

1 Prototype equation

1.1 Geometry

Let Ω be a bounded 2-dimensional region. Let Γ be a 1-dimensional curve that divides Ω into two subregions, denoted as Ω_1 and Ω_2 . Let $\partial\Omega_D$ and $\partial\Omega_N$ denote its Dirichlet and Neumann boundaries, respectively.

1.2 Diffusion-reaction equation for excitons

Consider the following equation defined on Ω :

$$-\nabla(a(x)\nabla u) + c(x)u = f(x) \quad (1)$$

and boundary conditions are:

$$u = u_D \quad \forall x \in \Omega_D \quad (2)$$

$$\frac{\partial u}{\partial n} = g_N \quad \forall x \in \Omega_N \quad (3)$$

In particular, u is not assumed to have continuous "flux" across $\Gamma = \overline{\Omega}_1 \cap \overline{\Omega}_2$, i.e.

$$u_1 = u_2 = u_\Gamma \quad \forall x \in \Gamma \quad (4)$$

$$\frac{\partial u_1}{\partial n_{12}} - \frac{\partial u_2}{\partial n_{12}} = -ku_\Gamma \quad \forall x \in \Gamma \quad (5)$$

The physical meaning here is that, if u is the density of excitons, then near the interface Γ , u has a drainage rate of " ku ".

To use Finite Element Method, one needs to write down the weak form. Let v be a test function on Ω , then

$$\int_{\Omega} a(x)\nabla u \cdot \nabla v \, dx + \int_{\Omega} c(x)uv \, dx + \int_{\Gamma} kuv \, ds = \int_{\Omega} fvd \, x \quad (6)$$

In FEM, the region Ω is decomposed into the union of disjoint triangles. Both u and v are approximated by piecewise linear polynomials.

2 Linear system assembly

In this section, we write down the detailed formula of assembling coefficient matrix. In particular, the 3 integrals above are replaced by the matrices A , C , and D .

2.1 Quadrature rule: 1D and 2D

Consider an integral over a triangular region with vertices $\{v_0, v_1, v_2\}$

$$I = \int_T f dx \quad (7)$$

Then first-order quadrature rule is

$$I \approx |T| * (f_0 + f_1 + f_2)/3 \quad (8)$$

where f_i denotes the function value $f(x_i, y_i)$.

We then consider an integral over a segment, for example an edge e of a triangle with end nodes being $\{v_0, v_1\}$. And we have 1st-order quadrature rule

$$I = \int_e f ds \quad (9)$$

$$\approx |e| \frac{f_0 + f_1}{2} \quad (10)$$

These quadrature rules are exact for linear polynomials and are used for assembling all the matrices below.

2.2 Matrices and vectors

2.2.1 Matrix A

Let ϕ_i be the linear basis Lagrange polynomial with 1 on the i -th node and 0 on all other nodes. Then

$$A_{ij} = \int_{\Omega} a(x) \nabla \phi_i \nabla \phi_j dx \quad (11)$$

In practice, this is done by looping through all triangular elements T 's. For each T , one identify its vertices i, j, k , and calculate the contribution of T to all 9 entries in A: $A_{\alpha\beta}$ where α and β can be either i, j , or k .

For example, assume triangle T has vertices 0,1,and 2. We denote the coordinates of these vertices by \vec{r}_0, \vec{r}_1 and \vec{r}_2 . We compute 2 quantities: $A_{00}(T)$ and $A_{01}(T)$. Other cases can be obtained by switching subscripts between 0, 1, and 2.

1. $A_{00}(T)$

First, note by a linear transformation, one can transform v_0, v_1 and v_2 easily to a reference triangle with vertices (0,0), (1,0), and (0,1). In this way, we compute

$$\nabla \phi_0 \cdot \nabla \phi_0 = \frac{|\vec{r}_2 - \vec{r}_1|^2}{4|T|^2} \quad (12)$$

And hence,

$$\int_T a(x) \nabla \phi_0 \cdot \nabla \phi_1 dx = |T| \cdot \frac{(a_0 + a_1 + a_2)}{3} \cdot \frac{|\vec{r}_2 - \vec{r}_1|^2}{4|T|^2} \quad (13)$$

$$= \frac{a_0 + a_1 + a_2}{12|T|} |\vec{r}_2 - \vec{r}_1|^2 \quad (14)$$

2. $A_{01}(T)$

Similar to $A_{01}(T)$, we again make use of the reference triangle and obtained

$$\nabla \phi_0 \cdot \nabla \phi_1 = - \frac{(\vec{r}_2 - \vec{r}_0) \cdot (\vec{r}_2 - \vec{r}_1)}{4|T|^2} \quad (15)$$

and thus

$$\int_T a(x) \nabla \phi_0 \cdot \nabla \phi_1 dx = - \frac{a_0 + a_1 + a_2}{12|T|} (\vec{r}_2 - \vec{r}_0) \cdot (\vec{r}_2 - \vec{r}_1) \quad (16)$$

Note the "-" sign on the right-hand side and symmetry in the formulation.

2.2.2 Matrix C

Applying the simple quadrature rule, we conclude the matrix C is diagonal, i.e. $A_{ij} = 0$ if $i \neq j$. We then compute the contribution to A_{ii} from element T if v_i is a vertex belonging to T .

$$A_{ii}(T) = \int_T c(x) \phi_i \phi_i dx \quad (17)$$

$$= |T| \frac{c_i}{3} \quad (18)$$

2.2.3 Matrix D

Similar structure takes place in D . In fact D is sparser than C , for it's only non-zero entries corresponds to the nodes on the interface Γ defined above. Furthermore, applying a 1st-order quadrature rule for line integral, we have

$$\int_e k \phi_i \phi_j ds = \delta_{ij} |e| \frac{k_i}{2} \quad (19)$$

where $|e|$ is the length of edge e .

2.2.4 Right-hand side vector

$$\int_T f(x) \phi_i dx = |T| \frac{f_i}{3} \quad (20)$$

2.3 Dirichlet boundary conditions

Finally, we identify the nodes on Dirichlet boundaries. We modify the linear system defined above on both sides (assuming v_i is a node on Dirichlet boundaries):

- Left side:
We re-write the i -th row as a unit vector with i -th entry being 1 and 0's elsewhere.
- Right side:
Corresponding to the modification to the coefficient matrix on the left side, we change the i -th entry of right-hand side vector to be the Dirichlet boundary value $u_D(v_i)$