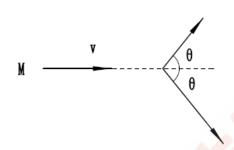
培尖教育 2018 年学科竞赛夏令营物理模拟卷 (九)

考试时间: 150 分钟 总分 320 分

(参考答案)

1 设一个质量为 α M,则另一个质量为 $(1-\alpha)$ M 取质心系 (向右以速度 ν_0 运动的参考系),设在质心系中, α M 物体速度 ν_1 ,则另一个质量为 $(1-\alpha)$ M



 α) **M** 的物体速度为 $\frac{\alpha v_1}{1-\alpha}$ (质心系动量为 0)

$$\Delta E = \frac{1}{2}\mu u^2$$
, $\mu = \frac{\alpha M \cdot (1-\alpha)M}{M} = \alpha(1-\alpha)M$ (2 $\%$)

$$u = v_1 + \frac{\alpha}{1-\alpha}v_1 = \frac{1}{1-\alpha}v_1 \ (1 \ \%)$$

$$\Delta E = \frac{1}{2} M v_1^2 \frac{\alpha}{1-\alpha} \quad (1 \text{ } \%)$$

由正弦定理
$$\frac{v_0}{\sin \theta_1} = \frac{v_1}{\sin \theta}$$
 (1分)

$$\frac{\alpha v_1}{(1-\alpha)\sin\theta} = \frac{v_0}{\sin\theta_2} = \frac{v_0}{\sin(\theta_1 + 2\theta)} \quad (2 \ \%)$$

得到 ν_1 的两个表达式 $\nu_1 = \frac{\nu_0 \sin \theta}{\sin \theta_1}$ (1分)

$$\pi \nabla_1 = \frac{v_0(1-\alpha)\sin\theta}{\alpha\sin(\theta_1+2\theta)}$$

有
$$\nu_1 \frac{\alpha}{1-\alpha} = \frac{\nu_0 \sin \theta}{\sin(\theta_1 + 2\theta)}$$
 (1分)

$$\Delta E = \frac{1}{2} M \nu_1 \frac{\alpha}{1-\alpha} \ \nu_1 \ (1 \ \%)$$

$$\Delta E = \frac{1}{2} M \frac{v_0 \sin \theta}{\sin \theta_1} \cdot \frac{v_0 \sin \theta}{\sin(\theta_1 + 2\theta)} (2 \%)$$

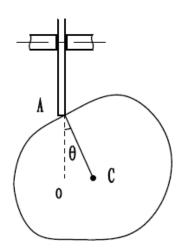
$$\theta_1 + \theta = \frac{1}{2}\pi$$
 时对应最小放出能量(1分)

$$\sin \theta_1 \sin(\theta_1 + 2\theta) = \cos^2 \theta \quad (1 \, \%)$$

$$\Delta E = \frac{I}{2}Mv_0^2 \tan^2 \theta \quad (1 \, \text{\frac{\beta}{2}})$$



2 (1) $wL \sin \theta = v \cos \theta$ (2 %) $v = \omega L \tan \theta$ (2 %)



(2)
$$a = \dot{v} = w\dot{L}\tan\theta + \frac{\omega L}{\cos^2\theta}\dot{\theta}$$
 (2 $\dot{\gamma}$)

$$\dot{L} = v$$
 (2分)

由余弦定理 $L^2 + \rho^2 - 2L\rho\cos\theta = \overline{00'}^2$ (2分)

对时间求导 AO 间距不变,对时间导数为 0 (1分)

ρ瞬时不变,对时间导数为0(1分)

余弦定理表达式对时间求导 $2L\dot{L}-2\dot{L}\rho\cos\theta+2L\rho\sin\theta\cdot\dot{\theta}=0$ (2分)

$$\dot{\theta} = \frac{-L\dot{L} + \rho\dot{L}\cos\theta}{L\rho\sin\theta} \ (2 \ \dot{\mathcal{D}})$$

代入
$$\dot{L} = v$$
 得到 $\dot{\theta} = \frac{(-L + \rho \cos \theta) v}{L\rho \sin \theta} = \frac{(-L + \rho \cos \theta) \omega}{\rho \cos \theta}$ (2分)

所以
$$a=\omega^2 L \tan^2 \theta + \frac{(-L+\rho\cos\theta)\omega^2 L}{\rho\cos^3\theta}$$
 (2 分)

再次相对速度为 0 时,相距为 x

取质心系则惯性力不做功

电荷 1 运动了距离
$$\frac{m_2}{m_1+m_2}(L-x)$$
 (2 分)

电荷 2 运动了距离
$$\frac{m_1}{m_1+m_2}(L-x)$$
 (2 分)

设



$$\frac{kq_1q_2}{x} + Eq_1 \frac{m_2}{m_1 + m_2} (L - x) - Eq_2 \frac{m_1}{m_1 + m_2} (L - x) = \frac{kq_1q_2}{L} \ (3 \ \%)$$

$$\frac{kq_1q_2}{xL}(L-x) = E \frac{m_1q_2-m_2q_1}{m_1+m_2}(L-x) (2 \%)$$

$$X = \frac{kq_1q_2(m_1+m_2)}{EL(m_1q_2-m_2q_1)} (2 \ \%)$$

$$X < L$$
 时,即 $\frac{kq_1q_2(m_1+m_2)}{E(m_1q_2-m_2q_1)} < L^2$ 时,最近为 $X = \frac{kq_1q_2(m_1+m_2)}{EL(m_1q_2-m_2q_1)}$ (2 分)

$$X>L$$
时,即 $\frac{kq_1q_2(m_1+m_2)}{E(m_1q_2-m_2q_1)}>L^2$ 时,最近为 L(2 分)

- 4(1)在外部放一电荷 Q, 等效电荷为
- ① 电荷 Q, 在距 O 为 a 处 (1分)
- ② 电荷 $-\frac{R}{a}Q$, 在 OQ 连线上距 O 为 $\frac{R^2}{a}$ (1分)
- ③ 电荷 $\frac{R}{a}Q$,在O点(1分)

在内部放一电荷 Q,在外部等效电荷为放在中心 Q 的 Q (1分) 在内部放一电荷 Q,在内部等效电荷为

- ① 电荷 Q, 在距 O 为 b 处 (1分)
- ② 电荷 $-\frac{R}{b}Q$,在 OQ 连线延长线上距 O 为 $\frac{R^2}{b}$ (1分)
- (2) 计算电能时要还原电荷原本位置 $E = \sum_{i=1}^{1} \varphi_i q_i$

球壳上净电荷为0且等势,造成静电能为0(1分)

计算点电荷电势时不考虑自己造成的 $\varphi = \frac{\kappa_a^R Q}{a} - \frac{k_a^R Q}{a - \frac{R^2}{a}} (1 分)$

$$\varphi = kRQ\left(\frac{1}{a^2} - \frac{1}{a^2 - R^2}\right) (2 \%)$$

E =
$$\frac{1}{2}kRQ^2\left(\frac{1}{a^2} - \frac{1}{a^2 - R^2}\right)$$
 (1 $\%$)

(3)计算电能时要还原电荷原本位置 $E = \sum_{i=1}^{1} \varphi_i q_i$

在内部仅考虑电荷 $-\frac{R}{b}Q$, 在 OQ 连线延长,线上距 O 为 $\frac{R^2}{b}$

从 O 到电荷 Q 处电势变化 $\varphi = -K\frac{R}{b}Q \left(\frac{1}{\frac{R^2}{b}-b} - \frac{1}{\frac{R^2}{b}}\right)$ (1 分)

$$\varphi = -kRQ \left(\frac{1}{R^2 - b^2} - \frac{1}{R^2} \right) (1 \%)$$

$$E = -\frac{1}{2}kRQ^2\left(\frac{1}{R^2-h^2} - \frac{1}{R^2}\right) (1 \%)$$

(4)此时壳上净电荷不为0,但壳电势为0,故壳上电荷静电能仍为0



此时壳内壁有电荷-Q,造成O点电势为 $\varphi_0 = -k\frac{Q}{R}$ (1分)

造成额外电势能 $\triangle E = -\frac{1}{2}k\frac{Q^2}{R}$ (1分)

E =
$$-\frac{1}{2}kRQ^2\left(\frac{1}{R^2-h^2} - \frac{1}{R^2}\right) - \frac{1}{2}k\frac{Q^2}{R}$$
 (1 $\%$)

$$E = -\frac{1}{2} \frac{kRQ^2}{R^2 - h^2} (2 \%)$$

(5) 球壳上净电荷为0且等势,造成静电能为0内壳有电荷-Q,外壳有电荷+Q

O 点电势为
$$\varphi_0 = -k\frac{\varrho}{r_1} + k\frac{\varrho}{r_2}$$
 (1分)

点电荷 Q 处电势

$$\varphi = -kr_1Q\left(\frac{1}{r_1^2 - b^2} - \frac{1}{r_1^2}\right) - k\frac{Q}{r_1} + k\frac{Q}{r_2} (1 \%)$$

$$\varphi = -kr_1Q \left(\frac{1}{r_1^2 - b^2} - \frac{1}{r_1r_2} \right) (2 \%)$$

$$E = -\frac{1}{2}kr_1Q^2 \left(\frac{1}{r_1^2 - b^2} - \frac{1}{r_1r_2}\right) (1 \%)$$

(6)电荷分布的本质为外壳+Q 均匀分布,内壳包含均匀分布的部分和非均匀分布的部分,总量-0

外壳电荷对球内任意一点造成电势

$$\varphi = k \frac{Q}{r_2} \ (1 \ \text{β})$$

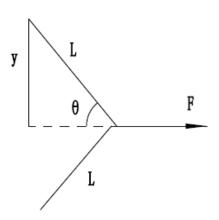
对空腔球心 o_1 ,电势为 $\varphi_{o_1} = k \frac{Q}{r_2} - k \frac{Q}{r_1}$ (1分)

故结果与上问相同

$$E = -\frac{1}{2}kr_1Q^2\left(\frac{1}{r_1^2 - b^2} - \frac{1}{r_1r_2}\right) (2 \, \cancel{/})$$

5 (1)
$$v_x = a_c t \ (1 \, \%)$$

$$F = 2ma_c, \ a_c = \frac{F}{2m} \ (1 \, \text{\%})$$



仅需求 t 即可,取质心系,由动能定理(惯性力不做功)



FLcos $\theta = \frac{1}{2} m v_y^2 \times 2 \ (1 \ \%)$

$$v_y = \sqrt{\frac{FL\cos\theta}{m}}$$
 (1 $\%$)

$$\cos\theta = \frac{\sqrt{L^2 - y^2}}{L} \quad (1 \, \text{\refta})$$

$$y = L \sin \theta \quad (1 \, \text{分})$$

 $dt = \frac{dy}{v_y}$ (此处未关注符号,默认各物理量为正)

$$dt = \sqrt{\frac{m}{F}} \frac{d(\frac{y}{L})}{\left(1 - \frac{y^2}{L^2}\right)^{\frac{1}{4}}} L^{\frac{1}{2}} \quad (1 \; \text{$\frac{1}{2}$})$$

$$t = \sqrt{\frac{m}{F}} \int_0^1 \frac{d(\frac{y}{L})}{\left(1 - \frac{y^2}{L^2}\right)^{\frac{1}{4}}} L^{\frac{1}{2}} = 1.197 \sqrt{\frac{mL}{F}} (1 \%)$$

$$v_x = 0.60 \sqrt{\frac{FL}{m}} \quad (1 \, \text{分})$$

(2) 由
$$F_x = -ax^2$$
, 积分得势能

$$Ep = \int_0^x ax^2 dx = \frac{1}{3}\alpha x^3$$
 (1 $\%$)

$$E = \frac{1}{3} \alpha A^3 \ (1 \ \%)$$

$$E_k = \frac{1}{3}\alpha(A^3 - x^3)$$
 (1 $\%$)

故 v=
$$\sqrt{\frac{2}{3}\frac{\alpha}{m}(A^3-x^3)}$$
 (1分)

$$dt = \frac{dx}{v} = \frac{A d\left(\frac{x}{A}\right)}{\sqrt{\frac{2\alpha}{3m}(A^3 - x^3)}} = \frac{A d\left(\frac{x}{A}\right)}{\sqrt{\frac{2\alpha}{3m}A^3\left(1 - \left(\frac{x}{A}\right)^3\right)}} (2 / T)$$

$$T = 4\sqrt{\frac{3m}{2\alpha A}} \int_0^1 \frac{d\left(\frac{x}{A}\right)}{\sqrt{1-\left(\frac{x}{A}\right)^3}} (2 \%)$$

由于计算器的限制,1不能取到

$$T = 4\sqrt{\frac{3m}{2\alpha A}} \int_{0}^{0.99999} \frac{d\left(\frac{x}{A}\right)}{\sqrt{1 - \left(\frac{x}{A}\right)^{3}}}$$

得 T=6.85
$$\sqrt{\frac{m}{\alpha A}}$$
 (1分)

6 由对称性, 易知光路对称



经凹透镜 L_3 时 $u_3 = \nu_3$ (2分)

①
$$ext{d} \frac{1}{u_3} + \frac{1}{v_3} = \frac{1}{-f} (1 \%)$$

得
$$u_3 = v_3 = -2f(1分)$$

经凸透镜时
$$\frac{1}{x} + \frac{1}{2f+d} = \frac{1}{f}$$
 (2分)

得
$$x = \frac{f(2f+d)}{f+d}$$
 (1分)

②
$$u_3 = v_3 = 0 \ (1 \ \%)$$

即光线在凹透镜中心汇聚,此时仍满足对称光路

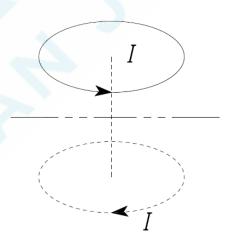
$$\frac{1}{x} + \frac{1}{d} = \frac{1}{f} (2 \%)$$

$$x = \frac{fd}{d-f} \ (1 \ \text{分})$$

讨论当 d>f 时,
$$x = \frac{f(2f+d)}{f+d}$$
或 $x = \frac{fd}{d-f}$ (2分)

7、(1) 通电圆环悬浮在 z=h 处,超导体的内部磁感应强度为零而表面外侧磁感应强度与表面平行,这可等效为通电圆环与它的像电流——在 z=-h 虚设一个相同的通以反向电流的环——共同产生的结果,如图所示,通电圆环必须有其所受重力与像电流施予的磁场力相平衡,由 r<<h 这个条件,将两环形电流近似为反向平行电流:

$$F = \frac{\mu_0}{2\pi} * \frac{I}{2h} * I * 2\pi r$$
 (2分)

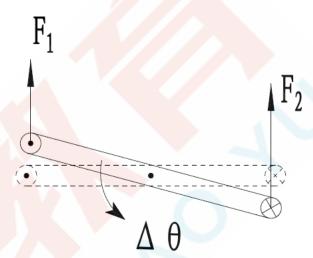


$$I = \sqrt{\frac{2mgh}{\mu_0 r}} \ (2 \ \%)$$



(2) 若令圆环水平地上下振动,当与平衡位置有任一位移(如向下 x)时, Σ F=mg- $\frac{\mu_0 r I^2}{2(h-x)}$ \approx mg- $\frac{\mu_0 r I^2}{2h} \left(1 + \frac{x}{h}\right) = -\frac{\mu_0 r I^2}{2h^2}$ x (3分),可见为一简谐运动,回复力常数 $\frac{\mu_0 r I^2}{2h^2} = \frac{mg}{h}$ (2分) 则 $T_1 = 2\pi \sqrt{\frac{h}{g}}$ (2分)

(3) 当以 P_1P_2 为轴作小幅摆动时,圆环转动惯量 $J=\frac{1}{2}mr^2$ (1分),当有一小量角位移 Δ θ 时,如图所示,有合力矩:



$$\begin{split} & \sum \mathbf{M} = \ (\mathbf{F}_1 - \mathbf{F}_2) \ \mathbf{r} \ (1 \, \dot{\mathcal{T}}) \\ & = \frac{\mu_0 I^2 r^2}{2} \left(\frac{1}{2 \hbar + r * \Delta \theta} - \frac{1}{2 \hbar - r * \Delta \theta} \right) \\ & = \ \frac{\mu_0 I^2 r^2}{4 \hbar} \left[\ (1 - \frac{r}{2 \hbar} * \Delta \theta) - \ (1 + \frac{r}{2 \hbar} * \Delta \theta) \right] \\ & = - \ \frac{\mu_0 I^2 r^2}{4 \hbar^2} \ \Delta \theta \ = - \ \frac{m g r^2}{2 \hbar} \ \Delta \theta, \ (4 \, \dot{\mathcal{T}}) \end{split}$$

$$& \boxed{\mathbb{M}} \ \mathbf{T}_2 = \ 2 \pi \sqrt{\frac{I}{K}} = \ 2 \pi \sqrt{\frac{\frac{1}{2} m r^2}{2 \hbar}} = \ 2 \pi \sqrt{\frac{\hbar}{g}} \ (2 \, \dot{\mathcal{T}}) \end{split}$$

8 (1),
$$q\nu_0 B = \frac{m\nu_0^2}{r_0}$$
 (2 $\frac{h}{m}$) $r_0 = \frac{m\nu_0}{q_B} = 3.41 \times 10^{-3} m(3 \frac{h}{m})$

(2) 设
$$\frac{1}{4\pi\epsilon_0} \frac{2q^2}{3c^3} = k$$
,近似有 $a = \omega^2 r$ (1分)

则辐射功率 $P=k(\omega^2 r)^2(1分)$

磁场中
$$\omega = \frac{\nu}{r} = \frac{qB}{m}$$
为一常量(1分)

$$dE = k (\omega^2 r)^2 dt = d\left(\frac{1}{2}mv^2\right) (1 \%)$$

$$v = \omega r \quad (1 \, \text{分})$$

dE=k
$$(ω^2 r)^2 dt = d\left(\frac{1}{2}mω^2 r^2\right) (2 \%)$$



 $dE = k\omega^4 r^2 dt = m\omega^2 r dr$

$$\frac{k\omega^2}{m}dt = \frac{dr}{r} \ (1 \ \%)$$

$$\stackrel{\text{to}}{=} \frac{\text{Joc}}{dt} \frac{d(\omega r)}{dt} / \omega^2 r = \frac{q^3 B^1}{6\pi \varepsilon_0 C^3 m_e^2} = 1.1 \times 10^{-14} \ (1 \ \text{fb})$$

故忽视切向加速度的近似有效(1分)

(3)代入 r=0.5r₀

t=3.59x10⁴s(3分)