

Review.

1. Chirping effect

* Linear chirping: GVD $-(1+aq)\frac{t^2}{t_0^2}$

$$E(t, 0) \rightarrow e$$

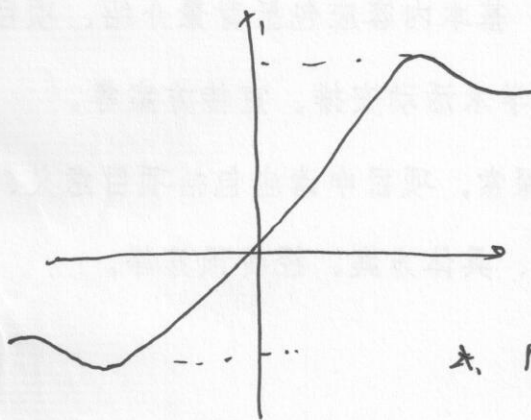
after dispersive propagation z

$$\text{new chirping: } b = a + \frac{2k''z}{t_0^2(1+q^2)}$$

* Nonlinear chirping:

$$\text{SPM} \sim \text{Kerr effect: } \Delta n = n_2 I$$

by using Gaussian pulse:



$$\rightarrow \frac{\Delta \omega}{\omega} \sim \frac{n_2 I_0 z}{c t_0}$$

* New frequency

* Complex spectral structure

Instant phase change $\phi(t)$ ~~is~~ ^{is} dependence on pulse shape.

Comments:

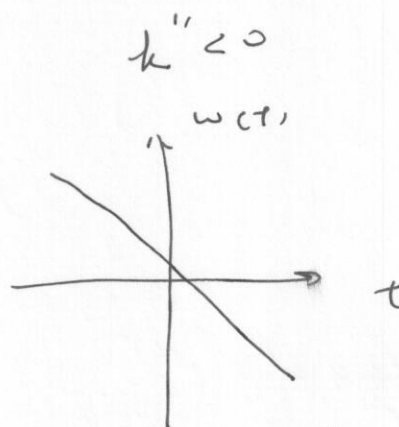
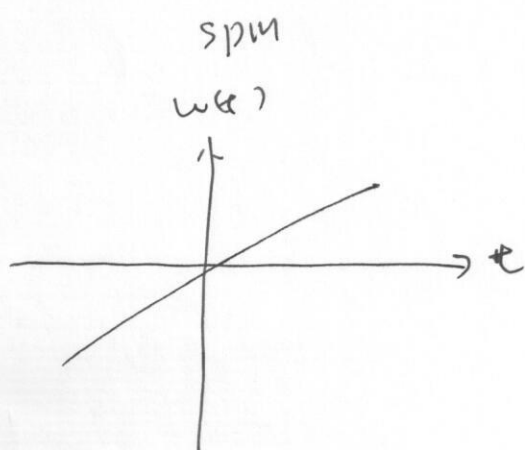
2.5 (2)

- ① $\text{Re}[\phi_{NL}(t)]$ does not change pulse profile
only changes spectrum
- ② Combined with GVD or higher order dispersion

* Nonlinear + Dispersion

Entangled together.

useful example: Soliton.



Make it possible \rightarrow balance between
SPM and Dispersion

2. Spatial-temporal problem.

Dispersion:

$$\frac{\partial A}{\partial z} + \frac{1}{2} j k'' \frac{\partial^2 A}{\partial t^2} = 0$$

$$\begin{cases} H(\omega) = e^{-j\pi \Lambda \omega^2 z} \\ h(t) = \sqrt{\frac{1}{j\Lambda z}} e^{j\frac{\pi}{\Lambda z} t^2} \end{cases}$$

$$\Lambda = -2\pi k''$$

Diffraction:

$$\frac{\partial A}{\partial z} - j \frac{1}{2k} \frac{\partial^2 A}{\partial x^2} = 0$$

$$-j\pi \Lambda \omega^2 z$$

$$\begin{cases} H(\omega) = e^{-j\frac{\pi}{\Lambda z} \omega^2} \\ h(x) = \sqrt{j\Lambda z} e^{j\frac{1}{2k} x^2} \end{cases}$$

$$\begin{aligned} \Lambda &\leftrightarrow \lambda \\ t &\leftrightarrow x \\ z &\leftrightarrow k_x \end{aligned}$$

$$-k'' \leftrightarrow \frac{1}{k}$$

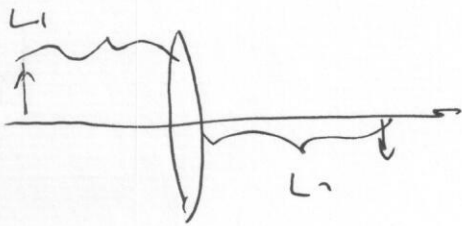
This Lecture

Lab 5 (3)

* finishing up phase modulation of linear system

* Fermi send optics

Solving dispersion problem by performing diffraction experiments: (time lens)



$$\frac{1}{L_1} + \frac{1}{L_2} = \frac{1}{f}$$

give lens phase function:

$$g_1, g_2 e^{-j \frac{k}{2f} (x^2 + y^2)} = t(x, y)$$

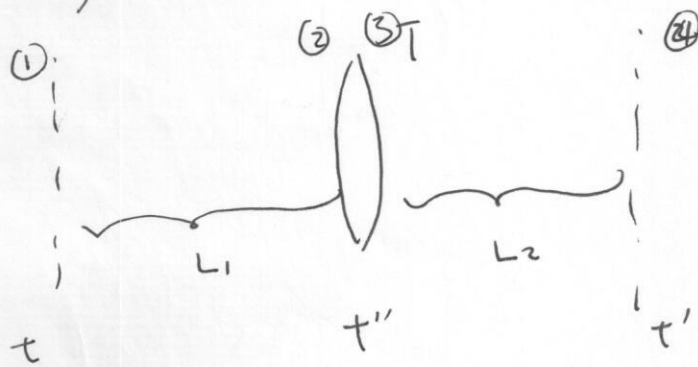


$$g_2 = g_1 \cdot t(x, y)$$

Assume we have certain system (Dispersion)

$$T = e^{-j \frac{\omega_0}{2f_t} t^2}$$

give initial pulse $f(t)$



$$\begin{aligned} (1) & f(t) \\ (2) & f(t) \otimes \sqrt{\frac{1}{j\lambda_1 L_1}} e^{j \frac{\pi}{\lambda_1 L_1} t^2} \\ (3) & \left[\downarrow \right] \otimes e^{j \frac{\pi}{\lambda_2 L_2} t^2} \\ (4) & \left[\downarrow \right] \otimes \sqrt{\frac{1}{\lambda_2 L_2}} \end{aligned}$$

$$(1) f(t)$$

$$j \frac{\pi}{\lambda_1 L_1} (t - t'')^2$$

\Rightarrow

$$\begin{aligned} (2) & \sqrt{\frac{1}{j\lambda_1 L_1}} \cdot \int dt \cdot f(t) \cdot e^{j \frac{\pi}{\lambda_1 L_1} t^2} \cdot e^{j \frac{\pi}{\lambda_1 L_1} t''^2} \cdot e^{-j \frac{\pi}{\lambda_1 L_1} 2tt''} \\ & = \frac{1}{\sqrt{j\lambda_1 L_1}} \cdot \int dt f(t) \cdot e \cdot e \cdot e \end{aligned}$$

L25 (4)

$$\begin{aligned}
 (3) \quad & \frac{1}{\sqrt{j\omega L_1}} \int dt f(t) \cdot e^{j\frac{\pi}{\omega L_1} t^2} \cdot e^{j\frac{\pi}{\omega L_1} t'^2} \cdot e^{-j\frac{\pi}{\omega L_1} 2tt'} \cdot e^{-j\frac{\omega_0}{2f_c} t'^2} \\
 & = \frac{1}{\sqrt{j\omega L_1}} \int dt f(t) \cdot e^{j\frac{\pi}{\omega L_1} t^2} \cdot e^{-j\frac{\pi}{\omega L_1} 2tt'} \cdot e^{j(\frac{\pi}{\omega L_1} - \frac{\omega_0}{2f_c}) t'^2}
 \end{aligned}$$

$$\begin{aligned}
 (4) \quad & \frac{1}{\sqrt{j\omega L_1}} \cdot \frac{1}{\sqrt{j\omega L_2}} \cdot \int dt f(t) e^{j\frac{\pi}{\omega L_1} t^2} \int dt'' e^{-j\frac{\pi}{\omega L_2} 2tt''} \cdot e^{j(\frac{\pi}{\omega L_1} - \frac{\omega_0}{2f_c}) t'^2} \\
 & = \frac{1}{\sqrt{-\omega L_1 \omega L_2}} \cdot \int dt f(t) e^{j\frac{\pi}{\omega L_1} t^2} \int dt'' e^{-j\frac{\pi}{\omega L_2} 2tt''} \cdot e^{j(\frac{\pi}{\omega L_1} - \frac{\omega_0}{2f_c}) t'^2} \\
 & = \frac{1}{\sqrt{C}} \cdot \int dt f(t) \cdot e^{j\frac{\pi}{\omega L_1} t^2} \cdot e^{j\frac{\pi}{\omega L_2} t'^2} \cdot e^{-j\frac{\pi}{\omega L_2} 2tt''} \cdot e^{j(\frac{\pi}{\omega L_1} + \frac{\pi}{\omega L_2} - \frac{\omega_0}{2f_c}) t'^2} \cdot e^{-j\frac{\pi}{\omega L_2} (\frac{t}{\omega L_1} + \frac{t'}{\omega L_2}) \cdot t''}
 \end{aligned}$$

if we can find a time lens.

$$\begin{aligned}
 \frac{\pi}{\omega L_1} + \frac{\pi}{\omega L_2} &= \frac{\omega_0}{2f_c} \quad j\frac{\pi}{\omega L_1} t^2 \quad j\frac{\pi}{\omega L_2} t'^2 \\
 (4) \Rightarrow & \frac{1}{\sqrt{-M}} \cdot \int dt \cdot f(t) \cdot e^{j\frac{\pi}{\omega L_1} t^2} \cdot e^{j\frac{\pi}{\omega L_2} t'^2} \cdot e^{-j\frac{\pi}{\omega L_2} (\frac{t}{\omega L_1} + \frac{t'}{\omega L_2}) \cdot t''} \\
 & \propto f(-\frac{t'}{M}) \quad \text{where } M = \frac{\omega L_2}{\omega L_1}
 \end{aligned}$$

Comments on $e^{-j \frac{\omega_0}{2ft} \cdot t^2}$

Dispersion ————— Diffraction

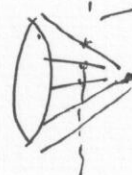
$\omega_0 \sim$ frequency

$k \sim$ spatial frequency (angular)

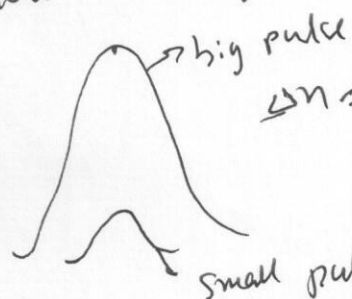
chirp: $\frac{f}{t}$
 \Rightarrow at different time position, ω is different

f (focal lens)

at different spatial position, k is different.



How to implementation of time lens



$\omega \sim \omega_0 I$

$$e^{j \frac{\omega}{c} n I_0 \frac{t^2}{t_0^2} z} \sim e$$

$$e^{j \frac{\omega}{c} n I_0 \frac{t^2}{t_0^2} z}$$

Temporal and Spatial evolution of a pulse
 $\vec{E}(x, y, 0, t) \xrightarrow{\text{F.T.}} \vec{E}(x, y, 0, \omega) \rightarrow \textcircled{x} \frac{e^{j\omega L}}{j\omega L} e^{j \frac{\omega}{\lambda L} (x^2 + y^2)}$

$z=0$

$$\Rightarrow \vec{E}(x, y, L, \omega)$$

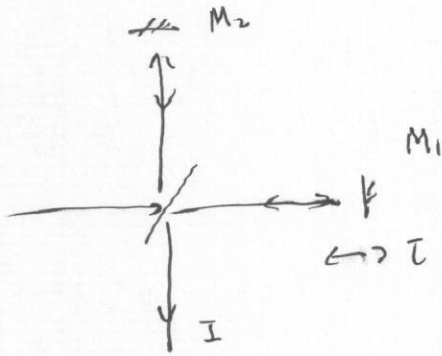
I.F.T

$$\rightarrow \vec{E}(x, y, L, t)$$

$z=L$

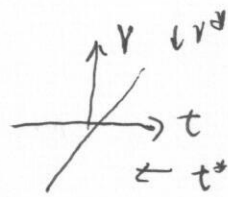
Notice phase term, here.

1. Short Pulse Interferometry



$$I(\tau) = |\bar{E}_1(t-\tau) + \bar{E}_2(t)|^2$$

about BS (Beamsplitter)



$r^*, t^* \Rightarrow$ time reversal

$$r^* t + t^* r = 0$$

base on causality
down port should be zero

$$\Rightarrow r^* t + (t \cdot r^*)^* = 0$$

$$\Rightarrow r^* t = i(\quad) \text{ pure imaginary}$$

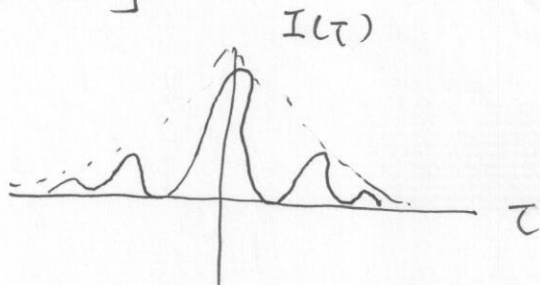
$$I(\tau) = \langle |\bar{E}_1|^2 \rangle + \langle |\bar{E}_2|^2 \rangle + \langle \bar{E}_1^*(t-\tau) \cdot \bar{E}_2(t) \rangle + \langle \bar{E}_1(t-\tau) \cdot \bar{E}_2^*(t) \rangle$$

Let $\bar{E}_1(t) = \xi_1(t) \cdot e^{-j\omega_0 t}$
 $\bar{E}_2(t) = \xi_2(t) \cdot e^{-j\omega_0 t}$

$$I(\tau) = \langle |\xi_1|^2 \rangle + \langle |\xi_2|^2 \rangle + \langle \xi_1^*(t-\tau) \xi_2(t) \rangle e^{-j\omega_0 \tau} + \langle \xi_1(t-\tau) \xi_2^*(t) \rangle e^{j\omega_0 \tau}$$

$$\propto 2 |A(\tau)| \cos[\omega_0 \tau - \phi(\tau)]$$

$$A(\tau) = |A| \cdot e^{j\phi}$$



perform F.T. to $I(t) \sim \tilde{I}(\omega)$

Lab 7

$$= \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} \tilde{I}_1(t-\tau) \cdot \tilde{I}_2(t) \cdot dt \right) \cdot e^{-j\omega\tau} \cdot d\tau$$

$$= \int_{-\infty}^{\infty} \left(\int_{-\infty}^{\infty} \tilde{I}_1(t-\tau) \cdot \tilde{I}_2(t) \cdot dt \right) \cdot e^{-j\omega\tau} \cdot d\tau$$

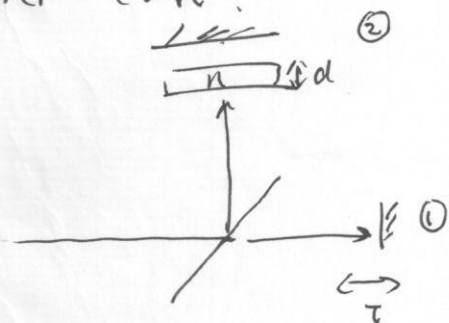
$$= \tilde{I}_1(\omega) \cdot \tilde{I}_2(\omega)$$

$$\Rightarrow \tilde{I}(\omega) = \tilde{I}_1(\omega) \cdot \tilde{I}_2(\omega)$$

if two pulses are identical: $\Rightarrow \tilde{I}(\omega) = |\tilde{I}(\omega)|^2$

Now we get the spectrum of the pulse

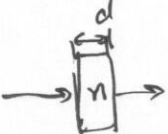
Another case:



$$\tilde{I}_2(\omega) = \tilde{I}_1(\omega) \cdot e^{j\phi(\omega)}$$

$$\tilde{I}_1(\omega) = \tilde{I}_2(\omega)$$

$$\phi(\omega) = \frac{\omega}{c} (n-1) \cdot d \cdot 2$$

why:  $\frac{\omega}{c} \cdot nd - \frac{\omega}{c} \cdot d$

$$\phi(\omega) = \phi(\omega_0) + \phi'(\omega_0) \cdot \Delta\omega + \frac{1}{2} \phi''(\omega_0) \cdot \Delta\omega^2 + \dots$$

$$\phi(\omega_0) = \frac{\omega_0}{c} \cdot (n-1) \cdot d$$

$$\phi'(\omega_0) = \frac{d}{d\omega} \left[\frac{\omega}{c} \cdot n \cdot d - \frac{\omega}{c} \cdot 2d \right] \Big|_{\omega=\omega_0}$$

$$= \frac{2d}{c/n_g} - \frac{2d}{c} = \Delta\tau$$

$$n_g = n + \omega_0 \cdot \frac{dn}{d\omega} \Big|_{\omega=\omega_0}$$

→ Stop here on 20/1.3.22

group index. $c/n_g \rightarrow$ group velocity

$$\phi''(\omega_0) = \frac{d^2}{d\omega^2} \left[\frac{\omega_0 2d n}{c} - \frac{\omega}{c} \cdot 2d \right]$$

$$= \frac{d^2}{d\omega^2} \left[\frac{\omega}{c} \cdot n \cdot 2d \right] = k'' \cdot 2d, \quad k'' = \frac{\omega}{c} \cdot n(\omega)$$

$$\tilde{E}_2(\omega) = \tilde{E}_1(\omega) \cdot e^{j\phi(\omega_0)} \cdot e^{j\Delta\tau \cdot (\omega - \omega_0)} \cdot e^{-\frac{1}{2} j k'' 2d (\omega - \omega_0)^2}$$

First let ignore k'' (GVD, from material)

take I.F.T

$$E_2(t) = E_1(t - \Delta\tau) \cdot e^{j\phi(\omega_0)} \cdot e^{j\Delta\tau \cdot \omega_0}$$

envelope

define. $\Delta\phi = \phi(\omega_0) - \omega_0 \cdot \Delta\tau = \frac{\omega_0}{c} (n - n_g) \cdot 2d$

$$= E_1(t - \Delta\tau) \cdot e^{j\Delta\phi}$$

Envelope is delayed by $\Delta\tau = \phi'(\omega_0)$