

The electron gas in two and three dimensions

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Outline

Motivation

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Setting up the system

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Setting up the system

What could we do?

- ▶ The first study of the 2D electron gas with SRG, FCIQMC, and CC
- ▶ How well can the methods describe correlations?
- ▶ At what densities do SRG and CC break down?
- ▶ Theoretical explanation of breakdown?
- ▶ Ring and ladder approximations vs. CCD

What could we do?

- ▶ What are differences between finite and periodic systems? (QD vs. 2DEG)
- ▶ Compare 2D with finite thickness 2D+1D (or 3D?)
- ▶ Virial theorem (Jørgen's thesis)?
- ▶ More ideas?

Books

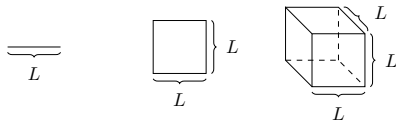
- ▶ Basics: Fetter and Walecka, Quantum theory of many-particle systems
- ▶ E-book: G. Giuliani and G. Vignale, Quantum theory of the electron liquid (library will buy it)

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The two-dimensional electron gas



- ▶ Interacting electrons in a quadratic box
- ▶ Homogeneous system (constant density)
- ▶ Constant positive background charge

Single-particle states

One electron in a quadratic box:

$$-\frac{\hbar^2}{2m}\nabla^2\varphi(\mathbf{x}) = \varepsilon\varphi(\mathbf{x}) \quad (1)$$

Solutions:

$$\varphi_{\mathbf{k}}(\mathbf{x}) = \frac{1}{\sqrt{A}} \exp(i\mathbf{k} \cdot \mathbf{x}) \quad (2)$$

Single-particle states

Periodic boundary conditions

$$\begin{aligned}\varphi_{\mathbf{k}}(x + L, y) &= \varphi_{\mathbf{k}}(x, y) \\ \varphi_{\mathbf{k}}(x, y + L) &= \varphi_{\mathbf{k}}(x, y)\end{aligned}\tag{3}$$

gives

$$\begin{aligned}\mathbf{k} &= \frac{2\pi}{L}(n_x, n_y), \\ n_x, n_y &\in \{0, \pm 1, \pm 2, \dots\}.\end{aligned}\tag{4}$$

Single-particle states

Periodic boundary conditions

$$\begin{aligned}\varphi_{\mathbf{k}}(x + L, y) &= \varphi_{\mathbf{k}}(x, y) \\ \varphi_{\mathbf{k}}(x, y + L) &= \varphi_{\mathbf{k}}(x, y)\end{aligned}\tag{3}$$

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$$\begin{aligned}\mathbf{k} &= \frac{2\pi}{L}(n_x, n_y), \\ n_x, n_y &\in \{0, \pm 1, \pm 2, \dots\}.\end{aligned}\tag{4}$$

Single-particle energies:

$$\begin{aligned}\varepsilon_{n_x, n_y} &= \frac{\hbar^2 k^2}{2m} \\ &= \frac{\hbar^2}{2m} \left(\frac{2\pi}{L} \right)^2 (n_x^2 + n_y^2)\end{aligned}\tag{5}$$

Hamiltonian operator

$$\hat{H} = \hat{H}_0 + \hat{H}_I + \hat{H}_{bb} + \hat{H}_{eb}$$

$$\hat{H}_0 = \sum_{i=1}^N -\frac{\hbar^2 \nabla_i^2}{2m}$$

$$\hat{H}_I = \frac{1}{2} e^2 \sum_{i < j}^N \frac{e^{-\mu |\mathbf{r}_i - \mathbf{r}_j|}}{|\mathbf{r}_i - \mathbf{r}_j|}$$

$$\hat{H}_{bb} = \frac{1}{2} e^2 \int d\mathbf{r} \int d\mathbf{r}' \rho(\mathbf{r}) \rho(\mathbf{r}') \frac{e^{-\mu |\mathbf{r} - \mathbf{r}'|}}{|\mathbf{r} - \mathbf{r}'|}$$

$$\hat{H}_{eb} = -e^2 \sum_{i=1}^N \int d\mathbf{r} \rho(\mathbf{r}) \frac{e^{-\mu |\mathbf{r} - \mathbf{r}_i|}}{|\mathbf{r} - \mathbf{r}_i|}$$

$$\rho(\mathbf{r}) = \text{particle density} = N/A$$

Hamiltonian operator

Can be shown (see Fetter and Walecka for 3D):

$$\hat{H} = \hat{H}_0 + \hat{H}_I - \frac{\pi e^2 N^2}{A_\mu} \quad (6)$$

The $\mathbf{q} = 0$ terms in the interaction operator give

$$\frac{\pi e^2}{A_\mu} (N^2 - N) \quad (7)$$

Let first $A \rightarrow \infty$ and then $\mu \rightarrow 0$

Interaction matrix elements

- ▶ Fetter and Walecka
- ▶ Hamaguchi: Basic Semiconductor Physics (e-book)

$$\begin{aligned}& \langle \mathbf{k}_1 m_{s_1} \mathbf{k}_2 m_{s_2} | v | \mathbf{k}_3 m_{s_3} \mathbf{k}_4 m_{s_4} \rangle \\&= \frac{e^2}{A^2} \int d\mathbf{r}_1 \int d\mathbf{r}_2 e^{-i\mathbf{k}_1 \cdot \mathbf{r}_1} e^{-i\mathbf{k}_2 \cdot \mathbf{r}_2} \frac{e^{-\mu|\mathbf{r}_1 - \mathbf{r}_2|}}{|\mathbf{r}_1 - \mathbf{r}_2|} e^{i\mathbf{k}_3 \cdot \mathbf{r}_1} e^{i\mathbf{k}_4 \cdot \mathbf{r}_2} \\& \quad \times \langle m_{s_1} | m_{s_3} \rangle \langle m_{s_2} | m_{s_4} \rangle \\&= \dots \\&= \frac{e^2}{A} \delta_{m_{s_1} m_{s_3}} \delta_{m_{s_2} m_{s_4}} \delta_{\mathbf{k}_1 + \mathbf{k}_2, \mathbf{k}_3 + \mathbf{k}_4} \int d\mathbf{r} e^{i(\mathbf{k}_3 - \mathbf{k}_1) \cdot \mathbf{r}} \frac{e^{-\mu r}}{r}\end{aligned}$$

Interaction matrix elements

From Abramowitz:

$$\int_0^{2\pi} e^{ipr \cos \phi} = 2\pi J_0(pr) \quad (8)$$

$J_0(x)$ = Bessel function of first kind

From Gradshteyn and Ryzhik:

$$\int_0^\infty dr e^{-\mu r} J_0(pr) = \frac{1}{\sqrt{\mu^2 + p^2}} \quad (9)$$

Antisymmetrized matrix elements

$$\begin{aligned} & \langle \mathbf{k}_1 m_{s_1} \mathbf{k}_2 m_{s_2} | v | \mathbf{k}_3 m_{s_3} \mathbf{k}_4 m_{s_4} \rangle_{AS} \\ &= \frac{e^2}{A} \delta_{\mathbf{k}_1 + \mathbf{k}_2, \mathbf{k}_3 + \mathbf{k}_4} \left\{ \delta_{m_{s_1} m_{s_3}} \delta_{m_{s_2} m_{s_4}} (1 - \delta_{\mathbf{k}_1 \mathbf{k}_3}) \frac{2\pi}{\sqrt{\mu^2 + |\mathbf{k}_3 - \mathbf{k}_1|^2}} \right. \\ & \quad \left. - \delta_{m_{s_1} m_{s_4}} \delta_{m_{s_2} m_{s_3}} (1 - \delta_{\mathbf{k}_1 \mathbf{k}_4}) \frac{2\pi}{\sqrt{\mu^2 + |\mathbf{k}_4 - \mathbf{k}_1|^2}} \right\} \end{aligned}$$

- ▶ Atomic units: $\hbar = m = e = 1$
- ▶ Classical electron radius: $N\pi r_s^2 = A$
- ▶ N and r_s input parameters
- ▶ Quantum numbers: (n_x, n_y, m_s)

Orbitals

$n_x^2 + n_y^2$	n_x	n_y	$N_{\uparrow\downarrow}$	$N_{\uparrow\uparrow}$
0	0	0	2	1
1	-1	0	10	5
	1	0		
	0	-1		
	0	1		
2	-1	-1	18	9
	-1	1		
	1	-1		
	1	1		
4	-2	0	26	13
	2	0		
	0	-2		
	0	2		
5	-2	-1	42	21
	2	-1		
	-2	1		
	2	1		
	-1	-2		
	-1	2		
	1	-2		
	1	2		

The three-dimensional electron gas

Single-particle energies:

$$\varepsilon_{n_x, n_y, n_z} = \frac{\hbar^2}{2m} \left(\frac{2\pi}{L} \right)^2 (n_x^2 + n_y^2 + n_z^2)$$

Antisymmetrized interaction matrix elements:

$$\begin{aligned} & \langle \mathbf{k}_1 m_{s_1} \mathbf{k}_2 m_{s_2} | v | \mathbf{k}_3 m_{s_3} \mathbf{k}_4 m_{s_4} \rangle_{AS} \\ &= \frac{e^2}{V} \delta_{\mathbf{k}_1 + \mathbf{k}_2, \mathbf{k}_3 + \mathbf{k}_4} \left\{ \delta_{m_{s_1} m_{s_3}} \delta_{m_{s_2} m_{s_4}} (1 - \delta_{\mathbf{k}_1 \mathbf{k}_3}) \frac{4\pi}{\mu^2 + (\mathbf{k}_3 - \mathbf{k}_1)^2} \right. \\ & \quad \left. - \delta_{m_{s_1} m_{s_4}} \delta_{m_{s_2} m_{s_3}} (1 - \delta_{\mathbf{k}_1 \mathbf{k}_4}) \frac{4\pi}{\mu^2 + (\mathbf{k}_4 - \mathbf{k}_1)^2} \right\} \end{aligned}$$