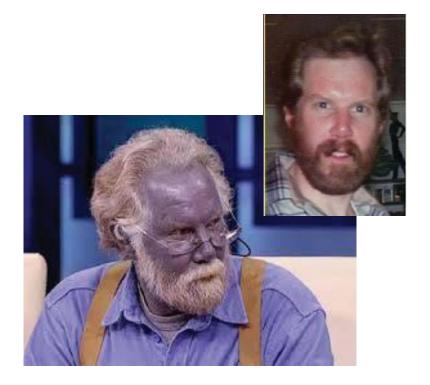
Blue skin man



Avatar



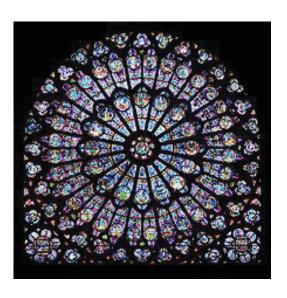
In 2007 press reports described Paul Karason, an American man whose entire skin gradually turned blue after consuming what he believed was colloidal silver made by himself with distilled water, salt and silver, and using a silver salve on his face in an attempt to treat problems with his sinus, dermatitis, acid reflux and other issues.

Think before class:

- What is the color of small (nm-size) gold?
- Why does it have these colors?





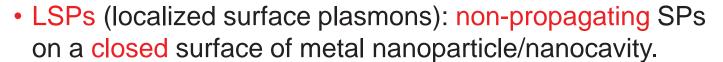


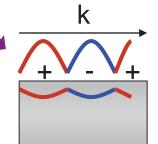
Content of this lecture

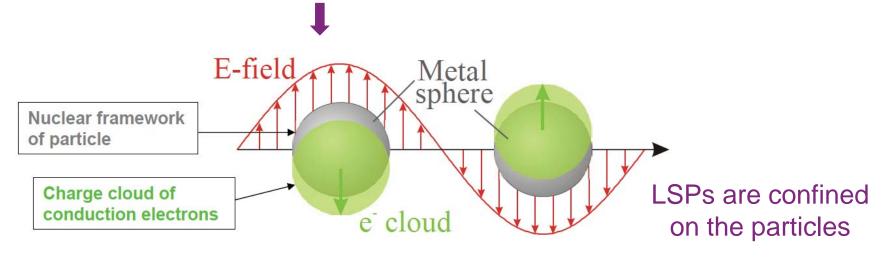
- 1. Localized surface plasmons (LSPs) on metal nanoparticles
 - Difference of LSPs with SPPs
 - Color effect of metal nanoparticles
 - Various metal nanoparticles
- 2. Resonance condition of LSPs
 - Review of dipole radiation
 - LSPs of nanoshperes (quasi-static approximation)
 - Size- and shape-dependence of LSPR (Mie theory)
 - LSPs of nanorods (Gans theory)
- 3. Coupling of LSPs between nanoparticles
- 4. LSPs of complex nanostructures nanoshell
- 5. Comparison of volume plasmons, SPPs, and LSPs

1. LSPs of metal nanoparticles

 SPPs: propagating SPs (photons coupled to SPs) on an extended metal-dielectric interface.





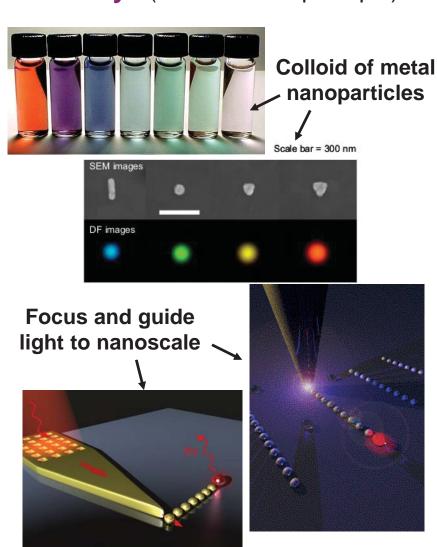


 The excitation of LSPs influences the extinction (i.e., absorption + scattering) of light from the nanoparticles → color effect

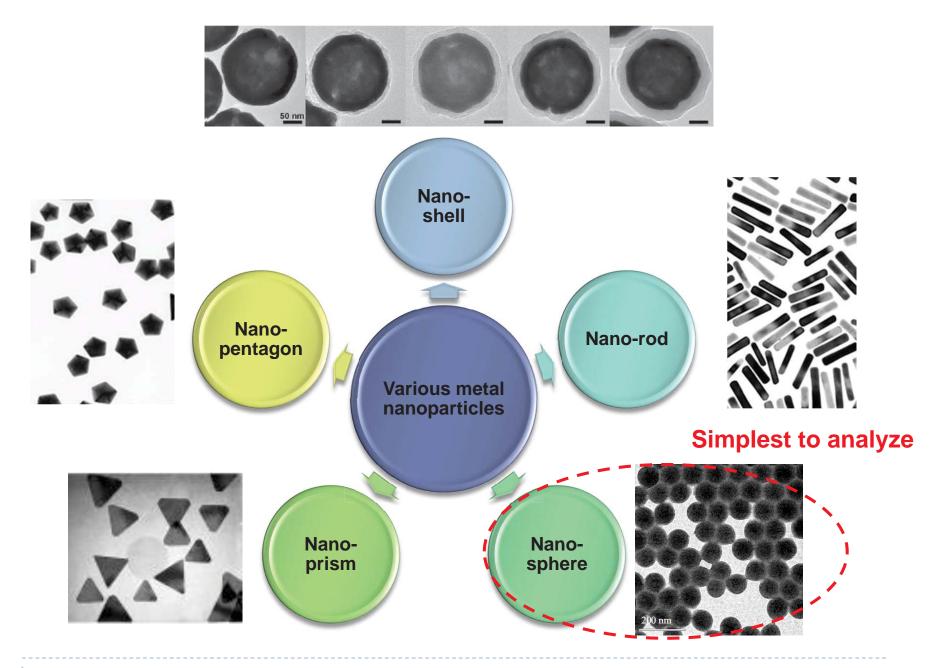
Color effect of metal nanoparticles

In the past: **Different shapes & sizes Lycurgus Cup** → distinct colors Stained glass in Au colloids in water church window

Nowadays (but the same principle):



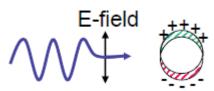
(M. Faraday ~1856)



2. Resonance condition of LSPs

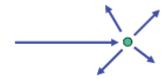
• When metal nanoshpere is small enough ($D << \lambda$), it can be regarded as an effective electric dipole.

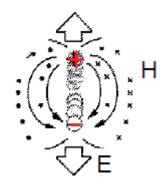
Harmonic E field drives harmonic oscillation of free electrons



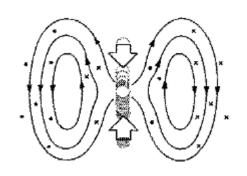


Oscillating dipole radiates (scattering of light)

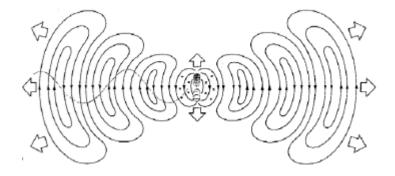




Start oscillation



Generation of EM wave



After several periods

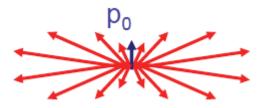
(radiation mainly ⊥oscillation direction)

Review of dipole radiation

• Radiated intensity: $I = \frac{p^2 \omega^4}{32\pi^2 \varepsilon_0 c^3 r^2} \sin^2 \theta$ (angle dependent)

Derivation: see any textbook on electromagnetism or electrodynamics

• Radiation pattern:



- Consider Lorentz model of the dipole moment: $\mathbf{p} = \frac{e^2}{m} \frac{1}{\omega_0^2 \omega^2 i\gamma\omega} \mathbf{E}_L$
- We get: $I_{S} = \frac{e^{4}\omega^{4}}{32\pi^{2}m^{2}\varepsilon_{0}c^{3}r^{2}} \left(\frac{1}{\omega_{0}^{2} \omega^{2} i\gamma\omega}\right)^{2} E_{L}^{2}\sin^{2}\theta$

See Lecture 2

Conclusions:

Rayleigh scattering

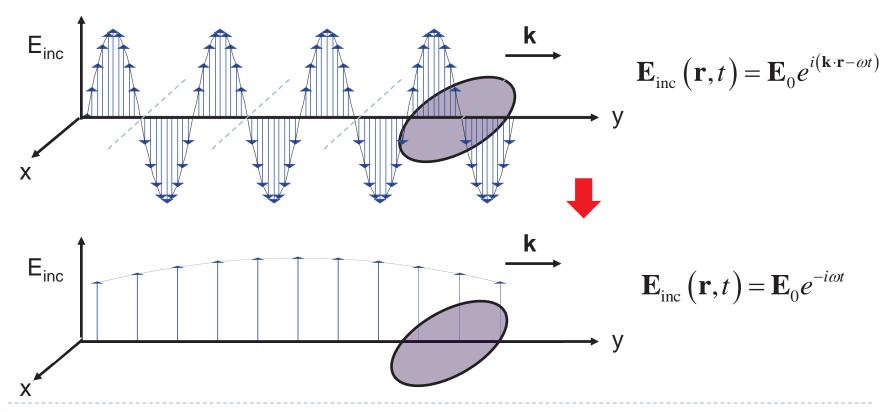
- Stronger scattering occurs at higher ω (shorter λ)
- Scattering occurs in both forward and backward directions

Incoming intensity

LSPs of nanospheres

- If particle small enough $(2a << \lambda)$
 - → treated as an electric dipole

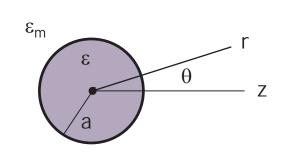
- $\epsilon_{\rm m}$ $\epsilon_{\rm m}$ $\epsilon_{\rm m}$
- quasi-static approximation: constant phase over the particle volume



$$\mathbf{E}_{\rm inc} = E_0 \hat{\mathbf{z}}$$

scalar potentials $\mathbf{E}_{1,2} = -\nabla \Phi_{1,2}$

• The electric fields inside (**E**₁) and outside (\mathbf{E}_2) the sphere may be found from the



The scalar potentials satisfy Laplace's equation:

$$\nabla^2 \Phi_1 = 0 \ (r < a), \ \nabla^2 \Phi_2 = 0 \ (r > a)$$

Boundary conditions:

$$\Phi_1 = \Phi_2 \ (r = a), \quad \varepsilon \frac{\partial \Phi_1}{\partial r} = \varepsilon_m \frac{\partial \Phi_2}{\partial r} \ (r = a), \quad \lim_{r \to \infty} \Phi_2 = -E_0 z$$

The following solutions satisfy the above equations:

$$\Phi_{1} = -\left(\frac{3\varepsilon_{m}}{\varepsilon + 2\varepsilon_{m}}\right) E_{0} r \cos \theta, \quad \Phi_{2} = -E_{0} r \cos \theta + a^{3} \left(\frac{\varepsilon - \varepsilon_{m}}{\varepsilon + 2\varepsilon_{m}}\right) E_{0} \frac{\cos \theta}{r^{2}}$$
applied field

Effective dipole at the sphere center

Introduce the dipole moment **p** as:

$$\Phi_{\text{out}} = -E_0 r \cos \theta + \frac{\mathbf{p} \cdot \mathbf{r}}{4\pi \varepsilon_0 \varepsilon_m r^3} \qquad \mathbf{p} = 4\pi \varepsilon_0 \varepsilon_m a^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m} \mathbf{E}_0$$

• Therefore, the polarizability (defined via $\mathbf{p} = \varepsilon_0 \varepsilon_m \alpha \mathbf{E}_0$):

$$\alpha = 4\pi a^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m}$$

$$\varepsilon = \varepsilon_1(\omega) + i\varepsilon_2(\omega)$$

$$\varepsilon_m, \text{ real permittivity of the embedding dielectric}$$

Resonant enhancement under the condition:

$$\left| \varepsilon(\omega) + 2\varepsilon_m \right| = \text{minimum} \leftarrow \text{called "Fröhlich condition"}$$

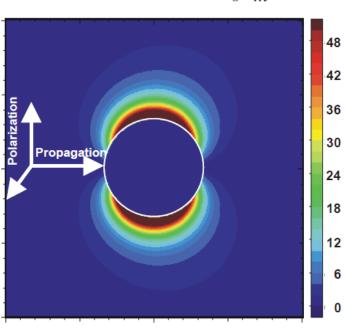
For Drude metal located in lossless dielectric: $\omega_{lsp} = \frac{\omega_p}{\sqrt{1+2\varepsilon_m}}$ can be used for sensing

$$\omega_{lsp} = \frac{\omega_p}{\sqrt{1 + 2\varepsilon_m}}$$

• The electric field can be obtained by $\mathbf{E} = -\nabla \Phi$:

$$\mathbf{E_{in}} = \frac{3\varepsilon_m}{\varepsilon + 2\varepsilon_m} \mathbf{E_0} \qquad \mathbf{E_{out}} = \mathbf{E_0} + \frac{3\mathbf{n} (\mathbf{n} \cdot \mathbf{p}) - \mathbf{p}}{4\pi \,\varepsilon_0 \varepsilon_m} \frac{1}{r^3}$$

Calculated **E** Field:
20 nm Ag nanosphere
in vacuum



- At resonance, the extinction (scattering + absorption) is large
- Near field enhancement → many prominent applications such as sensing, SERS, enhanced nonlinearity, data storage, ...

For full derivation, see J. D. Jackson, *Classical Electrodynamics* (Wiley, NY, 1999) and S. A. Maier, *Plasmonics: Fundamentals and Applications* (Springer, NY, 2007).

Application example 1: 80 • Sensing with 35 nm 60 Intensity Ag nanospheres $\Delta \lambda_{\text{max}} = 40.7 \text{ nm}$ a monolayer of smallmolecule adsorbates on the Ag nanoparticle 450 650 700 600 Wavelength (nm) 100-640 В -600.8 588.0-510.2 574.2 611.9 80 . λ_{max} = 203.1*RI + 306.5 600 -Resonance shift with 60 λ_{max} (nm) Intensity respect to refractive 560 index (RI) change of 40 surrounding medium 520 20 480 500 650 450 550 600 700 1.2

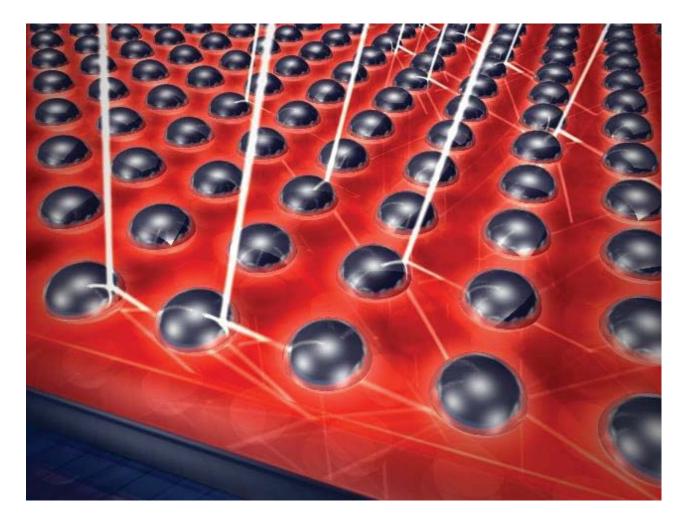
McFarland and Duyne, Nano Lett. 3, 1057 (2003).

Refractive Index

Wavelength (nm)

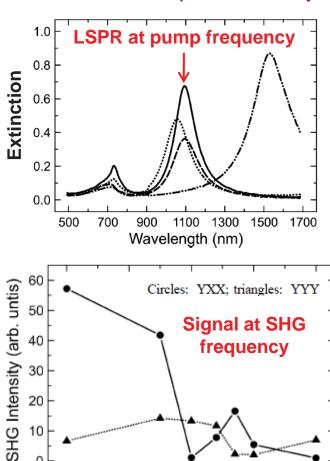
Application example 2:

Light trapping in solar cells by metallic nanoparticles



Application example 3:

Enhanced second-harmonic generation in metal nanoparticle array

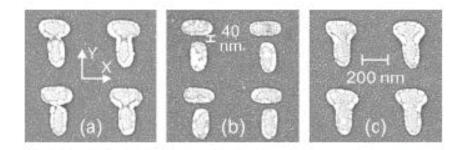


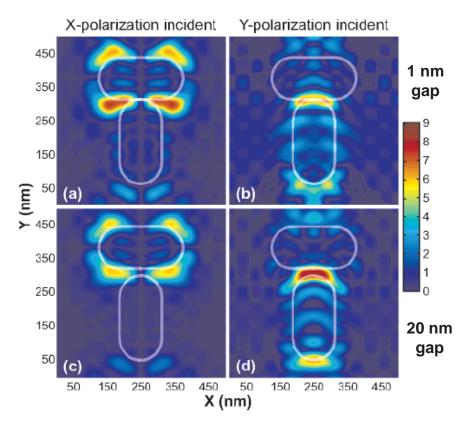
20

Nanogap Size (nm)

30

40

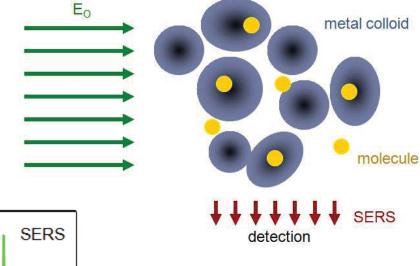


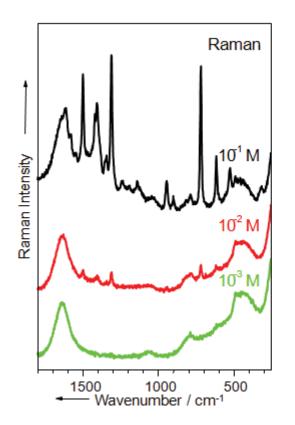


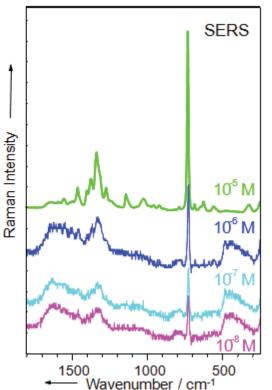
Canfield et al., Nano Lett. 7, 1251 (2007)

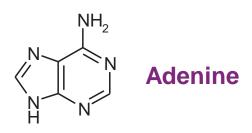
Application example 4:

Surface enhanced Raman spectroscopy (SERS)









SERS improves the detection limit

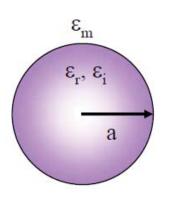
Size- and shape-dependence of LSPR

- Quasi-static approximation is valid only for nanoparticles < 100 nm at visible and near-infrared frequencies; size-dependence is not captured.
- For rigorous analysis of nanosphere Mie theory



Mie theory (1908): size dependence shape dependence

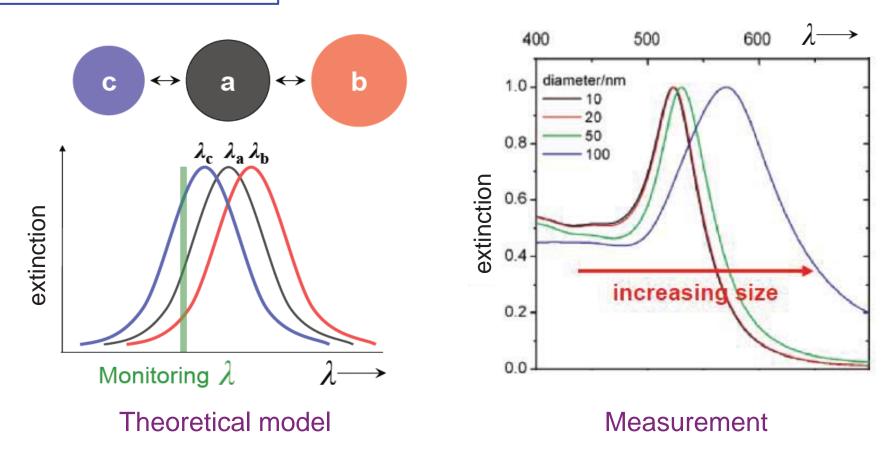
$$E(\lambda) = \frac{24\pi^2 N a^3 \varepsilon_m^{3/2}}{\lambda \ln(10)} \left[\frac{\varepsilon_i}{\left(\varepsilon_r + \chi \varepsilon_m\right)^2 + \varepsilon_i^2} \right]$$



 $E(\lambda)$ = Extinction spectrum = absorption + scattering χ = shape factor (2 for sphere, > 2 for spheroid) $\varepsilon_{\rm m}$ = external dielectric constant $\varepsilon_{\rm r}$ = real metal dielectric constant $\varepsilon_{\rm r}$ = imaginary metal dielectric constant

Mie, Ann. Phys. 25, 377 (1908).

Size dependence



Understand the size dependence qualitatively:

nanosphere size $\uparrow \Rightarrow$ charge distance $\uparrow \Rightarrow$ restoring force $\downarrow \Rightarrow$ resonance frequence $\omega \downarrow \Rightarrow$ resonance wavelength \uparrow

Shape dependence: LSPs of nanorods

Response of nanospheroid/nanorod – Gans theory (extension of Mie theory)

$$\sigma_{ext}(\lambda) = \frac{2\pi V \varepsilon_{med}^{3/2}}{3\lambda} \sum_{j} \frac{\left(1/P_{j}^{2}\right) \varepsilon''}{\left(\varepsilon' + \frac{1 - P_{j}}{P_{j}} \varepsilon_{med}\right)^{2} + \left(\varepsilon''\right)^{2}}$$

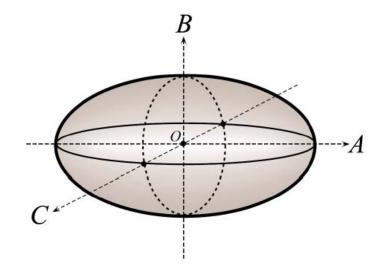
$$(A > B = C)$$

$$P_{A} = \frac{1 - e^{2}}{e^{2}} \left[\frac{1}{2e} \ln \left(\frac{1 + e}{1 - e} \right) - 1 \right]$$

$$P_B = P_C = \frac{1 - P_A}{2}$$

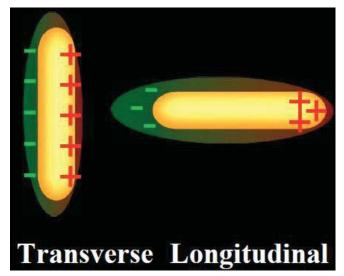
$$e = \sqrt{1 - \left(\frac{B}{A}\right)^2}$$

aspect ratio
$$R = \frac{A}{B}$$

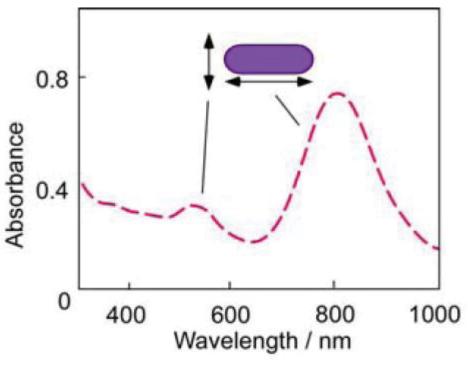


Link et al., J. Phys. Chem. B **103**, 3073 (1999).

- There are two maxima in the extinction spectra, corresponding to two resonance modes:
 - a longitudinal mode (dipole oscillation along the long axis)
 - a transverse mode (dipole oscillation along the short axis)

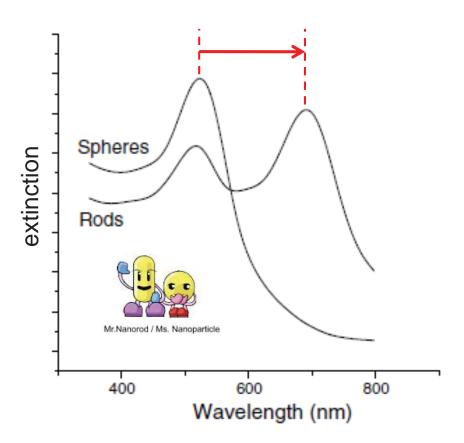


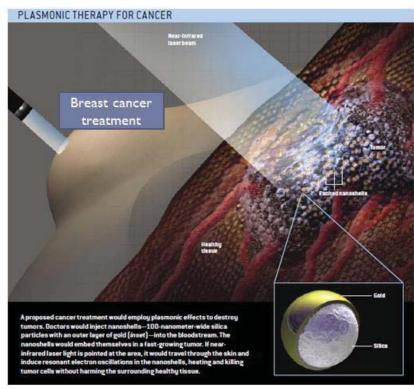




Gold nanorods show two absorption peaks; Visible region: 520-530 nm, Transverse Mode Near-infrared region: 700-1500 nm, Longitudinal Mode

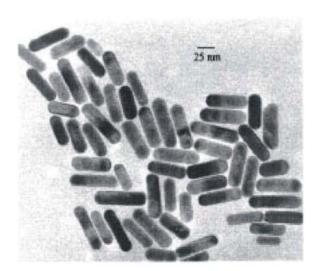
- The longitudinal mode has a significant red-shift compared to the resonance of a nanoshpere of the same volume
- Importance: aspect ratio ↑ ⇒ red-shift of longitudinal mode to near-IR
 ⇒ beneficial for biomedical applications (e.g., treatment of tumors)

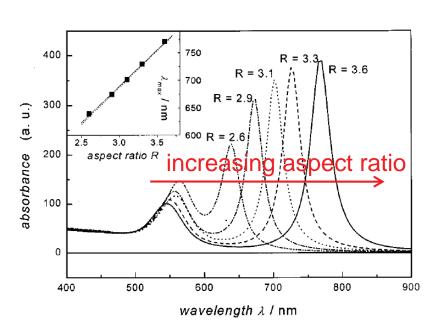




El-Brolossy et al., Eur. Phys. J. Special Topics **153**, 361–364 (2008).

Gold nanorod experiment





1.2 Longitudinal mode

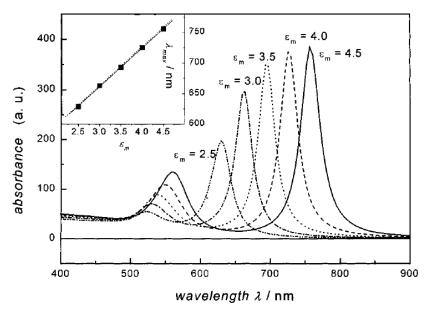
Transverse mode

0.4

0.2

400 500 600 700 800 900 1000

Wavelength (nm)



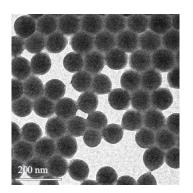
Dependence on aspect ratio

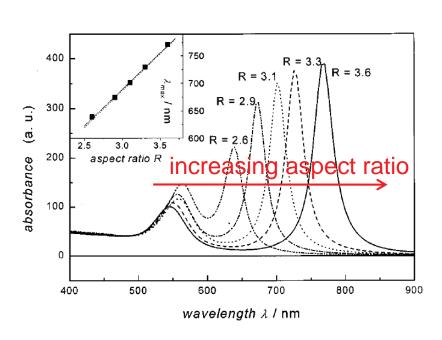
Dependence on surrounding medium

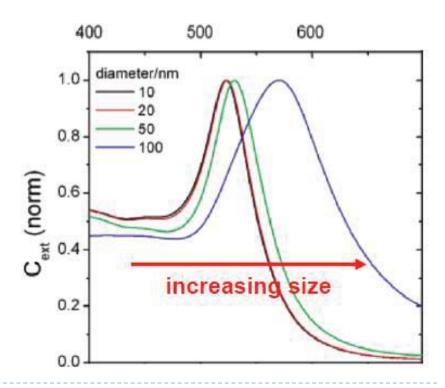
Nanorod vs. nanosphere: better tunability of resonance wavelength

nanorod

nanoshpere

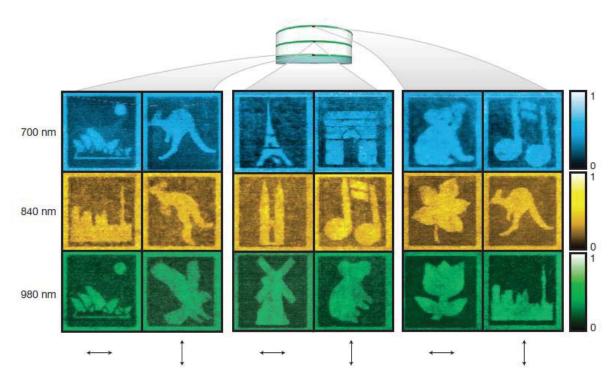




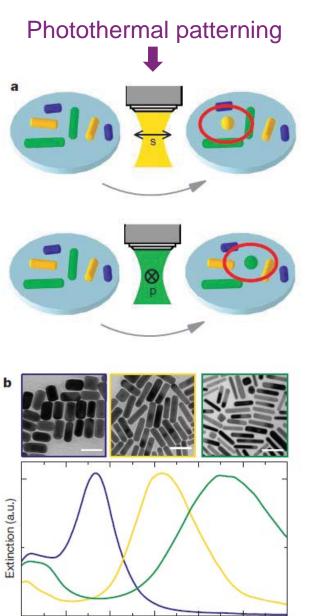


Application example:

Five-dimensional optical data storage with gold nanorods



Zijlstra et al., Nature 459, 410 (2009).



600

700

800

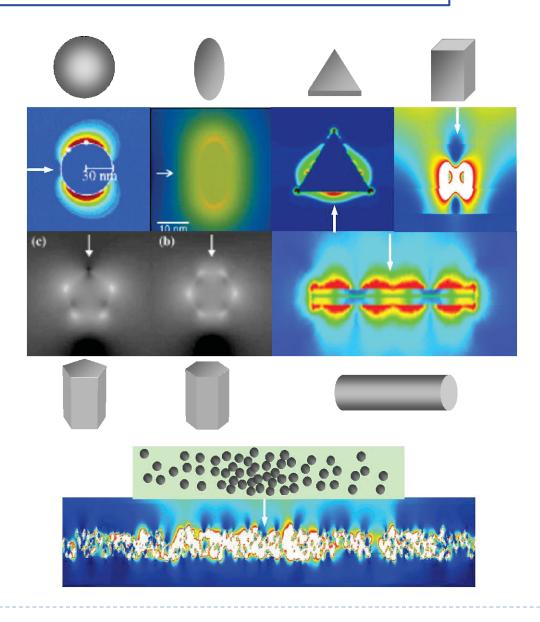
Wavelength (nm)

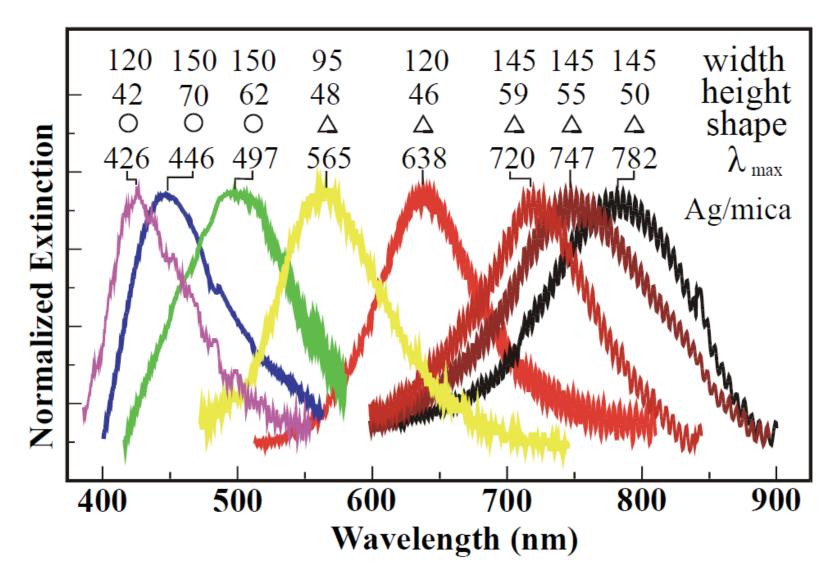
900

1,000

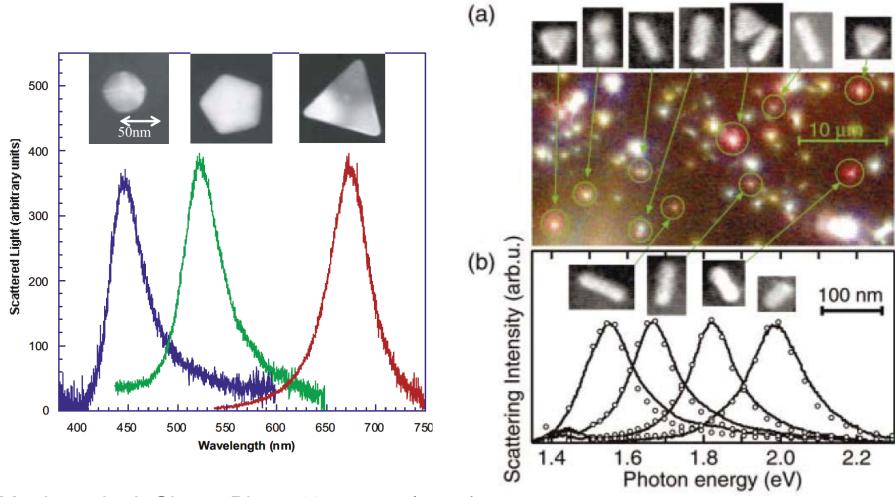
1,100

More about size and shape dependence





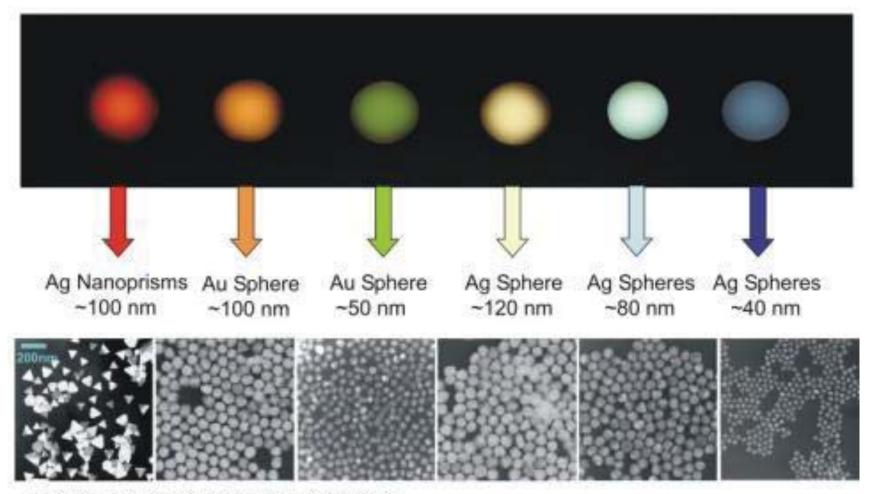
Jensen et al., J. Phys. Chem. B 104, 10549 (2000).



Mock et al., J. Chem. Phys. 116, 6755 (2002).

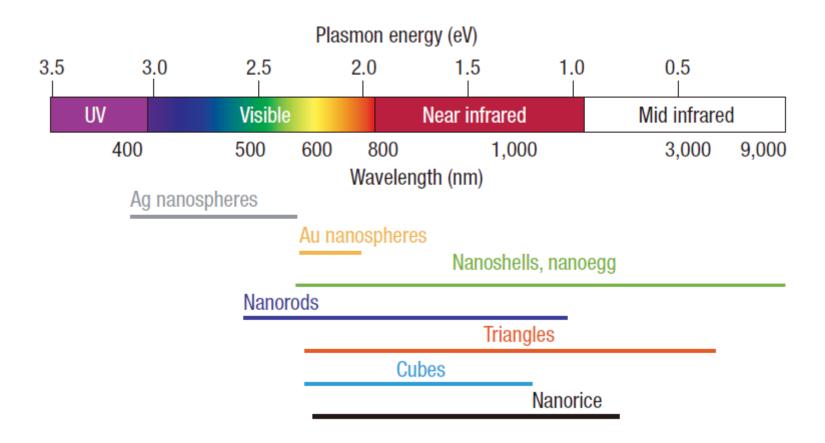
Kuwata et al., APL 83, 4625 (2003).

Rayleigh Light-Scattering of Nanocrystals: Shape, Size, and Composition Matter



^{*} The scale bar is the same for all the images.

A range of LSPR for a variety of particle shapes:



Lal et al., Nature Photon. 1, 641 (2007).

3. Coupling of LSPs between nanoparticles

For single nanoparticle:



An isolated sphere is symmetric, so the polarization direction does not matter.

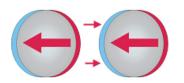
For closely spaced nanoparticles – near-field intercoupling:



TRANSVERSE:

restoring force *increased* by coupling to neighbor

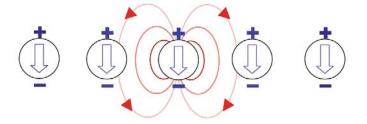
→ Resonance shifts to higher frequency



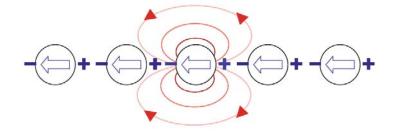
LONGITUDINAL:

restoring force *reduced* by coupling to neighbor → Resonance shifts to lower frequency

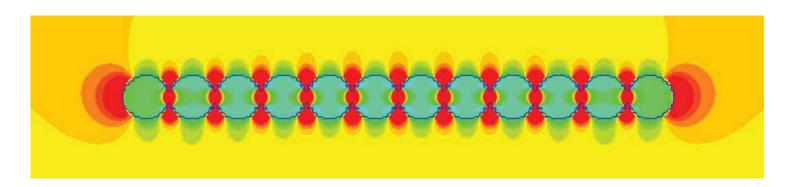
Increased restoring force Higher (blue-shifted) frequency



Reduced restoring force Lower (red-shifted) frequency



Another consequence: near-field enhancement in gaps

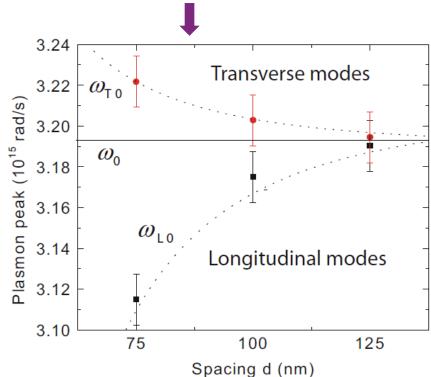


Experimental verification:

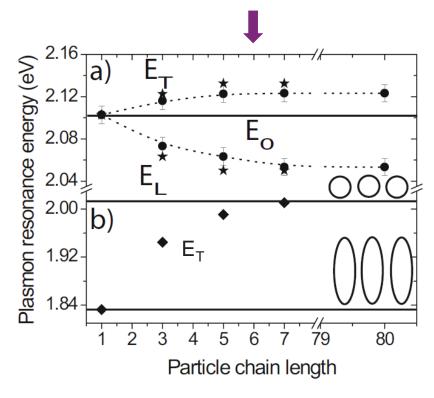
50 nm 800022 10.0kV ×15.0k' 2.00' m

Dependence on particle chain length



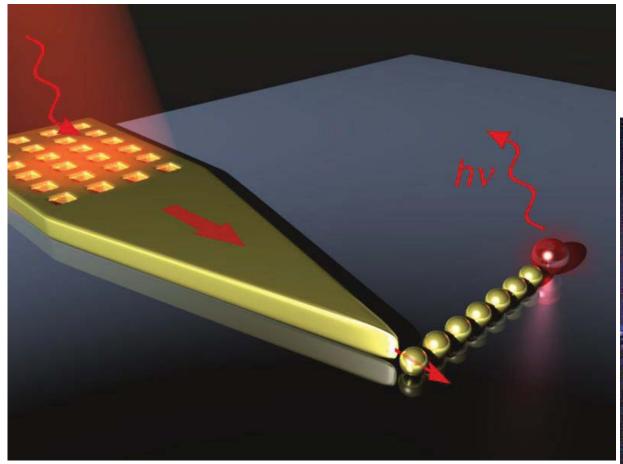


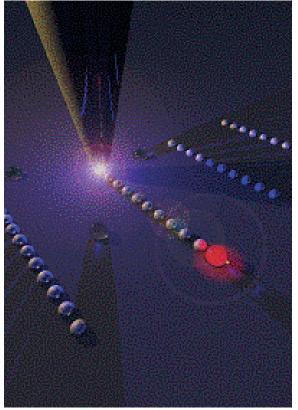
Maier et al., PRB **65**, 193408 (2002).



Maier et al., APL 81, 1714 (2002).

Application example: nanoparticle chain used as SPP waveguide

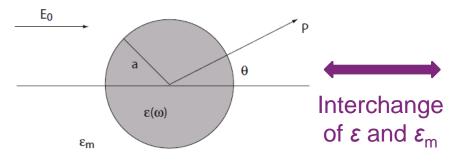




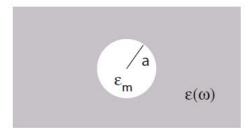
4. LSPs of complex nanostructures – nanoshell

Let's first consider nanosphere vs. nanocavity:

nanosphere



nanocavity



polarizability

$$\alpha = 4\pi a^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m}$$

 $\text{Re}\left[\varepsilon\left(\omega\right)\right] = -2\varepsilon_{m}$

Fröhlich condition

$$\omega_{\mathrm{lsp}}$$
 for Drude metal

$$\omega_{lsp} = \frac{\omega_p}{\sqrt{1 + 2\varepsilon_m}} \stackrel{\text{in air }}{=} \frac{\omega_p}{\sqrt{3}}$$

$$\alpha = 4\pi a^3 \frac{\varepsilon_m - \varepsilon}{\varepsilon_m + 2\varepsilon}$$

$$\mathbf{Re}\left[\varepsilon\left(\omega\right)\right] = -\frac{1}{2}\varepsilon_{m}$$

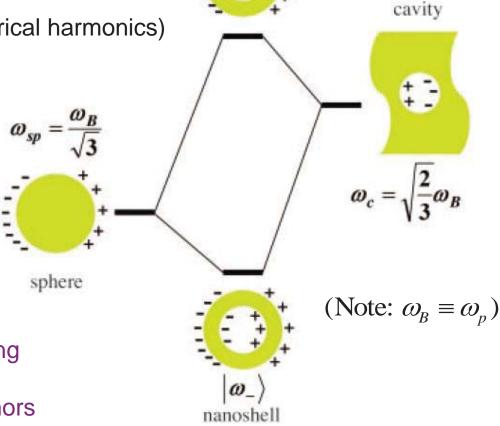
$$\omega_{lsp} = \frac{\omega_p}{\sqrt{1 + \frac{1}{2}\,\mathcal{E}_m}} \stackrel{\text{in air}}{=} \sqrt{\frac{2}{3}} \omega_p$$

Nanoshell – plasmon hybridization

$$\omega_{l,\pm}^2 = \frac{\omega_p^2}{2} \left[1 \pm \frac{1}{2l+1} \sqrt{1 + 4l (l+1) \left(\frac{a^{2l+1}}{b}\right)} \right]$$

(l: orders of spherical harmonics)

- Nanoshell modes = shpere mode + cavity mode
 - \rightarrow an antibonding ω_{+} mode
 - \rightarrow a bonding ω_{-} mode
- Consequences:
 - → resonance shifts to near-IR
 - → reduced resonance linewidth
- Superior to nanospheres for sensing and biomedical applications (e.g., treatment of nanoparticle-filled tumors via absorption-induced heating)



nanoshell

Prodan et al., Science **302**, 419 (2003).

5. Comparison of volume plasmons, SPPs, and LSPs

| | Volume plasmons | SPPs | LSPs (nanosphere) |
|--------------------------|--------------------------------------------------|-----------------------------------------------------------|-------------------------------------------------------------|
| Schematic | + + + | k + | |
| Mode property | Propagating mode in bulk metal | Propagating mode on metal surface | Non-propagating confined mode |
| Wave property | Longitudinal | Transverse & Longitudinal | - |
| Characteristic frequency | $\omega_p = \sqrt{\frac{Ne^2}{\varepsilon_0 m}}$ | $\omega_{sp} = \frac{\omega_p}{\sqrt{1 + \varepsilon_d}}$ | $\omega_{lsp} = \frac{\omega_p}{\sqrt{1 + 2\varepsilon_d}}$ |
| Interaction with light | No interaction (non-EM wave) | Coupled with photon to form polariton | Resonant extinction (scattering +absorption) |

Summary

- Localized surface plasmons (LSPs):
 LSPs: non-propagating SPs confined on nanoparticle/nanocavity,
 color effect, various metal nanoparticles
- Resonance condition of LSPs: metal nanoparticle treated as effective electric dipole, LSPs of nanospheres (quasi-static approximation), Fröhlich condition, sizeand shape-dependence (Mie theory), LSPs of nanorods (Gans theory), sensing and biomedical application of LSPs
- Coupling of LSPs between nanoparticles: transverse and longitudinal modes, near-field enhancement in gaps
- LSPs of complex nanostructures nanoshell
 LSPs of nanosphere & nanocavity, plasmon hybridization in nanoshell
- Comparison of volume plasmons, SPPs, and LSPs