

The Watrous Post-Quantum Zero-Knowledge Proof

A Crypto Reading Group Talk

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Post-Quantum ZK for NP

The model:

- ▶ Classical P and V
- ▶ ZK system for NP languages
- ▶ V^* can be quantum.
 - ▶ Modeled as a quantum polynomial-time (QPT) Turing machine.
 - ▶ equivalently (and more preferred in quantum-computing literature), poly-size quantum circuits.
 - ▶ Non-uniformity: V^* has an auxiliary quantum state that depends only on the security para. n . More accurately,

$$V^* = \{\text{QC}_n, |\psi_n\rangle\}_{n \in \mathbb{N}}$$

Post-Quantum (Black-Box) ZK Is Hard

Why's **rewinding** hard?

- ▶ information gain VS state disturbance
- ▶ the no-cloning theorem

The major result in [[Wat06](#)]: a quantum rewinding lemma

Some Historical Notes

Techniques inspired by Marriot-Watrous [[MW04](#)]

- ▶ error-gap amplification for QMA using only 1 witness state

First published at STOC'06 [[Wat06](#)]

- ▶ Explicit connection to [[MW04](#)]
- ▶ Simple, ad hoc proof
- ▶ This talk mainly focus on this version
- ▶ The notation herein is consistent with this version

Then, on SIAM Journal of Computing in 2009 [[Wat09](#)]

- ▶ Abstracts out a general quantum rewinding lemma
- ▶ Hides the connection with Marriot-Watrous
- ▶ we'll also see the high-level idea of this version

Agenda for Today

- ▶ Prove quantum ZK for the Graph Isomorphism protocol [GMW86] (in detail)
 - ▶ Originally ad hoc [Wat06]
 - ▶ We'll take a general perspective
- ▶ Extends to the Graph-3-coloring Protocol [GMW86] in the ideal Com model (simple)
 - ▶ General quantum rewinding lemma
- ▶ G3C ZK with computationally-secure Com (simple-yet-tedious)
 - ▶ Rewinding lemma in its most general form — allowing small perturbations
 - ▶ the widely-used version in crypto literature

GMW ZK for Graph Isomorphism (GI)

Some Remarks:

- ▶ GI is not known to be NP-complete.
- ▶ the 1st message of the GMW GI protocol is perfectly uniform.

Input for P : statement $(G_0, G_1) \in \mathcal{G}_n \times \mathcal{G}_n$, witness $w = \sigma$ s.t. $\sigma(G_1) = G_0$

Input for V : (G_0, G_1)

1. P samples $\pi \leftarrow S_n$, sends $H = \pi(G_0)$
2. V sends $a \leftarrow \{0, 1\}$
3. P sends $\tau = \pi \circ \sigma^a$

V 's decision: accept iff $\tau(G_a) = H$

Classical Sim: guess the bit a

Modeling in Quantum Way

Model a Quantum V^* : circuit family $\{\mathbf{V}_H\}_{H \in \mathcal{G}_n}$, auxiliary input $|\psi\rangle$

- ▶ receives H from P
- ▶ Perform $\mathbf{V}_H |\psi\rangle_W |0\rangle_V |0\rangle_A = \alpha_0 |\psi_0\rangle_{WV} |0\rangle_A + \alpha_1 |\psi_1\rangle_{WV} |1\rangle_A$
 - ▶ V : work space
 - ▶ A : single-qubit register to store V^* 's challenge.
 - ▶ Note that \mathbf{V}_H operates on space $W \otimes V \otimes A$

Modeling in Quantum Way

View the protocol through a quantum lens:

- ▶ The full space $W \otimes X$, where $X = V \otimes A \otimes Y \otimes B \otimes Z$
- ▶ P performs

$$\mathbf{T} |0\rangle_{YBZ} = \frac{1}{\sqrt{2n!}} \sum_{b \in \{0,1\}} \sum_{\pi \in S_n} |\pi(G_b)\rangle_Y |b\rangle_B |\pi\rangle_Z$$

- ▶ V apply $\mathbf{V} = \sum_{H \in \mathcal{G}} \mathbf{V}_H \otimes |H\rangle\langle H|_Y \otimes \mathbb{1}_{BZ}$ on the full space $W \otimes X$
 - ▶ recall that \mathbf{V}_H operates on $|\psi\rangle_W |0\rangle_V |0\rangle_A$
 - ▶ corresponding to the exec. in super-position
 - ▶ Output format:

$$\alpha_{00} |\psi_{00}\rangle |00\rangle_{AB} + \alpha_{01} |\psi_{01}\rangle |01\rangle_{AB} + \alpha_{10} |\psi_{10}\rangle |10\rangle_{AB} + \alpha_{11} |\psi_{11}\rangle |11\rangle_{AB}$$

In summary, the protocol up to step 2 is:

$$\underbrace{\mathbf{VT}}_{\text{on } W \otimes X} (|\psi\rangle_W |0\rangle_{X=VAYBZ}) \Leftrightarrow \underbrace{\mathbf{VT}(\mathbb{1}_W \otimes |0\rangle_X)}_{\text{only on } W} |\psi\rangle$$

Measuring the Guess

Define a binary-outcome measurement on the full space $W \otimes X$:

- ▶ $\Pi_0 = |0\rangle\langle 0|_{AB}$, $\Pi_1 := \mathbb{1}_{AB} - |0\rangle\langle 0|_{AB} = |1\rangle\langle 1|_{AB}$
- ▶ work on the full space $W \otimes X$. Just tensor identities on registers other than AB

Performing $\{\Pi_0, \Pi_1\}$ on $\mathbf{V}\mathbf{T} |\psi\rangle_W |0\rangle_X$ boils down to:

- ▶ the probability for outcome 0 is $\text{Tr}(\langle\psi| \mathbf{Q} |\psi\rangle)$
- ▶ the probability for outcome 1 is $\text{Tr}(\langle\psi| (\mathbb{1}_W - \mathbf{Q}) |\psi\rangle)$

where $\mathbf{Q} = (\mathbb{1}_W \otimes \langle 0|_X) \mathbf{T}^\dagger \mathbf{V}^\dagger \Pi_0 \mathbf{T} \mathbf{V} (\mathbb{1}_W \otimes |0\rangle_X)$.

Two important facts:

- ▶ $\{\mathbf{Q}, \mathbb{1}_W - \mathbf{Q}\}$ form a POVM
- ▶ $\text{Tr}(\langle\psi| \mathbf{Q} |\psi\rangle) = \text{Tr}(\langle\psi| (\mathbb{1}_W - \mathbf{Q}) |\psi\rangle) = \frac{1}{2}$, independent of $|\psi\rangle$. (Cuz 1st msg. of GI prot. is perfectly uniform.)

$$\Rightarrow \mathbf{Q} = \mathbb{1}_W - \mathbf{Q} = \frac{1}{2} \mathbb{1}_W$$

An Important Lemma

Let $\Delta_0 := \mathbb{1}_W \otimes |0\rangle\langle 0|_X$.

- ▶ Δ_0 projects register X to all-0 qubits.
- ▶ $\Delta_0 = \Delta_0^\dagger$
- ▶ $\Delta_1 := \mathbb{1}_{WX} - \Delta_0$. The $\{\Delta_0, \Delta_1\}$ form a POVM.

LEMMA 1:

For all $|\psi\rangle \in \mathcal{H}(W)$, $|\gamma_0\rangle = |\psi\rangle_W |0\rangle_X$ is an eigenvector of $\underbrace{\Delta_0^\dagger \mathbf{T}^\dagger \mathbf{V}^\dagger \Pi_0 \mathbf{T} \mathbf{V} \Delta_0}_{:=\mathbf{M}}$ with corresponding eigenvalue $\lambda = 1/2$.

Proof. Recall $\mathbf{Q} = (\mathbb{1}_W \otimes \langle 0|_X) \mathbf{T}^\dagger \mathbf{V}^\dagger \Pi_0 \mathbf{T} \mathbf{V} (\mathbb{1}_W \otimes |0\rangle_X) = \frac{1}{2} \mathbb{1}_W$.

$$\Rightarrow \Delta_0^\dagger \mathbf{T}^\dagger \mathbf{V}^\dagger \Pi_0 \mathbf{T} \mathbf{V} \Delta_0 = (\mathbb{1}_W \otimes |0\rangle_X) \mathbf{Q} (\mathbb{1}_W \otimes \langle 0|_X) = \frac{1}{2} \mathbb{1}_W \otimes |0\rangle\langle 0|_X$$

$$\Rightarrow \forall |\psi\rangle, \Delta_0^\dagger \mathbf{T}^\dagger \mathbf{V}^\dagger \Pi_0 \mathbf{T} \mathbf{V} \Delta_0 \underbrace{|\psi\rangle_W |0\rangle_X}_{|\gamma_0\rangle} = \left(\frac{1}{2} \mathbb{1}_W \otimes |0\rangle\langle 0|_X \right) \underbrace{|\psi\rangle_W |0\rangle_X}_{|\gamma_0\rangle} = \frac{1}{2} \underbrace{|\psi\rangle_W |0\rangle_X}_{|\gamma_0\rangle}$$

Marriot-Watrous Lemma

LEMMA 2: MARRIOT-WATROUS [MW04]

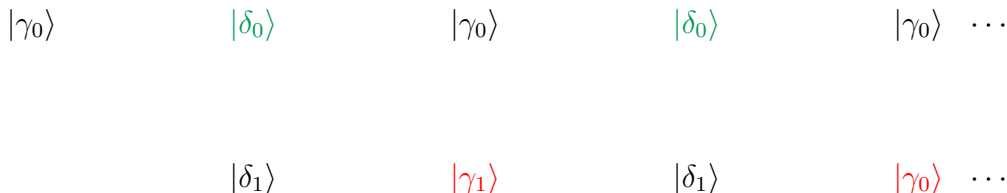
Given unitary \mathbf{U} , proj. mnt. $\{\Pi_0, \Pi_1\}$ and $\{\Delta_0, \Delta_1\}$. Assume $|\gamma_0\rangle$ is an evec. of $\Delta_0 \mathbf{U}^\dagger \Pi_0 \mathbf{U} \Delta_0$ with eval. λ . Define

$$|\delta_0\rangle := \frac{\Pi_0 \mathbf{U} |\gamma_0\rangle}{\sqrt{\lambda}}, \quad |\delta_1\rangle := \frac{\Pi_0 \mathbf{U} |\gamma_0\rangle}{\sqrt{1-\lambda}}, \quad |\gamma_1\rangle := \frac{\Delta_1 \mathbf{U}^\dagger |\delta_0\rangle}{\sqrt{1-\lambda}}.$$

Then, $\langle \gamma_0 | \gamma_1 \rangle = \langle \delta_0 | \delta_1 \rangle = 0$ and

$$\begin{aligned} \mathbf{U} |\gamma_0\rangle &= \sqrt{\lambda} |\delta_0\rangle + \sqrt{1-\lambda} |\delta_1\rangle & \mathbf{U}^\dagger |\delta_0\rangle &= \sqrt{\lambda} |\gamma_0\rangle + \sqrt{1-\lambda} |\gamma_1\rangle \\ \mathbf{U} |\gamma_1\rangle &= \sqrt{1-\lambda} |\delta_0\rangle - \sqrt{\lambda} |\delta_1\rangle & \mathbf{U}^\dagger |\delta_1\rangle &= \sqrt{1-\lambda} |\gamma_0\rangle - \sqrt{\lambda} |\gamma_1\rangle \end{aligned}$$

(draw the evolution diagram)



In Our Setting: Marriot-Watrous + Post-Mnt. Selection

In our setting, we have $\mathbf{U} = \mathbf{V}\mathbf{T}$, $\lambda = 1/2$, and $|\gamma_0\rangle = |\psi\rangle_{\mathbf{W}} |0\rangle_{\mathbf{X}}$

Lemma 2 $\Rightarrow |\gamma_0\rangle = \frac{1}{\sqrt{2}} |\delta_0\rangle + \frac{1}{\sqrt{2}} |\delta_1\rangle$, and the following:

$$|\delta_0\rangle = \sqrt{2}\mathbf{\Pi}_0\mathbf{V}\mathbf{T}|\gamma_0\rangle, \quad \mathbf{T}^\dagger\mathbf{V}^\dagger|\delta_1\rangle = \frac{1}{\sqrt{2}}|\gamma_0\rangle - \frac{1}{\sqrt{2}}|\gamma_1\rangle, \quad \mathbf{V}\mathbf{T}\left(\frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle\right) = |\delta_0\rangle$$

Starting with $|\gamma_0\rangle \rightarrow \mathbf{V}\mathbf{T}|\gamma_0\rangle \rightarrow \text{measurement } \{\mathbf{\Pi}_0, \mathbf{\Pi}_1\}$:

- ▶ w.p. 1/2, it is $|\delta_0\rangle$ — we are done!
- ▶ w.p. 1/2, it is $|\delta_1\rangle$
 - ▶ Key observation: $\mathbf{T}^\dagger\mathbf{V}^\dagger|\delta_1\rangle = \frac{1}{\sqrt{2}}|\gamma_0\rangle - \frac{1}{\sqrt{2}}|\gamma_1\rangle$
 - ▶ If we can flip the phase of the 2nd term $\Rightarrow \frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle$.
 - ▶ Then, simply do $\mathbf{V}\mathbf{T}\left(\frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle\right) = |\delta_0\rangle$

Yes, we can! (next slide)

Phase Flip for the 2nd Term

We want: $\frac{1}{\sqrt{2}} |\gamma_0\rangle - \frac{1}{\sqrt{2}} |\gamma_1\rangle \rightarrow \frac{1}{\sqrt{2}} |\gamma_0\rangle + \frac{1}{\sqrt{2}} |\gamma_1\rangle$

Recall the following

- ▶ $|\gamma_0\rangle = |\psi\rangle_W |0\rangle_X$ and $\Delta_0 = \mathbb{1}_W \otimes |0\rangle\langle 0|_X$
- ▶ $\Rightarrow \Delta_0 |\gamma_0\rangle = |\gamma_0\rangle$
- ▶ Lemma 2 says $|\gamma_1\rangle = \sqrt{2} \Delta_1 \mathbf{T}^\dagger \mathbf{V}^\dagger |\delta_0\rangle \Rightarrow \Delta_0 |\gamma_1\rangle = 0$

Therefore, it is not hard to come up with the following idea:

$$\begin{aligned} \underbrace{(2\Delta_0 - \mathbb{1}_{WX})}_{=\Delta_0 - \Delta_1} \left(\frac{1}{\sqrt{2}} |\gamma_0\rangle - \frac{1}{\sqrt{2}} |\gamma_1\rangle \right) &= \frac{2}{\sqrt{2}} \Delta_0 |\gamma_0\rangle - \frac{2}{\sqrt{2}} \Delta_0 |\gamma_1\rangle - \frac{1}{\sqrt{2}} |\gamma_0\rangle + \frac{1}{\sqrt{2}} |\gamma_1\rangle \\ &= \frac{2}{\sqrt{2}} |\gamma_0\rangle - 0 - \frac{1}{\sqrt{2}} |\gamma_0\rangle + \frac{1}{\sqrt{2}} |\gamma_1\rangle \\ &= \frac{1}{\sqrt{2}} |\gamma_0\rangle + \frac{1}{\sqrt{2}} |\gamma_1\rangle \end{aligned}$$

Summarizing the Watrous Simulator

- ▶ Start with $|\gamma_0\rangle_{\text{XW}} = |\psi\rangle_{\text{X}} |0\rangle_{\text{W}}$
- ▶ Perform \mathbf{VT} $|\gamma_0\rangle_{\text{XW}}$
- ▶ Perform measurement $\{\mathbf{\Pi}_0, \mathbf{\Pi}_1\}$
 - ▶ If outcome is 0 — guessed correctly (in $|\delta_0\rangle$). Go next step.
 - ▶ Otherwise, we are in $|\delta_1\rangle = \sqrt{2}\mathbf{\Pi}_1\mathbf{VT}|\gamma_0\rangle$.
 - ▶ Perform $\mathbf{T}^\dagger\mathbf{V}^\dagger|\delta_1\rangle = \frac{1}{\sqrt{2}}|\gamma_0\rangle - \frac{1}{\sqrt{2}}|\gamma_1\rangle$
 - ▶ Perform $(2\mathbf{\Delta}_0 - \mathbb{1}_{\text{WX}})(\frac{1}{\sqrt{2}}|\gamma_0\rangle - \frac{1}{\sqrt{2}}|\gamma_1\rangle) = \frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle$
 - ▶ Perform $\mathbf{VT}(\frac{1}{\sqrt{2}}|\gamma_0\rangle + \frac{1}{\sqrt{2}}|\gamma_1\rangle) = |\delta_0\rangle$. Go next step.
- ▶ Sim can finish the last round as the honest prover.

Extending to G3C—Idealized Com Model (1/3)

- ▶ The graph-3-coloring (G3C) problem is NP-complete
- ▶ will use the G3C classical ZK proof from [GMW86]
- ▶ $\Pr[\text{Guess correctly}] = \frac{1}{m}$, where $m = \# \text{ edges}$.
- ▶ $\Pr[\text{Guess correctly}] \perp |\psi\rangle?$
 - ▶ 1st msg using perfect-binding (PB) Com
 - ▶ What about binding? — Collapse-binding suffices [Unr16]
 - ▶ We assume an ideal Com for simplicity: perfect-hiding and perfectly-binding
 - ▶ Extends to PB and CH Com later

Extending to G3C—Idealized Com Model (2/3)

Key ingredients for the GI simulator:

- ▶ Define an operator: $\Delta_0^\dagger \mathbf{T}^\dagger \mathbf{V}^\dagger \Pi_0 \mathbf{T} \mathbf{V} \Delta_0$ ($=: \mathbf{M}$)
- ▶ An technical Lemma 1: $\lambda = \frac{1}{2}$ ($\perp |\psi\rangle$)
- ▶ Invoke Marriot-Watrous Lemma 2 with $\lambda = \frac{1}{2}$:
 - ▶ Voilà 😊! We can get $|\delta_0\rangle$ within ≤ 2 steps

What will change for the G3C protocol?

- ▶ \mathbf{M} defined as before (\mathbf{T} modified in the natural way)
- ▶ Lemma 1: $\lambda = \frac{1}{m}$ ($\perp |\psi\rangle$)
- ▶ Invoke Marriot-Watrous Lemma 2 with $\lambda = \frac{1}{m}$:
 - ▶ 😞! no guarantee for $|\delta_0\rangle$ within ≤ 2 steps
- ▶ Solution: use the full power of Matriott-Watrous analysis (next slide).

Extending to G3C—Idealized Com Model (3/3)

(draw the evolution diagram in the current setting)

$|\gamma_0\rangle$ $|\delta_0\rangle$ $|\gamma_0\rangle$ $|\delta_0\rangle$ $|\gamma_0\rangle \dots$

$|\delta_1\rangle$ $|\gamma_1\rangle$ $|\delta_1\rangle$ $|\gamma_0\rangle \dots$

The main take-away:

- ▶ $\mathbf{U} |\gamma_0\rangle = \sqrt{\lambda} |\delta_0\rangle + \sqrt{1-\lambda} |\delta_1\rangle$ $\mathbf{U}^\dagger |\delta_0\rangle = \sqrt{\lambda} |\gamma_0\rangle + \sqrt{1-\lambda} |\gamma_1\rangle$, where $\lambda = 1/m$.
 $\mathbf{U} |\gamma_1\rangle = \sqrt{1-\lambda} |\delta_0\rangle - \sqrt{\lambda} |\delta_1\rangle$ $\mathbf{U}^\dagger |\delta_1\rangle = \sqrt{1-\lambda} |\gamma_0\rangle - \sqrt{\lambda} |\gamma_1\rangle$
- ▶ Measure $\{\Pi_0, \Pi_1\}$ at each $|\delta\rangle$, if results in $|\delta_1\rangle$:
 - ▶ $\mathbf{U}(2\Delta - \mathbb{1})\mathbf{U}^\dagger |\delta_1\rangle = 2\sqrt{p(1-p)} |\delta_0\rangle + (1-2p) |\delta_1\rangle$
- ▶ Measure $\{\Pi_0, \Pi_1\}$. Go to $|\delta_1\rangle$ w.p. $(1-2p)$.
- ▶ Keep failing after t iteration: $1 - (1-p)(1-2p)^t$. Then Chernoff bound.

The General Quantum Rewinding Lemma (Exact)

LEMMA 3: EXACT QUANTUM REWINDING [Wat09]

\mathbf{Q} is a QC works on $|\psi\rangle$ and with $\Pr[\text{success}] = p \ (\perp |\psi\rangle)$ outputs $|\delta_0\rangle$. Then, for any $\varepsilon > 0$, there exists another QC \mathbf{R} of size

$$O\left(\frac{\log(1/\varepsilon)}{p(1-p)} \cdot \text{size}(\mathbf{Q})\right)$$

such that for every input $|\psi\rangle$, the output ρ of \mathbf{R} satisfies $\text{Tr}(\rho |\delta_0\rangle\langle\delta_0|) \geq 1 - \varepsilon$.

- ▶ “Exact” refers to the fact that $p \perp |\psi\rangle$.
- ▶ The $\frac{\log(1/\varepsilon)}{p(1-p)}$ is due to Chernoff bound.
- ▶ Only need poly-size for a negligible ε .
- ▶ For the trace, the closer to 1, the better.

The Version Allowing Small Perturbations

LEMMA 4: QUANTUM REWINDING WITH SMALL PERTURBATIONS [Wat09, Sec. 4.2]

Let \mathbf{Q} , $|\psi\rangle$, and $|\delta_0\rangle$ as before. But $\Pr[\text{success}] = p(|\psi\rangle)$ now depends on $|\psi\rangle$. Let $p_0, q \in (0, 1)$ and $\varepsilon \in (0, 1/2)$ be real numbers such that

$$(1). |p(\psi) - q| < \varepsilon \quad (2). p_0 \leq p(\psi) \quad (3). p_0(1 - p_0) \leq q(1 - q)$$

Then, for any $\varepsilon > 0$, there exists another QC \mathbf{R} of size $O\left(\frac{\log(1/\varepsilon)}{p_0(1-p_0)} \cdot \text{size}(\mathbf{Q})\right)$ such that for every input $|\psi\rangle$, the output ρ of \mathbf{R} satisfies:

$$\text{Tr}(\rho |\delta_0\rangle\langle\delta_0|) \geq 1 - 16\varepsilon \frac{\log^2(1/\varepsilon)}{p_0^2(1 - p_0)^2}.$$

Proof at a high-level:

- ▶ Consider each eigen-space separately (next slide).
- ▶ For detailed calculation, see [Wat09, Sec. 4.2].

Proof Sketch for Lemma 4

Proof Sketch:

- ▶ In Lemma 1, $|\gamma_0\rangle = |\psi\rangle_{\mathcal{W}} |0\rangle_{\mathcal{X}}$ is no longer an evec. of \mathbf{M}
 - ▶ The reason: $|\psi\rangle_{\mathcal{W}}$ is not an evec. of \mathbf{Q}
- ▶ (mental exper.) Thus, decomp. $|\psi\rangle$ in the evecs $\{|\psi_i\rangle\}_{i \in [\text{dim}]}$ of \mathbf{Q}
- ▶ (mental exper.) For each i , we obtain Lemmas 1 and 2
- ▶ (mental exper.) In the Marriot-Watrous procedure, in each egin space:
$$\mathbf{V}\mathbf{T} |\psi_i\rangle_{\mathcal{W}} |0\rangle_{\mathcal{X}} = \sqrt{p(|\psi_i\rangle)} |\delta_0(|\psi_i\rangle)\rangle + \sqrt{1 - p(|\psi_i\rangle)} |\delta_1(|\psi_i\rangle)\rangle$$
- ▶ (mental exper.) Define a unitary \mathbf{N} such that for all $i \in [\text{dim}]$:
$$\sqrt{p(|\psi_i\rangle)} |\delta_0(|\psi_i\rangle)\rangle + \sqrt{1 - p(|\psi_i\rangle)} |\delta_1(|\psi_i\rangle)\rangle \rightarrow \sqrt{q} |\delta_0(|\psi_i\rangle)\rangle + \sqrt{1 - q} |\delta_1(|\psi_i\rangle)\rangle$$
- ▶ (mental exper.) Ready to apply the Exact Rewinding Lemma 3 (w/ p_0 as we don't know p .) Need $p_0(1 - p_0) \leq q(1 - q)$

In summary, this is a Sim w/ an imaginary operator \mathbf{N} , giving the same trace bound as in Lemma 3. But for the real Sim, there is no \mathbf{N} .

- ▶ Doesn't matter. \mathbf{N} only affects the trace bound negligibly.
- ▶ By tedious-yet-elementary linear algebra (see [Wat09, Sec. 4.2]).

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