

MNM Study Guide Coursework 2

Notation:

Scalars – a or a

Vectors – \mathbf{v} or $\vec{\mathbf{v}}$ or $\bar{\mathbf{v}}$, vector length $|\mathbf{v}|$

Tensors – \mathbf{T} or (if rank 2) $\underline{\mathbf{T}}$

Unit vector along direction of \mathbf{v} : $\hat{\mathbf{e}}_v = \frac{\mathbf{v}}{|\mathbf{v}|}$

Unit outward normal for a plane: $\hat{\mathbf{n}}$

Equations/concepts you are expected to know and be able to apply:

Examples given here all for 3-D, orthonormal Cartesian reference frame

- Index notation: vector or tensor components written as v_i or T_{ij} with $i,j=1,2,3$ or $i,j=x,y,z$
- Einstein convention – implied summation of the same index repeated twice within a single term, e.g. $v_i w_i = \sum_{i=1}^3 v_i w_i$
- Vector and tensor products:
 - dot product: $\mathbf{v} \cdot \mathbf{w} = v_i w_i$ or $\mathbf{T} \cdot \mathbf{v} = T_{ij} v_j$
 - multiple contraction, e.g. $\boldsymbol{\sigma} = \mathbf{C} : \boldsymbol{\varepsilon} = C_{ijkl} \varepsilon_{kl}$
 - cross product: $\mathbf{v} \times \mathbf{w} = \varepsilon_{ijk} v_i w_j \hat{\mathbf{e}}_k$
 - tensor product: $\mathbf{vw} = v_i w_j$
- Transpose: $T_{ji} = T_{ij}^T$
- Tensor symmetry:
 - Symmetric in i,j : $T_{ji} = T_{ij}$,
 - Antisymmetric in i,j : $T_{ji} = -T_{ij}$
- Tensor trace: for rank 2 tensor $\text{tr}(\mathbf{T}) = T_{11} + T_{22} + T_{33} = T_{ii}$.
- Kronecker delta $\delta_{ij} = 1$ if $i=j$, $=0$ if $i \neq j$
- Levi-Civita tensor $\varepsilon_{ijk} = 1$ for even permutations of 1,2,3, $\varepsilon_{ijk} = -1$ for odd permutations of 1,2,3, $\varepsilon_{ijk} = 0$ if any i,j,k are equal
- Lagrangian or material – description of motion following a ‘particle’, all fields described as a function of position $\boldsymbol{\xi}$ at a reference time t_0 and time t .
- Eulerian or spatial – description of motion from a fixed observation point. All fields described as a function of position \mathbf{x} and time t .
- Material Derivative
 - in spatial description, the full time derivative of a field $P(\mathbf{x},t)$ becomes: $\frac{DP}{Dt} = \frac{\partial P}{\partial t} + \mathbf{v} \cdot \nabla P$, i.e., contains a time and advective term
 - in material description, the time derivative of the field $P(\boldsymbol{\xi},t)$ is $\frac{DP}{Dt} = \frac{\partial P}{\partial t}$

- Divergence: $\nabla \cdot \mathbf{v} = \frac{\partial v_1}{\partial x_1} + \frac{\partial v_2}{\partial x_2} + \frac{\partial v_3}{\partial x_3}$
 - represents source/sink of a \mathbf{v} field
 - can also be applied to tensors, e.g. $(\nabla \cdot \mathbf{T})_i = \frac{\partial T_{1i}}{\partial x_1} + \frac{\partial T_{2i}}{\partial x_2} + \frac{\partial T_{3i}}{\partial x_3}$
- Curl: $\nabla \times \mathbf{v} = \left(\frac{\partial v_3}{\partial x_2} - \frac{\partial v_2}{\partial x_3}, \frac{\partial v_1}{\partial x_3} - \frac{\partial v_3}{\partial x_1}, \frac{\partial v_2}{\partial x_1} - \frac{\partial v_1}{\partial x_2} \right)$
 - represents vorticity of a \mathbf{v} field
- Gradient:
 - of a scalar $\nabla f = \left(\frac{\partial f}{\partial x_1}, \frac{\partial f}{\partial x_2}, \frac{\partial f}{\partial x_3} \right)$
 - or of a vector $(\nabla \mathbf{v})_{ij} = \frac{\partial v_j}{\partial x_i}$
- Laplacian: $\nabla \cdot \nabla f = \nabla^2 f = \Delta f = \frac{\partial^2 f}{\partial x_i \partial x_i}$
- Cauchy Stress tensor:
 - stress tensor component σ_{ij} represents a force in $\hat{\mathbf{e}}_j$ direction on a plane with normal in $\hat{\mathbf{e}}_i$ direction. Positive normal stress corresponds to extension.
 - stress tensor is symmetric: $\sigma_{ij} = \sigma_{ji}$ (conservation of angular momentum)
 - traction \mathbf{t} on a plane with normal $\hat{\mathbf{n}}$ is $\mathbf{t} = \boldsymbol{\sigma} \cdot \hat{\mathbf{n}}$
 - the stress tensor can be diagonalised, with principal components $\sigma_1, \sigma_2, \sigma_3$ which include maximum and minimum normal stress
 - Can be decomposed into isotropic stress (pressure $p = -\sigma_{kk}/3$) and deviatoric stress $\boldsymbol{\sigma}'$ such that $\sigma_{ij} = -p\delta_{ij} + \sigma'_{ij}$
- Conservation of linear momentum (per unit volume): $\rho \frac{D^2 \mathbf{u}}{Dt^2} = \mathbf{f} + \nabla \cdot \boldsymbol{\sigma}$, where ρ is density, \mathbf{u} is displacement and \mathbf{f} is body force.
- Infinitesimal strain tensor:
 - Infinitesimal strain tensor component $\varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_j}{\partial x_i} + \frac{\partial u_i}{\partial x_j} \right)$, where \mathbf{u} is the displacement field. Applicable if $\nabla \mathbf{u} \ll 1$.
 - An original line segment described by vector \mathbf{x} deforms to a new line segment \mathbf{x}' as: $\mathbf{x}' = \boldsymbol{\varepsilon} \cdot \mathbf{x}$
 - Diagonal components of ε_{ij} represent fractional length changes, i.e., if \mathbf{x} is a vector in $\hat{\mathbf{e}}_1$ direction then $\varepsilon_{11} = \frac{|\mathbf{x}'| - |\mathbf{x}|}{|\mathbf{x}|}$. Similarly, for a given vector \mathbf{s} , the product $\mathbf{s} \cdot \boldsymbol{\varepsilon} \cdot \mathbf{s}$ corresponds to the fractional change in $|\mathbf{s}|$ by the strain $\boldsymbol{\varepsilon}$.
 - Off-diagonal components represent changes in angle (i.e., shape), such that $2\varepsilon_{12}$ equals the change in angle between a line segment originally in $\hat{\mathbf{e}}_1$ direction and one originally in $\hat{\mathbf{e}}_2$ direction. Given two originally perpendicular vectors \mathbf{s} and \mathbf{p} , $2 \times$ the product $\mathbf{p} \cdot \boldsymbol{\varepsilon} \cdot \mathbf{s}$ corresponds to the change in the angle between \mathbf{s} and \mathbf{p} by the strain $\boldsymbol{\varepsilon}$.
 - $tr(\boldsymbol{\varepsilon}) = \nabla \cdot \mathbf{u}$ and represents the fractional change in volume.
 - ε_{ij} is symmetric and can be diagonalised, such that principal strain components $\varepsilon_1, \varepsilon_2, \varepsilon_3$ include the maximum and minimum fractional length changes in the strain field described by $\boldsymbol{\varepsilon}$.

- Can be decomposed into isotropic and deviatoric strain, like the stress tensor
- Strain rate tensor
 - Strain rate tensor $\mathbf{D} = \mathbf{D}\mathbf{\epsilon}/Dt$ has same kind of properties as the infinitesimal strain tensor, but depends on the velocity field \mathbf{v} : $D_{ij} = \frac{1}{2} \left(\frac{\partial v_j}{\partial x_i} + \frac{\partial v_i}{\partial x_j} \right)$.
 - $\text{tr}(\mathbf{D}) = \nabla \cdot \mathbf{v} = 0$ means no change in volume and is the conservation of mass equation for an incompressible material.
- Energy equation – if given the equation, understand the different terms (conduction, advection, heat production, power dissipated by deformation) and be able to use (e.g., to solve for temperature for simple case)
- Rheology – know difference between elastic and viscous rheology. Be able to apply. If more complex equations are necessary, they will be given.
 - Elasticity – $\boldsymbol{\sigma} = \mathbf{C}:\boldsymbol{\epsilon}$, linear relationship between stress and infinitesimal strain. For an isotropic medium, only two independent parameters, e.g. Lamé parameters (λ, μ) , bulk and shear moduli (K, μ) , or Young's modulus and Poisson's ratio (E, ν) . In terms of Lamé parameters: $\boldsymbol{\sigma} = \lambda \theta \mathbf{I} + 2\mu \boldsymbol{\epsilon}$, where $\theta = \text{tr}(\boldsymbol{\epsilon})$. Bulk modulus: $-p = K\theta$, $K = \lambda + \frac{2}{3}\mu$; In uniaxial stress: Young's modulus $E = \sigma_{11}/\epsilon_{11}$, and Poisson's ratio $\nu = -\epsilon_{33}/\epsilon_{11}$
 - Newtonian Viscosity – linear relationship between deviatoric stress $\boldsymbol{\sigma}'$ and strain rate \mathbf{D} . If isotropic \Rightarrow bulk viscosity ζ and shear viscosity η as the two material parameters: $\boldsymbol{\sigma} = (-p + \zeta \Delta) \mathbf{I} + 2\eta \mathbf{D}$
- Equations of motion – wave equation for elastic media and Navier Stokes for fluids, understand terms and derive simple solutions if equations given.
- Dimensional analysis
 - For non-dimensionalising equations, understand how to produce dimensionless versions of the dependent and independent variables and to use those to form the dimensionless the dimensionless version of the equations.
 - Understand how to use Buckingham Pi theory to determine the number of dimensionless groups required to describe a system and to produce an appropriate set of dimensionless groups