilde{\phi} = \mathcal{F}[\phi]$ . Taking the form

$$\begin{aligned}\psi_{j\mathbf{x}}^{(k)} &= \bar{\Pi}_F \left( \phi_{j\mathbf{x}}^{(0)}, \phi_{j\mathbf{x}}^{(1)}, \dots, \phi_{j\mathbf{x}}^{(r=M_p)} \right) \\ &= \mathcal{F}^{-1} \left[ \frac{\sqrt{I_{j\mathbf{q}}} \tilde{\phi}_{j\mathbf{q}}^{(k)}}{\sqrt{\sum_k |\tilde{\phi}_{j\mathbf{q}}^{(k)}|^2}} \right]\end{aligned}$$

again with

$$\tilde{\phi}_{j\mathbf{q}}^{(k)} = \mathcal{F} \left[ P_{\mathbf{x}}^{(k)} O_{\mathbf{x}-\mathbf{x}_j} \right]$$

## Step2: Overlap projection

$$O_{\mathbf{x}} = \frac{\sum_k \sum_j P_{\mathbf{x}+\mathbf{x}_j}^{(k)*} \psi_{j,\mathbf{x}+\mathbf{x}_j}^{(k)}}{\sum_k \sum_j \left| P_{\mathbf{x}+\mathbf{x}_j}^{(k)} \right|^2}$$
$$P_{\mathbf{x}}^{(k)} = \frac{\sum_j O_{\mathbf{x}-\mathbf{x}_j}^* \psi_{j\mathbf{x}}^{(k)}}{\sum_j \left| O_{\mathbf{x}-\mathbf{x}_j} \right|^2}$$

solved numerically by applying them sequentially for a few iterations.

# Outcome

Denote  $O_j \in \mathbb{C}^{\bar{m} \times \bar{m}}$  as a (diagonal) matrix to represent linear transform to  $w$ , s.t.  $\mathcal{S}_j u \circ \omega = O_j w$ . Denote  $f_q^* \in \mathbb{C}^{1 \times \bar{m}}$  as a row vector constructed from Fourier transform  $\mathcal{F}$ , to represent projection on frequency element. Construct measurement matrix  $\mathcal{I}_{j\mathbf{q}} = O_j^* f_q f_q^* O_j$  and density matrix  $\rho$ , we get another form (actually a natural one in quantum state tomography) of the model:

Find  $O_j, \rho$ , s.t.

$$f_{pc,j}(q) = \text{Tr}(\mathcal{I}_{j\mathbf{q}} \rho) \quad (3)$$

$\rho$  is positive semi-definite, with  $\text{rank} \leq r$



# Derivation

Simple calculation process:

$$\begin{aligned} f_{pc,j}(q) &= |f_q^* O_j w|^2 = (f_q^* O_j w)^* (f_q^* O_j w) = w^* (O_j^* f_q f_q^* O_j) w \\ &= \text{Tr}[w^* (O_j^* f_q f_q^* O_j) w] = \text{Tr}[(O_j^* f_q f_q^* O_j)(w w^*)] \\ &= \text{Tr}(\mathcal{I}_{j\mathbf{q}} \rho) \end{aligned}$$

When  $w$  is in pure state (a vector in Hilbert space),  $\rho = w w^*$  is a rank-one matrix. In partially coherent case, **we use mixed state to model  $w$** . Mixed state is represented by **generalizing the density matrix to one with higher rank**:

$$\rho = \sum_{k=1}^r w_k w_k^*$$

$$f_{pc,j}(q) = \text{Tr} \mathcal{I}_{j\mathbf{q}} \rho = \text{Tr}[\mathcal{I}_{j\mathbf{q}} \sum_{k=1}^r w_k w_k^*] = \sum_{k=1}^r w_k^* \mathcal{I}_{j\mathbf{q}} w_k = \sum_{k=1}^r |f_q^* O_j w_k|^2$$

And that is exactly (2)  $f_{pc,j} = \sum_{k=1}^r |\mathcal{F}(\mathcal{S}_j u \circ (\omega_k))|^2$ .

# Further explanation about mixed state

For example, with probability 0.5 in state  $\psi_1$  and 0.5 in  $\psi_2$  ( $\psi_1$  and  $\psi_2$  are not necessarily orthogonal here). Now  $w$  can no longer be represented by a vector (ps.  $w \neq p_1\psi_1 + p_2\psi_2$ , the latter is still a determined pure state vector).

$$\rho = \sum_k p_k \psi_k \psi_k^*$$

Easy to find  $\rho$  is a positive semi-definite matrix, we can decompose  $\rho$  using spectral theorem, with  $r$  (rank of  $\rho$ ) orthogonal state  $w_k$ :

$$\rho = \sum_{k=1}^r w_k w_k^*$$

## Related work

Fannjiang, Albert, and Thomas Strohmer. "The numerics of phase retrieval." *Acta Numerica* 29 (2020): 125-228. 6.3. Low-rank phase retrieval problems.

The phase retrieval problem has a natural generalization to recovering low-rank positive semidefinite matrices. More precisely, we are concerned with the task of reconstructing a finite-dimensional quantum mechanical system which is fully characterized by its density operator  $\rho$  (an  $n \times n$  positive semidefinite matrix with trace one.)