### The Watrous Post-Quantum Zero-Knowledge Proof

A Crypto Reading Group Talk

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### Post-Quantum ZK for NP

#### The model:

- ► Classical P and V
- ZK system for NP languages
- $ightharpoonup V^*$  can be quantum.
  - ▶ Modeled as a quantum polynomial-time (QPT) Turing machine.
  - equivalently (and more preferred in quantum-computing literature), poly-size quantum circuits.
  - Non-uniformity:  $V^*$  has an auxiliary quantum state that depends only on the security para. n. More accurately,

$$V^* = \{ \mathsf{QC}_n, |\psi_n\rangle \}_{n \in \mathbb{N}}$$

### Post-Quantum (Black-Box) ZK Is Hard

### Why's **rewinding** hard?

- ▶ information gain VS state disturbance
- ▶ the no-cloning theorem

The major result in [Wat06]: a quantum rewinding lemma

### Some Historical Notes

Techniques inspired by Marriot-Watrous [MW04]

error-gap amplification for QMA using only 1 witness state

### First published at STOC'06 [Wat06]

- ► Explicit connection to [MW04]
- Simple, ad hoc proof
- ► This talk mainly focuses on this version
- ► The notation herein is consistent with this version

#### Then, on SIAM Journal of Computing in 2009 [Wat09]

- ► Abstracts out a general quantum rewinding lemma
- ► Hides the connection with Marriot-Watrous
- ▶ We'll also see the high-level idea of this version

# Agenda for Today

- ▶ Prove quantum ZK for the Graph Isomorphism protocol [GMW86] (in detail)
  - ► Originally ad hoc [Wat06]
  - ► We'll take a general perspective
- ► Extends to the Graph-3-coloring Protocol [GMW86] in the ideal Com model (simple)
  - General quantum rewinding lemma
- ► G3C ZK with computationally-secure Com (simple-yet-tedious)
  - Rewinding lemma in its most general form allowing small perturbations
  - the widely-used version in crypto literature

# GMW ZK for Graph Isomorphism (GI)

#### Some Remarks:

- ► GI is not known to be NP-complete.
- ▶ the 1st message of the GMW GI protocol is perfectly uniform.

Input for P: statement  $(G_0, G_1) \in \mathcal{G}_n \times \mathcal{G}_n$ , witness  $w = \sigma$  s.t.  $\sigma(G_1) = G_0$ Input for V:  $(G_0, G_1)$ 

- 1. P samples  $\pi \leftarrow S_n$ , sends  $H = \pi(G_0)$
- 2. V sends  $a \leftarrow \{0, 1\}$
- 3. P sends  $\tau = \pi \circ \sigma^a$

V's decision: accept iff  $\tau(G_a) = H$ 

**Classical Sim:** guess the bit b. Set  $H = \pi(G_b)$ . Win if b == a.

# Modeling in Quantum Way

### **Model a Quantum** $V^*$ : circuit family $\{V_H\}_{H \in \mathcal{G}_n}$ , auxiliary input $|\psi\rangle$

- ▶ Receives *H* from P
- ► Perform  $V_H |\psi\rangle_W |0\rangle_V |0\rangle_A = \alpha_0 |\psi_0\rangle_{WV} |0\rangle_A + \alpha_1 |\psi_1\rangle_{WV} |1\rangle_A$ 
  - ► V: work space
  - ightharpoonup A: single-qubit register to store  $V^*$ 's challenge.
  - ▶ Note that  $V_H$  operates on space  $W \otimes V \otimes A$

# Modeling in Quantum Way

View the protocol through a quantum lens:

- ▶ The full space  $W \otimes X$ , where  $X = V \otimes A \otimes Y \otimes B \otimes Z$
- ► Sim performs (classical Sim in superposition)

$$\mathbf{T} |0\rangle_{\mathsf{YBZ}} = \frac{1}{\sqrt{2n!}} \sum_{b \in \{0,1\}} \sum_{\pi \in S_n} |\pi(G_b)\rangle_{\mathsf{Y}} |b\rangle_{\mathsf{B}} |\pi\rangle_{\mathsf{Z}}$$

- ▶ V apply  $\mathbf{V} = \sum_{H \in \mathcal{G}} \mathbf{V}_H \otimes |H\rangle\langle H|_{\mathsf{Y}} \otimes \mathbb{1}_{\mathsf{BZ}}$  on the full space  $\mathsf{W} \otimes \mathsf{X}$ 
  - ightharpoonup recall that  $V_H$  operates on  $|\psi\rangle_W\,|0\rangle_V\,|0\rangle_A$
  - corresponding to the exec. in super-position
  - Output format:

$$\alpha_{00}\left|\psi_{00}\right\rangle\left|00\right\rangle_{\mathsf{AB}}+\alpha_{01}\left|\psi_{01}\right\rangle\left|01\right\rangle_{\mathsf{AB}}+\alpha_{10}\left|\psi_{10}\right\rangle\left|10\right\rangle_{\mathsf{AB}}+\alpha_{11}\left|\psi_{11}\right\rangle\left|11\right\rangle_{\mathsf{AB}}$$

In summary, the protocol up to step 2 is:

$$\underbrace{\mathbf{VT}}_{\text{on } \mathbb{W} \otimes \mathsf{X}} (|\psi\rangle_{\mathsf{W}} |0\rangle_{\mathsf{X}=\mathsf{VAYBZ}}) \quad \Leftrightarrow \quad \underbrace{\mathbf{VT}(\mathbb{1}_{\mathsf{W}} \otimes |0\rangle_{\mathsf{X}})}_{\text{only on } \mathbb{W}} |\psi\rangle \tag{1}$$

## Measuring the Guess

Define a binary-outcome measurement on the full space  $W \otimes X$ :

- $\blacksquare \mathbf{\Pi}_0 = |0\rangle\langle 0|_{\mathsf{AB}}, \quad \mathbf{\Pi}_1 \coloneqq \mathbb{1}_{\mathsf{AB}} |0\rangle\langle 0|_{\mathsf{AB}} = |1\rangle\langle 1|_{\mathsf{AB}}$
- $\blacktriangleright$  work on the full space W  $\otimes$  X. Just tensor identities on registers other than AB

### Performing $\{\Pi_0, \Pi_1\}$ on $\mathbf{VT} |\psi\rangle_{\mathsf{W}} |0\rangle_{\mathsf{X}}$ :

- w.p. Tr  $(\langle \psi | \mathbf{Q} | \psi \rangle)$ , the outcome is 0.
- w.p.  $\operatorname{Tr}\left(\left\langle \psi\right|\left(\mathbb{1}_{\mathsf{W}}-\mathbf{Q}\right)|\psi\right)$ , the outcome is 1.

where 
$$\mathbf{Q}=(\mathbb{1}_{\mathsf{W}}\otimes\langle 0|_{\mathsf{X}})\mathbf{T}^{\dagger}\mathbf{V}^{\dagger}\mathbf{\Pi}_{0}\mathbf{T}\mathbf{V}(\mathbb{1}_{\mathsf{W}}\otimes|0\rangle_{\mathsf{X}})$$
. (See Expression (1).)

#### Two important facts:

- $ightharpoonup \{Q, \mathbb{1}_W Q\}$  form a POVM
- ► Tr  $(\langle \psi | \mathbf{Q} | \psi \rangle)$  = Tr  $(\langle \psi | (\mathbb{1}_{\mathsf{W}} \mathbf{Q}) | \psi \rangle)$  =  $\frac{1}{2}$ , independent of  $|\psi\rangle$ . (Cuz 1st msg. of GI prot. is perfectly uniform.)

$$\Rightarrow \mathbf{Q} = \mathbb{1}_{\mathsf{W}} - \mathbf{Q} = \frac{1}{2} \mathbb{1}_{\mathsf{W}}$$

### An Important Lemma

Let  $\Delta_0 := \mathbb{1}_{\mathsf{W}} \otimes |0\rangle\langle 0|_{\mathsf{X}}$ .

- $ightharpoonup \Delta_0$  projects register X to all-0 qubits.
- $lackbox{\Delta}_0 = oldsymbol{\Delta}_0^{\dagger}$
- $ightharpoonup \Delta_1 := \mathbb{1}_{\mathsf{WX}} \Delta_0$ . The  $\{\Delta_0, \Delta_1\}$  form a POVM.

#### LEMMA 1:

For all  $|\psi\rangle \in \mathcal{H}(W)$ ,  $|\gamma_0\rangle = |\psi\rangle_W |0\rangle_X$  is an eigenvector of  $\underbrace{\Delta_0^{\dagger} \mathbf{T}^{\dagger} \mathbf{V}^{\dagger} \mathbf{\Pi}_0 \mathbf{T} \mathbf{V} \Delta_0}_{:=\mathbf{M}}$  with corresponding eigenvalue  $\lambda = 1/2$ .

**Proof.** Recall  $\mathbf{Q} = (\mathbb{1}_{\mathsf{W}} \otimes \langle 0|_{\mathsf{X}}) \mathbf{T}^{\dagger} \mathbf{V}^{\dagger} \mathbf{\Pi}_{0} \mathbf{T} \mathbf{V} (\mathbb{1}_{\mathsf{W}} \otimes |0\rangle_{\mathsf{X}}) = \frac{1}{2} \mathbb{1}_{\mathsf{W}}.$ 

$$\Rightarrow \quad \mathbf{\Delta}_0^{\dagger} \mathbf{T}^{\dagger} \mathbf{V}^{\dagger} \mathbf{\Pi}_0 \mathbf{T} \mathbf{V} \mathbf{\Delta}_0 = (\mathbb{1}_{\mathsf{W}} \otimes |0\rangle_{\mathsf{X}}) \mathbf{Q} (\mathbb{1}_{\mathsf{W}} \otimes \langle 0|_{\mathsf{X}}) = \frac{1}{2} \mathbb{1}_{\mathsf{W}} \otimes |0\rangle \langle 0|_{\mathsf{X}}$$

$$\Rightarrow \forall |\psi\rangle, \mathbf{\Delta}_{0}^{\dagger}\mathbf{T}^{\dagger}\mathbf{V}^{\dagger}\mathbf{\Pi}_{0}\mathbf{T}\mathbf{V}\mathbf{\Delta}_{0}\underbrace{|\psi\rangle_{\mathsf{W}}|0\rangle_{\mathsf{X}}}_{|\gamma_{0}\rangle} = (\frac{1}{2}\mathbb{1}_{\mathsf{W}}\otimes|0\rangle\langle 0|_{\mathsf{X}})\underbrace{|\psi\rangle_{\mathsf{W}}|0\rangle_{\mathsf{X}}}_{|\gamma_{0}\rangle} = \frac{1}{2}\underbrace{|\psi\rangle_{\mathsf{W}}|0\rangle_{\mathsf{X}}}_{|\gamma_{0}\rangle}$$

### Marriot-Watrous Lemma

#### Lemma 2: Marriot-Watrous [MW04]

Given unitary  $\mathbf{U}$ , proj. mnt.  $\{\mathbf{\Pi}_0, \mathbf{\Pi}_1\}$  and  $\{\boldsymbol{\Delta}_0, \boldsymbol{\Delta}_1\}$ . Assume  $|\gamma_0\rangle$  is an evec. of  $\boldsymbol{\Delta}_0\mathbf{U}^{\dagger}\mathbf{\Pi}_0\mathbf{U}\boldsymbol{\Delta}_0$  with eval.  $\lambda$ . Define

$$|\delta_0\rangle \coloneqq \frac{\mathbf{\Pi}_0\mathbf{U}\,|\gamma_0\rangle}{\sqrt{\lambda}}, \ |\delta_1\rangle \coloneqq \frac{\mathbf{\Pi}_0\mathbf{U}\,|\gamma_0\rangle}{\sqrt{1-\lambda}}, \ |\gamma_1\rangle \coloneqq \frac{\mathbf{\Delta}_1\mathbf{U}^\dagger\,|\delta_0\rangle}{\sqrt{1-\lambda}}.$$

Then,  $\langle \gamma_0 | \gamma_1 \rangle = \langle \delta_0 | \delta_1 \rangle = 0$  and

$$\mathbf{U} |\gamma_0\rangle = \sqrt{\lambda} |\delta_0\rangle + \sqrt{1-\lambda} |\delta_1\rangle \qquad \qquad \mathbf{U}^{\dagger} |\delta_0\rangle = \sqrt{\lambda} |\gamma_0\rangle + \sqrt{1-\lambda} |\gamma_1\rangle$$

$$\mathbf{U} |\gamma_1\rangle = \sqrt{1-\lambda} |\delta_0\rangle - \sqrt{\lambda} |\delta_1\rangle \qquad \qquad \mathbf{U}^{\dagger} |\delta_1\rangle = \sqrt{1-\lambda} |\gamma_0\rangle - \sqrt{\lambda} |\gamma_1\rangle$$

### (draw the evolution diagram)

$$|\gamma_0\rangle$$

$$|\delta_0\rangle$$

$$|\gamma_0\rangle$$

$$|\delta_0\rangle$$

$$|\gamma_0\rangle$$
 ...

$$|\delta_1\rangle$$

$$|\gamma_1\rangle$$

$$|\delta_1\rangle$$

$$\gamma_0\rangle$$

## In Our Setting: Marriot-Watrous + Post-Mnt. Selection

In our setting, we have  $\mathbf{U}=\mathbf{VT},$   $\lambda=1/2,$  and  $|\gamma_0\rangle=|\psi\rangle_{\mathsf{W}}\,|0\rangle_{\mathsf{X}}$ 

Lemma 2  $\Rightarrow$   $|\gamma_0\rangle = \frac{1}{\sqrt{2}} |\delta_0\rangle + \frac{1}{\sqrt{2}} |\delta_1\rangle$ , and the following:

$$|\delta_0\rangle = \sqrt{2}\mathbf{\Pi}_0\mathbf{V}\mathbf{T}|\gamma_0\rangle, \quad \mathbf{T}^{\dagger}\mathbf{V}^{\dagger}|\delta_1\rangle = \frac{1}{\sqrt{2}}|\gamma_0\rangle - \frac{1}{\sqrt{2}}|\gamma_1\rangle, \quad \mathbf{V}\mathbf{T}(\frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle) = |\delta_0\rangle$$

Starting with  $|\gamma_0\rangle \to VT |\gamma_0\rangle \to \text{measurement } \{\Pi_0, \Pi_1\}$ :

- $\blacktriangleright$  w.p. 1/2, it is  $|\delta_0\rangle$  we are done!
- w.p. 1/2, it is  $|\delta_1\rangle$ 
  - Key observation:  $\mathbf{T}^{\dagger}\mathbf{V}^{\dagger} |\delta_1\rangle = \frac{1}{\sqrt{2}} |\gamma_0\rangle \frac{1}{\sqrt{2}} |\gamma_1\rangle$
  - ► If we can flip the phase of the 2nd term  $\Rightarrow \frac{1}{\sqrt{2}} |\gamma_0\rangle + \frac{1}{\sqrt{2}} |\gamma_1\rangle$ .
  - ► Then, simply do  $VT(\frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle) = |\dot{\delta}_0\rangle$

Yes, we can! (next slide)

### Phase Flip for the 2nd Term

We want:  $\frac{1}{\sqrt{2}} |\gamma_0\rangle - \frac{1}{\sqrt{2}} |\gamma_1\rangle \rightarrow \frac{1}{\sqrt{2}} |\gamma_0\rangle + \frac{1}{\sqrt{2}} |\gamma_1\rangle$ 

Recall the following

$$|\gamma_0\rangle = |\psi\rangle_W |0\rangle_X$$
 and  $\Delta_0 = \mathbb{1}_W \otimes |0\rangle\langle 0|_X$ 

$$\blacktriangleright \Rightarrow \Delta_0 |\gamma_0\rangle = |\gamma_0\rangle$$

Lemma 2 says 
$$|\gamma_1\rangle = \sqrt{2}\Delta_1 \mathbf{T}^{\dagger} \mathbf{V}^{\dagger} |\delta_0\rangle \implies \Delta_0 |\gamma_1\rangle = 0$$

Therefore, it is not hard to come up with the following idea:

$$\underbrace{\frac{(2\Delta_0 - \mathbb{1}_{\text{WX}})}{=\Delta_0 - \Delta_1}} (\frac{1}{\sqrt{2}} | \gamma_0 \rangle - \frac{1}{\sqrt{2}} | \gamma_1 \rangle) = \frac{2}{\sqrt{2}} \Delta_0 | \gamma_0 \rangle - \frac{2}{\sqrt{2}} \Delta_0 | \gamma_1 \rangle - \frac{1}{\sqrt{2}} | \gamma_0 \rangle + \frac{1}{\sqrt{2}} | \gamma_1 \rangle$$

$$= \frac{2}{\sqrt{2}} | \gamma_0 \rangle - 0 - \frac{1}{\sqrt{2}} | \gamma_0 \rangle + \frac{1}{\sqrt{2}} | \gamma_1 \rangle$$

$$= \frac{1}{\sqrt{2}} | \gamma_0 \rangle + \frac{1}{\sqrt{2}} | \gamma_1 \rangle$$

## Summarizing the Watrous Simulator

- Start with  $|\gamma_0\rangle_{XW} = |\psi\rangle_{X} |0\rangle_{W}$
- ightharpoonup Perform  $VT |\gamma_0\rangle_{XW}$
- ▶ Perform measurement  $\{\Pi_0, \Pi_1\}$ 
  - ▶ If outcome is 0 guessed correctly (in  $|\delta_0\rangle$ ). Go next step.
  - Otherwise, we are in  $|\delta_1\rangle = \sqrt{2}\Pi_1 \mathbf{V} \mathbf{T} |\gamma_0\rangle$ .
    - Perform  $\mathbf{T}^{\dagger}\mathbf{V}^{\dagger} |\delta_1\rangle = \frac{1}{\sqrt{2}} |\gamma_0\rangle \frac{1}{\sqrt{2}} |\gamma_1\rangle$
    - ▶ Perform  $(2\Delta_0 \mathbb{1}_{WX})(\frac{1}{\sqrt{2}}|\gamma_0\rangle \frac{1}{\sqrt{2}}|\gamma_1\rangle) = \frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle$
    - Perform  $VT(\frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle) = |\delta_0\rangle$ . Go next step.
- Sim can finish the last round as the honest prover.

# Extending to G3C—Idealized Com Model (1/3)

- ► The graph-3-coloring (G3C) problem is NP-complete
- ► Start point: the G3C classical ZK proof from [GMW86]

#### Caveats:

- ▶  $\Pr[\text{Guess correctly}] = \frac{1}{m}$ , where m = # edges.
- ▶  $\Pr[Guess correctly] \perp |\psi\rangle$ ?
  - ▶ 1st msg using perfect-binding (PB) Com
  - ▶ What about binding? Collapse-binding suffices [Unr16]
  - We assume an ideal Com for simplicity: perfect-hiding and perfectly-binding
  - Extends to comp.-hiding Com later

# Extending to G3C—Idealized Com Model (2/3)

#### Key ingredients for the GI simulator:

- ▶ Define an operator:  $\Delta_0^{\dagger} \mathbf{T}^{\dagger} \mathbf{V}^{\dagger} \mathbf{\Pi}_0 \mathbf{T} \mathbf{V} \Delta_0$  (=: M)
- ► An technical Lemma 1:  $\lambda = \frac{1}{2} \quad (\bot \quad |\psi\rangle)$
- ► Invoke Marriot-Watrous Lemma 2 with  $\lambda = \frac{1}{2}$ :
  - ▶ Voilà  $\Theta$ ! We can get  $|\delta_0\rangle$  within  $\leq 2$  steps

#### What will change for the G3C protocol?

- ▶ M defined as before (T modified in the natural way)
- ▶ Lemma 1:  $\lambda = \frac{1}{m} \; (\bot \; |\psi\rangle)$
- ► Invoke Marriot-Watrous Lemma 2 with  $\lambda = \frac{1}{m}$ :
- ► Solution: use the full power of Matrriot-Watrous analysis (next slide).

# Extending to G3C—Idealized Com Model (3/3)

(draw the evolution diagram in the current setting)

$$|\gamma_0\rangle$$
  $|\delta_0\rangle$   $|\gamma_0\rangle$   $|\delta_0\rangle$   $|\gamma_0\rangle$  ...

$$|\delta_1\rangle$$
  $|\gamma_1\rangle$   $|\delta_1\rangle$   $|\gamma_0\rangle$  ...

The main take-away:

$$\begin{array}{c} \mathbf{U} \left| \gamma_0 \right> = \sqrt{\lambda} \left| \delta_0 \right> + \sqrt{1 - \lambda} \left| \delta_1 \right> & \mathbf{U}^{\dagger} \left| \delta_0 \right> = \sqrt{\lambda} \left| \gamma_0 \right> + \sqrt{1 - \lambda} \left| \gamma_1 \right> \\ \mathbf{U} \left| \gamma_1 \right> = \sqrt{1 - \lambda} \left| \delta_0 \right> - \sqrt{\lambda} \left| \delta_1 \right> & \mathbf{U}^{\dagger} \left| \delta_1 \right> = \sqrt{1 - \lambda} \left| \gamma_0 \right> - \sqrt{\lambda} \left| \gamma_1 \right> , \text{ where } \lambda = 1/m. \end{array}$$

- ▶ Measure  $\{\Pi_0, \Pi_1\}$  at each  $|\delta\rangle$ , if results in  $|\delta_1\rangle$ :
  - $\mathbf{U}(2\Delta_0 1)\mathbf{U}^{\dagger} |\delta_1\rangle = 2\sqrt{p(1-p)} |\delta_0\rangle + (1-2p) |\delta_1\rangle$
- ► Measure  $\{\Pi_0, \Pi_1\}$ . Go to  $|\delta_1\rangle$  w.p. (1-2p).
- ▶ Keep failing after t iteration:  $1 (1 p)(1 2p)^t$ . Then Chernoff bound.

## The General Quantum Rewinding Lemma (Exact)

### Lemma 3: Exact Quantum Rewinding [Wat09]

 ${f Q}$  is a QC works on  $|\psi\rangle$  and with  $\Pr[{\sf success}]=p\ (\perp\ |\psi\rangle)$  outputs  $|\delta_0\rangle$ . Then, for any  $\varepsilon>0$ , there exists another QC  ${f R}$  of size

$$O\left(\frac{\log(1/\varepsilon)}{p(1-p)} \cdot \mathsf{size}(\mathbf{Q})\right)$$

such that for every input  $|\psi\rangle$ , the output  $\rho$  of  $\mathbf{R}$  satisfies  $\mathrm{Tr}(\rho|\delta_0\rangle\langle\delta_0|) \geq 1 - \varepsilon$ .

- "Exact" refers to the face that  $p \perp |\psi\rangle$ .
- ► The  $\frac{\log(1/\varepsilon)}{p(1-p)}$  is due to Chernoff bound.
- ▶ Only need poly-size for a negligible  $\varepsilon$ .
- ► For the trace, the closer to 1, the better.

# G3C ZK with Comp.-Hiding Com

- ► Sim's 1st msg.  $\stackrel{c}{\approx}$  Prover's 1st msg.
- ▶  $V^*$ 's challenge  $a \not\perp$  the 1st msg.
- ▶ In Lemma 3,  $\Pr[\mathsf{success}] = p(|\psi\rangle)$ .
  - $ightharpoonup p(|\psi\rangle)$  jiggles within an negl. small interval.
- ▶ Need a version of Lemma 3 allowing small perturbations

### The Version Allowing Small Perturbations

### Lemma 4: Quantum Rewinding with Small Perturbations [Wat09, Sec. 4.2]

Let  $\mathbf{Q}$ ,  $|\psi\rangle$ , and  $|\delta_0\rangle$  as before. But  $\Pr[\mathrm{success}] = p(|\psi\rangle)$  now depends on  $|\psi\rangle$ . Let  $p_0, q \in (0,1)$  and  $\varepsilon \in (0,1/2)$  be real numbers such that

(1). 
$$|p(\psi) - q| < \varepsilon$$
 (2).  $p_0 \le p(\psi)$  (3).  $p_0(1 - p_0) \le q(1 - q)$ 

Then, for any  $\varepsilon > 0$ , there exists another QC  $\mathbf{R}$  of size  $O\left(\frac{\log(1/\varepsilon)}{p_0(1-p_0)} \cdot \operatorname{size}(\mathbf{Q})\right)$  such that for every input  $|\psi\rangle$ , the output  $\rho$  of  $\mathbf{R}$  satisfies:

$$\operatorname{Tr}(\rho|\delta_0\rangle\langle\delta_0|) \ge 1 - 16\varepsilon \frac{\log^2(1/\varepsilon)}{p_0^2(1-p_0)^2}.$$

#### Proof at a high-level:

- Consider each eigen-space separately (next slide).
- ► For detailed calculation, see [Wat09, Sec. 4.2].

### Proof Sketch for Lemma 4

Proof Sketch:

- ► In Lemma 1,  $|\gamma_0\rangle = |\psi\rangle_W |0\rangle_X$  is no longer an evec. of M
  - ► The reason:  $|\psi\rangle_{W}$  is not an evec. of **Q**
- (mental exper.) Thus, decomp.  $|\psi\rangle$  in the evecs  $\{|\psi_i\rangle\}_{i\in[\text{dim}]}$  of Q
- $\blacktriangleright$  (mental exper.) For each i, we obtain Lemmas 1 and 2
- ▶ (mental exper.) In the Marriot-Watrous procedure, in each egein space:

$$\mathbf{VT} |\psi_i\rangle_{\mathsf{W}} |0\rangle_{\mathsf{X}} = \sqrt{p(|\psi_i\rangle)} |\delta_0(|\psi_i\rangle)\rangle + \sqrt{1 - p(|\psi_i\rangle)} |\delta_1(|\psi_i\rangle)\rangle$$

- (mental exper.) Define a unitary N such that for all  $i \in [dim]$ :
  - $\sqrt{p(|\psi_i\rangle)} |\delta_0(|\psi_i\rangle)\rangle + \sqrt{1 p(|\psi_i\rangle)} |\delta_1(|\psi_i\rangle)\rangle \rightarrow \sqrt{q} |\delta_0(|\psi_i\rangle)\rangle + \sqrt{1 q} |\delta_1(|\psi_i\rangle)\rangle$
- ▶ (mental exper.) Ready to apply the Exact Rewinding Lemma 3 (w/  $p_0$  as we don't know p.) (Need  $p_0(1-p_0) \le q(1-q)$ .)

In summary, this is a Sim w/ an imaginary operator N, giving the same trace bound as in Lemma 3. But for the real Sim, there is no N.

- ▶ Doesn't matter. N only affects the trace bound negligibly.
- ▶ By tedious-yet-elementary linear algebra (see [Wat09, Sec. 4.2]).

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