Scaling rules for the ionization of biological molecules by highly charged ions

A. M. P. Mendez, C. C. Montanari, J. E. Miraglia Instituto de Astronomía y Física del Espacio (CONICET-UBA), Buenos Aires, Argentina. (Dated:)

In the present work, we We investigate scaling rules for the ionization cross sections of multicharged ions on molecules of biological interest. The cross sections are obtained from distorted wave calculations for atomic targets combined with a stoichiometric model for the molecules proposed using a methodology presented in [Mendez et al. J. Phys B (2020)], which considers distorted-wave calculations for atomic targets combined with a molecular stoichiometric model. We examine ions with nuclear charges Z from +1 to +8 in-impacting on five nucleobases -adenine, cytosine, guanine, thymine, uracil-, tetrahydrofuran, pyrimidine, and also in-water. We propose a scaling with the ion charge , which is valid in the intermediate to high energy range, i.e., 0.2-5 MeV/amu for oxygen impact. We extend our work to a universal scaling for any ion and molecule, merging the forty ion molecule Furthermore, we introduce a modified scaling rule, which takes into account the number of active electrons per molecule. This scaling rule is independent of the ion charge and the nature of the molecular target, and it allows us to merge the forty ion-molecule systems analyzed here into a single band. Furthermore, our model proved to be valid for other molecules toe We confirm the generality of our independent scaling law with several collisional systems.

31

PACS numbers: 34.50Gb, 34.80Gs, 34.80Dp

Keywords: ionization, scaling, molecules, charged-ions, DNA, multicharged ions

I. INTRODUCTION

1

10

11

12

13

14

15

16

17

18

19

20

21

23

24

27

28

The interest in the ionization of biological ³⁴ molecules by multicharged ions has gained increasing ³⁵ interest increased due to medical and environmental ³⁶ reasonsimplementations [1], from including medical ³⁷ treatments [2-4] to [2-4] and contaminant recognition in ³⁸ biological materials [5, 6]. Many semiempirical [7] and ³⁹ theoretical efforts are currently being undertaken [8-13] ⁴⁰ [8-13] to get reliable values for the ionization cross ⁴¹ sections of these molecules molecular systems.

Recently In recent work [8], we combined the continue 44 distorted wave continuum distorted-wave calculations 45 (CDW) for atoms and the simple stoichiometric model 46 (SSM), also known as the Bragg sum rule, to approximate the ionization cross sections of complex molecular targets by charged ions [8]. The CDW SSM approximation showed reasonable the impact of charged ions. The molecular ionization cross section σ_M was expressed as a linear combination of atomic CDW 48 calculations σ_A , weighted with the number of atoms for each specie n_A , i.e., $\sigma_M = \sum_A n_A \sigma_A$. The CDW-SSM 49 approximation showed consistent results for over a hun-50 dred of ion molecule biologically relevant ion-molecule 51 systems. As expected, in the high energy range (i.e., 52 above 5 MeV/amu), the ionization cross sections present 53 of the molecular systems follow the Z^2 dependence pre- 54 dicted by the first Born approximation. However, at 55 intermediate energies, the dependence with Z is more 56 complex, and non-perturvative not straightforward since 57 non-perturbative models are mandatory.

The intention of this letter is to give a follow up This contribution constitutes a follow-up of our previous work [8] by proposing a scaling. We introduce here a two-folded scaling rule for the ionization cross sections of complex molecules by charged ions. Our approach considers the dependence of the cross section with the ion charge Z and incorporates the scaling of the ionization cross sections of complex molecules, valid at intermediate energies. In general, scaling rules are used as first order tests with the number of active electron $n_{\rm g}$ of the molecular targets. Scaling rules are in generally very useful since they can be used as first-order approximations in experimental measurements and multipurpose codes. Based on [8], we propose a universal scaling for any ion target system.

At intermediate impact energies,

II. SCALING RULES

A. Scale with the ion charge

In the development of our scaling rule, we examine forty collisional systems. The target-ion systems are composed of eight targets: the DNA and RNA nucleobases –adenine, cytosine, guanine, thymine, uracil–, tetrahydrofuran (THF), pyrimidine, and water; and five ion species: H⁺, He⁺², Be⁺⁴, C⁺⁶, and O⁺⁸. We consider these systems as a benchmark for the present rule.

We found two types of Z-scaling laws in the literature applicable to the intermediate impact energy range. The

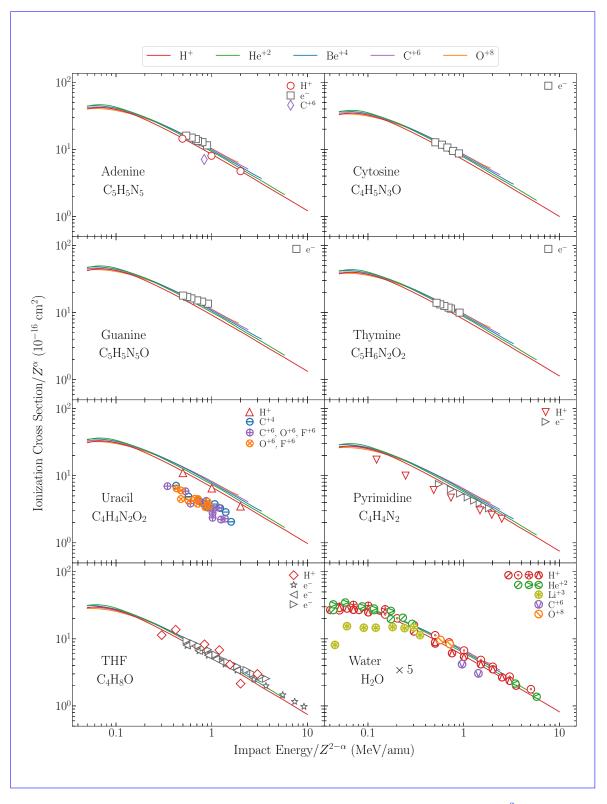


FIG. 1. (Color online) Scaled ionization cross section σ/Z^{α} as a function of ion impact energy $E/Z^{2-\alpha}$ with $\alpha=1.2$. Colors are associated with the incident ion labeled on top of the figure. Curves: present CDW-SSM theoretical results. Symbols: experimental impact of $\overset{\bullet}{\circ}$ H⁺ [18] and $\overset{\bullet}{\circ}$ C⁺⁶ [21] on adenine; H⁺ on $\overset{\bullet}{\circ}$ uracil [17], $\overset{\bullet}{\circ}$ C⁺⁴, $\overset{\bullet}{\circ}$ C⁺⁶, O⁺⁶, F⁺⁶, and $\overset{\bullet}{\circ}$ O⁺⁸, F⁺⁸ on uracil [22, 23]; H⁺ on $\overset{\bullet}{\circ}$ pyrimidine [19], and $\overset{\bullet}{\circ}$ THF [20]; $\overset{\bullet}{\circ}$ [24], $\overset{\bullet}{\circ}$ [25], $\overset{\bullet}{\circ}$ [26], $\overset{\bullet}{\circ}$ [27] H⁺, $\overset{\bullet}{\circ}$ [28], $\overset{\bullet}{\circ}$ [29], $\overset{\bullet}{\circ}$ [27] He⁺², $\overset{\bullet}{\circ}$ $\overset{\bullet}{\circ}$ C⁺⁶ [31, 32], and $\overset{\bullet}{\circ}$ O⁺⁸ [33] on water. Markers $\overset{\bullet}{\circ}$ [34], $\overset{\bullet}{\circ}$ [35], $\overset{\bullet}{\circ}$ [36], and $\overset{\bullet}{\circ}$ [37] correspond to electron impact ionization with the equi-velocity conversion.

sus E/Z as to be the natural reduced form of the ionization cross section σ and the incident ion energy E. Much more More recently, Montenegro and co-workers [15, 16] proposed co-workers [15, 16] suggested an alternative scaling by taking into account that σ the cross section is a function of Z^2/E at high energies. Their scaling, given by

$$\sigma/Z^{\alpha} = f(E/Z^{2-\alpha}),\tag{1}$$

keeps the Z^2/E relationship for any value of the parameter α . In Ref. [15], the authors propose The authors proposed $\alpha = 4/3$ for the ionization of He and H₂ by different charged ions.

(Color online) Scaled ionization cross section σ/Z^{α} as 120 a function of ion impact energy $E/Z^{2-\alpha}$ with $\alpha=1.2.^{121}$ Colors are associated with the incident ion labeled on 122 top of the figure. Curves: present CDW SSM theoretical 123 results. Symbols: experimental impact of O H+ [18] 124 and OC+6 [21] on adenine; H+ on OC+4 125 OC+6, OC+6, OC+6, OC+6, OC+6, OC+7 and OC+8, OC+8 on uracil [22, 23]; H+ on 126 OC+7 pyrimidine [19], and OC+8 [20], OC-12] He+2, OC-15 [31, 32], and OC+8 [33] on water. Markers OC-16 [31, 32], OC-17 and OC-18 [37] correspond to electron impact ionization with the equi-velocity conversion.

Combining our recent CDW SSM results [8] and ¹²⁹ Eq. (??), we propose here a Z-sealing and implement ¹³⁰ it for forty collisional systems. The ion molecule ¹³¹ systems are composed of eight targets: the DNA and ¹³² RNA nucleobases—adenine, cytosine, guanine, thymine, ¹³³ uracil, tetrahydrofuran (THF), pyrimidine, and water; ¹³⁴ and five charged ions: H⁺, He⁺², Be⁺⁴, C⁺⁶, and O⁺⁸. We ¹³⁵ considered these systems as a benchmark for the present ¹³⁶ sealing. differently charged ions [15].

We—Following the work of Montenegro and collaborators, we found that the parameter α from 139 Eq. (??) that fits the CDW-SSM scaled cross sections₁₄₀ for all the ions that best converges the CDW-SSM cross₁₄₁ sections of the forty collisional systems over the broadest energy range is $\alpha = 1.2$. The validity of the theoretical sealing with the ion charge is very clear this particular 142 scaling is evident in Fig. 1, where -for each targetthe CDW-SSM curves lays corresponding to different 143 ions lay one over the other. Our It is worth noting,144 that our theoretical results are valid for impact energies, 145 around and above the maximum of the cross sections, 146 which means above corresponds to an impact energy, 147 range from 50 keV for impact of H⁺ to 250 keV/amu for₁₄₈ impact of O⁺⁸.

The scaling was tested with-We also examined the ex- $_{150}$ perimental data available by impact of different charged $_{151}$ ions [17–25, 29–31, 33], and also by electron impact at $_{152}$ sufficiently high velocity [34–37] for the forty ion-target $_{153}$ systems [17–33] with the Z^{α} -scaling rule. As can be $_{154}$ noted, most of the data in Fig. 1 confirms confirm the $_{155}$ present scaling, even for O+8 in water [33]. However, $_{156}$ the Only two data sets are off our predictions: the $_{157}$ ionization data of uracil by swift C, O, and F ions in-from $_{156}$

Molecule	n_e	Molecule	n_e	Molecule	n_e
H_2O	6	CO ₂ —	12	C ₄ H ₅ N ₃ O —	37
N_2	8	C_4H_8O —	28	$C_5H_6N_2O_2$ —	42
O_2	8	$C_4H_4N_2$ —	28	$C_5H_5N_5$ —	45
CH_4	8	$C_4H_4N_2O_2$ —	36	$C_5H_5N_5O$ —	49

TABLE I. Number of active electrons per target at intermediate to high energies obtained from the CDW calculations [8].

Refs. [22, 23] are too low compared with our CDW SSM results, but also as compared with Itoh et al. data [17], and with the CTMC calculations by Sarkadi [38]. The CDW-SSM results, and the data for Li⁺³ in water from Ref. [30] also spreads spread out from the present theoretical curves for E < 600 keV/amu. For targets with none or little experimental data, we included electron impact ionization results [34–37] at high velocity with the corresponding equivelocity conversion.

B. Scale with the molecular target

The good results obtained in the scaling with the ion charge challenged us to look for a more general scaling rule encouraged us to further investigate a scaling law that could predict values for ionization cross sections of any ion in any molecule. To that this end, we resorted to considered the number of active electrons in each atom molecule n_e proposed in [8], Ref. [8] and combined it with the ZZ^{α} -scaling displayed in Fig. 1from Section II A.

The CDW ionization cross sections σ^{CDW} of atomic In our previous work, we noticed that the CDW atomic ionization cross section σ_A of H, C, N, O targets scale as $\sigma_e = \sigma^{\text{CDW}}/n_e$ with n_e being with the number of active electron per atom ν_A , where ν_A is 1 for H, and 4 for C, N, and O, i. e.

$$\frac{\sigma_{\rm H}^{\rm CDW}}{1} \sim \frac{\sigma_{\rm C}^{\rm CDW}}{4} \sim \frac{\sigma_{\rm N}^{\rm CDW}}{4} \sim \frac{\sigma_{\rm O}^{\rm CDW}}{4}$$

The SSMleads to the molecular numbers. By considering the SSM, we define the number of active electrons included in TableI per molecule $n_e = \sum_A n_A \nu_A$, which are given in Table I for the molecular targets considered throughtout this work. The scaling with the molecular number of active electrons (see Fig. 6 of Ref. [8]) gave excellent results and proved to be a good scaling.

(Color online) Universal scaling with the ions charge Z and the number of active electrons in the molecule n_e given by Eq. (??). Curves: present CDW-SSM theoretical results. Symbols: experimental impact of H⁺ on \bigcirc adenine [18], \triangle uracil [17], ∇ pyrimidine [19] and \triangle THF [20]; \triangle C⁺⁶ on adenine [21]; \bigcirc [24], \bigcirc [25], \bigcirc [26], \bigcirc [27] H⁺, \bigcirc [28], \bigcirc [29], \bigcirc [27] He⁺², \bigcirc C⁺⁶ [31, 32], and \bigcirc O⁺⁸ [33] on water. H⁺ impact on \bigcirc N₂, \bigcirc O₂, \bigcirc CO₂, and \bigcirc CH₄; and He⁺² impact on \bigcirc N₂.

■ O_2 , \square CO, \square CO_2 , and \square CH_4 [29, 39], \bigcirc H^+ on 195 CH_4 [40]; and electron impact on \triangleright pyrimidine [35], and 196 \triangleleft , \Rightarrow [36, 37] THF.

160

161

162

163

164

165

166

167

168

169

170

171

172

173

175

176

177

178

179

181

182

184

185

186

187

188

189

190

192

193

235

236

237

238

239

240

241

242

243

244

245

246

247

248

C. Scale with the ion charge and the molecular target

The universal scalingwe propose here By incorporating₂₀₃ the Z^{α} reduction and the molecular number of active₂₀₄ electron scaling, we introduce the scaled and reduced₂₀₅ ionization cross section of molecules $\tilde{\sigma}$, which is expressed₂₀₆ as σ_{U} as a function of $E/Z^{2-\alpha}$, with and it is given by 207

$$\underline{\sigma_U}\tilde{\sigma} = \frac{\sigma_e}{Z^{\alpha}} = \frac{\sigma/n_e}{Z^{\alpha}} \frac{\sigma_M/n_e}{Z^{\alpha}}, \qquad (2)_{210}^{209}$$

199

201

202

 σ where σ_M is the ionization cross section for the molecular target, $\alpha = 1.2$ and n_e is the number of active electrons per molecule given in Table I. In I and the parameter²¹³ is $\alpha = 1.2$. Fig. 2, we test the universal scaling of₂₁₄ Eq. (??) for all the shows the reduced-scaled theoret-215 ical and experimental values displayed in Fig.1. We also 216 included a gray area representing the 30% deviation of 217 our theoretical curvesdata from Fig. 1 obtained using218 Eq. (2). As can be noted, the universal scaling works very 219 well. All the curves and data lays and is independent of 220 the ion charge or the complexity of the molecular target. 221 Our theoretical curves lay in a narrow band valid for₂₂₂ any charged ion (scaled with Zreduced with Z^{α}) in any₂₂₃ molecule (scaled with the number of active electrons) -224 We decided not to include in this figure with a dispersion₂₂₅ of about xx%. If we consider the experimental data, the 226 uncertainty of our scaling grows to 30%, which is shown₂₂₇ with the gray area. It is worth noting that we did not₂₂₈ include the data for uracil from Refs. [22, 23], and for Li⁺³₂₂₉ on water [30]. The discussion about these experimental₂₃₀ values exceeds the present work.

In principle, the *universal* scaling should

We consider this scaling robust enough to be valid₂₃₃ for different ion-molecule combinations. We proved₂₃₄

this statement by including tested the generality of our scaling by including several data sets of molecular targets that were not considered previously. We included in Fig. 2 the measurements by Rudd *et al.* [29, 39] for H⁺ and He⁺² in N₂, O₂, CH₄, CO and CO₂, and the recent values by Luna *et al.* [40] for H⁺ in CH₄.

The good agreement shown in Fig. 2 summaries the main result of this work, and holds the validity of our universal scaling. Although the theoretical CDW-SSM results are valid for energies above the maximum of the cross sections, it is worth noting from Fig. 2 that the scaling of the experimental data extends even to lower impact energies, as can be noted in Fig. 2. The importance of scaling rules lays in their predictive capability. New measurements for other ions and molecules are expected to reinforce the present proposal.

III. CONCLUSIONS

In this letter, we present a scaling

We present scaling rules for the ionization cross sections of highly charged ions in biological targets. The scaling rule states the cross sections divided by first scaling reduces the nature of the projectile by scaling the cross section with the ion charge, Z^{α} , as a function of the reduced impact energy $E/Z^{2-\alpha}$, with $\alpha = 1.2$. The scaling was second scaling considers the molecular characteristic of the target by taking into account the number of active electrons per molecule, n_e . The last scaling law combines the Z-reduction with the n_e -scaling of cross section, and it becomes independent of the ion charge and the molecular target. The scalings were obtained by means of the CDW-SSM calculations for five different charged ions in eight targets and tested with the available experimental data. A universal scaling rule is also proposed, which reduced the cross sections with the number of active electrons of the molecule. The universal scaling The generality of our independent scaling is proved to be valid for a large in a wide energy range by considering a significant number of experimental data -sets for other collisional systems.

 ^[1] T. Liamsuwan and H. Nikjoo, Phys. Med. Biol. 58 641

–249

672 (2013).

^[2] O. Mohamad, B. J. Sishc, J. Saha, A. Pompos, A.251 Rahimi, M. D. Story, A. J. Davis, D. N. Kim, Cancers 9,252 66 (2017).

 ^[3] A. V. Solov'yov, E. Surdutovich, E. Scifoni, I. Mishustin, 254
 and W. Greiner, Phys. Rev. E 79, 011909 (2009); 255

^[4] Denifl S., Mrk T.D., Scheier P. Eds: García Gómez-256 Tejedor G., Fuss M. Springer, Dordrecht (2012)

^[5] N. A. Gafur, M. Sakakibara, S. Sano, K. A. Sera, Water₂₅₈10, 1507 (2018); doi:10.3390/w10111507.

^[6] D. Benedetti, E. Nunes, M. Sarmento, C. Porto, C. E.260 Iochims dos Santos, J. Ferraz Dias, J. da Silva, Mutation261 Research/Genetic Toxicology and Environmental Muta-262

genesis, Volume 752, 28-33 (2013);

^[7] P. de Vera, R. Garcia-Molina, I. Abril, and A. V. Solovyov, Phys. Rev. Lett. 110, 148104 (2013).

^[8] A. M. P. Mendez, C. C. Montanari, and J. E. Miraglia, J. Phys. B: At. Mol. Opt. Phys. 53, 055201 (2020).

^[9] M. A. Quinto, J. M. Monti, C. A. Tachino, P. F. Weck, O. A. Fojón, C. Champion, R. D. Rivarola, Rad. Phys. Chem. 167, 108337 (2020);

^[10] H. J. Lüdde, M. Horbatsch and T. Kirchner, J. Phys. B: At. Mol. Opt. Phys. 52, 195203 (2019).

^[11] H. J. Lüdde, M. Horbatsch and T. Kirchner, Eur. Phys. J. B 91, 99 (2018).

^[12] H. J. Lüdde, A. Achenbach, T. Kalkbrenner, H.-C. Jankowiak and T. Kirchner, Eur. Phys. J. D 70, 82

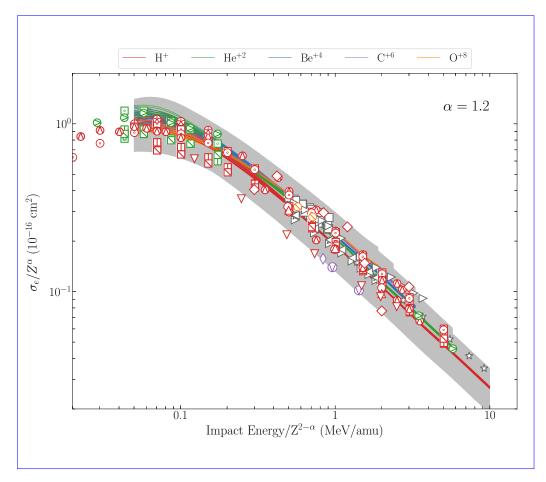


FIG. 2. (Color online) Universal scaling with the ions charge Z and the number of active electrons in the molecule n_{ϵ} given by Eq. (2) with $\alpha = 1.2$. Curves: present CDW-SSM theoretical results. Symbols: experimental impact of H^+ on \bigcirc adenine [18], \triangle uracil [17], ∇ pyrimidine [19] and \bigcirc THF [20]; \bigcirc C⁺⁶ on adenine [21]; \bigcirc [24], \bigcirc [25], \bigcirc [26], \bigcirc [27] H^+ , \bigcirc [28], \bigcirc [29], \bigcirc [27] $H^{e^{+2}}$, \bigcirc C⁺⁶ [31, 32], and \bigcirc O⁺⁸ [33] on water. H^+ impact on \bigcirc N₂, \bigcirc O₂, \bigcirc CO₂, and \bigcirc CH₄; and \bigcirc CH₄ [29, 39], \bigcirc H⁺ on CH₄ [40]; and electron impact on \triangleright pyrimidine [35], and \bigcirc , \bigcirc [36, 37] THF.

(2016).

263

264

265

266

267

268

269

270

271

272

273

274

275

276

277

278

279

280

281

282

283

- [13] C. Champion, ME. E. Galassi, O. Fojón, H. Lekadir, 285 J. Hanssen, RD. D. Rivarola, PF. F. Weck, AN. N. 286 Agnihotri, S. Nandi, and LC. C. Tribedi. J. Phys.: Conf. 287 Ser. 373, 012004 (2012).
- [14] R. K. Janev and L. P. Presnyakov J. Phys. B: At. Mol. 289 Opt. Phys. 13, 4233 (1980).
- [15] R. D. DuBois, E. C. Montenegro and G. M. Sigaud, AIP₂₉₁ Conference Proceeding 1525, 679 (2013).
- [16] E. C. Montenegro, G. M. Sigaud, and R. D. DuBois, 293 Phys. Rev. A 87 012706 (2013).
- [17] A. Itoh, Y. Iriki, M. Imai, C. Champion, and R. D. Ri-295 varola, Phys. Rev. A 88, 052711 (2013).
- [18] Y. Iriki, Y. Kikuchi, M. Imai, and A. Itoh Phys. Rev. A₂₉₇ 84, 052719 (2011).
- [19] W. Wolff, H. Luna, L. Sigaud, A. C. Tavares, and E. C. 299 Montenegro J. Chem. Phys. 140, 064309 (2014).
- [20] M. Wang, B. Rudek, D. Bennett, P. de Vera, M. Bug, T.301 Buhr, W. Y. Baek, G. Hilgers, H. Rabus, Phys. Rev. A302 93, 052711 (2016).
- [21] S. Bhattacharjee, C. Bagdia, M. R. Chowdhury, A. Man-304

- dal, J. M. Monti, R. D. Rivarola, and L. C. Tribedi, Phys. Rev. A **100**, 012703(2019).
- [22] A. N. Agnihotri, S. Kasthurirangan, S. Nandi, A. Kumar, M. E. Galassi, R. D. Rivarola, O. Fojón, C. Champion, J. Hanssen, H. Lekadir, P. F. Weck, and L. C. Tribedi. Phys. Rev. A 85, 032711 (2012).
- [23] A. N. Agnihotri, S. Kasthurirangan, S. Nandi, A. Kumar, C. Champion, H. Lekadir, J. Hanssen, P. F. Weck, M. E. Galassi, R. D. Rivarola, O. Fojón and L. C. Tribedi, J. Phys. B: At. Mol. Opt. Phys. 46, 185201 (2013).
- [24] H. Luna, A. L. F. de Barros, J. A. Wyer, S. W. J. Scully, J. Lecointre, P. M. Y. Garcia, G. M. Sigaud, A. C. F. Santos, V. Senthil, M. B. Shah, C. J. Latimer, and E. C. Montenegro, Phys. Rev. A 75, 042711 (2007).
- [25] M. A. Bolorizadeh and M. E. Rudd, Phys. Rev. A 33, 888 (1986).
- [26] M. E. Rudd, T. V. Goffe, R. D. DuBois, L. H. Toburen, Phys. Rev. A 31, 492 (1985).
- [27] L. H. Toburen, W. E. Wilson and R. J. Popowich, Radiat. Res. 82, 27–44 (1980).
- [28] D. Ohsawa, Y. Sato, Y. Okada, V. P. Shevelko, and F.

Soga Phys. Rev. A **72**, 062710 (2005).

305

306

307

311

312

313

[29] M. E. Rudd, T. V. Goffe, and A. Itoh, Phys. Rev. A 32,322 2128 (1985).

321

- [30] H. Luna, W. Wolff, E. C. Montenegro, Andre C. Tavares, 324
 H. J. Ludde, G. Schenk, M. Horbatsch, and T. Kirchner, 325
 Phys. Rev. A 93, 052705 (2016).
 - [31] C. Dal Cappello, C. Champion, O. Boudrioua, H.327 Lekadir, Y. Sato, D. Ohsawa, Nuclear Instruments and Methods in Physics Research B 267 (2009) 781–790.
- 314 [32] Shamik Bhattacharjee, S. Biswas, J. M. Monti, R. D.330
 315 Rivarola, and L. C. Tribedi, Phys. Rev A 96, 052707331
 316 (2017).
- [33] S. Bhattacharjee, S. Biswas, C. Bagdia, M. Roychowd-333
 hury, S. Nandi, D. Misra, J. M. Monti, C. A. Tachino, R. 334
 D. Rivarola, C. Champion and L. C. Tribedi, J. Phys. B:335
 At. Mol. Opt. Phys. 49, 065202 (2016).

- [34] M. A. Rahman and E. Krishnakumar, Electron ionization of DNA bases, J. Chem. Phys. 144, 161102 (2016).
- [35] M. U. Bug, W. Y. Baek, H. Rabus, C. Villagrasa, S. Meylan, A. B. Rosenfeld, Rad. Phys. Chem. 130, 459– 479 (2017).
- [36] W. Wolff, B. Rudek, L. A. da Silva, G. Hilgers, E. C. Montenegro, M. G. P. Homem, J. Chem. Phys. 151, 064304 (2019).
- [37] M. Fuss, A. Muoz, J. C. Oller, F. Blanco, D. Almeida, P. Limo-Vieira, T. P. D. Do, M. J. Brunger, G. García, Phys. Rev. A 80, 052709 (2009).
- [38] L. Sarkadi, J. Phys. B: At. Mol. Opt. Phys. 49 185203 (2016)
- [39] M. E. Rudd, R. D. DuBois, L. H. Toburen, and C. A. Ratcliffe, T. V. Goffe, Phys. Rev. A 28, 3244 (1983).
- [40] H. Luna, W. Wolff, and E. C. Montenegro, L. Sigaud, Phys. Rev. A 99, 012709 (2019).