Question:

When t = 0 a switch closes connecting a voltage source to a series RL circuit. The source is characterised by: $v(t) = V_m \sin(\omega t + \phi)$. Find the time domain expression of the current, i(t).

$$I(s) = \frac{V(s)}{Z(s)} \tag{1}$$

Finding V(s)

$$\mathcal{L}\left\{v(t)\right\} = \mathcal{L}\left\{V_m \sin(\omega t + \phi)\right\} = V_m \mathcal{L}\left\{\sin(\omega t + \phi)\right\}$$
 (2)

$$V_m \mathcal{L}\left\{\sin(\omega t)\cos(\phi) + \cos(wt)\sin(\phi)\right\}$$
(3)

$$V_m \left(\mathcal{L} \left\{ \sin(\omega t) \cos(\phi) \right\} + \mathcal{L} \left\{ \cos(wt) \sin(\phi) \right\} \right)$$
(4)

Note that $\cos(\phi)$ and $\sin(\phi)$ are constants with respect to time can subsequently be removed from the argument of the Laplace transform.

$$V_m \left(\cos(\phi) \mathcal{L} \left\{ \sin(\omega t) \right\} + \sin(\phi) \mathcal{L} \left\{ \cos(wt) \right\} \right)$$
 (5)

$$\therefore V(s) = V_m \left(\cos(\phi) \frac{\omega}{s^2 + \omega^2} + \sin(\phi) \frac{s}{s^2 + \omega^2} \right)$$
 (6)

Finding Z(s) by simple s-domain series impedance:

$$Z(s) = sL + R \quad \Rightarrow \quad \frac{1}{Z(s)} = \frac{1}{sL + R} = \frac{\frac{1}{L}}{s + \frac{R}{L}}$$
 (7)

Then to current, I(s)

$$I(s) = \frac{V(s)}{Z(s)} = V_m \left(\cos(\phi) \frac{\omega}{s^2 + \omega^2} + \sin(\phi) \frac{s}{s^2 + \omega^2}\right) \frac{\frac{1}{L}}{s + \frac{R}{\tau}}$$
(8)

$$=V_m \left(\frac{\cos(\phi)\omega + \sin(\phi)s}{s^2 + \omega^2}\right) \frac{\frac{1}{L}}{s + \frac{R}{L}}$$
(9)

$$= \frac{V_m}{L} \left(\frac{\cos(\phi)\omega + \sin(\phi)s}{(s^2 + \omega^2)(s + \frac{R}{L})} \right)$$
 (10)

We want to decompose the bottom into separate individual poles so we can use the residue theorem. It's helpful to decompose $s^2 + \omega^2$ into $(s + j\omega)(s - j\omega)$.

$$= \left(\frac{V_m}{L}\right) \frac{\left(\cos(\phi)\omega + \sin(\phi)s\right)}{\left(s + j\omega\right)\left(s - jw\right)\left(s + \frac{R}{L}\right)} \tag{11}$$

Recall the residue theorem:

$$f(t) = 2\pi i \sum_{k=1}^{m} \frac{1}{(n-1)!} \lim_{s \to s_k} \frac{d^{n-1}}{ds^{n-1}} (s - s_k)^n F(s)$$
 (12)

s.t.

m = number of poles n = order of the pole $s_k = \text{the } k\text{-th pole in } F(s)$

Using the residue theorem to solve inverse Laplace transforms:

$$\mathcal{L}^{-1}\left\{F(s)\right\} = f(t) = \sum_{k=1}^{m} \operatorname{Res}(e^{st}F(s), s_k) = \sum_{k=1}^{m} \frac{1}{(n-1)!} \lim_{s \to s_k} \frac{d^{n-1}}{ds^{n-1}} (s - s_k)^n e^{st}F(s)$$
(13)

Looking at the denominator of the frequency domain function of current, we can see the set of poles,

$$\{s\} = \left\{ -j\omega, \ j\omega, \ -\frac{R}{L} \right\}$$

Now we have to insert the poles and evaluate Eq.13 but we can first do some simplification. For our poles, they are all of order 1.

$$\sum_{k=1}^{m} \frac{1}{(n-1)!} \lim_{s \to s_k} \frac{d^{n-1}}{ds^{n-1}} (s - s_k)^n e^{st} F(s) \bigg|_{n=1} = \sum_{k=1}^{m} \lim_{s \to s_k} (s - s_k)^n e^{st} F(s)$$
 (14)

$$\sum_{k=1}^{m} \lim_{s \to s_k} (s - s_k) e^{st} F(s) = \lim_{s \to j\omega} (s - s_k) e^{st} F(s) + \lim_{s \to -j\omega} (s - s_k) e^{st} F(s) + \lim_{s \to \alpha} (s - s_k) e^{st} F(s)$$
(15)

$$\sum_{k=1}^{m} \lim_{s \to s_k} (s - s_k) e^{st} F(s) = (s - j\omega) e^{j\omega t} F(j\omega) + (s + j\omega) e^{-j\omega t} F(-j\omega) + (s - \alpha) e^{\alpha t} F(\alpha)$$
(16)

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