

培尖教育 2018 年学科竞赛夏令营物理模拟卷 (七)

考试时间: 150 分钟 总分 320 分

(参考答案)

一、 1. 设小球、大球的加速度分别为 a_1, a_2 , 细杆的角加速度为 β
列出动力学方程

$$mg - N_1 \sin 30 = ma_1$$

$$N_1 \cos 30 - N_2 \sin 30 = 4ma_2$$

$$N_2 \cdot 2\sqrt{3}r - 3mg \cdot 2r \sin 30 = \frac{1}{3} \cdot 3m(4r)^2 \beta$$

列出加速度关系: 法向加速度相同

$$a_1 \sin 30 = a_2 \cos 30$$

$$2\sqrt{3}r\beta = a_2 \cos 30$$

联立以上五式, 解得

$$a_1 = \frac{9\sqrt{3}g}{4(7\sqrt{3}+1)} \quad a_2 = \frac{9g}{4(7\sqrt{3}+1)} \quad \beta = \frac{9g}{16(7\sqrt{3}+1)r}$$

2. 据法向速度相同

$$v_2 \cos 30 = v_0 \sin 30 \quad v_2 = v_0 \tan 30$$

取小球为系, 则大球以 $v' = \frac{v_0}{\cos 30}$ 绕小球旋转, 故

$$a_n = \frac{v'^2}{3r} = \frac{4v_0^2}{9r}$$

$$a_n \text{ 为 } a_2 \text{ 在法向的投影, } \therefore a_2 = \frac{a_n}{\cos 30} = \frac{8\sqrt{3}v^2}{27r}$$

$$\text{杆的角速度 } 2\sqrt{3}\omega r = v_2 \cos 30 \quad \omega = \frac{\sqrt{3}v_0}{12r}$$

据加速度关联可得 $a_c = -a_2 \cos 30 - \frac{v_3^2}{2r}$, 其中 $v_3 = \frac{1}{2}v_2$ 为杆相对球的切向速度

以 A 点为原点写出 C 点加速度的极坐标表达式 $a_c = 2\ddot{l}\theta + l\ddot{\theta}$

$$AC=AB, \text{ 故 } \dot{l} = -v_2$$

$$\text{联立(10)(11)两式可得 } \beta = -\frac{23\sqrt{3}v^2}{432r^2}$$

$$\text{改写(1)式为 } F + N_1 \sin 30 = mg$$

联立(2)(3)(13)式可得

$$F = \frac{3}{4}mg + \left(\frac{32}{27} + \frac{23\sqrt{3}}{324}\right)\frac{mv_0^2}{r}$$

二、 显然 $kv_0^2 = 0.44mg$

运用动能定理, 上升过程中

$$-kv^2 dy - mgdy = d\left(\frac{1}{2}mv^2\right)$$

$$-(kv^2 + mg)dy = mv dv$$

$$-dy = \frac{mv dv}{kv^2 + mg}$$

$$\text{积分得 } \ln \frac{mg}{kv^2 + mg} = -\frac{2k}{m} H$$

$$\text{代入数据得 } H = \frac{m}{k} \ln \frac{6}{5}$$

下落过程中

$$-kv^2 dy + mg dy = mv dv \quad (\text{两次 } y \text{ 的方向取得相反})$$

$$dy = \frac{mv dv}{mg - kv^2}$$

$$\text{积分, 得 } -\frac{2k}{m} H = \ln \frac{mg - kv^2}{mg}$$

$$\text{得 } v = \sqrt{\frac{11mg}{36k}}$$

三、 1. 设下滑 x 时, 对整体分析, 列出 $F = \frac{dp}{dt}$

$$(m + \lambda x) g \sin \theta = \frac{d(m + \lambda x)v}{dt} = \frac{d(m + \lambda x)v}{dx} \frac{dx}{dt}$$

$$(m + \lambda x)^2 g \sin \theta dx = \frac{1}{2} d[(m + \lambda x)v]^2$$

$$\text{积分得 } \frac{g \sin \theta}{3\lambda} [(m + \lambda x)^3 - m^3] = \frac{1}{2} [(m + \lambda x)v]^2$$

$$v = \sqrt{\frac{2g \sin \theta [(m + \lambda x)^3 - m^3]}{3\lambda(m + \lambda x)^2}}$$

$$2. \text{设绳子长为 } l, \text{ 则下降 } x \text{ 时的质心位置为 } x_c = \frac{mx + \frac{1}{2}\lambda x^2}{m + \lambda l}$$

$$\text{对其求导, 可得 } dx_c = \frac{m + \lambda x}{m + \lambda l} dx$$

据质心运动定理, 可得

$$\int_0^x (m + \lambda x) g \sin \theta dx_c = \frac{1}{2} (m + \lambda l) \left(\frac{m + \lambda x}{m + \lambda l} v \right)^2$$

代入式, 得

$$\int_0^x (m + \lambda x)^2 g \sin \theta dx = \frac{1}{2} (m + \lambda l)^2 \left(\frac{m + \lambda x}{m + \lambda l} v \right)^2$$

$$\text{得 } v = \sqrt{\frac{2g \sin \theta [(m + \lambda x)^3 - m^3]}{3\lambda(m + \lambda x)^2}}$$

四、 取管系研究, 设 P_0 不变可得

$$\frac{T}{V} = \text{Const} \quad \frac{T_0 + \Delta T}{V_0 + \Delta V} = \frac{T_0}{V_0} \text{ 得 } T_0 \Delta V = V_0 \Delta T \dots\dots\dots(1)$$

$$\text{理想气体状态方程 } P_0 V = \nu R T_0 \text{ 其中 } V = \nu S \Delta t \dots\dots\dots(2)$$

据功能原理 $\frac{5}{2} \nu R \Delta T = q \Delta t - P_0 \Delta V \dots\dots\dots(3)$

$$(1)(3) \text{ 联立得 } \frac{7}{2} P_0 \Delta V = q \Delta t \quad \frac{\Delta V}{\Delta t} = \frac{2q}{7P_0} \quad \Delta v \cdot S = \frac{2q}{7P_0} \dots\dots\dots(4)$$

$$\text{代入(2)式得 } \Delta v = \frac{2qv\Delta t}{7\nu RT_0}$$

$$\therefore m \frac{\Delta v}{\Delta t} = \frac{2qv}{7\nu RT_0} \cdot \nu \mu = \frac{2qv\mu}{7RT_0} = \frac{\Delta P}{\Delta t} = M \frac{dv}{dt}$$

积分可得 $v = v_0 e^{\frac{2q\mu}{7MRT_0} t}$

$$x = \frac{7MRT_0 v_0}{2q\mu} (e^{\frac{2q\mu}{7MRT_0} t} - 1)$$

五、 1. 旋转时，体电流密度为

$$\text{六、 } j = \rho v = \rho r(\omega_0 + \beta t)$$

在 r 处 $B(r)$ 可积分为为（运用无限长直螺线管公式）

$$B = \int_r^a \mu_0 n i = \int_r^a \mu_0 j r dr = \frac{1}{2} \rho \mu_0 (\omega_0 + \beta t)(a^2 - r^2)$$

或直接运用

$$\nabla \times \vec{B} = \mu_0 \vec{j} \quad -\frac{\partial B}{\partial r} = \mu_0 \rho (\omega_0 + \beta t) r$$

利用边界条件 $r = a, B = 0$

$$\text{得 } B = \frac{1}{2} \rho \mu_0 (\omega_0 + \beta t) (a^2 - r^2)$$

$$\text{综上 } r < a \text{ 时 } B = \frac{1}{2} \rho \mu_0 (\omega_0 + \beta t) (a^2 - r^2)$$

$r > a$ 时 $B = 0$

2. 据高斯定理 $r < a$ 时

$$E_r = \frac{\rho \pi r^2 l}{\varepsilon_0 2 \pi r l} = \frac{\rho r}{2 \varepsilon_0}$$

$$\text{磁通量 } \varphi = \int_0^r B(r) 2 \pi r dr$$

$$= \pi \rho \mu_0 (\omega_0 + \beta t) \left(\frac{1}{2} a^2 r^2 - \frac{r^4}{4} \right)$$

$$E_\theta = - \frac{\partial \varphi}{\partial t} \div 2 \pi r = - \frac{\rho \mu_0 \beta}{8} (2 a^2 r - r^3)$$

$r > a$ 时

$$E_r = \frac{\rho a^2}{2 r}$$

$$E_\theta = \frac{\rho \mu_0 \beta a^4}{8 r}$$

3. 外力矩为电磁力矩与使其加速的力矩之和

$$M_1 = \frac{1}{2} M a^2 \beta$$

$$\begin{aligned} M_2 &= \int r \times F = \int_0^a r \cdot \frac{\rho \mu_0 \beta}{8} (2 a^2 r - r^3) \rho 2 \pi r l dr \\ &= \frac{\mu_0 \rho^2 \beta l \pi a^6}{12} \end{aligned}$$

$$M = M_1 + M_2 = \frac{1}{2} M a^2 \beta + \frac{\mu_0 \rho^2 \beta l \pi a^6}{12}$$

$$\text{六、1. 火星的运行速度为 } v_0 = \sqrt{\frac{GM}{R}} \dots\dots\dots(1)$$

无限远罗向太阳的彗星可视为能量为零，且垂直进入火星轨道

$$\text{故 } v = \sqrt{\frac{2GM}{R}}, \text{ 相对速度为 } \sqrt{\frac{3GM}{R}} = \sqrt{3} v_0 \dots\dots\dots(2)$$

据引力弹弓效应，转换到原参考系可得

$$v' = (\sqrt{3} - 1) v_0 \dots\dots\dots(3)$$

据能量关系

$$\frac{1}{2} m v'^2 - \frac{GMm}{R} = - \frac{GMm}{2a} \dots\dots\dots(4)$$

$$\text{得 } a = \frac{\sqrt{3} + 1}{4} R \dots\dots\dots(5)$$

$$\therefore r_1 = 2a - R = \frac{\sqrt{3} - 1}{2} R \dots\dots\dots(6)$$

2.角动量守恒

$$v_1 r_1 = (\sqrt{3} - 1)v_0 R \quad v_1 = 2v_0 \dots\dots\dots(7)$$

$$\text{恰好圆周运动时 } v_2 = \sqrt{\frac{GM}{r_1}} = \sqrt{(\sqrt{3} + 1)v_0} \dots\dots\dots(8)$$

分裂过程中动量守恒

$$m \cdot 2v_0 = \frac{m}{2}v_2 + \frac{m}{2}v_3 \quad v_3 = (4 - \sqrt{\sqrt{3} + 1})v_0 \dots\dots\dots(9)$$

据(4)式可得

$$a = \frac{R}{8\sqrt{\sqrt{3} + 1} + \sqrt{3} - 15} \dots\dots\dots(10)$$

七、先推导余弦型电荷分布在球内产生的场强，取球心分析

$$\sigma_{(\theta)} = \sigma_0 \cos \theta$$

$$dE = \frac{2\pi R \sin \theta R d\theta \cdot \sigma_{(\theta)}}{4\pi \epsilon_0 R^2} \cos \theta \dots\dots\dots(1)$$

$$E = \int_0^\pi dE = \frac{\sigma_0}{3\epsilon_0} \dots\dots\dots(2)$$

设内表面 $\sigma_{(\theta)} = \sigma_1 \cos \theta$ ，外表面 $\sigma_{(\theta)} = \sigma_2 \cos \theta$

则内表面电荷等效电偶极子为

$$p = \frac{4}{3}\pi b^3 \sigma_1 \quad E_{(\theta)} = \frac{2p \cos \theta}{4\pi \epsilon_0 r^3} \dots\dots\dots(3)$$

在外表面处研究

$$E_{\text{外}} = \frac{2p \cos \theta}{4\pi \epsilon_0 a^3} + \frac{2\sigma_2}{3\epsilon_0} \cos \theta + E_0 \cos \theta \dots\dots\dots(4)$$

$$E_{\text{内}} = \frac{2p \cos \theta}{4\pi \epsilon_0 a^3} - \frac{\sigma_2}{3\epsilon_0} \cos \theta + E_0 \cos \theta \dots\dots\dots(5)$$

据电位移连续，有 $\epsilon_r E_{\text{内}} = \epsilon_0 E_{\text{外}} \dots\dots\dots(6)$

$$\text{得 } \epsilon_r \left(\frac{2p \cos \theta}{4\pi \epsilon_0 a^3} - \frac{\sigma_2}{3\epsilon_0} \cos \theta + E_0 \cos \theta \right) = \epsilon_0 \left(\frac{2p \cos \theta}{4\pi \epsilon_0 a^3} + \frac{2\sigma_2}{3\epsilon_0} \cos \theta + E_0 \cos \theta \right) \dots\dots\dots(7)$$

在内表面处研究

$$E_{\text{外}} = \frac{2\sigma_2 \cos \theta}{3\epsilon_0} - \frac{\sigma_1 \cos \theta}{3\epsilon_0} + E_0 \cos \theta \dots\dots\dots(8)$$

$$E_{\text{内}} = -\frac{\sigma_2 \cos \theta}{3\epsilon_0} - \frac{\sigma_1 \cos \theta}{3\epsilon_0} + E_0 \cos \theta \dots\dots\dots(9)$$

据电位移连续，有 $\epsilon_0 E_{\text{内}} = \epsilon_r E_{\text{外}} \dots\dots\dots(10)$

$$\text{得 } \epsilon_0 \left(-\frac{\sigma_2 \cos \theta}{3\epsilon_0} - \frac{\sigma_1 \cos \theta}{3\epsilon_0} + E_0 \cos \theta \right) = \epsilon_r \left(\frac{2\sigma_2 \cos \theta}{3\epsilon_0} - \frac{\sigma_1 \cos \theta}{3\epsilon_0} + E_0 \cos \theta \right) \dots\dots\dots(11)$$

(3)(7)(11)联立，解得

$$\sigma_1 = -\frac{9\epsilon_0 E_0 a^3 (\epsilon_r - 1)}{a^3 (\epsilon_r + 2)(2\epsilon_r + 1) - 2b^3 (\epsilon_r - 1)^2}$$

$$\sigma_2 = \frac{3\varepsilon_0 E_0 (\varepsilon_r - 1) [a^3 (2\varepsilon_r + 1) - 2b^3 (\varepsilon_r - 1)]}{a^3 (\varepsilon_r + 2)(2\varepsilon_r + 1) - 2b^3 (\varepsilon_r - 1)^2}$$

题八:

$v_0 = \sqrt{\frac{2E_0}{M}}$; 因水平面方向受合外力为 0, 所以分裂过程动量守恒. 有

$$\begin{cases} Mv_0 = \lambda Mv_1 \cos \theta + (1-\lambda)Mv_2 \cos \theta, \\ \lambda Mv_1 \sin \theta = (1-\lambda)Mv_2 \sin \theta. \end{cases}$$

$$\text{得} \begin{cases} v_0 = [\lambda v_1 + (1-\lambda)v_2] \cos \theta, \\ \lambda v_1 = (1-\lambda)v_2. \end{cases}$$

$$\text{得 } v_1 = \frac{v_0}{2\lambda \cos \theta} = \frac{\sec \theta}{2\lambda} \sqrt{\frac{2E_0}{M}},$$

$$v_2 = \frac{\lambda}{1-\lambda} v_1 = \frac{\sec \theta}{1-\lambda} \sqrt{\frac{E_0}{2M}}.$$

(2) 分裂后两裂块的总能量

$$E = \frac{1}{2} \lambda M v_1^2 + \frac{1}{2} (1-\lambda) M v_2^2 = \frac{E_0}{4\lambda(1-\lambda)} \sec^2 \theta.$$

$$\text{内部机构提供的能量 } E' = E - E_0 = E_0 \left(\frac{\sec^2 \theta}{4\lambda(1-\lambda)} - 1 \right).$$

求 E' 的极小值即求 $\lambda(1-\lambda)$ 的极大值, $\frac{d\lambda(1-\lambda)}{d\lambda} = 0$,

$$\text{解得 } \lambda = \frac{1}{2}, \text{ 因 } \frac{d^2[\lambda(1-\lambda)]}{d\lambda^2} \Big|_{1/2} = -2 < 0.$$

表明 $\lambda = \frac{1}{2}$ 时, E' 有极小值, $E'_{\min} = E_0 (\sec^2 \theta - 1) = E_0 \tan^2 \theta$.

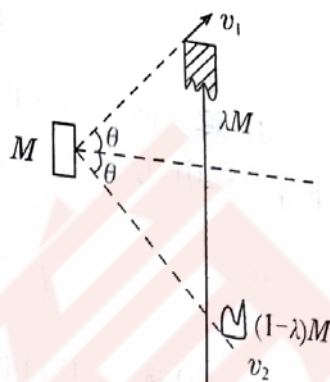


图 5-35