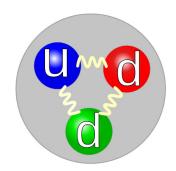
on the magnetic field profile for the nEDM search

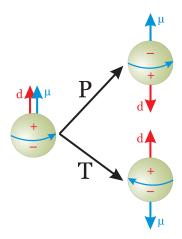
Aritra Biswas Kellogg Radiation Laboratory Mentors: Brad Filippone, Simon Slutsky

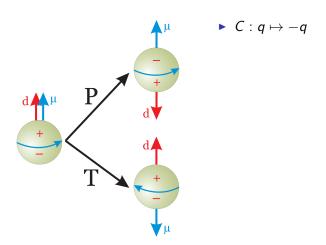
November 18, 2014

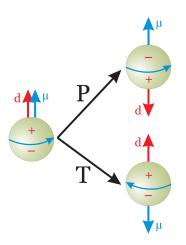
nEDM = neutron electric dipole moment

- distributed + and charges inside neutron
- electric dipole moment (EDM) measures separation between centers of + and charge

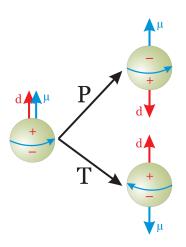




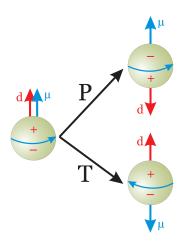




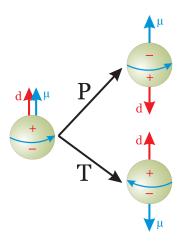
- $ightharpoonup C: q \mapsto -q$
- $P: (t, x, y, z) \mapsto (t, -x, -y, -z)$



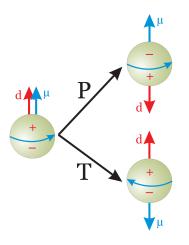
- \triangleright C: $q \mapsto -q$
- $P: (t, x, y, z) \mapsto (t, -x, -y, -z)$
- $ightharpoonup T: (t, x, y, z) \mapsto (-t, x, y, z)$



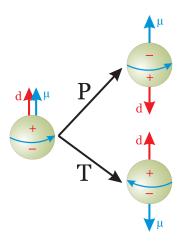
- \triangleright $C: q \mapsto -q$
- $P:(t,x,y,z)\mapsto(t,-x,-y,-z)$
- $ightharpoonup T: (t, x, y, z) \mapsto (-t, x, y, z)$
- CPT symmetry



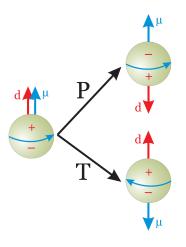
- \triangleright $C: q \mapsto -q$
- $P:(t,x,y,z)\mapsto(t,-x,-y,-z)$
- $ightharpoonup T: (t, x, y, z) \mapsto (-t, x, y, z)$
- CPT symmetry+ P violation



- \triangleright $C: q \mapsto -q$
- $P:(t,x,y,z)\mapsto(t,-x,-y,-z)$
- $ightharpoonup T: (t, x, y, z) \mapsto (-t, x, y, z)$
- CPT symmetry
 - + P violation
 - + T violation



- \triangleright $C: q \mapsto -q$
- $P:(t,x,y,z)\mapsto(t,-x,-y,-z)$
- \vdash $T:(t,x,y,z)\mapsto(-t,x,y,z)$
- CPT symmetry
 - + P violation
 - + T violation
 - \Rightarrow *CP* violation



- $ightharpoonup C: q \mapsto -q$
- $P: (t, x, y, z) \mapsto (t, -x, -y, -z)$
- $ightharpoonup T: (t, x, y, z) \mapsto (-t, x, y, z)$
- CPT symmetry
 - + P violation
 - + T violation
 - \Rightarrow *CP* violation
- reformulations of Standard Model
- matter-antimatter asymmetry

▶ put ultra-cold neutrons (UCN) in **E** and **B** fields

Pendlebury et. al. Geometric-phase-induced false electric dipole moment signals for particles in traps. Phys. Rev. A. 70, 032102 (2004).

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar}$$

Pendlebury et. al. Geometric-phase-induced false electric dipole moment signals for particles in traps. Phys. Rev. A. 70, 032102 (2004).

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

▶
$$\frac{\partial \mathbf{B}}{\partial (\mathbf{x}, \mathbf{y}, \mathbf{z})} \neq 0 \Rightarrow \frac{\partial \mathbf{B}}{\partial t} \neq 0 \Rightarrow \mathbf{E}$$
 field

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

▶
$$\frac{\partial \mathbf{B}}{\partial (x,y,z)} \neq 0 \Rightarrow \frac{\partial \mathbf{B}}{\partial t} \neq 0 \Rightarrow \mathbf{E} \text{ field } \Rightarrow \Delta \omega_{geo}$$

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

▶ geometric phase ⇒ false measurement!

Pendlebury et. al. Geometric-phase-induced false electric dipole moment signals for particles in traps. Phys. Rev. A. 70, 032102 (2004).

- put ultra-cold neutrons (UCN) in E and B fields
- lacktriangle neutron's spin will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
 (1)

$$\Delta\omega = \pm \frac{2d_n E}{J\hbar} \pm \Delta\omega_{geo} \tag{2}$$

▶
$$\frac{\partial \mathbf{B}}{\partial (x,y,z)} \neq 0 \Rightarrow \frac{\partial \mathbf{B}}{\partial t} \neq 0 \Rightarrow \mathbf{E} \text{ field } \Rightarrow \Delta \omega_{geo}$$

- ▶ geometric phase ⇒ false measurement!
- engineering challenge: creating an uniform magnetic field

Pendlebury et. al. Geometric-phase-induced false electric dipole moment signals for particles in traps. Phys. Rev. A. 70, 032102 (2004).





▶ about 2 meters tall



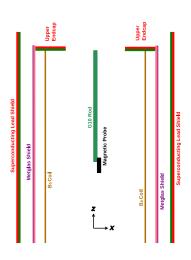
- ▶ about 2 meters tall
- inside a cryostat (cools to 4 K)

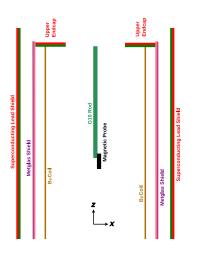


- about 2 meters tall
- inside a cryostat (cools to 4 K)
- only creates B field; no measurement cells

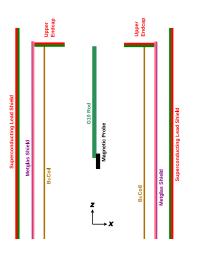


- about 2 meters tall
- inside a cryostat (cools to 4 K)
- only creates B field; no measurement cells
- final experiment at Oak Ridge National Laboratory

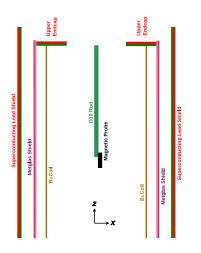




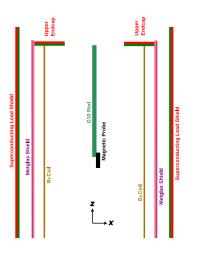
▶ B_0 coil: $\cos \theta$ coil geometry



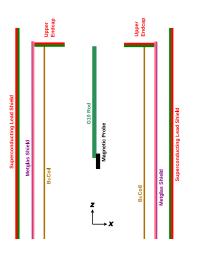
- ▶ B_0 coil: $\cos \theta$ coil geometry
 - ▶ **B** field in *x* direction



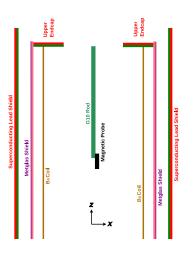
- ▶ B_0 coil: $\cos \theta$ coil geometry
 - ▶ **B** field in *x* direction
- ferromagnetic Metglas shield



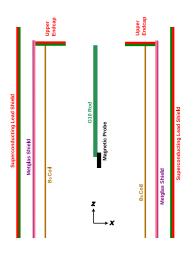
- ▶ B_0 coil: $\cos \theta$ coil geometry
 - ▶ **B** field in *x* direction
- ferromagnetic Metglas shield
 - $\blacktriangleright \ \ \mathsf{high} \ \mu$



- ▶ B_0 coil: $\cos \theta$ coil geometry
 - ▶ **B** field in *x* direction
- ferromagnetic Metglas shield
 - high μ
- superconducting axial shield

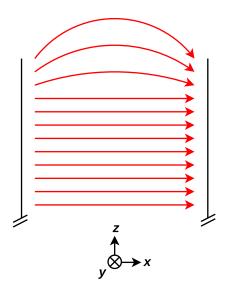


- ▶ B_0 coil: $\cos \theta$ coil geometry
 - ▶ **B** field in *x* direction
- ferromagnetic Metglas shield
 - \blacktriangleright high μ
- superconducting axial shield
 - $\mu = 0$

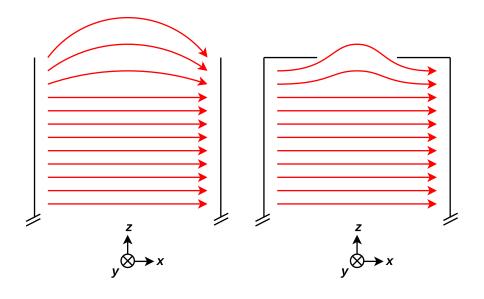


- ▶ B_0 coil: $\cos \theta$ coil geometry
 - ▶ **B** field in *x* direction
- ferromagnetic Metglas shield
 - ▶ high μ
- superconducting axial shield
 - $\blacktriangleright \ \mu = 0$
- superconducting upper endcap

edge effects and the superconducting endcap

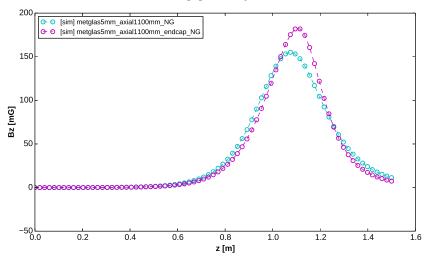


edge effects and the superconducting endcap

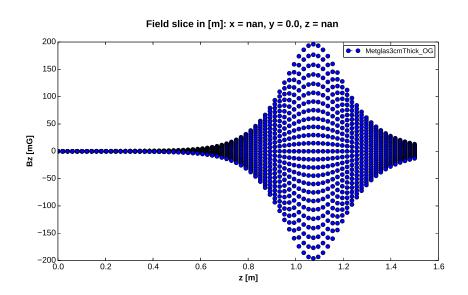


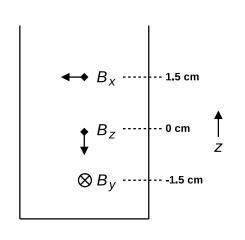
simulations of endcap effect

Field slice in [m]: x = 0.1, y = 0.0, z = nan

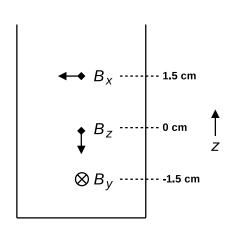


correction: probe *x* centering

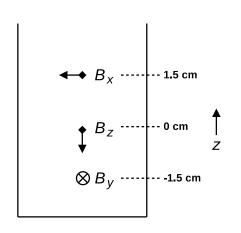




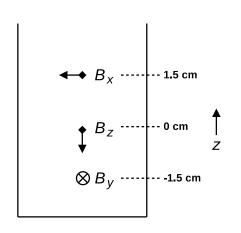
▶ 3 separate 1-axis probes



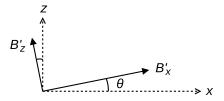
- ▶ 3 separate 1-axis probes
- ▶ incomplete vector map

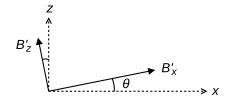


- ▶ 3 separate 1-axis probes
- ▶ incomplete vector map
- need to store z-axis offset vector along with z array

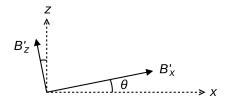


- ▶ 3 separate 1-axis probes
- ▶ incomplete vector map
- need to store z-axis offset vector along with z array



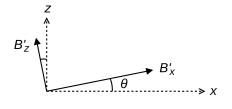


$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta$$
 (3)



$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta$$
 (3)

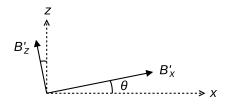
1. θ is small:



$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta$$
 (3)

1. θ is small:

$$B_x = B_x' - B_z'\theta, \quad B_z = B_z' + B_x'\theta$$

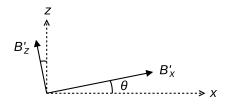


$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta$$
 (3)

1. θ is small:

$$B_x = B_x' - B_z'\theta, \quad B_z = B_z' + B_x'\theta$$

2. $B_z = 0$ at center:



$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta$$
 (3)

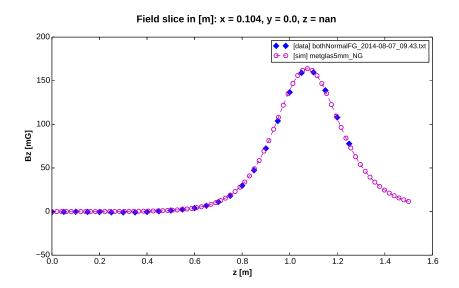
1. θ is small:

$$B_x = B_x' - B_z'\theta, \quad B_z = B_z' + B_x'\theta$$

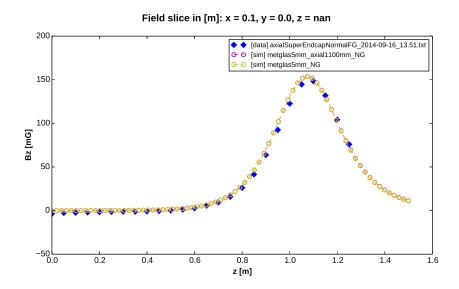
2. $B_z = 0$ at center:

$$\theta = -\frac{B_z'}{B_z'}$$

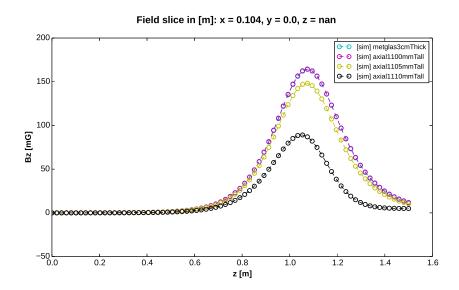
comparison: axial shield normal, endcap normal



comparison: axial shield SC, endcap normal

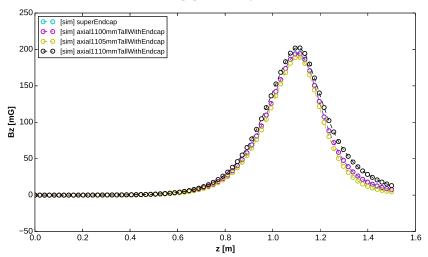


simulation: varying axial shield height

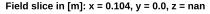


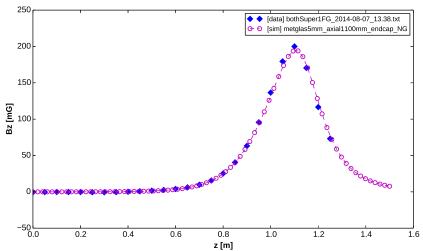
simulation: varying axial shield height, with endcap

Field slice in [m]: x = 0.104, y = 0.0, z = nan



comparison: axial shield SC, endcap SC





simulations are effective in predicting endcap behaviors

- simulations are effective in predicting endcap behaviors
- motivates further simulated studies with different endcap geometries

- simulations are effective in predicting endcap behaviors
- motivates further simulated studies with different endcap geometries
- our endcap seems to shift the B_z peak away from magnet center

- simulations are effective in predicting endcap behaviors
- motivates further simulated studies with different endcap geometries
- our endcap seems to shift the B_z peak away from magnet center
- axial shield effect is stronger when more of it is "uncovered" by the Metglas

- simulations are effective in predicting endcap behaviors
- motivates further simulated studies with different endcap geometries
- our endcap seems to shift the B_z peak away from magnet center
- axial shield effect is stronger when more of it is "uncovered" by the Metglas
- SC endcap hides axial shield influence, even over small variation in height



ongoing and future work

endcap will likely be effective in final experiment

ongoing and future work

- endcap will likely be effective in final experiment
- new model with top and bottom endcaps

ongoing and future work

- endcap will likely be effective in final experiment
- new model with top and bottom endcaps
- analysis of field gradients in measurement cell volumes

acknowledgments

- Mentor: Brad Filippone
- Co-mentor: Simon Slutsky
- Chris Swank, Chub Osthelder, Bob Carr
- Arthur R. Adams SFP Fellowship
- Caltech SURF Program