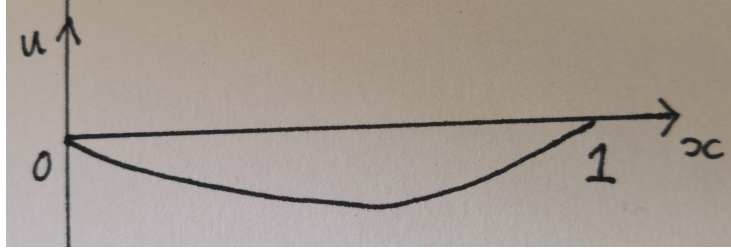


# 1 Notes on Green's Functions

## 1.1 Loaded String

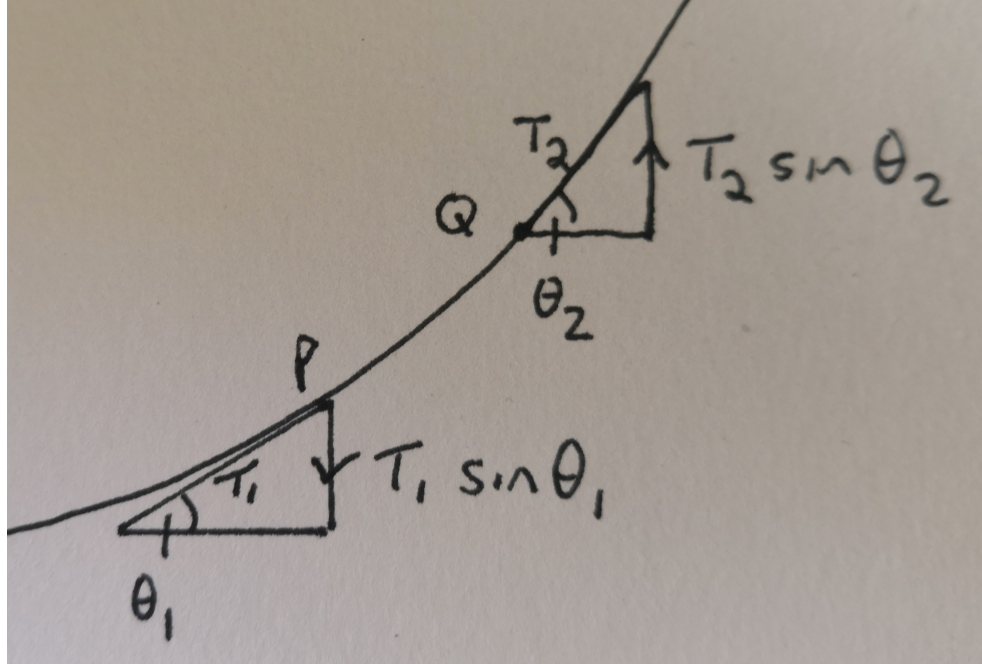


The equation for the displacement  $u(x)$  of a string under tension that is subject to a force per unit length  $\phi(x)$  is

$$u''(x) = \phi(x) \quad (1)$$

There are a number of available proofs of this but essentially the argument is as follows. Let us consider a string that is displaced downwards by a force.

Consider a small piece of the string  $dx$  between points  $P$  and  $Q$ . The tension in the string at  $P$  is  $T_1$  and the tension in the string at  $Q$  is  $T_2$ . The string is under tension so the forces at either end of  $dx$  pull in different directions.



Since the string element  $dx$  does not move horizontally, the horizontal component of the tension at point  $P$  and point  $Q$  must be equal (let us call this constant  $C$ ) i.e.

$$T_1 \cos \theta_1 = T_2 \cos \theta_2 = C \quad (2)$$

Additionally, because the string is stationary, the net vertical component upwards of the tension on element  $dx$  must be equal to the net force downwards on  $dx$  i.e.

$$T_2 \sin \theta_2 - T_1 \sin \theta_1 = \phi(x) dx \quad (3)$$

Dividing 3 by the constant  $C$  but using the appropriate term in 2

$$\frac{T_2 \sin \theta_2}{T_2 \cos \theta_2} - \frac{T_1 \sin \theta_1}{T_1 \cos \theta_1} = \frac{\phi(x) dx}{C}$$

$$\tan \theta_2 - \tan \theta_1 = \frac{\phi(x) dx}{C}$$

But the tangent is just the gradient.

$$\frac{du}{dx} \big|_Q - \frac{du}{dx} \big|_P = \frac{\phi(x) dx}{C}$$

$$u'_Q - u'_P = \frac{\phi(x) dx}{C}$$

$$\frac{(u'_Q - u'_P)}{dx} = \frac{\phi(x)}{C}$$

Including the constant  $C$  into the function  $\phi$ ,

$$\frac{(u'_Q - u'_P)}{dx} = \phi(x) \quad (4)$$

we get equation 1.

Note the embedded constant  $C$  on the RHS. Intuitively this must exist as the displacement of the string must be related to the tension in the string.

For the following sections I'll use the form

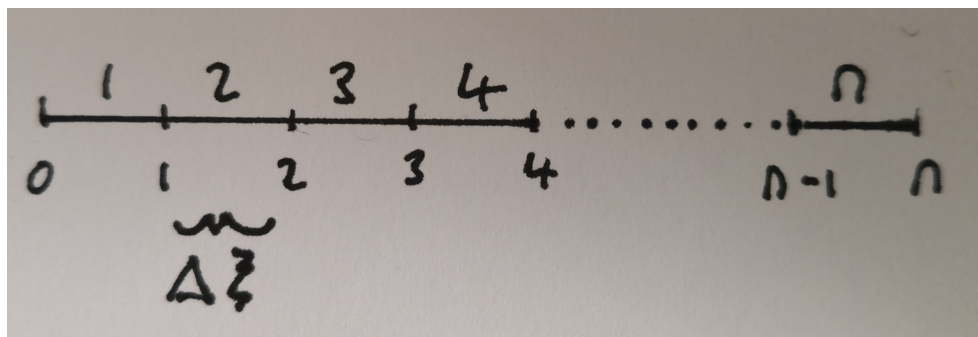
$$u'_Q - u'_P = \phi(x) dx \quad (5)$$

## 1.2 Unit Force

Let the length of the string be 1 and be fixed at both ends.

$$u(0) = u(1) = 0$$

We'll consider  $n$  segments of the string. Each segment then has length  $= \frac{1}{n}$ . Let us call this length  $\Delta\xi$ . A segment is identified using the index  $i$  where  $i = 1..n$  and is the interval  $[i-1, i]$

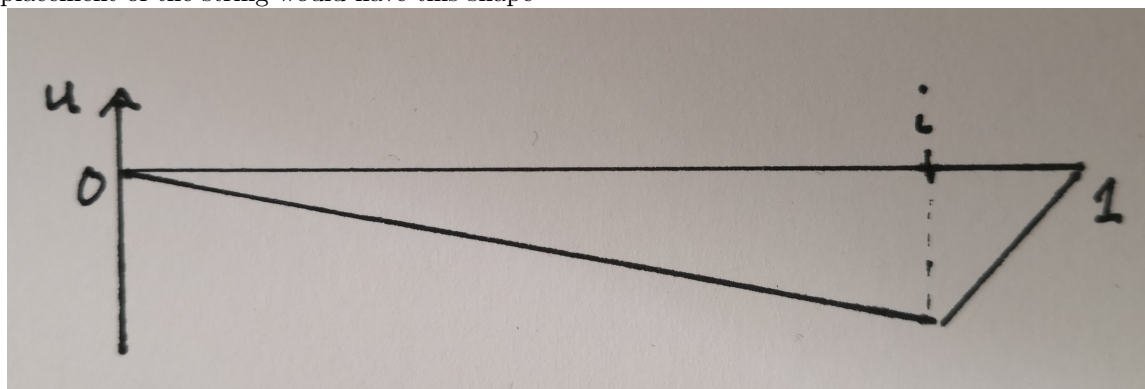


From 5 we have for the  $i^{th}$  segment

$$\Delta u'_i = u'_i - u'_{i-1} = \phi[i-1, i] \Delta \xi$$

From this and the derivation above we can see the change in gradient at segment  $i$  only depends on the force per unit length  $\phi$  over the interval  $[i-1, i]$ .

If we were to make  $\phi$  equal to 0 *outside* the interval  $[i-1, i]$  then the RHS becomes 0 and  $\Delta u' = 0$  outside the interval  $[i-1, i]$ . This would imply the displacement of the string would have this shape



From 5 we have

$$u'_i - u'_{i-1} = \phi[i-1, i] \Delta \xi = F \quad (6)$$

So if the gradient at  $i-1$  is  $m_1$  and the gradient at  $i$  is  $m_2$  then from 6 we must have

$$m_2 - m_1 = F$$

$$\implies m_1 = m_2 - F$$

We also know that the 2 slopes meet at a point in the interval  $[i-1, i]$ . Let us call this point  $\xi$ . Also the length of the string is 1.  $u$  is displaced in the -ve direction so

$$m_1 = \frac{\Delta u}{\xi}$$

$$\begin{aligned}
m_2 &= -\frac{\Delta u}{1-\xi} \\
\Rightarrow \frac{m_1}{m_2} &= -\frac{1-\xi}{\xi} = \frac{\xi-1}{\xi} \\
\Rightarrow \frac{m_2-F}{m_2} &= \frac{\xi-1}{\xi} \\
\Rightarrow -\frac{F}{m_2} &= \frac{\xi-1}{\xi} - 1 = \frac{\xi-1-\xi}{\xi} = -\frac{1}{\xi} \\
m_2 &= F\xi
\end{aligned}$$

and so

$$m_1 = F\xi - F = F(\xi - 1)$$

For the first slope

$$u_1(x) = F(\xi - 1)x + c_1$$

but  $u(0) = 0$  so  $c_1 = 0$  so

$$u_1(x) = F(\xi - 1)x$$

For the second slope

$$u_2(x) = F\xi x + c_2$$

but  $u(1) = 0$  so  $c_2 = -F\xi$  so

$$u_2(x) = F\xi x - F\xi = F\xi(x - 1)$$

which gives

$$u(x) = F \begin{cases} (\xi - 1)x & x \leq \xi \\ (x - 1)\xi & x \geq \xi \end{cases}$$

For the special case  $F = \phi(\xi) \Delta\xi = 1$  then  $F$  is a unit force. Let us use  $G$  for the this case, rather than  $u$ .

$$G(x, \xi) = \begin{cases} (\xi - 1)x & x \leq \xi \\ (x - 1)\xi & x \geq \xi \end{cases}$$

so

$$u(x) = FG(x, \xi) = G(x, \xi) \phi(\xi) \Delta\xi$$

### 1.3 Symmetry of $G(\xi, x)$

$G(x, \xi)$  is the displacement at  $x$  due to a unit force at  $\xi$ .

$$G(x, \xi) = \begin{cases} (\xi - 1)x & x \leq \xi \\ (x - 1)\xi & x \geq \xi \end{cases}$$

$G(\xi, x)$  is the displacement at  $\xi$  due to a unit force at  $x$ .

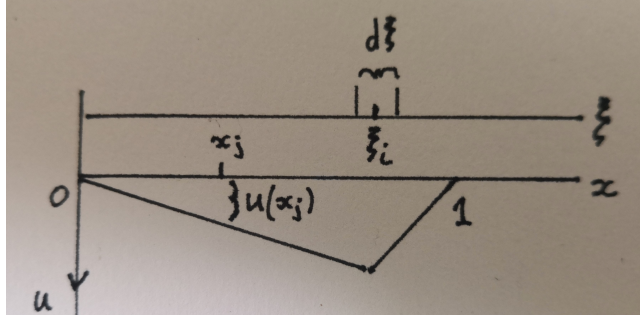
$$G(\xi, x) = \begin{cases} (x - 1)\xi & \xi \leq x \\ (\xi - 1)x & \xi \geq x \end{cases}$$

We can see from the above that  $G(\xi, x) = G(x, \xi)$  i.e. the displacement at  $\xi$  due to unit force at  $x$  is the same as the displacement at  $x$  due to a unit force at  $\xi$ . Basically because if  $\xi < x$  then  $x > \xi$ .

### 1.4 Superposition

$$G(x, \xi) = \begin{cases} (\xi - 1)x & x \leq \xi \\ (x - 1)\xi & x \geq \xi \end{cases}$$

We can envisage this as below. A unit force  $f$  at  $\xi_i$  over an interval  $d\xi$  so that  $f(\xi_i) d\xi = 1$  and outside the interval  $d\xi$ ,  $f = 0$ . This causes a deflection  $u$  at a point  $x_j$



Consider applying a force to the  $i^{th}$  segment. Let the force on this segment be  $\phi(\xi_i) \Delta\xi$ . The deflection of the string at the point  $x_j$  is given by

$$u(x_j) = G(x_j, \xi_i) \phi(\xi_i) \Delta\xi$$

and because  $G(x_j, \xi_i) = G(\xi_i, x_j)$

$$u(x_j) = G(\xi_i, x_j) \phi(\xi_i) \Delta\xi$$

The total deflection at point  $x_j$  due to forces being applied over the complete range of  $[0, 1]$  is

$$u(x_j) = \sum_{i=1}^n G(\xi_i, x_j) \phi(\xi_i) \Delta\xi$$

In the limit

$$u(x_j) = \int_0^1 G(\xi, x_j) \phi(\xi) d\xi$$

And more generally

$$u(x) = \int_0^1 G(\xi, x) \phi(\xi) d\xi$$