

CMP工作示意图

## 具有对流效应的CMP润滑方程的推导

连续方程为：

$$\nabla \cdot \mathbf{u} = 0 \quad (1)$$

该方程等价于

$$\frac{\partial u_1}{\partial x} + \frac{\partial u_2}{\partial y} + \frac{\partial u_3}{\partial z} = 0 \quad (2)$$

运动方程为：

$$\rho \frac{\partial \mathbf{u}}{\partial t} + \rho \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \eta \Delta \mathbf{u} + \rho \mathbf{F} \quad (3)$$

当假设  $\mathbf{F} = \mathbf{0}$  时，该方程等价于

$$\begin{cases} \rho \left( \frac{\partial u_1}{\partial t} + u_1 \frac{\partial u_1}{\partial x} + u_2 \frac{\partial u_1}{\partial y} + u_3 \frac{\partial u_1}{\partial z} \right) = -\frac{\partial p}{\partial x} + \eta \left( \frac{\partial^2 u_1}{\partial x^2} + \frac{\partial^2 u_1}{\partial y^2} + \frac{\partial^2 u_1}{\partial z^2} \right) \\ \rho \left( \frac{\partial u_2}{\partial t} + u_1 \frac{\partial u_2}{\partial x} + u_2 \frac{\partial u_2}{\partial y} + u_3 \frac{\partial u_2}{\partial z} \right) = -\frac{\partial p}{\partial y} + \eta \left( \frac{\partial^2 u_2}{\partial x^2} + \frac{\partial^2 u_2}{\partial y^2} + \frac{\partial^2 u_2}{\partial z^2} \right) \\ \rho \left( \frac{\partial u_3}{\partial t} + u_1 \frac{\partial u_3}{\partial x} + u_2 \frac{\partial u_3}{\partial y} + u_3 \frac{\partial u_3}{\partial z} \right) = -\frac{\partial p}{\partial z} + \eta \left( \frac{\partial^2 u_3}{\partial x^2} + \frac{\partial^2 u_3}{\partial y^2} + \frac{\partial^2 u_3}{\partial z^2} \right) \end{cases}$$

### 假设条件

1. CMP过程中载荷完全由抛光液来承载，晶片由于抛光液的隔离完全与抛光垫分离，因为膜厚比较薄，忽略质量力。
2. 假设抛光液是牛顿性流体，在CMP过程中是稳态不可压的，流体的黏性系数是常数  $\eta$ 。由于稳态所以  $\frac{\partial u_i}{\partial t} = 0$ 。
3. 满足无量纲润滑条件。即  $\frac{\partial^2 u_i}{\partial x^2} + \frac{\partial^2 u_i}{\partial y^2}$  比  $\frac{\partial^2 u_i}{\partial z^2}$  量级小，可以忽略。
4. 在  $z$  轴方向总压力梯度为0， $z$  轴方向的速度忽略不计。即  $u_3 = 0$

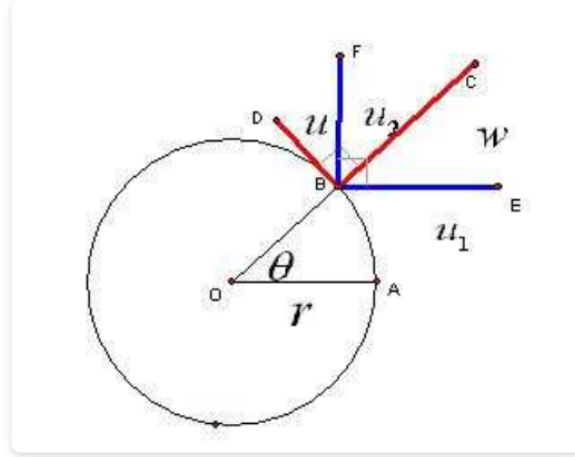
所以连续方程简化为

$$\frac{\partial u_1}{\partial x} + \frac{\partial u_2}{\partial y} = 0 \quad (5)$$

运动方程简化为

$$\begin{cases} \frac{\partial p}{\partial x} + \rho u_1 \frac{\partial u_1}{\partial x} + \rho u_2 \frac{\partial u_1}{\partial y} = \eta \frac{\partial^2 u_1}{\partial z^2} \\ \frac{\partial p}{\partial y} + \rho u_1 \frac{\partial u_2}{\partial x} + \rho u_2 \frac{\partial u_2}{\partial y} = \eta \frac{\partial^2 u_2}{\partial z^2} \\ \frac{\partial p}{\partial z} = 0 \end{cases} \quad (6)$$

上式的坐标系是直角坐标系，将其转化为柱坐标，则速度转换关系如下图



CMP速度关系式

变换关系式为:

$$\begin{cases} x = r \cos \theta \\ y = r \sin \theta \\ z = z \end{cases} \quad (7)$$

$$\begin{cases} \frac{\partial x}{\partial r} = \cos \theta & \frac{\partial x}{\partial \theta} = -r \cdot \sin \theta \\ \frac{\partial y}{\partial r} = \sin \theta & \frac{\partial y}{\partial \theta} = r \cdot \cos \theta \end{cases} \quad (8)$$

$$\begin{cases} \frac{\partial r}{\partial x} = \cos \theta & \frac{\partial \theta}{\partial x} = -\frac{\sin \theta}{r} \\ \frac{\partial r}{\partial y} = \sin \theta & \frac{\partial \theta}{\partial y} = \frac{\cos \theta}{r} \end{cases} \quad (9)$$

速度关系式

$$\begin{cases} u_1 = \omega \cdot \cos \theta - u \cdot \sin \theta \\ u_2 = \omega \cdot \sin \theta + u \cdot \cos \theta \end{cases} \quad (10)$$

于是可得(5)与(6)式的柱坐标形式如下:

$$\begin{aligned}
\frac{\partial rw}{\partial r} + \frac{\partial u}{\partial \theta} &= 0 \\
\frac{\partial p}{\partial r} + \rho w \frac{\partial w}{\partial r} + \rho \frac{u}{r} \frac{\partial w}{\partial \theta} - \rho \frac{u^2}{r} &= \eta \frac{\partial^2 w}{\partial z^2} \\
\frac{1}{r} \frac{\partial p}{\partial \theta} + \rho w \frac{\partial u}{\partial r} + \rho \frac{u}{r} \frac{\partial u}{\partial \theta} + \rho \frac{wu}{r} &= \eta \frac{\partial^2 u}{\partial z^2} \\
\frac{\partial p}{\partial z} &= 0
\end{aligned}$$

(10)

p只与x,y相关, 即 $p=p(r,\theta)$ , 记

$$\begin{aligned}
F(w, u) &= \rho w \frac{\partial w}{\partial r} + \rho \frac{u}{r} \frac{\partial w}{\partial \theta} - \rho \frac{u^2}{r} \\
G(w, u) &= \rho w \frac{\partial u}{\partial r} + \rho \frac{u}{r} \frac{\partial u}{\partial \theta} + \rho \frac{wu}{r}
\end{aligned}$$

(11)

对(10)的中间两项在 $[0, z]$ 上进行两次积分可得

$$\left. \begin{aligned}
w &= \frac{z^2}{2\eta} \frac{\partial p}{\partial r} + \frac{1}{\eta} \int_0^z \int_0^z F(w, u) dz dz + \\
&\quad \left. \frac{\partial w}{\partial z} \right|_{z=0} z + w_0 \\
u &= \frac{z^2}{2r\eta} \frac{\partial p}{\partial \theta} + \frac{1}{\eta} \int_0^z \int_0^z G(w, u) dz dz + \\
&\quad \left. \frac{\partial u}{\partial z} \right|_{z=0} z + u_0
\end{aligned} \right\}$$

(12)

对一阶对z的偏导数进行泰勒展式:

$$\left. \begin{aligned}
\frac{\partial w}{\partial z} \Big|_{z=0} &\approx \frac{w_h - w_0}{h} - \frac{h}{2} \frac{\partial^2 w}{\partial z^2} \Big|_{z=0} \\
\frac{\partial u}{\partial z} \Big|_{z=0} &\approx \frac{u_h - u_0}{h} - \frac{h}{2} \frac{\partial^2 u}{\partial z^2} \Big|_{z=0}
\end{aligned} \right\}$$

(13)

将(10)中间两项代入(13), (13)代入(12), 然后再将(12)代入到(10)的第一个公式中, 则有

$$\begin{aligned}
&\frac{\partial}{\partial r} \left[ r \left( \frac{z^2 - hz}{2\eta} \frac{\partial p}{\partial r} + \frac{w_h - w_0}{h} z + w_0 + \right. \right. \\
&\quad \left. \left. \frac{1}{\eta} \int_0^z \int_0^z F(w, u) dz dz - \frac{hz}{2\eta} F(w_0, u_0) \right) \right] + \\
&\frac{\partial}{\partial \theta} \left[ \frac{z^2 - hz}{2r\eta} \frac{\partial p}{\partial r} + \frac{u_h - u_0}{h} z + u_0 + \right. \\
&\quad \left. \frac{1}{\eta} \int_0^z \int_0^z G(w, u) dz dz - \frac{hz}{2\eta} G(w_0, u_0) \right] = 0
\end{aligned}$$

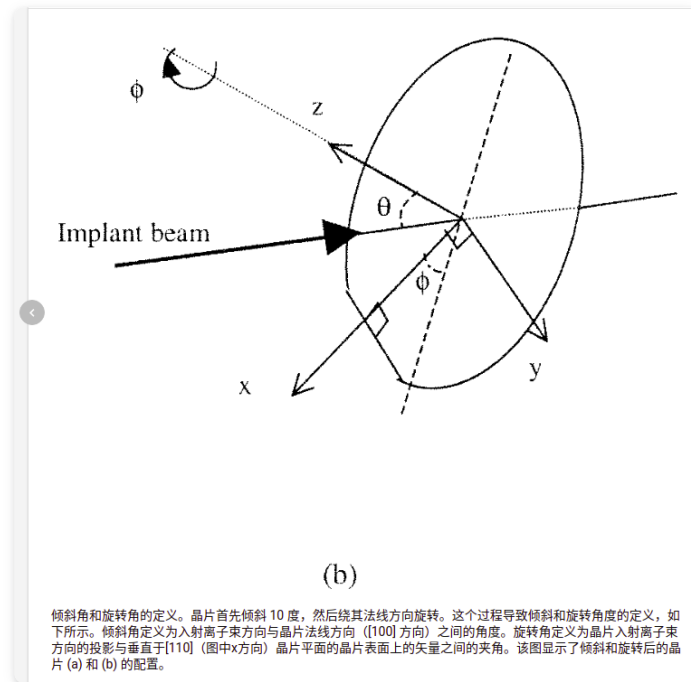
(14)

将(14)沿z轴[0,h]积分，积分过程中采用中矩形公式近似，于是得到了CMP润滑方程

$$\begin{aligned} \frac{\partial}{\partial r} \left( r h^3 \frac{\partial p}{\partial r} \right) + \frac{1}{r} \frac{\partial}{\partial \theta} \left( h^3 \frac{\partial p}{\partial \theta} \right) = \\ 6\eta \frac{\partial}{\partial r} [(w_h + w_0)h] + 6\eta \frac{\partial}{\partial \theta} [(u_h + u_0)h] + \\ \frac{\partial}{\partial r} [r h^3 (F(w_h, u_h) - 2F(w_0, u_0))] + \\ \frac{\partial}{\partial \theta} [h^3 (G(w_h, u_h) - 2G(w_0, u_0))] \end{aligned}$$

记晶片中心厚度为 $h_{piv}$ ，晶片转角和倾角分别为 $\alpha$   $\beta$ ，则在 $(r,\theta)$ 处膜厚 $h$ 满足

$$h = h_{piv} - r \sin \alpha \cos \theta - r \sin \beta \sin \theta$$

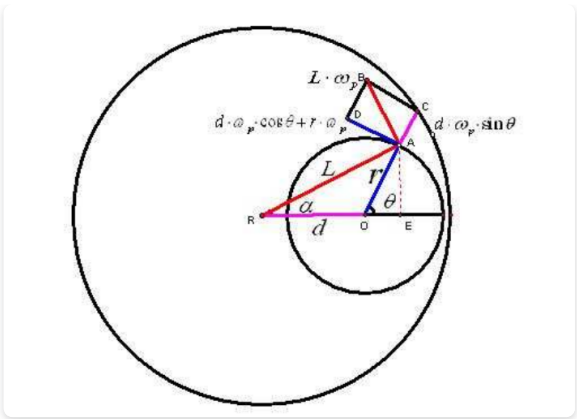


转角倾角示意图

速度边界条件满足如下：

$$\left. \begin{aligned} u &= (r + d \cos \theta) \omega_p \\ w &= d \sin \theta \omega_p, \quad z = 0 \\ u &= r \omega_w \\ w &= 0, \quad z = h \end{aligned} \right\}$$

其中 $\omega_w$ 为晶片的转速， $\omega_p$ 为抛光垫的转速， $d$ 为抛光垫中心到晶片中心的距离。



速度边界关系图

无量纲

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将一个物理导出量用若干个基本量的乘方之积表示出来的表达式，称为该物理量的量纲式，简称量纲。它是在选定了单位制之后，由[基本物理量](#)单位表达的式子。有量纲的物理量都可以进行无量纲化处理。

中文名	无量纲	意思	没有单位的 <a href="#">物理量</a>
外文名	dimensionless	成因	两个 <a href="#">量纲</a> 相同的物理量的比值

无量纲百度解释

形成流体动力润滑的条件是

1.两工件之间的间隙必须有楔形间隙。

2.两工件表面之间必须连续充满润滑油或其它液体。

3.两工件表面必须有相对滑动速度。其运动方向必须保证润滑油从大截面流进，从小截面出来

润滑条件百度解释