

$$u = u(x, t), D > 0$$

(passive trans.)

$$u$$

$$\forall A, B \in \mathbb{R}$$

(i) Heat Eq. (Diffusing Particles) $U_t = DU_{xx}$; (ii) Wave Eq. (wave prop) $U_{tt} = c^2 U_{xx}$ (iii) Transport Eq. $U_t + cU_x = 0$; $u(x, y)$. $U_x = 0 \Rightarrow u(x, y) = f(y)$; $y'' + y = 0 \Rightarrow y = A \cos(ky) + B \sin(ky)$

$\therefore U_x + U_t = 0 \Rightarrow u(x, y) = A \cos(kx) + B \sin(kx)$ Order: Highest derivative, Transformation of linear if $c^2 f(u+v) = f(u)$, $f(cw) = c f(w)$; PDE Linear if $f(u) = f$, f : Lin. tran. f.ind. of u . non-lin.

$U_t - DU_{xx} = 0 : f(u) = U_t - DU_{xx}, f = 0, f(cu+v) = f(cu), f(cw) = cf(w)$; (ii) $U_t + DU_{xx} + DU_{xx} = 0 : f(cw) \neq c f(w)$, (iii) $U_{tt} + c^2 U_{xx} = 0 : \text{Linear}$; $c^2 U_{xx} + bU_{xy} + cU_{yy} + dU_{xz} + eU_{yz} + fU = g$, $D = b^2 - 4ac$.

(i) D₀: Elliptic (ii) D>0: Hyperbolic (iii) D=0: parabolic. $U_t = U_x : U(x, 0) = u_0(x)$ LHS: $bU_0(x+b\tau) = U_0(x+b\tau)$, RHS: $\frac{\partial U_0(x, \tau)}{\partial x} = \frac{\partial U_0(x+b\tau)}{\partial x} = U_0(x+b\tau)$. LHS: $\frac{\partial U(x, \tau)}{\partial t} = \frac{\partial U_0(x+b\tau)}{\partial t} = U_0(x+b\tau)$.

(Gradient): (Directional D. of u in direction \vec{v}) $u = u(x, y)$, $\nabla u = (U_x, U_y)$, if $\|\vec{v}\| = 1 \Rightarrow \nabla u \cdot \vec{v} = U_x \cdot \vec{v} + U_y \cdot \vec{v} = 0 \Rightarrow u \text{ constant in direction } \vec{v} \cdot \vec{e}_1 < 0, \vec{e}_2 > 0$ R.P. $\vec{v} \rightarrow \vec{v}$

Heat: $U_t = DU_{xx}$. Let $U(x, t) = u(x, t)h + U_{xx}(x, t)h^2 + \dots$, $U(x, t+\tau) = u(x, t)h + U_{xx}(x, t)h^2 + \dots \therefore U(x, t+\tau) - U(x, t) = u(x, t) + O(h^2); U(x-h, t) + U(x+h, t) - 2u(x, t) = U_{xx}(x, t)h^2 + O(h^3)$

$\frac{h^2}{2\tau}, \text{ choose } \tau = 2D \therefore \frac{h^2}{2\tau} = D$. Take $h \rightarrow 0 \therefore U_t = DU_{xx}$ Euler F. $e^{i\varphi} = \cos(\varphi) + i\sin(\varphi) \therefore e^{ikx} = \cos(kx) + i\sin(kx)$. Fourier Transform: $\hat{u}(k)$ of $u(x)$

$(F_u)(k) = \hat{u}(k) = \int_{-\infty}^{\infty} e^{ikx} u(x) dx$. Let $U(x) = e$ $\therefore \hat{U}(k) = \int_{-\infty}^{\infty} e^{ikx} e^{-Dk^2x} dx = \frac{1}{2} \int_{-\infty}^{\infty} e^{ikx - Dk^2x} dx = \int_{-\infty}^{\infty} e^{ikx - Dk^2x} dx = \frac{i}{2\pi} \int_{-\infty}^{\infty} e^{ikx - Dk^2x} dk = \frac{i}{2\pi} \int_{-\infty}^{\infty} e^{ikx - Dk^2x} dk$

$= \frac{i}{2\pi} \int_{-\infty}^{\infty} e^{ikx - Dk^2x} dk = \frac{i}{2\pi} \int_{-\infty}^{\infty} e^{ikx - \frac{D}{4} - \frac{D}{4}} dk = \frac{i}{2\pi} e^{-\frac{D}{4}} \int_{-\infty}^{\infty} e^{ikx - \frac{D}{4}} dk = \frac{i}{2\pi} e^{-\frac{D}{4}} \int_{-\infty}^{\infty} e^{ikx} dk = \frac{i}{2\pi} e^{-\frac{D}{4}}$ (5) Schurz Class

$u(x) \in \mathcal{S} \forall n, m \in \mathbb{N} \exists U^{(n)}(x), |x| \rightarrow \infty \text{ I.F.C. } u(x) \in \mathcal{S} \text{ and } u(x) = (\hat{F}^{-1} \hat{u})(x) = 2\pi \int_{-\infty}^{\infty} e^{-ikx} \hat{u}(k) dk$. IBP: $\int_{-\infty}^{\infty} UV' dx = - \int_{-\infty}^{\infty} U' V dx + U(-\infty)V(-\infty) = 0$. $U_t = DU_{xx}, U(x, 0) = f(x)$

$\Rightarrow \hat{U}_t(k, t) = \hat{U}(k, 0)e^{-Dk^2t} = \hat{f}(k)e^{-Dk^2t} \Rightarrow U(x, t) = \int_{-\infty}^{\infty} e^{ikx} e^{-Dk^2t} dk$. Convolution: $u(x), v(x) \in \mathcal{S}, (u \ast v)(x) := \int_{-\infty}^{\infty} u(x-y)v(y) dy$; $\hat{(u \ast v)}(k) = (\hat{u} \ast \hat{v})(k) \iff$

$[\hat{(u \ast v)}(k)](x) = \hat{u}(k) \hat{v}(k) \therefore u(x, t) = \hat{F}^{-1}(e^{-Dt}) \hat{f}(k)x, v(x, t) = \hat{F}^{-1}(e^{-Dt}) \hat{g}(k)x$. Recall: $\hat{F}(e^{-\alpha x^2}) = \sqrt{\frac{\pi}{\alpha}} e^{-x^2/\alpha} \Rightarrow \hat{F}(e^{-k^2/4a}) = \sqrt{\frac{\pi}{4a}} e^{-k^2/4a}$, need: $\frac{1}{4a} = Dt \Rightarrow a = \frac{1}{4Dt}$

$\Rightarrow G(x, t) = \frac{\sqrt{\pi}}{4\sqrt{Dt}} e^{-x^2/(4Dt)} \Rightarrow u(x, t) = (G \ast f)(x, t), v(x, t) = (G \ast g)(x, t)$. Sol. Heat Eq. $\frac{\partial u}{\partial t} = \frac{1}{4\sqrt{Dt}} \int_{-\infty}^{\infty} e^{-x^2/(4Dt)} f(y) dy$. Derivatives: $(\hat{u}(k))_{(k)} = -ik\hat{u}(k), (\hat{u}_t(k))_{(k)} = \frac{1}{2} \hat{u}(k)$

PDE in $(x, t) \xrightarrow{\text{PDE in } (k, t)}$ Prop: $u(x, t), v(x, t)$ Conserv. of mass: $\int u(x, t) dx = \int v(x, t) dx$, (i) $\exists! u(x, t),$ (ii) Bounded. Let $M = \max|f(x)| \Rightarrow |u(x, t)| \leq M$ Infinite S. of. Prop. $\forall x \in \mathbb{R}, t > 0$

Sol. $u(x, t) \xrightarrow{\text{Sol. } \hat{u}(k, t)}$ (i) Decay: $|u(x, t)| \leq \sqrt{\frac{1}{4\pi Dt}} \int_{-\infty}^{\infty} |f(y)| dy \therefore |u(x, t)| \rightarrow 0 \text{ uniformly}; (V) If } \exists a, b \in \mathbb{R} \text{ s.t. } f(x) > 0 \text{ in } (a, b), f(x) = 0 \text{ elsewhere} \Rightarrow u(x, t) \rightarrow 0 \forall x \in \mathbb{R}, t > 0$

Proof: (i) $\int_a^b u(x, t) dx = D \int_a^b u_0(x, t) dx = -D \int_a^b x(1) U_0(x, t) dx = 0 \therefore \int_a^b U_0(x, t) dx = 0 \Rightarrow \int_a^b u(x, t) dx = \int_a^b u_0(x, 0) dx = \int_a^b f(x) dx$ (ii) Let

$u(x, t), u_0(x, t) \text{ sol. } V(x, t) := u_0(x, t) - u_0(x, 0) \therefore V \text{ sol. } \left\{ \begin{array}{l} \int_a^b V_0(x, t) dx = D \int_a^b V_0(x, 0) dx = D \int_a^b f(x) dx \\ \int_a^b V(x, t) dx = D \int_a^b V(x, 0) dx = D \int_a^b V_0(x, 0) dx = D \int_a^b f(x) dx \end{array} \right. \text{ LHS} = \frac{1}{2} \int_a^b \int_a^b V(x, t)^2 dx, \text{ RHS} = -D \int_a^b \int_a^b V(x, t)^2 dx \leq 0$

$\therefore \int_a^b V(x, t)^2 dx \text{ non-inc.} \Rightarrow 0 \leq \int_a^b V(x, t)^2 dx \leq \int_a^b V(x, 0)^2 dx \Rightarrow V(x, t) = 0 \forall x \in \mathbb{R}$. Energy Method: Show $E(t) := \int_a^b V(x, t)^2 dx$ is not decreasing in time. (iii) Given (ii) $\exists! U_0(x, t) = \frac{1}{\sqrt{4\pi Dt}} \int_{-\infty}^{\infty} e^{-x^2/(4Dt)} f(y) dy, g = y-x \therefore |U(x, t)| \leq M \frac{1}{\sqrt{4\pi Dt}} \int_{-\infty}^{\infty} e^{-y^2/(4Dt)} dy, |V(x, t)| \leq M \frac{1}{\sqrt{4\pi Dt}} \int_{-\infty}^{\infty} e^{-y^2/(4Dt)} dy$

$\frac{1}{4\pi Dt} \int_{-\infty}^{\infty} e^{-x^2/(4Dt)} f(y) dy > 0$. Wave Eq. $U_{tt} = c^2 U_{xx}$, $U(x, t)$ displacement of a vib. at (x, t) . Damped Wave Eq. $U_{tt} = c^2 U_{xx} - k^2 U$. $\frac{\partial^2}{\partial t^2} U - c^2 \frac{\partial^2}{\partial x^2} U = (\frac{\partial^2}{\partial t^2} - c^2 \frac{\partial^2}{\partial x^2}) U$, change var.

$\left\{ \begin{array}{l} x = x + \frac{1}{2}kt \\ t = \frac{x}{2c} \end{array} \right. \therefore \frac{\partial}{\partial x} = \frac{1}{2} \frac{\partial}{\partial t} + \frac{1}{2c} \frac{\partial}{\partial t}, \frac{\partial}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} - \frac{1}{2c} \frac{\partial}{\partial x}$. Chain Rule: $\frac{\partial}{\partial x} = \frac{\partial}{\partial x} + \frac{\partial}{\partial t}$. Let $V(\frac{x}{2}, \frac{t}{2c}) = U(x, t) = U(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c})$. C.R. $\frac{\partial}{\partial x} = U_x \frac{\partial}{\partial x} + U_t \frac{\partial}{\partial t}$

$= \frac{1}{2} U_x(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c}) + \frac{1}{2c} U_t(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c}) \therefore V_x(\frac{x}{2}, \frac{t}{2c}) = \frac{1}{2} U_{xx}(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c}) + \frac{1}{2c} U_{xt}(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c}) + \frac{1}{2} U_{tt}(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c}) \frac{\partial t}{\partial x} = 0 \Rightarrow$

$V_{tt} = \frac{1}{4} U_{ttt} - \frac{1}{4c} U_{ttx} + \frac{1}{4c} U_{txx} - \frac{1}{4c^2} U_{ttt} = \frac{1}{4c^2} (c^2 U_{xx}(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c}) - U_{tt}(\frac{x}{2} + \frac{1}{2}kt, \frac{t}{2c})) = 0 \therefore V_{tt} = 0 \Rightarrow V_t = 0 \Rightarrow V_t = f(\frac{x}{2}) + G(\frac{t}{2}) = F(\frac{x}{2} + \frac{1}{2}kt) + G(\frac{t}{2})$

$\therefore u(x, t) = V(\frac{x}{2}, \frac{t}{2c}) = F(x + ct) + G(x - ct)$ AF, G, gen. sol. wave Eq. Wave-like Eq: $U_{tt} + AU_{xx} + BU_{xx} = 0$ If $\exists a, b \in \mathbb{R}$ st $a \neq b$ and $(\frac{\partial^2}{\partial t^2} - a^2 \frac{\partial^2}{\partial x^2}) U = 0$

\Rightarrow gen. sol. is: $F(x+at) + G(x+bt)$. $\left\{ \begin{array}{l} U(x, 0) = f(x) \\ U_t(x, 0) = g(x) \end{array} \right. \text{ init. position}$ $\therefore U(x, 0) = F(x) + G(x), U_t(x, 0) = C_F(x+ct) - C_G(x-ct) \Rightarrow$

$U_t(x, 0) = C_F(x) + C_G(x) = g(x) \therefore f(x) = C_F(x), g(x) = C_G(x) \therefore F(x) = \frac{1}{2} \int_{-\infty}^x g(s) ds, G(x) = \frac{1}{2} \int_x^{\infty} g(s) ds \therefore F(x) - G(x) = \frac{1}{2} \int_x^{\infty} g(s) ds = \frac{1}{2} \int_x^{\infty} g(s) ds$

$F(x) = \frac{1}{2} f(x) + \frac{1}{2c} \int_{-\infty}^x g(s) ds + \frac{A}{2}, G(x) = \frac{1}{2} f(x) - \frac{1}{2c} \int_x^{\infty} g(s) ds - \frac{A}{2} \Rightarrow U(x, t) = F(x+ct) + G(x-ct)$. Sel. $U(x, t) = \frac{1}{2} f(x+ct) + \frac{1}{2c} \int_{-\infty}^x g(s) ds + \frac{1}{2} f(x-ct) + \frac{1}{2c} \int_x^{\infty} g(s) ds$ (D'Alembert's Formula)

If init veloci. $g(x) = 0$, $f(x)$ splitted evenly into 2: $\left\{ \begin{array}{l} \leftarrow, \rightarrow, U(x, t) = \frac{1}{2} f(x+ct) - \frac{1}{2} f(x-ct) \\ \text{K.E.} \quad \text{P.E.} \end{array} \right.$

Conservation of Energy: Define $E(t) = \frac{1}{2} \int_{-\infty}^{\infty} (U(x, t))^2 dx + \frac{C}{2} \int_{-\infty}^{\infty} (U_t(x, t))^2 dx + \frac{C}{2} \int_{-\infty}^{\infty} (U_{xx}(x, t))^2 dx$. If $U(x, t) \in \mathcal{S} \Rightarrow E(t) = E(0) \forall t \in \mathbb{R}$. Proof: $\left\{ \begin{array}{l} (-, U_t) \\ (+, U_{xx}) \end{array} \right. \therefore \int_{-\infty}^{\infty} U_{tt} U_t dx + C \int_{-\infty}^{\infty} U_{xx} U_{xx} dx = 0$

LHS: $\int_{-\infty}^{\infty} \frac{2}{\partial t^2} \left(\frac{1}{2} (U_t)^2 \right) dx + C \int_{-\infty}^{\infty} (U_{xx})^2 dx = \frac{d}{dt} \left(\frac{1}{2} \int_{-\infty}^{\infty} (U_t)^2 dx + \frac{C}{2} \int_{-\infty}^{\infty} (U_{xx})^2 dx \right) = \frac{d}{dt} (E(t)) = 0 \Rightarrow E(t) = E(0)$. Remark: For energy method for (i) Heat Eq. (ii) u

(ii) Wave Eq. $(-, U_t)$. Theorem: $f: a, b \in \mathbb{R}$. If $g(x) = f(x-a) \Rightarrow g(k) = e^{ika} f(k)$

(i) Sketch characteristic lines, solve Transport Eq. $\left\{ \begin{array}{l} U(x, 0) = x^3 \\ U_t + 2U_x = 0 \end{array} \right. \Leftrightarrow (U_x, U_t) \cdot \langle 2, 1 \rangle = 0 \therefore \text{Characteristic lines II} \langle 2, 1 \rangle$

$\therefore U(x, t) = U(x-2t, 0) = (x-2t)^3 \therefore U(x, 0) = x^3, U(x, 2) = (x-2)^3, U(x, 4) = (x-4)^3$

(i) Derive transp. Eq. Assume $u(x, t) : \#P \mid_{x, t}$. Let $\alpha, \beta \in \mathbb{R}$. $\left\{ \begin{array}{l} x = x - 2t \\ t = t \end{array} \right. \therefore x = x - 2t, x = x - 2t, x = x - 2t, \dots$ Assume $\forall P \in [t, t+\tau]$

$\tau \rightarrow \text{nearest interval to the right. (a) Balance Eq. } \left\{ \begin{array}{l} \leftarrow, \rightarrow, U(x, t) = h(x, t) + h(x-2t, t) - h(x, t) \\ \text{P.E.} \end{array} \right. \therefore h(x, t+\tau) = h(x, t) + h(x-2t, t) - h(x, t) + \dots$

1^o: $\# p \in [x-\frac{h}{2}, x+\frac{h}{2}] \Big|_{t+\tau}, 2^{\circ}: \# p \in [x-\frac{3h}{2}, x+\frac{h}{2}] \Big|_t, 3^{\circ}: \# p \in [x-\frac{3h}{2}, x-\frac{h}{2}], 4^{\circ}: \# p \text{ going out}$. (b) Rearrange Terms and use Taylor Expansions:
 $\frac{\partial u(x, t+\tau)}{\tau} = \frac{h}{\tau} \left(\frac{u(x-h, t) - u(x, t)}{h} \right)$. Taylor Eq: $u(x, t+\tau) = u(x, t) + u_t(x, t)\tau + O(\tau^2); u(x-h, t) = u(x, t) - u_x(x, t)\tau + O(h^2) \Rightarrow \text{plug into } \frac{\partial u(x, t+\tau)}{\tau}$
 $\frac{\partial u_t(x, t)}{\tau} + O(\tau) = \frac{h}{\tau} (-U_x(x, t) + O(h))$ (Choose h, τ s.t. $h, \tau \rightarrow 0$ and $\frac{\partial}{\partial t} \rightarrow \text{Transp. Eq: } \tau = c, h \rightarrow 0 \therefore u_t(x, t) = -c u_x(x, t) \iff u_t + cu_x = 0$
 (i) If $g(x) = f(x-a) \Rightarrow \hat{g}(k) = e^{ik(a-t)} \hat{f}(k)$. Use F and theorem to solve Transp. $\begin{cases} u(x, 0) = g(x) \\ u_t - ik\hat{u} = 0, \hat{u}(k, 0) = \hat{f}(k) \end{cases}$
 $\therefore \hat{u}(k, t) = \hat{u}(k, 0) e^{ikt} = \hat{f}(k) e^{ikt} \Rightarrow u(x, t) = F^{-1}(\hat{f}(k) e^{ikt}) = f(x-ct)$. (ii) Consider $u_c = u_{xxx}$ for $x \in \mathbb{R}$ with $u(x, 0) = f(x) \in \mathcal{D}$.
 Derive an expression for F.T. $\hat{u}(k, t)$ of the solution $u(x, t)$ in terms of k, t and $f(k)$: $\hat{u}_t = \hat{u}_x = \widehat{u_{xxx}} = (-ik) \widehat{u_{xx}} = -ik^2 \widehat{u_{xx}} = -ik^3 \widehat{u_x} = (-ik^3) \hat{u} = k^3 \hat{u}$
 $\therefore \hat{u}(k, t) = \hat{u}(k, 0) e^{ikt} = \hat{f}(k) e^{ikt}$ (viii) For $D > 0$, $a \in \mathbb{R}$, consider $u_c = Du_{xx} + cu_x - au$ for $x \in \mathbb{R}, t > 0$. Assume $u(x, t)$ sol. Verify
 $v(y, t) := u(y-ct, t) e^{at}$ sol. to $V_t = DV_{yy}$. $\therefore V_t = \frac{\partial}{\partial t} (u(y-ct, t) e^{at}) = -cu_x(y-ct, t) e^{at} + u_t(y-ct, t) e^{at} + acu_y(y-ct, t) e^{at}$
 $= -cu_x e^{at} + (Du_{xx} + cu_x e^{at} - au) e^{at} + au e^{at} = Du_{xx}(y-ct, t) e^{at} \quad DV_{yy} = D \frac{\partial^2}{\partial y^2} (u(y-ct, t) e^{at}) = D \frac{\partial^2}{\partial y^2} (u_x(y-ct, t) e^{at}) = Du_{xx}(y-ct, t) e^{at}$
 (ix) $v(y, t) := \frac{1}{\sqrt{4\pi t}} e^{-\frac{|y|^2}{4t}}$ sol. $V_t = V_{yy}$. Use (viii) to find $u(x, t)$ sol. to $u_t = u_{xx} + 2u_x - 0.5u$. $a = \frac{1}{2}, c = 2 \therefore v(y, t) = u(y-2, t) e^{-\frac{|y|^2}{4t}}$
 let $y = x+2t \therefore V(x+2t, t) = u(x, t) e^{-\frac{|x+2t|^2}{4t}}$. $u(x, t) = V(x+2t, t) e^{\frac{|x+2t|^2}{4t}} = \frac{1}{\sqrt{4\pi t}} e^{\frac{|x+2t|^2}{4t}}$. (x) Solve $u(x, 0) = e^x$
 By sol. formula $u(x, t) = \frac{1}{\sqrt{4\pi D t}} \int_{\mathbb{R}} e^{\frac{-|x-y|^2}{4Dt}} dy$. Let $z = y-x \therefore \int_{\mathbb{R}} e^{\frac{-|x-z|^2}{4Dt}} dz = e^x \int_{\mathbb{R}} e^{\frac{-|z|^2}{4Dt}} dz = e^x \frac{1}{\sqrt{4\pi Dt}} \int_{\mathbb{R}} e^{\frac{-|z|^2}{4Dt}} dz$
 let $s = z-2\sqrt{Dt} \Rightarrow e^{\frac{x+Dt}{\sqrt{4Dt}}} \int_{\mathbb{R}} e^{\frac{-s^2}{4Dt}} ds = e^{\frac{x+Dt}{\sqrt{4Dt}}} \quad$ (xi) Use F to find gen. sol to wave Eq $u_{tt} = c^2 u_{xx}$: $\hat{u}_{tt} = c^2 \hat{u}_{xx} \therefore \hat{u}_{tt} = c^2 c i k^2 \hat{u}(k, t)$
 $\therefore \hat{u}_{tt}(k, t) = -c^2 k^2 \hat{u}(k, t) \therefore \hat{u}(k, t) = A(k) e^{ikt} + B(k) e^{-ikt}$. Let $F(k) = F^{-1}(A(k)), G(k) = F^{-1}(B(k)) \Rightarrow u(x, t) = F^{-1}(F(k) + G(k, t))$
 $= F^{-1}(A(k) e^{ikt}) + F^{-1}(B(k) e^{-ikt}) = F(x-ct) + G(x+ct)$. (xii) Solve $\begin{cases} \int_R u_{tt} - 3u_{xx} - 4u_{xx} = 0 \\ u(x, 0) = 2x^2 \\ u_t(x, 0) = 4x^2 \end{cases} \Rightarrow \begin{cases} \frac{\partial^2}{\partial t^2} - 4\frac{\partial^2}{\partial x^2} (\frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial x^2}) u = 0 \\ u(x, 0) = 2x^2 \\ u_t(x, 0) = 4x^2 \end{cases} \therefore \begin{cases} F(k) + G(k) = 2x^2 \\ F'(k) - G'(k) = 4e^{ikx} \end{cases} \therefore \begin{cases} F(k) + G(k) = 2x^2 \\ 4F(k) - G(k) = 4e^{ikx} \end{cases} \therefore \begin{cases} 5F(k) = 2x^2 + 4e^{ikx} \\ G(k) = \frac{2}{5}x^2 - \frac{4}{5}e^{ikx} \end{cases} \therefore F(k) = \frac{2}{5}x^2 + 2e^{ikx} + \frac{4}{5}e^{ikx}$
 $G(k) = \frac{2}{5}x^2 - \frac{4}{5}e^{ikx} \therefore u(x, t) = \frac{2}{5}(x+4t)^2 + 2e^{ikt} + \frac{4}{5}(x-t)^2 - \frac{4}{5}e^{ikt}$ (xiii) Choose $\omega > 0$. Let $u(x, t)$ be a sol. $\in \mathcal{D}$ to damped
 Wave Eq. $u_{tt} + \alpha u_t = u_{xx}$. Show that energy $E(t) = \frac{1}{2} \int_R [u_t(x, t)^2 + (u_x(x, t))^2] dx$ is non-increasing ($\frac{d}{dt} E(t) \leq 0$)
 1) $(\alpha, u_t) \Rightarrow \int_R u_{tt} u_t dx + \alpha \int_R (u_t)^2 dx = \int_R u_{xx} u_t dx$ $\stackrel{\text{LBBP}}{\Rightarrow} \frac{d}{dt} \left(\frac{1}{2} \int_R (u_t)^2 dx \right) = -\alpha \int_R (u_t)^2 dx \leq 0 \Rightarrow \frac{d}{dt} E(t) \leq 0$
 $\frac{d}{dt} \left(\frac{1}{2} \int_R (u_t)^2 dx \right) = -\int_R u_x u_{xt} dx = -\frac{d}{dt} \left(\frac{1}{2} \int_R (u_x)^2 dx \right)$

$a < k < b$

Heat Eq: Bounded Intervals: $u(x,t) : u_t = D u_{xx}$ $\Rightarrow u(a,t) = T_1, u(b,t) = T_2 \forall t > 0$ (Dirichlet Condition), T_1, T_2 constants.

(ii) $u(x,a) = u(x,b) = 0 \forall t > 0$ (Neumann Condition), $\frac{u(a-h,t) - u(a+h,t)}{h} = 0 \Rightarrow u_x(a,h) = 0$. Balance Eq. of x : $\# \rho \in [a-h, a] = \# \rho \in [a, a+h]$. Balance Eq. of x :

$h u(a-\frac{h}{2},t) = h u(a+\frac{h}{2},t) \therefore \lim_{h \rightarrow 0} \frac{h u(a-\frac{h}{2},t) - h u(a+\frac{h}{2},t)}{h^2} = 0 \Rightarrow u_{xx}(a,0) = 0$. **Homogeneous Dirichlet Conditions:** $[a,b] = [0,\pi]$, $D=1$

(1) $\int u(x,t) dx = \int u(\pi, t) dx = 0, t > 0$. Recall $V = A V \in \mathbb{R}^2$. (2) Sol: $v(t) = v_0 e^{kt}$ (iii) $(A - \lambda)v_0 = 0$ (2ii) $\exists \lambda_1, \lambda_2$ sol. $\det(A - \lambda I) = 0$ (iv) eigenvectors: \vec{v}_1, \vec{v}_2 sol. $(A - \lambda_i) \vec{v}_i = 0, \lambda_i \neq 0$ (V) $v(t) = A_1 e^{\lambda_1 t} \vec{v}_1 + A_2 e^{\lambda_2 t} \vec{v}_2$ (vi) A_1, A_2 from V(i). (1): Look solutions of form $u(x,t) = e^{kt} \sin(nx)$

LHS = $\frac{\partial}{\partial t} (e^{kt} \sin(nx)) = \lambda e^{kt} \sin(nx)$, RHS = $\frac{\partial^2}{\partial x^2} (e^{kt} \sin(nx)) = -n^2 e^{kt} \sin(nx) : \lambda = -n^2 \Rightarrow \text{LHS} = \text{RHS}$, $u(x,t) = e^{-n^2 t} \sin(nx) = 0 \forall t, u_{xx}(t) = e^{-n^2 t} \sin(n\pi) = 0 (n \in \mathbb{N})$

$\Rightarrow e^{-n^2 t} \sin(n\pi) = 0 \forall n \in \mathbb{N}$. **Superposition:** If u_1, u_2 sol $\Rightarrow (a, u_1 + u_2)$ sol. $\forall a_1, a_2 \therefore u(x,t) = \sum_{n=1}^{\infty} a_n e^{-n^2 t} \sin(nx)$. **Separation of Variables:**

$u(x,t) = X(x)T(t) \Rightarrow (1) \Rightarrow X'(x)T(t) = X''(x)T(t), T'(t) = \frac{X''(x)}{X(x)}T(t) = \lambda \in \mathbb{R}$, $u(x,0) = X(x)T(0) = 0 \Rightarrow X(0) = 0 \Rightarrow X(\pi) = 0 \therefore \begin{cases} X(0) = 0 \\ X(\pi) = 0 \end{cases}$ set $V(x) = X(x)$

$\therefore V'(x) = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} x' \\ 0 \end{pmatrix} \therefore V' = AV, A = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \therefore \text{eigenvalues: } \lambda_1 = i\sqrt{\lambda}, \lambda_2 = -i\sqrt{\lambda}$. **Case 1**: $\exists w \in \mathbb{R}$ s.t. $\lambda = w^2 \therefore \lambda_1 = w, \lambda_2 = -w$. $A - \lambda I = \begin{pmatrix} -w & 0 \\ 0 & -w \end{pmatrix}$

$(A - \lambda I) \vec{v}_i = \vec{0} \Rightarrow \lambda_1 w v_1 = \vec{0} \therefore \lambda_1 w v_1 = \vec{0} \therefore V_1 = \begin{pmatrix} w \\ 1 \end{pmatrix}; (A - \lambda_2 I) \vec{v}_2 = \begin{pmatrix} w \\ -1 \end{pmatrix} \therefore V_2 = \begin{pmatrix} -w \\ 1 \end{pmatrix} \therefore V(x) = \begin{pmatrix} x' \\ 0 \end{pmatrix} = A e^{wx} \begin{pmatrix} 1 \\ 1 \end{pmatrix} + B e^{-wx} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \therefore X(x) = Ae^{wx} + Be^{-wx} \forall A, B$

$X(0) = a_1 + a_2 = 0 \Rightarrow a_1 = -a_2, X(\pi) = a_1 e^{w\pi} - a_2 e^{-w\pi} = 0 \Rightarrow a_1 = 0 \therefore X(x) = 0$ Only solution. **Case 2**: $\lambda_1 = \lambda_2 = 0, A - \lambda_1 I = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \therefore \lambda_1 \text{ is not a generalized eigenvector}$

$\Rightarrow V_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \Rightarrow V_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \Rightarrow V_{gen} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} \Rightarrow V_{gen} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \therefore V(x) = \begin{pmatrix} x' \\ 0 \end{pmatrix} = a_1 e^{0x} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + a_2 (x e^{0x} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + e^{0x} \begin{pmatrix} 0 \\ 1 \end{pmatrix}) \Rightarrow X(x) = a_1 + a_2 x \forall a_1, a_2$

$\Rightarrow X(x) = 0$. **Case 3: $\lambda \neq 0$** : $\lambda = i\omega, \lambda_2 = -i\omega \Rightarrow X(x) = a_1 \cos(\omega x) + a_2 \sin(\omega x)$ **Ansatz**, $X(0) = a_1 = 0, X(\pi) = a_2 \sin(\omega\pi) = 0 \Rightarrow \omega = n\pi \in \mathbb{N}$. $X(x) = a_2 \sin(nx)$

$\lambda = -n^2 \Rightarrow T'(t) = -n^2 t \Rightarrow T(t) = C e^{-n^2 t} \therefore X(x)T(t) = C e^{-n^2 t} \sin(nx), n \in \mathbb{N}$. **By Superposition:** $u(x,t) = \sum_{n=1}^{\infty} a_n e^{-n^2 t} \sin(nx)$. As determined from

$u(x,0) = f(x) \Rightarrow \sum_{n=1}^{\infty} a_n \sin(nx) = f(x)$. **Sol:** $\sum_{n=1}^{\infty} a_n \sin(nx) + f(x) = 0 \Rightarrow a_1 = 0, a_2 = 0, a_n = 0 \text{ for } n \neq 1, 2, \dots$. $u(x,t) = \sum_{n=1}^{\infty} a_n e^{-n^2 t} \sin(nx) + f(x)$. **Solution is unique.**

Assume u_1, u_2 soln. to (1). Let $V = u_1 - u_2 \Rightarrow \begin{cases} V(0,t) = V(\pi,t) = 0 \\ V(x,0) = 0 \end{cases} \therefore \int V e^V dx + \int V_{xx} e^V dx, \frac{\partial^2}{\partial t^2} (\int V^2 dx) = V e^V \int_{x=0}^{\pi} \int_{t=0}^T \int_{x=0}^{\pi} \int_{t=0}^T V^2 dx dt \Rightarrow \partial S(E(t)) \subseteq E(0) = 0, E(t) = \int_0^T V(x,t) dx$

$\Rightarrow V(x,t) = 0 \forall t \Rightarrow u_1(x,t) = u_2(x,t) \therefore \exists! u(x,t)$. **Inhomogeneous B.C.** $\begin{cases} u(0,t) = T_1 \\ u(\pi,t) = T_2 \end{cases} \therefore u(x,t) = ax + b$ soln. $\therefore u_1(0,t) = T_1, u_2(\pi,t) = T_2 \Rightarrow b = T_1, a = \frac{T_2 - T_1}{\pi}$. $\therefore V(x,t) = u_1(x,t) - u_2(x,t) = \frac{T_2 - T_1}{\pi} x + T_1 + \sum_{n=1}^{\infty} a_n e^{-n^2 t} \sin(nx)$. **Wave Eq:** $\begin{cases} u(0,t) = 0 \\ u(\pi,t) = 0 \end{cases} \therefore \frac{T''(t)}{T(t)} = \frac{X''(x)}{X(x)} = \lambda \therefore X(x) = \lambda X(x), X'(0) = X'(\pi) = 0 \therefore X(x) = a_1 e^{wx} + a_2 e^{-wx}$

$X'(0) = a_1 w e^{wx} - a_2 w e^{-wx}, X'(0) = a_1 (w e^{wx} - w e^{-wx}) = 0 \Rightarrow a_1 = a_2 = 0$. **Case 2**: $\lambda = 0$: $X(x) = a_1 + a_2 x, X'(0) = a_2 = 0 \Rightarrow a_2 = 0, X'(0) = -a_1 w \sin(wx) = 0 \Rightarrow w = n\pi \in \mathbb{N}$

Case 3: $\lambda = -\omega^2$: $X(x) = a_1 \cos(\omega x) + a_2 \sin(\omega x)$, $X'(x) = -a_1 \omega \sin(\omega x) + a_2 \omega \cos(\omega x)$, $X'(0) = a_2 \omega = 0 \Rightarrow a_2 = 0, X'(0) = -a_1 \omega \sin(\omega x) = 0 \Rightarrow w = n\pi \in \mathbb{N}$

$\therefore X(x) = a_1 \cos(\omega x) + a_2 \sin(\omega x) + C_1 \cos(\omega t) + C_2 \sin(\omega t) \therefore X(x)T(x) = C_1 \cos(\omega x) \cos(\omega t) + C_2 \sin(\omega x) \sin(\omega t)$

$\therefore u(x,t) = a_1 + b_1 t + \sum_{n=1}^{\infty} a_n \cos(\omega x) \cos(\omega nt) + b_n \sin(\omega x) \sin(\omega nt) \therefore \text{Dirichlet B.C. } (X'' = \lambda X) \quad (X(0) = X(\pi) = 0) \Rightarrow X(x) = C \sin(nx) \forall c$.

Neumann B.C. ($X'(0) = X'(\pi) = 0$): $X(x) = C \cos(nx) \forall c$. From (1) $u(x,t) = \sum_{n=1}^{\infty} a_n e^{-n^2 t} \sin(nx)$, a_n s.t. $f(x) = \sum_{n=1}^{\infty} a_n \sin(nx) \therefore \int_0^{\pi} f(x) \sin(nx) dx = \int_0^{\pi} \sum_{n=1}^{\infty} a_n \sin(nx) \sin(nx) dx$

$= \sum_{n=1}^{\infty} a_n \int_0^{\pi} \sin(nx) \sin(nx) dx$. **Case (a+b) = $\cos(a)\cos(b) - \sin(a)\sin(b)$, $\cos(a-b) = \cos(a)\cos(b) + \sin(a)\sin(b) \Rightarrow \sin(a)\sin(b) = \frac{1}{2} [\cos(a-b) - \cos(a+b)]$** . $\therefore \sum_{n=1}^{\infty} a_n \sin(nx) \sin(nx) dx = \sum_{n=1}^{\infty} \frac{a_n}{2} \int_0^{\pi} (-\cos(2nx) - \cos((n+2)x)) dx + \sum_{n=1}^{\infty} \frac{a_n}{2} \int_0^{\pi} \cos((n+2)x) - \cos((n-2)x) dx = \frac{\pi}{2} a_n - \frac{a_n}{2} \frac{\pi}{2} \sin(n\pi) + \sum_{n=1}^{\infty} \frac{a_n}{2} \left[\frac{\sin(-n\pi)}{n-\pi} - \frac{\sin(n\pi)}{n+\pi} \right] = \frac{\pi}{2} a_n$

$\Rightarrow a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$. **Fourier Series:** $f^2(a,b) = \{f : [a,b] \rightarrow \mathbb{R} \mid f \in L^2[a,b]\}$, $\|f\|_2 = \sqrt{\int_a^b |f(x)|^2 dx} \text{ (norm of } f \in f^2(a,b))$, $\langle f, g \rangle \text{ scalar product}$

$\langle f, g \rangle = \int_a^b f(x)g(x) dx, f, g \in f^2(a,b) \Rightarrow f, g \text{ orthogonal if } \langle f, g \rangle = 0$. **Ex:** $V_n(x) = \sin(nx) \in f^2$, $\|V_n\|_2 = \sqrt{\int_0^{\pi} |\sin(nx)|^2 dx} = \sqrt{\frac{\pi}{2} \cos(0) - \frac{1}{2} \cos(n\pi)} = \sqrt{\frac{\pi}{2}}$

$\langle V_n, V_m \rangle = \int_0^{\pi} \sin(nx) \sin(mx) dx = 0 \text{ for } n \neq m \Rightarrow V_n \perp V_m \forall n \neq m \in \mathbb{Z}$; let $W_n(x) = \cos(nx) \in f^2$: $\|W_n\|_2 = \sqrt{\frac{\pi}{2}} \neq 0, \|W_n\|_2 = \sqrt{\pi}, W_n \perp W_m \forall n \neq m$

$V_n \perp W_m, n, m \in \mathbb{Z}$. **Definition:** Let $f, f_n \in f^2(a,b)$, $n = 1, 2, \dots$ we say $\frac{1}{n} \sum_{k=1}^n f_k$ if $\|f - \sum_{k=1}^n f_k\|_2 \rightarrow 0$ as $N \rightarrow \infty$. **Theorem:** $\forall f \in f^2(a, \pi)$, we have

(i) $\frac{1}{n} \sum_{k=1}^n f_k \in \text{ansin}(nx)$, $a_n = \frac{\|f_n\|_2}{\|\sin(nx)\|_2^2}$ (ii) $\frac{1}{n} \sum_{k=1}^n f_k \in \text{acos}(nx)$, $a_n = \frac{\|f_n\|_2}{\|\cos(nx)\|_2^2}$ **Remark:** for $a \neq 0$: $\|\sin(nx)\|_2^2 = \|\cos(nx)\|_2^2 = \frac{\pi}{2} \Rightarrow a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$ or $a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \cos(nx) dx$. For $n=0$: $\|\cos(nx)\|_2^2 = \int_0^{\pi} |\cos(nx)|^2 dx \rightarrow a_0 = \frac{1}{\pi} \int_0^{\pi} f(x) dx$. Ex: $f(x) = x$, $a_n = \frac{2}{\pi} \int_0^{\pi} x \sin(nx) dx = (-1)^{n+1} \frac{\pi}{n} \therefore x = \frac{2}{\pi} \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\pi}{n} \sin(nx)$. $\therefore \text{solution to } \{u(x,0) = x\}$

$u(x,t) = \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\pi}{n} e^{-n^2 t} \sin(nx)$. $f(x) = \sum_{n=1}^{\infty} a_n \sin(nx) \Rightarrow f(x) = \sum_{n=1}^{\infty} a_n \sin(nx) \forall x \in \mathbb{R}$. $x = \pi \neq \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\pi}{n} \sin(nx)$

(i) $\max_{0 \leq x \leq \pi} |S_n(x) - f(x)| \rightarrow 0, N \rightarrow \infty$: Uniform Convergence to $f(x)$ on $[0, \pi]$, (II): Fix $0 < x \leq \pi$, then $|S_n(x) - f(x)| \rightarrow 0, N \rightarrow \infty$: Pointwise convergence

(III) For $0 < x \leq \pi$, we have $S_n(x) - g(x) = \begin{cases} f(x), & x \neq d \\ \frac{1}{2} f(d) + \frac{1}{2} f(d^+) & x = d \end{cases}$. Def: $f(x)$ on $[a, b]$ called piecewise continuous, if has at most finitely many points of discontinuity, and $f(x)$ has left and right limits at each point. **Theorem:** $f(x)$ piecewise cont. on $[0, \pi]$. $f(x)$ piecewise

$\lim_{N \rightarrow \infty} \frac{\int_a^b |f(x) - S_n(x)|^2 dx}{2} \rightarrow 0$. Application: $x = \frac{\pi}{2} = \frac{1}{2} \sin(\frac{\pi}{2}) - \frac{1}{2} \sin(\frac{3\pi}{2}) + \frac{1}{2} \sin(\frac{5\pi}{2}) - \dots \Rightarrow \frac{\pi}{4} = -\frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots$. **Poisson's Id:** Assume $f(x) = \frac{1}{2} \sum_{n=1}^{\infty} a_n \sin(nx)$

$\|f(x)\|_2^2 = \langle f(x), f(x) \rangle = \sum_{n=1}^{\infty} \langle a_n \sin(nx), \sum_{m=1}^{\infty} a_m \sin(mx) \rangle = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \langle a_n \sin(nx), a_m \sin(mx) \rangle = \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} a_n a_m \int_0^{\pi} \sin(nx) \sin(mx) dx$ By orthogonality of sin(nx) and cos(mx)

$\|f(x)\|_2^2 = \sum_{n=1}^{\infty} a_n^2 \int_0^{\pi} |\sin(nx)|^2 dx = \sum_{n=1}^{\infty} a_n^2 \pi = \frac{\pi}{2} \sum_{n=1}^{\infty} a_n^2$

If $f(x) = \sum_{n=1}^{\infty} a_n \sin(nx)$, $\|f\|_2^2 = \pi a_1^2 + \sum_{n=2}^{\infty} a_n^2$. Heat Eq. $\{u(x,0) = f(x)\}, u(x,t) = \sum_{n=1}^{\infty} a_n \sin(nx), u_0(x) = f(x)$; use Eq. $\{u_t(x,0) = g(x)\}, f, g$ bounded p.c.

$u(x,t) = \sum_{n=1}^{\infty} (a_n \cos(nt) + b_n \sin(nt)) \sin(nx), u(x,t) = \sum_{n=1}^{\infty} (-b_n \sin(nt) + a_n \cos(nt)) \sin(nx) \Rightarrow u(x,0) = \sum_{n=1}^{\infty} a_n \sin(nx) = f(x) \Rightarrow a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$

$\Rightarrow b_n = \frac{2}{\pi} \int_0^{\pi} g(x) \sin(nx) dx$. Q: $\sum_{n=1}^{\infty} a_n e^{-nt} \sin(nx) \leq \infty$. $|a_n| = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx \leq 2 \pi \int_0^{\pi} |f(x)| |\sin(nx)| dx \leq 2 \max|f(x)| = M$. Ratio Test: $\exists \delta, b > 0, \exists c < 1$ s.t. $\frac{b}{\delta} \leq \epsilon$ \Rightarrow pointwise convergence

$\forall x: \lim_{n \rightarrow \infty} f_n(x) = f(x)$. Uniform C. $\forall \epsilon > 0 \exists N: |f_n(x) - f(x)| < \epsilon \forall n \geq N, \forall x$.

