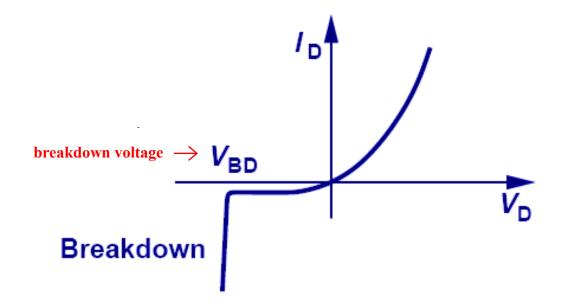


#### Reverse Breakdown

• As the reverse bias voltage increases, the electric field in the depletion region increases. Eventually, it can become large enough to cause the junction to break down so that a large reverse current flows:



## Reverse Breakdown Mechanisms

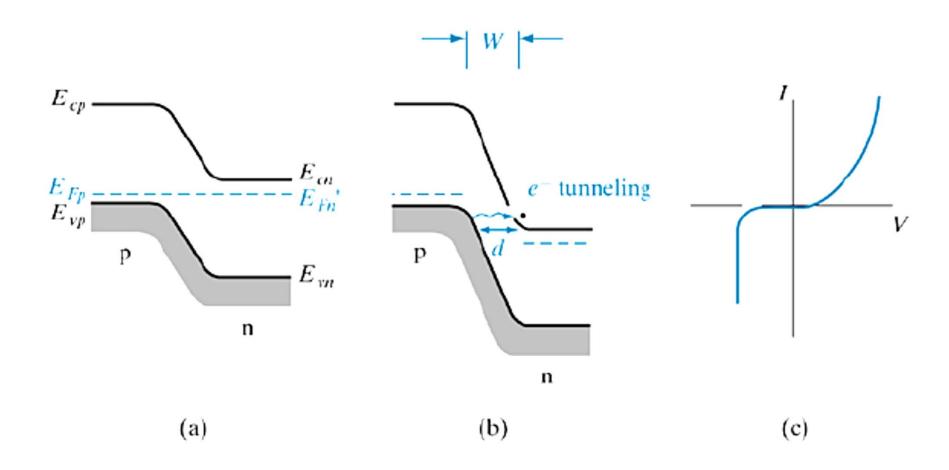
#### Zener breakdown

occurs when the electric field is sufficiently high to pull an electron out of a covalent bond (to generate an electron-hole pair).

#### Avalanche breakdown

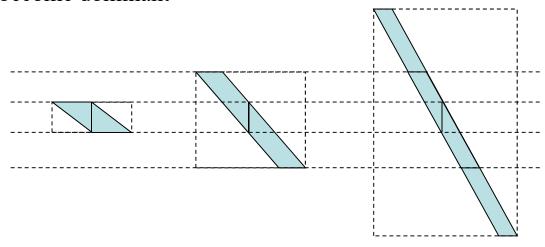
occurs when electrons and holes gain sufficient kinetic energy (due to acceleration by the E-field) in-between scattering events to cause electron-hole pair generation upon colliding with the lattice.

#### Zener breakdown

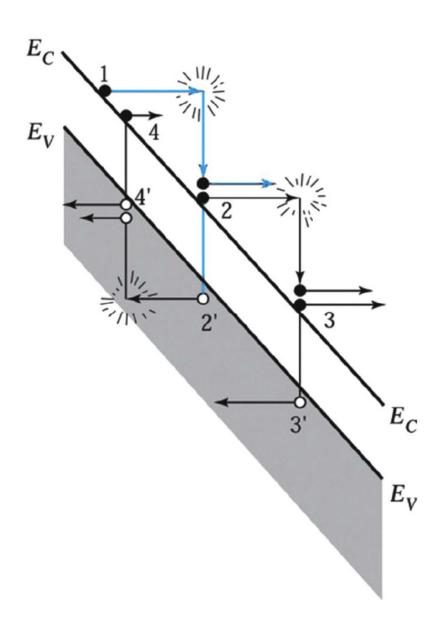


#### For Zener breakdown (tunneling)

- Extends only a very short distance W from each side of the junction
  - •The junction be sharp
  - The doping high
- The tunneling distance d may be too large for appreciable tunneling. However,
  - d becomes smaller as the reverse bias is increased, because the higher electric fields result in steeper slopes for the band edges
    - This assumes that the transition region width W does not increase appreciably with reverse bias
      - For low voltages and heavy doping on each side of the junction, this is a good assumption
- If Zener breakdown does not occur with reverse bias of a few volts, avalanche breakdown will become dominant



## Energy band diagram for the avalanche process

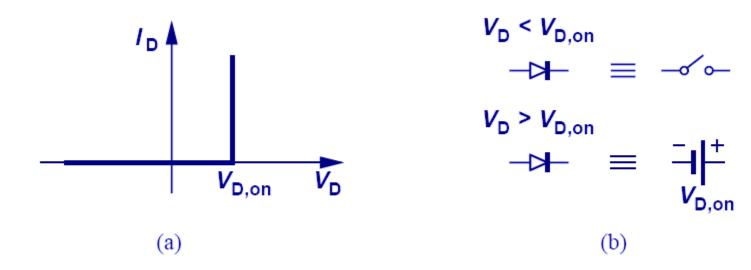


## For Avalanche breakdown

• Impact ionization rather than field ionization (Zener)

• Carrier multiplication

## Constant-Voltage Diode Model

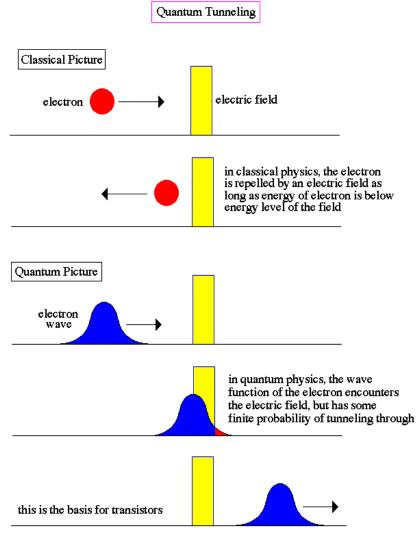


- If  $V_D < V_{D,on}$ : The diode operates as an open circuit.
- If  $V_D \ge V_{D,on}$ : The diode operates as a constant voltage source with value  $V_{D,on}$ .

# Tunneling

## The Tunneling of a Particle

- Tunneling refers to the ability of a particle to overcome and cross a potential energy barrier that it would not be able do based on classical understanding.
- It is only the particle's wave-nature that allows for this phenomenon.
- The probability wave describing the particle's position is an integral that overlaps into the energy barrier, allowing for some finite probability that the particle might actually "tunnel" through.



Images:http://4.bp.blogspot.com/,
abyss.uoregon.edu/.../ quantum tunneling.gif

#### <u>A Simple</u> Potential Step

 $\psi_A = Ae^{-jk_1x}$  $\psi_B = Be^{-jk_1x}$ 

CASE I :  $E_o > V$ 

$$E = 0$$
 Region 1 
$$x = 0$$
 Region 2 
$$x = 0$$

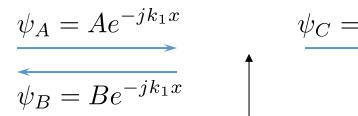
$$E_o\psi = -\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2}$$

$$\implies k_1^2 = \frac{2mE_o}{\hbar^2}$$

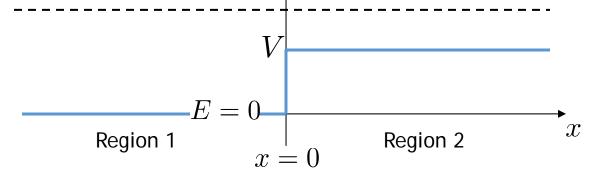
$$(E_o - V) \psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2}$$

$$(E_o - V)\psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} \qquad \Longrightarrow \quad k_2^2 = \frac{2m(E_o - V)}{\hbar^2}$$

## A Simple Potential Step



## CASE I : $E_o > V$



$$\psi_1 = Ae^{-jk_1x} + Be^{jk_1x}$$

$$\psi_2 = Ce^{-jk_2x}$$

 $\psi$  is continuous:

$$\psi_1(0) = \psi_2(0)$$

$$A + B = C$$

$$\frac{\partial \psi}{\partial x}$$
 is continuous:

$$\frac{\partial \psi}{\partial x} \text{ is continuous: } \qquad \frac{\partial}{\partial x} \psi(0) = \frac{\partial}{\partial x} \psi_2(0) \qquad \Longrightarrow \qquad A - B = \frac{k_2}{k_1} C$$

$$A - B = \frac{k_2}{k_1}C$$

## A Simple Potential Step

 $\psi_A = Ae^{-jk_1x}$   $\psi_B = Be^{-jk_1x}$ 

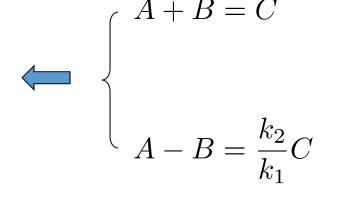
CASE I :  $E_o > V$ 

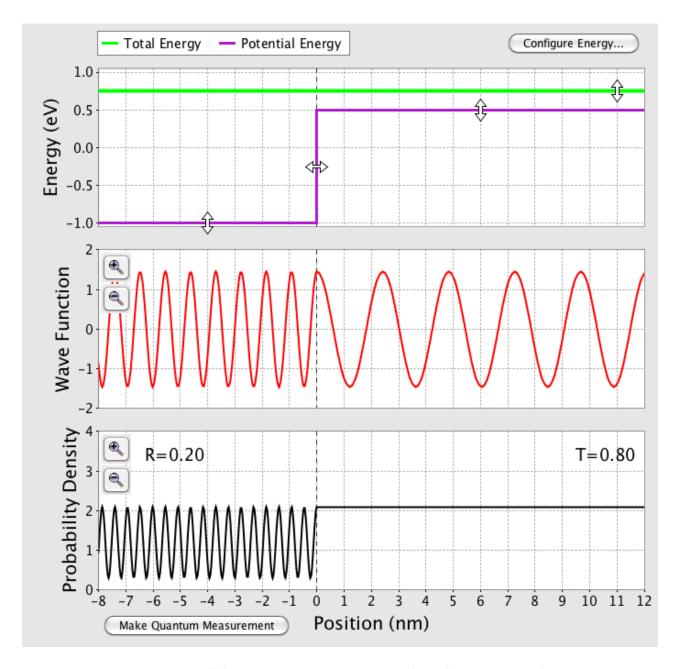
$$E = 0$$
 Region 1 
$$x = 0$$
 Region 2 
$$x = 0$$

$$\frac{B}{A} = \frac{1 - k_2/k_1}{1 + k_2/k_1}$$
$$= \frac{k_1 - k_2}{k_1 + k_2}$$

$$\frac{B}{A} = \frac{1 - k_2/k_1}{1 + k_2/k_1} \qquad \frac{C}{A} = \frac{2}{1 + k_2/k_1}$$

$$= \frac{k_1 - k_2}{k_1 + k_2} \qquad = \frac{2k_1}{k_1 + k_2}$$





Example from: <a href="http://phet.colorado.edu/en/get-phet/one-at-a-time">http://phet.colorado.edu/en/get-phet/one-at-a-time</a>

#### **Quantum Electron Currents**

Given an electron of mass m

that is located in space with charge density  $~ 
ho = q \left| \psi(x) \right|^2$  and moving with momentum ~ corresponding to  $~ < v > = \hbar k/m$ 

... then the current density for a single electron is given by

$$J = \rho v = q \left| \psi \right|^2 \left( \hbar k / m \right)$$

## <u>A Simple</u> <u>Potential Step</u>

 $\psi_A = Ae^{-jk_1x}$   $\psi_B = Be^{-jk_1x}$ 

 $\psi_C = Ce^{-jk_2 x}$ 

CASE I :  $E_0 > V$ 

$$E = 0$$
 Region 1 
$$x = 0$$
 Region 2 
$$x = 0$$

Reflection = 
$$R = \frac{J_{reflected}}{J_{incident}} = \frac{J_B}{J_A} = \frac{|\psi_B|^2(\hbar k_1/m)}{|\psi_A|^2(\hbar k_1/m)} = \left|\frac{B}{A}\right|^2$$

Transmission = 
$$T = \frac{J_{transmitted}}{J_{incident}} = \frac{J_C}{J_A} = \frac{|\psi_C|^2(\hbar k_2/m)}{|\psi_A|^2(\hbar k_1/m)} = \left|\frac{C}{A}\right|^2 \frac{k_2}{k_1}$$

$$\frac{B}{A} = \frac{1 - k_2/k_1}{1 + k_2/k_1} \qquad \frac{C}{A} = \frac{2}{1 + k_2/k_1}$$

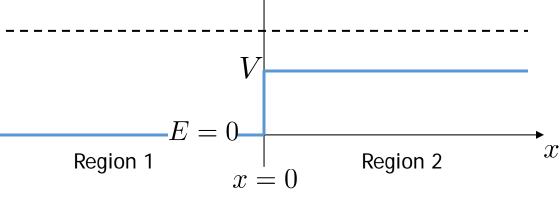
## <u>A Simple</u> <u>Potential Step</u>

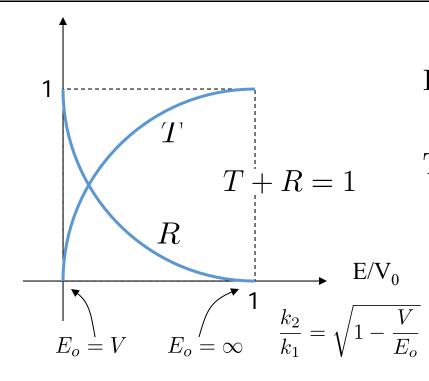
$$\psi_A = Ae^{-jk_1x}$$

$$\psi_B = Be^{-jk_1x}$$

$$\psi_C = Ce^{-jk_2 x}$$

CASE I : 
$$E_o > V$$





Reflection = 
$$R = \left| \frac{B}{A} \right|^2 = \left| \frac{k_1 - k_2}{k_1 + k_2} \right|^2$$

Transmission = 
$$T = 1 - R$$
  
=  $\frac{4k_1k_2}{|k_1 + k_2|^2}$ 

#### A Simple Potential Step

 $\psi_A = Ae^{-jk_1x}$   $\psi_B = Be^{-jk_1x}$ 

CASE II :  $E_o < V$ 

$$E_o\psi = -\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2}$$

$$\implies k_1^2 = \frac{2mE_o}{\hbar^2}$$

$$(E_o - V) \psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2}$$

$$(E_o - V) \psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} \qquad \Longrightarrow \quad \kappa^2 = \frac{2m (E_o - V)}{\hbar^2}$$

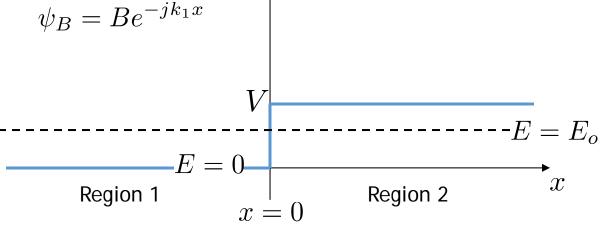
## A Simple Potential Step

$$\psi_A = Ae^{-jk_1x}$$

$$\psi_B = Be^{-jk_1x}$$

$$\psi_C = Ce^{-\kappa x}$$

CASE II : 
$$E_o < V$$



$$\psi_1 = Ae^{-jk_1x} + Be^{jk_1x}$$

$$\psi_2 = Ce^{-\kappa x}$$

$$\psi$$
 is continuous:

$$\psi_1(0) = \psi_2(0)$$

$$\longrightarrow$$
  $A+B=C$ 

$$\frac{\partial \psi}{\partial x}$$
 is continuous:

$$\frac{\partial \psi}{\partial x} \text{ is continuous: } \qquad \frac{\partial}{\partial x} \psi(0) = \frac{\partial}{\partial x} \psi_2(0) \qquad \Longrightarrow \qquad A - B = -j \frac{\kappa}{k_1} C$$

$$A - B = -j\frac{\kappa}{k_1}C$$

$$\psi_A = Ae^{-j\kappa_1 x}$$

$$\psi_B = Be^{-jk_1 x}$$

$$\psi_C = Ce^{-\kappa x}$$

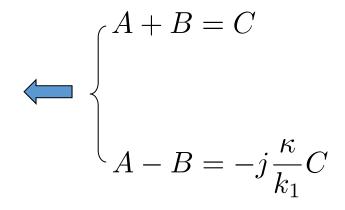
#### CASE II : $E_o < V$

$$\frac{B}{A} = \frac{1 + j\kappa/k_1}{1 - j\kappa/k_1} = \exp(2j\phi)$$
, where  $\phi = \kappa/k_1$ 

$$\frac{C}{A} = \frac{2}{1 - j\kappa/k_1}$$

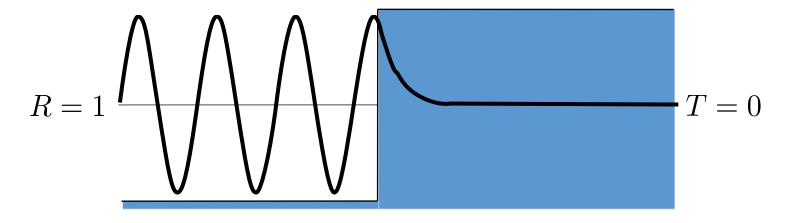
$$R = \left| \frac{B}{A} \right|^2 = 1 \qquad T = 0$$

Total reflection → Transmission must be zero

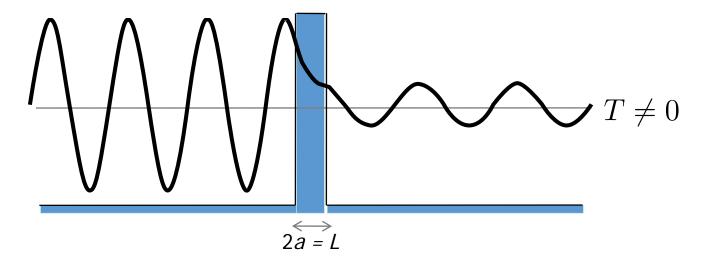


#### Quantum Tunneling Through a Thin Potential Barrier

#### Total Reflection at Boundary



#### Frustrated Total Reflection (Tunneling)



#### <u>A Rectangular</u> Potential Step

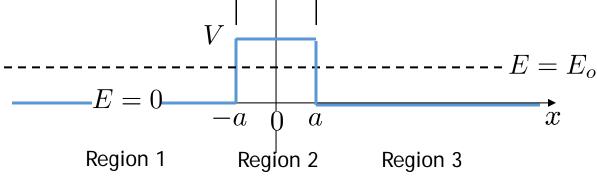
$$\psi_{A} = Ae^{-jk_{1}x} \qquad \psi_{C} = Ce^{-\kappa x}$$

$$\psi_{B} = Be^{jk_{1}x} \qquad \psi_{D} = De^{\kappa x}$$

$$\psi_F = Fe^{-jk_1x}$$

$$\psi_D = De^{\kappa x}$$

#### CASE II : $E_o < V$



In Regions 1 and 3:

$$E_o \psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} \quad \Longrightarrow \quad k_1^2 = \frac{2mE_o}{\hbar^2}$$

In Region 2:

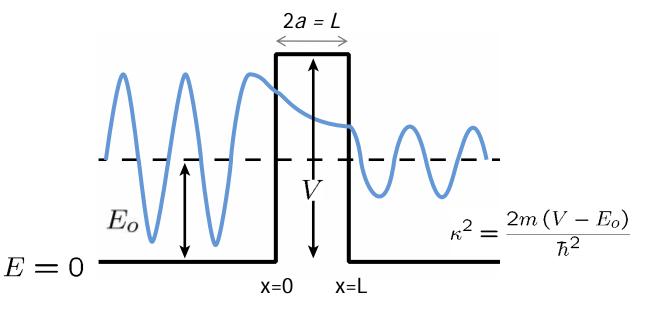
$$(E_o - V)\psi = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} \quad \Longrightarrow \quad \kappa^2 = \frac{2m(V - E_o)}{\hbar^2}$$

for 
$$E_o < V$$
:

$$T = \left| \frac{F}{A} \right|^2 = \frac{1}{1 + \frac{1}{4} \frac{V^2}{E_o(V - E_o)} \sinh^2(2\kappa a)}$$

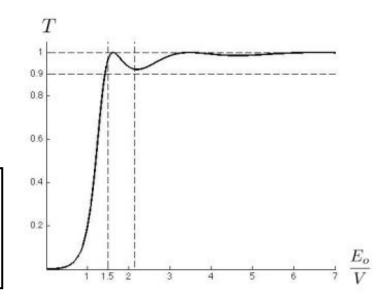
## A Rectangular Potential Step

Real part of  $\Psi$  for  $E_o < V$ , shows hyperbolic (exponential) decay in the barrier domain and decrease in amplitude of the transmitted wave.



for  $E_o < V$ :

$$T = \left| \frac{F}{A} \right|^2 = \frac{1}{1 + \frac{1}{4} \frac{V^2}{E_o(V - E_o)} \sinh^2(2\kappa a)}$$



Transmission Coefficient versus  $E_o/V$  for barrier with  $2m(2a)^2V/\hbar=16$