

培尖教育 2018 年学科竞赛夏令营物理模拟卷 (四)

考试时间: 150 分钟 总分 320 分 (参考答案)

1、解: (1)(i)在最低点有

$$mv_1 = mv_2$$

$$\frac{1}{2}mv_1^2 + \frac{1}{2}mv_2^2 = mgR$$

$$N - mg = \frac{m(v_1 + v_2)^2}{R}$$

$$\Rightarrow N = 5mg$$

(ii)设水平之间的相互作用冲量为 I,则有

$$I = mv_1$$

$$IR = (\frac{2}{3}mR^2 + mR^2)\omega$$

$$\frac{1}{2}mv_1^2 + \frac{1}{2} \times \frac{5}{3}mR^2\omega^2 = mgR$$

$$v_1 = \sqrt{\frac{5}{4}gR} \qquad \omega = \frac{3}{5}\sqrt{\frac{5g}{4R}}$$

有

$$N - mg = \frac{m(v_1 + R\omega)^2}{R}$$

$$\Rightarrow N = \frac{21}{5}mg$$

(2) (i)分离时设连线与竖直方向夹角为 θ

$$mv_x = mv$$

$$\tan \theta = \frac{v_y}{v_x + v}$$

$$mgR(1-\cos\theta) = \frac{1}{2}mv^2 + \frac{1}{2}m(v_x^2 + v_y^2)$$

且分离时,有条件

$$\frac{m(v_x + v)^2 + v_y^2}{R} = mg \cos \theta$$

$$\Rightarrow \theta \approx 42.94^{\circ} \left[\alpha = \arccos(\sqrt{3} - 1) \right]$$



 $_{
m (ii)}$ 仍设heta角,设水平冲量为 $I_{
m x}$ 。有

$$I_x = mv_x$$

$$I_{x}R = (\frac{2}{5}mR^{2} + mR^{2})\omega$$

且有
$$\frac{v_y}{v_x + \omega R} = \tan \theta$$

$$mgR(1-\cos\theta) = \frac{1}{2} \times \frac{7}{5} mR^2 \omega^2 + \frac{1}{2} m(v_x^2 + v_y^2)$$

分离时,有条件

$$mgR\cos\theta = \frac{m\left[(v_x + \omega R)^2 + v_y^2\right]}{R}$$

$$\Rightarrow \cos \theta = 44.1^{\circ}$$

2、解: (1) 由题意可得

$$2v_0\sin\theta = gt$$

$$d = v_0 \cos \theta t$$

可以得到

$$v_0^2 \sin 2\theta = gd$$

(2) 在竖直方向,速度变为了 $ev_0\sin\theta$,因而可得

$$I_N = (1+e)mv_0 \sin \theta$$

而
$$I_f \le \mu I_N = \mu (1+e) m v_0 \sin \theta$$

有关系

$$mv_0\cos\theta - I_f = mv_x$$

$$I_f r = \frac{2}{3} m r^2 \omega$$

因而可得

$$v_x = v_0 \cos \theta - \frac{I_f}{m}$$

$$\omega r = \frac{3I_f}{2m}$$



要求接触点速度

$$v_x - \omega r \ge 0$$

因而可得

$$I_f \le \frac{2}{5} m v_0 \cos \theta$$

因而,若
$$\mu \le \frac{2}{5(1+e)\tan\theta}$$
,则 $I_f = \mu(1+e)mv_0\sin\theta$,则

$$v_x = v_0 \left[\cos \theta - \mu (1+e) \sin \theta \right]$$

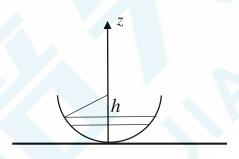
因而
$$\theta' = \arctan \left[\frac{\tan \theta}{1 - \mu(1 + e) \tan \theta} \right]$$

若
$$\mu > \frac{2}{5(1+e)\tan\theta}$$
,则 $I_f = \frac{2}{5}mv_0\cos\theta$,则

$$v_x = \frac{3}{5}v_0\cos\theta$$

因而
$$\theta' = \arctan\left(\frac{5}{3}\tan\theta\right)$$

3、解: (1) 如图所示



$$dq = 2\pi R\sigma dh$$

因而
$$d\varphi = \frac{dq}{4\pi\varepsilon_0\sqrt{R^2 - h^2 + (z - R + h)^2}} = \frac{\sigma R dh}{2\varepsilon_0\sqrt{2R^2 + z^2 - 2zR + 2(z - R)h}}$$

积分可得

$$\varphi = \int_{0}^{R} \frac{\sigma R dh}{2\varepsilon_0 \sqrt{2R^2 + z^2 - 2zR + 2(z - R)h}}$$
$$= \frac{\sigma R}{2(z - R)\varepsilon_0} (z - \sqrt{2R^2 + z^2 - 2zR})$$



这个是 $z \neq R$ 的情况, 当z = R时, 易得

$$\varphi = \frac{\sigma R}{2\varepsilon_0}$$

(2) 在球壳中央, 电场强度为

$$E = \int_{0}^{\pi/2} \frac{\sigma \cdot 2\pi R \sin \theta}{4\pi \varepsilon_0 R^2} \cos \theta R d\theta = \frac{\sigma}{4\varepsilon_0}$$

因而受力上满足

$$qE = mg$$

解得
$$m = \frac{\sigma q}{4\varepsilon_0 g}$$

(3) 能量上,满足方程

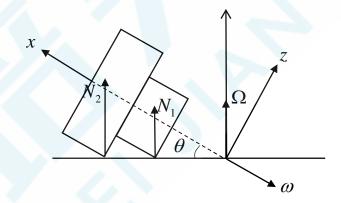
$$q\varphi_0 + 2mgR = q\varphi + mgz$$

其中,
$$\varphi_0 = \frac{2-\sqrt{2}}{2} \frac{\sigma R}{\varepsilon_0}$$
, $\varphi = \frac{\sigma R}{2(z-R)\varepsilon_0} (z-\sqrt{2R^2+z^2-2zR})$

联立可求得 z = 0.1012R

(4) 再补齐另一半球壳,则其电场必定为<mark>零</mark>,因而无论是向上微扰还是向下微扰,电场力的变化都是相同的,因而合力的方向也是相同的,运动趋势也相同,但是微扰方向不同,因而该点不可能是稳定平衡位置。

4、解:(1)画出侧视图



有几何关系
$$\tan \theta = \frac{R_2 - R_1}{h}$$

与地面接触点速度为零,因而有

$$\Omega l \cos \theta - \omega l \tan \theta = 0$$

解得
$$\Omega = \omega \frac{\tan \theta}{\cos \theta} = \omega \frac{(R_2 - R_1)\sqrt{(R_2 - R_1)^2 + h^2}}{h^2}$$

(2) 计算可得



转动惯量中

$$I_{xx} = \frac{1}{2} \rho \pi h (R_1^4 + R_2^4)$$

$$I_{zz} = \frac{1}{4} \rho \pi h (R_1^4 + R_2^4) + \frac{1}{3} (\frac{h}{R_2 - R_1})^3 \left[R_2^2 (2R_2 - R_1)^3 - R_2^5 + R_1^2 R_2^3 - R_1^5 \right]$$

因而 $L_z = I_{zz}\omega \tan \theta$

$$L_{x} = -I_{xx}(\omega - \Omega \sin \theta) = -I_{xx}\omega(1 - \tan^{2} \theta)$$

角动量的水平分量为

$$L = I_{xx}\omega(1 - \tan^2\theta)\cos\theta + I_{zz}\tan\theta\sin\theta$$

满足关系

$$N_{1} \frac{R_{1}}{\sin \theta} + N_{2} \frac{R_{2}}{\sin \theta} - \rho \pi R_{1}^{2} h g \left(\frac{R_{1}}{R_{2} - R_{1}} h + \frac{1}{2} h \right) \cos \theta - \rho \pi R_{2}^{2} h g \left(\frac{R_{2}}{R_{2} - R_{1}} h + \frac{1}{2} h \right) \cos \theta$$

$$= \omega L$$

又有关系

$$N_1 + N_2 = \rho \pi h g (R_1^2 + R_2^2)$$

解得

$$N_{1} = \frac{\omega L \sin \theta + m_{1}gl_{1} \sin \theta \cos \theta + m_{2}gl_{2} \sin \theta \cos \theta - (m_{1} + m_{2})gR_{2}}{R_{1} - R_{2}}$$

$$N_{2} = \frac{\omega L \sin \theta + m_{1}gl_{1} \sin \theta \cos \theta + m_{2}gl_{2} \sin \theta \cos \theta - (m_{1} + m_{2})gR_{1}}{R_{2} - R_{1}}$$

其中,
$$L = I_{xx}\omega(1-\tan^2\theta)\cos\theta + I_{zz}\tan\theta\sin\theta$$

$$I_{xx} = \frac{1}{2} \rho \pi h (R_1^4 + R_2^4)$$

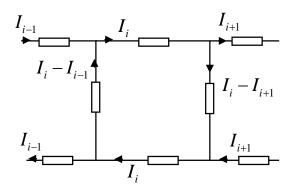
$$I_{zz} = \frac{1}{4} \rho \pi h (R_1^4 + R_2^4) + \frac{1}{3} (\frac{h}{R_2 - R_1})^3 \left[R_2^2 (2R_2 - R_1)^3 - R_2^5 + R_1^2 R_2^3 - R_1^5 \right]$$

$$l_1 = \frac{R_1}{R_2 - R_1} h + \frac{h}{2}; l_2 = \frac{R_2}{R_2 - R_1} h + \frac{h}{2}$$

$$m_1 = \rho \pi R_1^2 h; m_2 = \rho \pi R_2^2 h$$

5、对于第 i 个方框,如图所示





易得回路方程为

$$I_{i}\frac{2}{3}r+(I_{i}-I_{i+1})r+I_{i}\frac{2}{3}r+(I_{i}-I_{i-1})r=kL^{2}$$

化简为
$$I_{i+1} - \frac{10}{3}I_i + I_{i-1} = -\frac{kL^2}{r}$$

因而电流满足形式

$$I_i = A \cdot 3^i + B \cdot 3^{-i} + \frac{3kL^2}{4r}$$

代入
$$i = 0, i = N + 1$$
, 均有 $I = 0$

解得
$$A = -\frac{3kL^2}{4r} \frac{1 - 3^{-(N+1)}}{3^{N+1} - 3^{-(N+1)}}$$

$$B = -\frac{3kL^2}{4r} \frac{3^{N+1} - 1}{3^{N+1} - 3^{-(N+1)}}$$

因而可得

$$I_{i} = \frac{3kL^{2}}{4r} \left[1 - 3^{i} \frac{1 - 3^{-(N+1)}}{3^{N+1} - 3^{-(N+1)}} - 3^{-i} \frac{3^{N+1} - 1}{3^{N+1} - 3^{-(N+1)}} \right]$$

对于匀强磁场中的闭合电路, 其受力必定为零 功率满足

$$P = kL^2 \sum_{i=1}^{N} I_i$$

计算得

$$P = \frac{3k^2L^4}{4r} \left[N - \frac{3(3^N - 1)}{2(3^{N+1} + 1)} \right]$$

代入
$$N = 5$$
, 得 $P = \frac{9861}{2920} \frac{k^2 L^4}{r} = 3.377 \frac{k^2 L^4}{r}$



6、解:(1)内部压强为
$$p = p_0 + \frac{4\sigma}{r}$$

理想气体状态方程为 pV = vRT

其中
$$V = \frac{4}{3}\pi r^3$$

结合初态 $r = R_1$

联立可求得
$$T(r) = \frac{p_0 r^3 + 4\sigma r^2}{p_0 R_1^3 + 4\sigma R_1^2} T_0$$

(2) 初态,满足

$$dQ = -\kappa \frac{4\pi R_1^2 (T_1 - T_0)}{d} dt = C_V dT + (p_0 + \frac{4\sigma}{R_1}) \cdot 4\pi R_1^2 dr$$

$$\overline{m} \qquad dT = \frac{3p_0 R_1^2 + 8\sigma R_1}{p_0 R_1^3 + 4\sigma R_1^2} T_0 dr$$

代入可解得

$$\frac{dr}{dt} = -\frac{\kappa \cdot 4\pi R_1^2 (T_1 - T_0) (p_0 R_1^2 + 4\sigma R_1)}{d \left[C_V (3p_0 R + 8\sigma) T_0 + 4\pi (p_0 R_1^2 + 4\sigma R_1)^2 \right]}$$

(3)(i)肥皂泡带点后,由于带电产生的又一项附加压强为

$$p_e = -\frac{q^2}{32\pi\varepsilon_0 r^4}$$

因而内部压强表达式为

$$p = p_0 + \frac{4\sigma}{r} - \frac{q^2}{32\pi^2 \varepsilon_0 r^4}$$

结合
$$pV = vRT$$

$$V = \frac{4}{3}\pi r^3$$

可得
$$T(r) = \frac{p_0 r^3 + 4\sigma r^2 - \frac{q^2}{32\pi^2 \varepsilon_0 r}}{p_0 R_1^3 + 4\sigma R_1^2 - \frac{q^2}{32\pi^2 \varepsilon_0 R_1}} T_0$$

(ii) 由傅里叶热传导定律,可得

$$dQ = -\kappa \frac{4\pi R_1^2 (T_1 - T_0)}{d} dt = C_V dT + (p_0 + \frac{4\sigma}{R_0} - \frac{q^2}{32\pi^2 \varepsilon_0 R_0^4}) \cdot 4\pi R_1^2 dr$$



其中
$$dT = \frac{3p_0 R_1^2 + 8\sigma R_1 - \frac{q^2}{32\pi^2 \varepsilon_0 R_1^2}}{p_0 R_1^3 + 4\sigma R_1^2 - \frac{q^2}{32\pi^2 \varepsilon_0 R_1}} T_0 dr$$

联立可得

$$\frac{dr}{dt} = -\frac{\kappa \cdot 4\pi R_1^2 (T_1 - T_0)(p_0 R_1^2 + 4\sigma R_1 - \frac{q^2}{32\pi^2 \varepsilon_0 R_1^2})}{d \left[C_V \left(3p_0 R + 8\sigma - \frac{q^2}{32\pi^2 \varepsilon_0 R_1^3} \right) T_0 + 4\pi \left(p_0 R_1^2 + 4\sigma R_1 - \frac{q^2}{32\pi^2 \varepsilon_0 R_1^2} \right)^2 \right]}$$

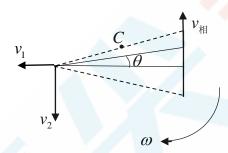
7、解:(1)对于一条边,相对于中心的转动惯量为

$$\Delta I = \frac{1}{12} \frac{m}{n} (2d \tan \frac{\pi}{n})^2 + \frac{m}{n} d^2 = \frac{m}{n} d^2 (1 + \frac{1}{3} \tan^2 \frac{\pi}{n})$$

因而总的转动惯量为

$$I = n \triangle I = md^2 \left(1 + \frac{1}{3} \tan^2 \frac{\pi}{n}\right)$$

(2) 在某一时刻,呈现如图位形时



可以表示出虫子在地面系中的速度

$$v_{\pm x} = -v_1 + \omega d \tan \theta$$

$$v_{\pm v} = -v_2 + \omega d + v_{\dagger l}$$

又有角动量守恒方程,即

$$mv_{\pm y} \frac{d}{2} + mv_2 \frac{d}{2} - mv_{\pm x} \frac{d \tan \theta}{2} - mv_1 \frac{d \tan \theta}{2} - I\omega = 0$$

又有
$$v_{\text{H}} = \frac{d}{\cos^2 \theta} \dot{\theta}$$

化简得到

$$\omega = \frac{\dot{\theta}}{1 + \cos^2 \theta (2 + \frac{2}{3} \tan^2 \frac{\pi}{n})}$$

因而可得



$$\alpha = 2n \int_{0}^{\frac{\pi}{n}} \frac{d\theta}{1 + \cos^2 \theta (2 + \frac{2}{3} \tan^2 \frac{\pi}{n})}$$

当n=3时,可得

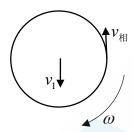
$$\alpha = 1.768 rad = 101.32^{\circ}$$

(3) 对于 $n \to \infty$,有

$$d\alpha = \frac{d\theta}{1 + \cos^2\theta (\frac{2}{3}\tan^2\frac{\pi}{n} + 2)} \approx \frac{d\theta}{3}$$

因而有 $\alpha = \frac{2\pi}{3}$

对于圆的情况来说,则有



角动量关系,满足

$$mv_{\pm}\frac{R}{2}+mv_1\frac{R}{2}=mR^2\omega$$

$$mv_1 = m(v_{\dagger \! \! 1} - v_1 - \omega R) = mv_{\pm}$$

解得

$$\omega = \frac{v_{\text{H}}}{3R}$$

因而可得

$$\alpha = \frac{2\pi}{3}$$

可见两者的结果是相同的。

8、解:

(1)
$$B_0 = \frac{\mu_0}{4\pi} \frac{2\pi RI}{R^2} = \frac{\mu_0 I}{2R}$$

(2)
$$B_z = \frac{\mu_0}{4\pi} \frac{2\pi RI}{(R^2 + Z^2)} \frac{R}{\sqrt{Z^2 + R^2}} = \frac{\mu_0 I R^2}{2(R^2 + Z^2)^{\frac{3}{2}}}$$

 $Z \gg R$

$$B_z = \frac{\mu_0 I R^2}{2 Z^2}$$



(3)

(i)
$$Li = \frac{\mu_0 I R^2}{2z_0^3} \pi r^2 = \frac{\pi \mu_0 I R^2 r^2}{2z_0^3}$$

$$i = \frac{\mu_0 I \pi R^2 r^2}{2z_0^3 L}$$

$$B_{zrdz}$$

$$B_r$$

$$B_z$$

有
$$B_z \pi r^2 = B_{z+dz} \pi r^2 + 2\pi r dz * B_r$$

$$\Rightarrow B_r = -\frac{r}{2} \frac{dB_z}{dz} = \frac{r}{2} \frac{\partial \mu_0 IR^2}{2z^4} = \frac{\partial \mu_0 IR^2 r}{4z^4}$$

$$\mathbb{B}_r = \frac{\partial \mu_0 I R^2 r}{4z^4}$$

因而有
$$2\pi riB_r = mg$$

$$\Rightarrow mg = 2\pi r * \frac{\mu_0 I \pi R^2 r^2}{2z_0^3 L} \frac{3\mu_0 I R^2 r}{4z_0^4} = \frac{3\pi^2 \mu_0^2 I^2 R^4 r^4}{4z_0^7 L}$$

$$\Rightarrow z_0 = \sqrt[7]{\frac{3\pi^2 \mu_0^2 I^2 R^4 r^4}{4mgL}}$$

因而
$$i = \frac{\mu_0 I \pi R^2 r^2}{2L} \left(\frac{3\pi^2 \mu_0^2 I^2 R^4 r^4}{4mgL} \right)^{-\frac{3}{7}}$$

(iii)
$$z_0 \rightarrow z_0 + \Delta z$$
 时

$$F = \frac{3\pi^2 \mu_0 I^2 R^4 r^4}{4(z_0 + \Delta z)^7 L} = mg(1 - \frac{7\Delta z}{z_0})$$

$$m\Delta z^{J} = -mg\,\frac{\Delta z}{z_0}$$

 $\Rightarrow \Delta z^{\prime} = -\frac{g}{z_0} \Delta z$

$$\Rightarrow \omega = \sqrt{\frac{g}{z_0}} = \sqrt{g\sqrt{\frac{4mgL}{3\pi^2\mu_0^2I^2R^4r^4}}} = \sqrt{\frac{4mg^8L}{3\pi^2\mu_0^2I^2R^4r^4}}$$

$$\Rightarrow z = z_0 + \Delta z_0 \cos \omega t$$