Ph22 Set 5 Aritra Biswas

1 Introduction

We build on the framework developed in Ph 22.3. Instead of the brute-force calculation of pairwise forces, we use the Barnes-Hut quadtree algorithm to clump together far-away gravitational sources.

1a Implementation choices

A QuadTree class was created to hold the Barnes-Hut tree. The tree is composed of QuadNode objects, which hold their member particles, their total mass, their effect center of mass, and pointers to their children if they exist. The RootNode object is a special subclass of QuadNode, containing special initialization routines since it is the only node not generated by calling QuadNode.split() on a parent node. The QuadNode.split() and QuadNode.calculate() methods are recursive, so calling them on the root node initializes the whole tree.

The tree is recreated at each evolution step, used for calculation of the net force on each particle, and then deleted for minimal memory usage. The repeated access to this tree is an important optimization concern. When implementing this algorithm in a lower-level language like C, cache-optimizing the QuadTree could lead to huge performance improvements.

1b Opening angle

The Barnes-Hut quadtree algorithm uses an opening angle θ to determine the level of approximation. When computing the force that the particles inside a given quadtree node exert on a test particle, we compare the size l of the quadtree node (an upper bound on the distances between the particles in this group) to the distance d from the group of particles to the test particle. If $l/d < \theta$, the group is considered far enough away that it can be approximated as a single particle.

Thus, $\theta = 0$ gives us the exact brute-force calculation, and higher θ means fewer calculations through more approximation. We aim to find a suitable θ that balances speed with accuracy, as well as study the behavior of the system with θ .

2 Qualitative observations

Qualitative observations were conducted on a system of N=30 particles visualized using vpython. Large variations in θ were used to emphasize visual differences; more careful study follows below.

 $\theta=1$ yielded a notable performance improvement over the brute-force approach, but the effects of clumping were quite visible. Particle motion seemed quite unnatural compared to the exact brute-force solution: it seemed as if there was an invisible larger mass in a small orbit near the center of the system influencing the motion of all the other particles. This behavior could be explained by considering the extreme $\theta\gg 1$ limit: all particles would be grouped, and each particle would be attracted to the center of mass of the entire system.

A short visualization with $\theta=5$ yielded complete disaster, with all particles quickly moving off-screen. This behavior is unsurprising given our examination of the $\theta\gg1$ limit: the particles' motion is highly coupled because all are attracted to the center of mass, so motion of the center of mass is amplified: particles follow the center of mass, moving the center of mass further in the direction in which it was originally going.

3 Quantitative analysis

Figure 1 shows the algorithm's wall-clock runtime for various opening angles. As expected, the performance increases dramatically with θ , although the log plot shows that the dependence is slower than exponential decay. Figure 2 shows the energy evolution for various θ .

As a measure of approximation error, we use an estimate of the average energy of the N=100 system over the computed period. This average is estimated by sampling at the timesteps plotted in Figure 2. Here, we find an optimal $\hat{\theta}=0.1$:

$$\langle E \rangle_{\theta=0} = -21500.289563 \qquad \langle E \rangle_{\theta=0.1} = -21400.060124 \qquad \% \text{ error} = 0.466\%.$$
 (1)

We have achieved a 13.4% performance boost with 0.5% error.

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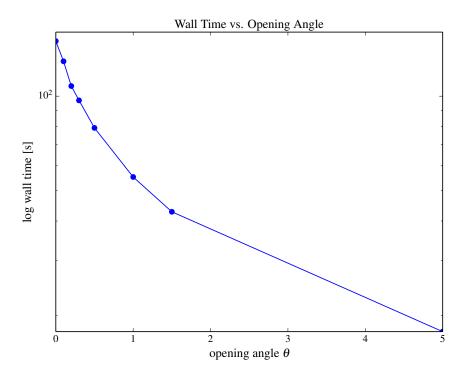


Figure 1: Wall time of the Barnes-Hut algorithm on N = 100 particles randomly distributed in an unit circle with unit speeds for various opening angles.

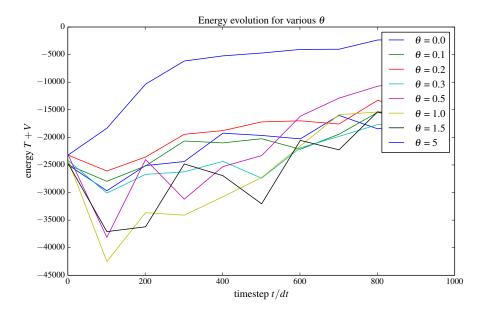


Figure 2: Evolution of the total energy of the gravitational system with N = 100 for various opening angles θ .