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In[60]:= Off[ClebschGordan::"tri"];
Off[SixJSymbol::"tri"];
(*PRD32,189-231(1985) Mesons in a relativized quark model with chromodynamics*)
(*计算所用参数*)
 $\alpha = \{0.25, 0.15, 0.20\}; \gamma = \{1/2, \sqrt{10}/2, \sqrt{1000}/2\};$ 
mu = 0.22;
ms = 0.419;
mc = 1.628;
mb = 4.977;
mt = 35;
bconst = 0.18;
cconst = -0.253;
 $\sigma_0 = 1.8;$ 
s = 1.55;
 $\epsilon_c = -0.168;$ 
 $\epsilon_t = 0.025;$ 
 $\epsilon_{sov} = -0.035;$ 
 $\epsilon_{sos} = 0.055;$ 
(*默认参数*)
r1 = 0.5; rmax = 15; nmax = 10;
(*自旋,轨道,张量项,smearing参数 $\sigma, \tau$ *)

$$dspin[S_] := \frac{1}{2} S(S+1) - \frac{3}{4};$$


$$dso[S_, L_, J_] := J(J+1) - L(L+1) - S(S+1);$$


$$dsqq[S_, L_, J_] :=$$


$$\text{If}[S == 0 \mid \mid L == 0, 0, \text{Which}[J == L+1, (J-1)/2, J == L, -1/2, J == L-1, (-J-2)/2]]];$$


$$\sigma_{ij}[m1_, m2_] := \sqrt{\sigma_0^2 * \left( \frac{1}{2} + \frac{1}{2} \left( \frac{4 m1 * m2}{(m1 + m2)^2} \right)^4 \right) + s^2 \left( \frac{2 m1 * m2}{m1 + m2} \right)^2};$$


$$\tau_{ij}[m1_, m2_, k_] := \sqrt{\frac{\gamma[[k]]^2 * \sigma_{ij}[m1, m2]^2}{\gamma[[k]]^2 + \sigma_{ij}[m1, m2]^2}};$$


$$\text{tensor}[S_, L_, J_] :=$$


$$\text{If}[S == 0 \mid \mid L == 0, 0, \text{Which}[J == L+1, -\frac{2(J-1)}{2J+1}, J == L, 2, J == L-1, -\frac{2(J+2)}{2J+1}]]];$$


$$Gtilde[m1_, m2_, r_] := - \sum_{k=1}^3 \frac{4 * \alpha[[k]]}{3 r} \text{Erf}[\tau_{ij}[m1, m2, k] * r];$$


$$Stilde[m1_, m2_, r_] := bconst * r * \left( \frac{\text{Exp}[-\sigma_{ij}[m1, m2]^2 r^2]}{\sqrt{\pi} \sigma_{ij}[m1, m2] * r} + \left( 1 + \frac{1}{2 \sigma_{ij}[m1, m2]^2 * r^2} \right) \text{Erf}[\sigma_{ij}[m1, m2] * r] \right) + cconst;$$

(*高斯基坐标和动量空间波函数*)

$$Nn1[i_, l_] := \left( \frac{2^{1+2} (2 \sqrt{n[i]})^{1+3/2}}{\sqrt{\pi} (2 l + 1) !!} \right)^{1/2};$$


$$\phi r[i_, l_, r_] := Nn1[i, l] * r^l * e^{-\sqrt{n[i]} * r^2};$$


$$\phi p[i_, l_, p_] := Nn1[i, l] * \frac{1}{(2 \sqrt{n[i]})^{1+3/2}} e^{-p^2/(4 \sqrt{n[i]})} p^l;$$

(*计算势能项相关矩阵元,动量空间和坐标空间分开计算*)

$$pij[i_, j_, l_, m1_, m2_, \epsilon x_] :=$$


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NIntegrate[ $\phi p[i, l, p] * \left( \frac{m1 * m2}{\sqrt{p^2 + m1^2} \sqrt{p^2 + m2^2}} \right)^{1/2 + \epsilon x} * \phi p[j, l, p] * p^2, \{p, 0, \infty\}];$ 
Gpij[i_, j_, l_, m1_, m2_] := NIntegrate[
 $\phi p[i, l, p] * \left( 1 + \frac{p^2}{\sqrt{p^2 + m1^2} \sqrt{p^2 + m2^2}} \right)^{1/2} * \phi p[j, l, p] * p^2, \{p, 0, \infty\}];$ 
Gtildeij[i_, j_, l_, m1_, m2_] := NIntegrate[
 $\phi r[i, l, r] * Gtilde[m1, m2, r] * \phi r[j, l, r] * r^2, \{r, 0, \infty\}];$ 
Stildeij[i_, j_, l_, m1_, m2_] := NIntegrate[
 $\phi r[i, l, r] * Stilde[m1, m2, r] * \phi r[j, l, r] * r^2, \{r, 0, \infty\}];$ 
Gij1[i_, j_, l_, m1_, m2_] := NIntegrate[ $\phi r[i, l, r] * \frac{1}{r} * (D[Gtilde[m1, m2, rp], rp] /. rp \rightarrow r) * \phi r[j, l, r] * r^2, \{r, 0, \infty\}];$ 
Gij2[i_, j_, l_, m1_, m2_] := NIntegrate[ $\phi r[i, l, r] * \frac{1}{r^2} (D[rp^2 * D[Gtilde[m1, m2, rp], rp], rp] /. rp \rightarrow r) * \phi r[j, l, r] * r^2, \{r, 0, \infty\}];$ 
Gij3[i_, j_, l_, m1_, m2_] := NIntegrate[ $\phi r[i, l, r] * (D[Gtilde[m1, m2, rp], \{rp, 2\}] /. rp \rightarrow r) * \phi r[j, l, r] * r^2, \{r, 0, \infty\}];$ 
Sij[i_, j_, l_, m1_, m2_] := NIntegrate[ $\phi r[i, l, r] * \frac{1}{r} (D[Stilde[m1, m2, rp], rp] /. rp \rightarrow r) * \phi r[j, l, r] * r^2, \{r, 0, \infty\}];$ 

mon[i_, j_, l_, m1_, m2_] :=
NIntegrate[ $\phi p[i, l, p] * (\sqrt{p^2 + m1^2} + \sqrt{p^2 + m2^2}) * \phi p[j, l, p] * p^2, \{p, 0, \infty\}];$ 

Eij[i_, j_, l_] :=  $\left( \frac{2 \sqrt{vn[i] * vn[j]}}{vn[i] + vn[j]} \right)^{1+3/2};$ 

(*求解本征方程:Sol函数,返回本征值.
*)
Sol[n1_, l1_, S1_, J1_, m11_, m21_, r11_, rmax1_, nmax1_, zz1_] :=
Module[{n = n1, l = l1, S = S1, J = J1, m1 = m11, m2 = m21, r1 = r11,
rmax = rmax1, nmax = nmax1, zz = zz1}, acon =  $\left( \frac{rmax}{r1} \right)^{1/(nmax-1)};$ 

vn[i_] :=  $\frac{1}{(r1 * acon^{i-1})^2};$ 

(*势能项矩阵*)
mat1 = Table[pij[i, j, l, m1, m1, esov], {i, 1, nmax}, {j, 1, nmax}];
mat2 = Table[pij[i, j, l, m2, m2, esov], {i, 1, nmax}, {j, 1, nmax}];
mat3 = Table[pij[i, j, l, m1, m2, esov], {i, 1, nmax}, {j, 1, nmax}];
mat4 = Table[pij[i, j, l, m1, m2, ec], {i, 1, nmax}, {j, 1, nmax}];
mat5 = Table[pij[i, j, l, m1, m2, et], {i, 1, nmax}, {j, 1, nmax}];
mat6 = Table[pij[i, j, l, m1, m1, esos], {i, 1, nmax}, {j, 1, nmax}];
mat7 = Table[pij[i, j, l, m2, m2, esos], {i, 1, nmax}, {j, 1, nmax}];
mat8 = Table[Gpij[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
mat9 = Table[Gtildeij[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
mat10 = Table[Stildeij[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
mat11 = Table[Gij1[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
mat12 = Table[Gij2[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
mat13 = Table[Gij3[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];

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mat14 = Table[Sij[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
(*动能矩阵*)
tt = Table[mon[i, j, l, m1, m2], {i, 1, nmax}, {j, 1, nmax}];
te = Table[Eij[i, j, l], {i, 1, nmax}, {j, 1, nmax}];
tein = Inverse[te];
(*H=T+V*)
th = tt + mat8.tein.mat9.tein.mat8 +  $\frac{dsqq[S, l, J]}{2 m1^2} * mat1.tein.mat11.tein.mat1 +$ 
 $\frac{dsqq[S, l, J]}{2 m2^2} * mat2.tein.mat11.tein.mat2 + \frac{2 dsqq[S, l, J]}{m1 * m2} *$ 
 $mat3.tein.mat11.tein.mat3 + \frac{2 dspin[S]}{3 m1 * m2} * mat4.tein.mat12.tein.mat4 -$ 
 $\frac{1}{12} * \frac{tensor[S, l, J]}{m1 * m2} (mat5.tein.mat13.tein.mat5 - mat5.tein.mat11.tein.mat5) +$ 
 $mat10 - \left( \frac{dsqq[S, l, J]}{2 m1^2} * mat6.tein.mat14.tein.mat6 + \right.$ 
 $\left. \frac{dsqq[S, l, J]}{2 m2^2} * mat7.tein.mat14.tein.mat7 \right);$ 
(*求解本征值和本征向量,得到波函数*)
{Ed, wave} = Eigensystem[{th, te}];
en = Ed[[-(n + 1)]];
tb = Table[Sqrt[2] vn[i], {i, nmax}];
coabc = Sign[Sum_{i=1}^{nmax} wave[[-(n + 1)]] [[i]] *  $\phi r[i, l, 0.1]$ ] *
(wave[[-(n + 1)]] . te.wave[[-(n + 1)]] )^{-1/2} * wave[[-(n + 1)]];
unr[n_, l_, r_] := Sum_{i=1}^{nmax} coabc[[i]] *  $\phi r[i, l, r]$  * r;

(*输出能量,画图*)
If[zz == 0, Print["Eigenvalue = ", en];
Print[Plot[unr[n, l, r], {r, 0, 10}]]]; Return[unr[n, l, r]];

```

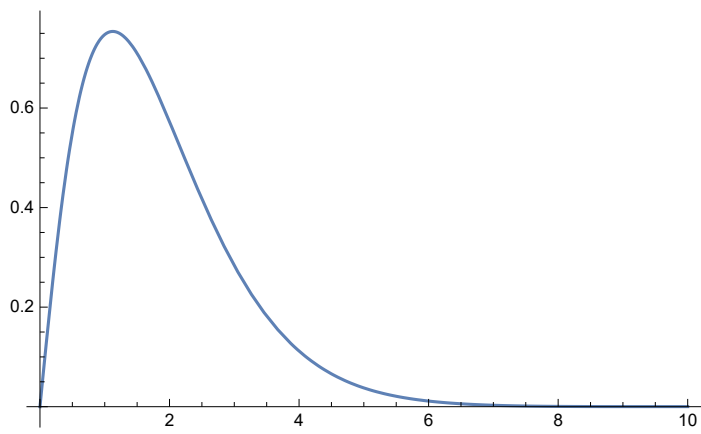
Results

In[131]:= << MaTeX`

J/ψ

In[153]:= n1 = 0; l1 = 0; S1 = 1; J1 = 0; m11 = mc; m21 = mc;
psi[r_] = Sol[n1, l1, S1, J1, m11, m21, r1, rmax, nmax, 0];

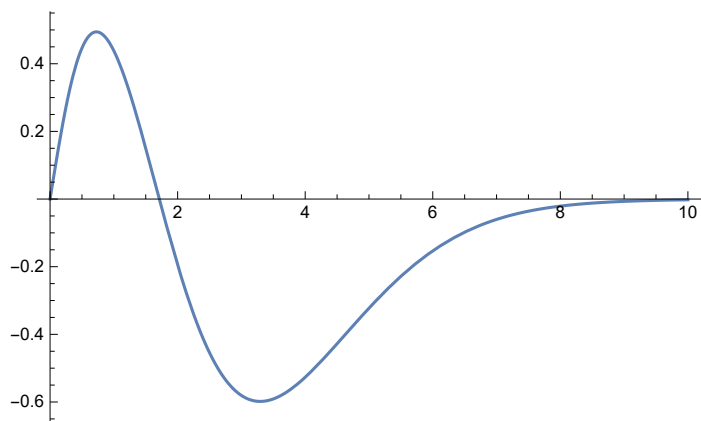
Eigenvalue = 3.09138



$\psi(2S)$

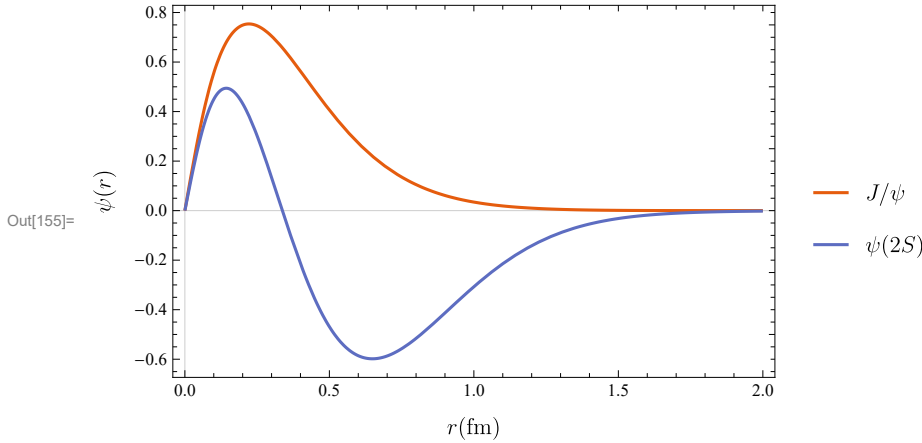
```
In[149]:= n1 = 1; l1 = 0; S1 = 1; J1 = 0; m11 = mc; m21 = mc;
psi2[r_] = Sol[n1, l1, S1, J1, m11, m21, r1, rmax, nmax, 0];
```

Eigenvalue = 3.67913



Wavefunction

```
In[155]:= Plot[{psi[r/0.197], psi2[r/0.197]}, {r, 0, 2}, Frame → True,
  PlotTheme → "Scientific", FrameLabel → {MaTeX["r (\\rm fm)"], MaTeX["\\psi(r)"]},
  PlotLegends → {MaTeX["J/\\psi"], MaTeX["\\psi(2S)"]}]
```



Check Orthonormal

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In[159]:= NIntegrate[psi[r]^2, {r, 0, Infinity}]
NIntegrate[psi2[r]^2, {r, 0, Infinity}]
NIntegrate[psi[r] * psi2[r], {r, 0, Infinity}] // Quiet
```

Out[159]= 1.

Out[160]= 1.

Out[161]= 1.3522×10^{-13}

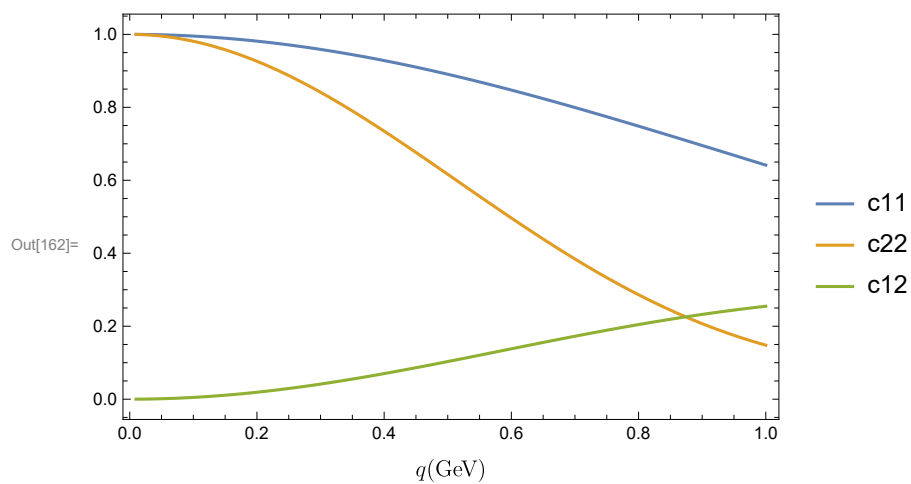
Momentum emission

Naively thinking, $\pi\pi$ emission with 3-momentum q from charmonia can be simply treated as the momentum loss of one c quark. The coupling of $\psi\psi\pi\pi$ should be proportional to the inner product of the initial and final states of charmonia. After the loss of momentum q , the state becomes

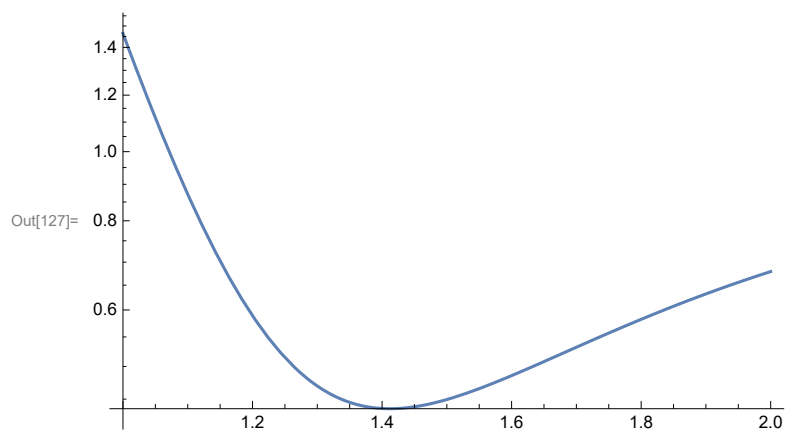
$$|\psi_f\rangle = e^{iqr} |\psi_i\rangle \text{ and then the coupling } \propto \langle \psi_f | \psi_i \rangle = \int dr \psi_f(r)^* \psi_i(r) \frac{\sin(qr)}{qr}$$

```
In[121]:= c12[q_] := NIntegrate[psi[r] * psi2[r] *  $\frac{\text{Sin}[q r]}{q r}$ , {r, 0, 20}]
c11[q_] := NIntegrate[psi[r] * psi[r] *  $\frac{\text{Sin}[q r]}{q r}$ , {r, 0, 20}]
c22[q_] := NIntegrate[psi2[r] * psi2[r] *  $\frac{\text{Sin}[q r]}{q r}$ , {r, 0, 20}]
```

```
In[162]:= Plot[{c11[q], c22[q], c12[q]}, {q, 0.01, 1}, Frame → True,
  FrameLabel → {MaTeX["q(\rm GeV)"]}, PlotLegends → {"c11", "c22", "c12"}]
```



```
In[127]:= LogPlot[c11[r] * c22[r] / c12[r]^2, {r, 1, 2}]
```



```
In[128]:= LogPlot[c11[r] / c12[r], {r, 2 * 0.138, 1}]
```

