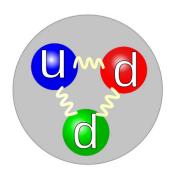
in minimizing magnetic field gradients for the nEDM search

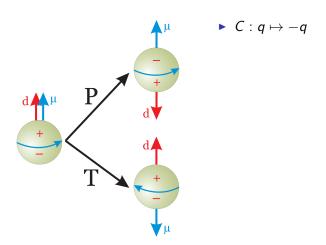
Aritra Biswas
Filippone Group, Kellogg Radiation Laboratory

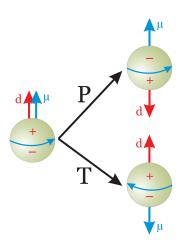
August 7, 2014

nEDM = neutron electric dipole moment

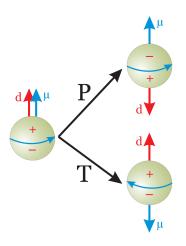
- up quark: $\frac{2}{3}e$
- ► each down quark: $-\frac{1}{3}e$
- electric dipole moment:
 vector measuring separation
 between + and charges
 and their orientation



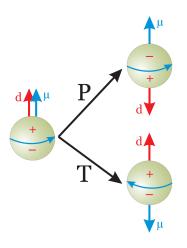




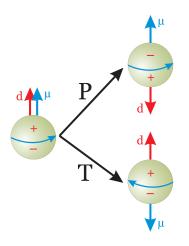
- $ightharpoonup C: q \mapsto -q$
- $P: (t, x, y, z) \mapsto (t, x, -y, z)$



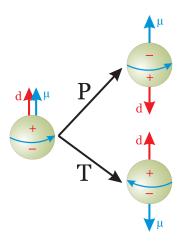
- \triangleright C: $q \mapsto -q$
- $P: (t, x, y, z) \mapsto (t, x, -y, z)$
- $T: (t, x, y, z) \mapsto (-t, x, y, z)$



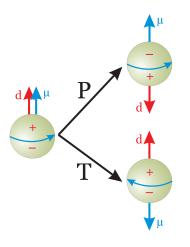
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- ► *CPT* symmetry



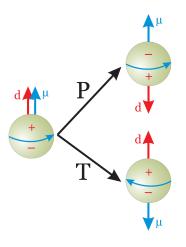
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- CPT symmetry+ P violation



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- CPT symmetry
 - + P violation
 - + T violation
 - \Rightarrow *CP* violation
- reformulations of Standard Model
- matter-antimatter asymmetry

▶ put ultra-cold neutrons (UCN) in **E** and **B** fields

Pendlebury et. al. Geometric-phase-induced false electric dipole moment signals for particles in traps. Phys. Rev. A. 70, 032102 (2004).

- put ultra-cold neutrons (UCN) in E and B fields
- ightharpoonup neutron will precess at frequency ω

$$\omega_{\uparrow\uparrow} = -\frac{\mu_n B + d_n E}{J\hbar}, \quad \omega_{\uparrow\downarrow} = -\frac{\mu_n B - d_n E}{J\hbar}$$
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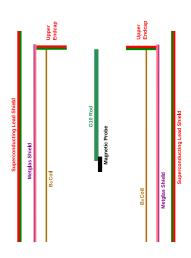
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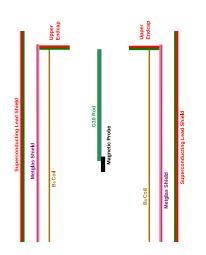
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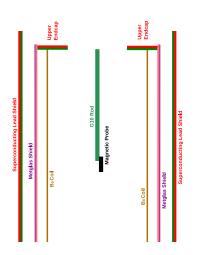
- ▶ $\frac{\partial \mathbf{B}}{\partial (x,y,z)} \neq 0 \Rightarrow \frac{\partial \mathbf{B}}{\partial t} \neq 0 \Rightarrow \text{effect of } \mathbf{E} \text{ field } \Rightarrow \Delta \omega_{geo}$
- ▶ geometric phase ⇒ false measurement!

Pendlebury et. al. Geometric-phase-induced false electric dipole moment signals for particles in traps. Phys. Rev. A. 70, 032102 (2004).

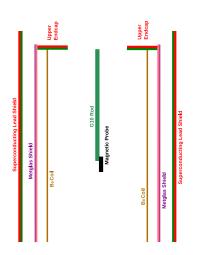




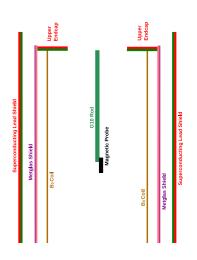
▶ B_0 coil: $\cos \theta$ coil geometry, emulates sheet current



- ▶ B_0 coil: $\cos \theta$ coil geometry, emulates sheet current
- ► ferromagnetic Metglas shield

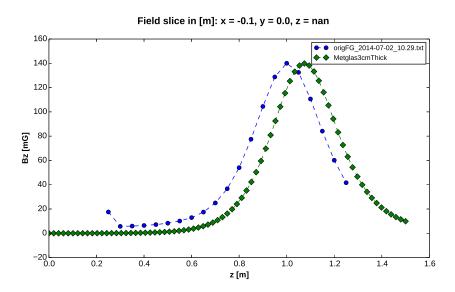


- ▶ B_0 coil: $\cos \theta$ coil geometry, emulates sheet current
- ferromagnetic Metglas shield
- superconducting axial shield



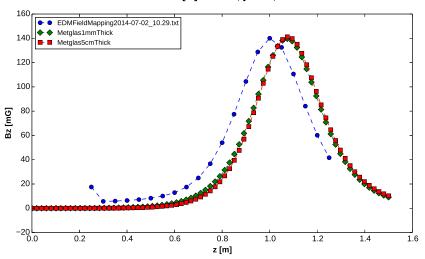
- ▶ B_0 coil: $\cos \theta$ coil geometry, emulates sheet current
- ► ferromagnetic Metglas shield
- superconducting axial shield
- superconducting upper endcap

original comparison, warm



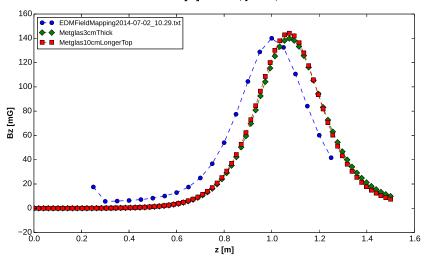
Metglas thickness

Field slice in [m]: x = -0.1, y = 0.0, z = nan

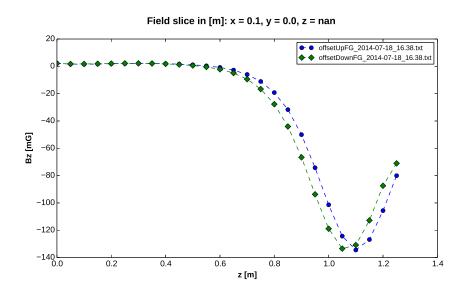


Metglas 10 cm longer on top

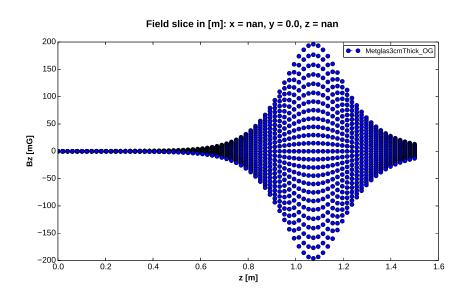
Field slice in [m]: x = -0.1, y = 0.0, z = nan



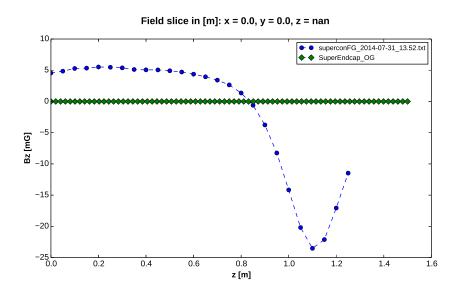
correction 1: probe measurement time



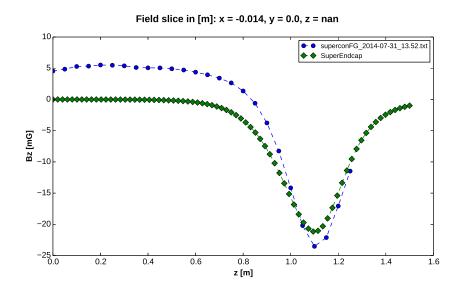
correction 2: x centering



correction 2: x centering, superconducting

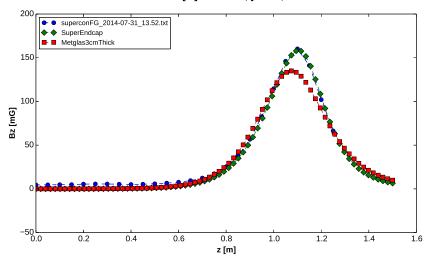


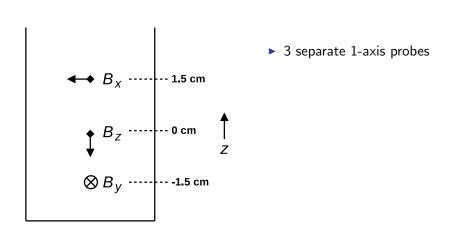
correction 2: x centering, superconducting, 1.4 cm offset

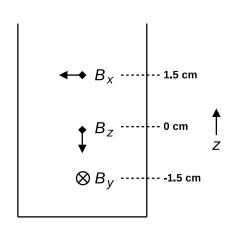


comparison, superconducting

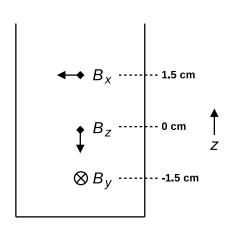
Field slice in [m]: x = 0.086, y = 0.0, z = nan



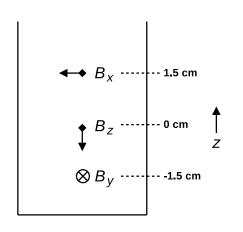




- ▶ 3 separate 1-axis probes
- ▶ incomplete vector map

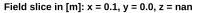


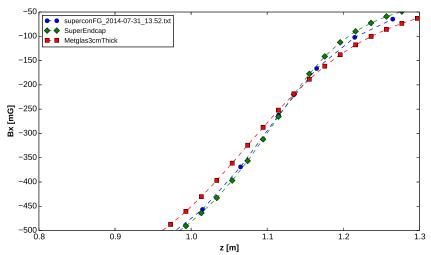
- ▶ 3 separate 1-axis probes
- incomplete vector map
- need to store z-axis offset vector along with z array



- 3 separate 1-axis probes
- incomplete vector map
- need to store z-axis offset vector along with z array
- OffsetAxis class to return proper spatial axis array based on desired vector component

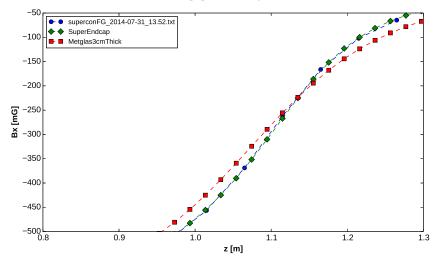
comparison, superconducting, no offsets

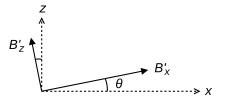


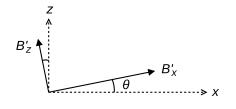


comparison, superconducting, x-centering and axis offset

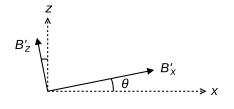
Field slice in [m]: x = 0.086, y = 0.0, z = nan





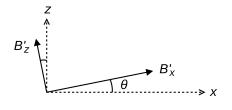


$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta$$
 (3)



$$B_{x} = B'_{x} \cos \theta - B'_{z} \sin \theta, \quad B_{z} = B'_{z} \cos \theta + B'_{x} \sin \theta$$
 (3)

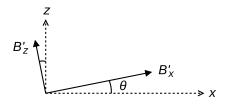
1. θ is small:



$$B_x = B_x' \cos \theta - B_z' \sin \theta, \quad B_z = B_z' \cos \theta + B_x' \sin \theta \tag{3}$$

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$$B_x = B_x' - B_z'\theta, \quad B_z = B_z' + B_x'\theta$$

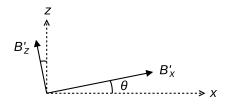


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2. $B_z = 0$ at center:



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$$\theta = -\frac{B_z'}{B_x'}$$

comparison, superconducting

