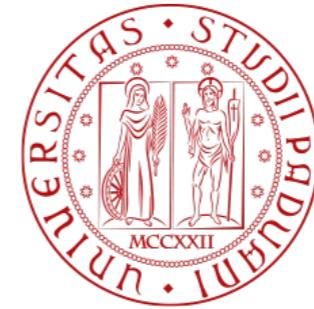
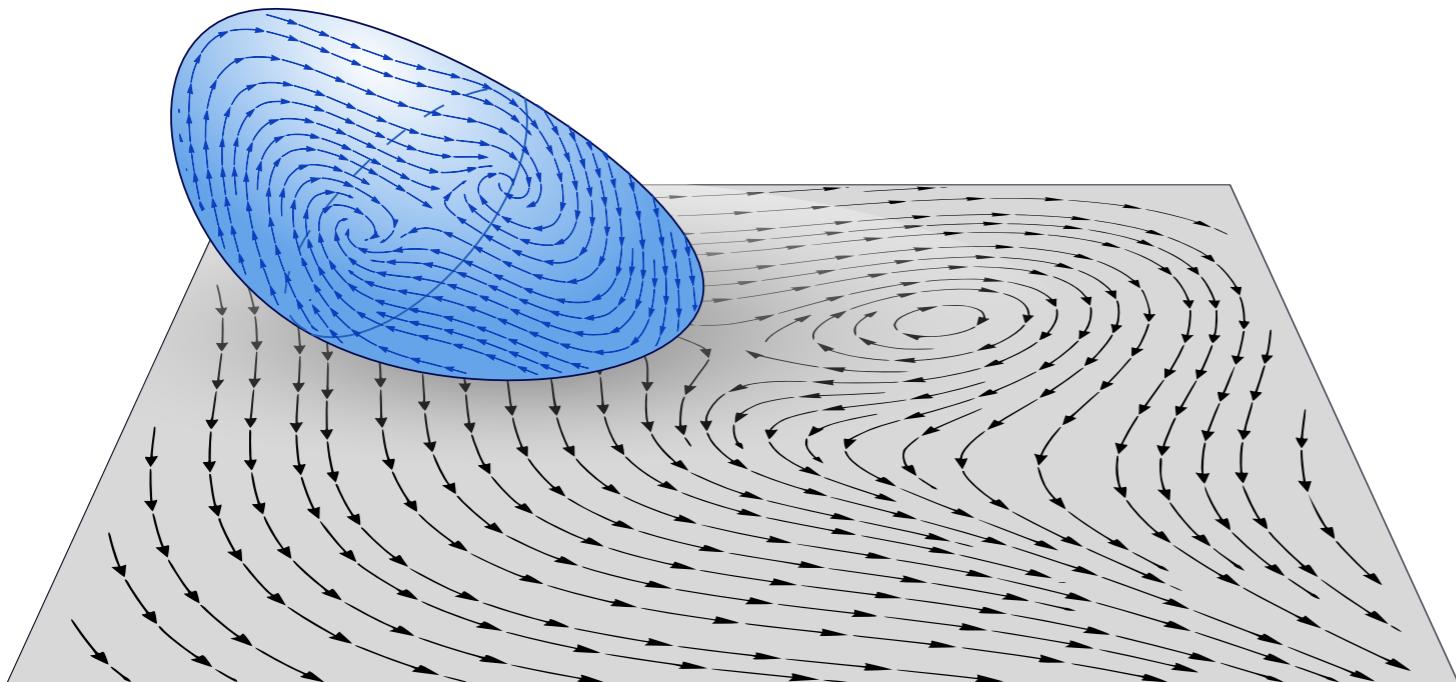


# Integrable and chaotic generalizations of the rolling of a convex body without slipping on a plane

Mariana Costa Villegas  
Joint work with Luis García Naranjo



UNIVERSITÀ  
DEGLI STUDI  
DI PADOVA

DIPARTIMENTO  
**MATEMATICA**  
DIPARTIMENTO DI MATEMATICA “TULLIO LEVI-CIVITA”

XVIII International Young  
Researchers Workshop in Geometry,  
Dynamics and Field Theory

Warsaw, Poland, February 21-23, 2024

# Classic convex body rolling without slipping on a plane

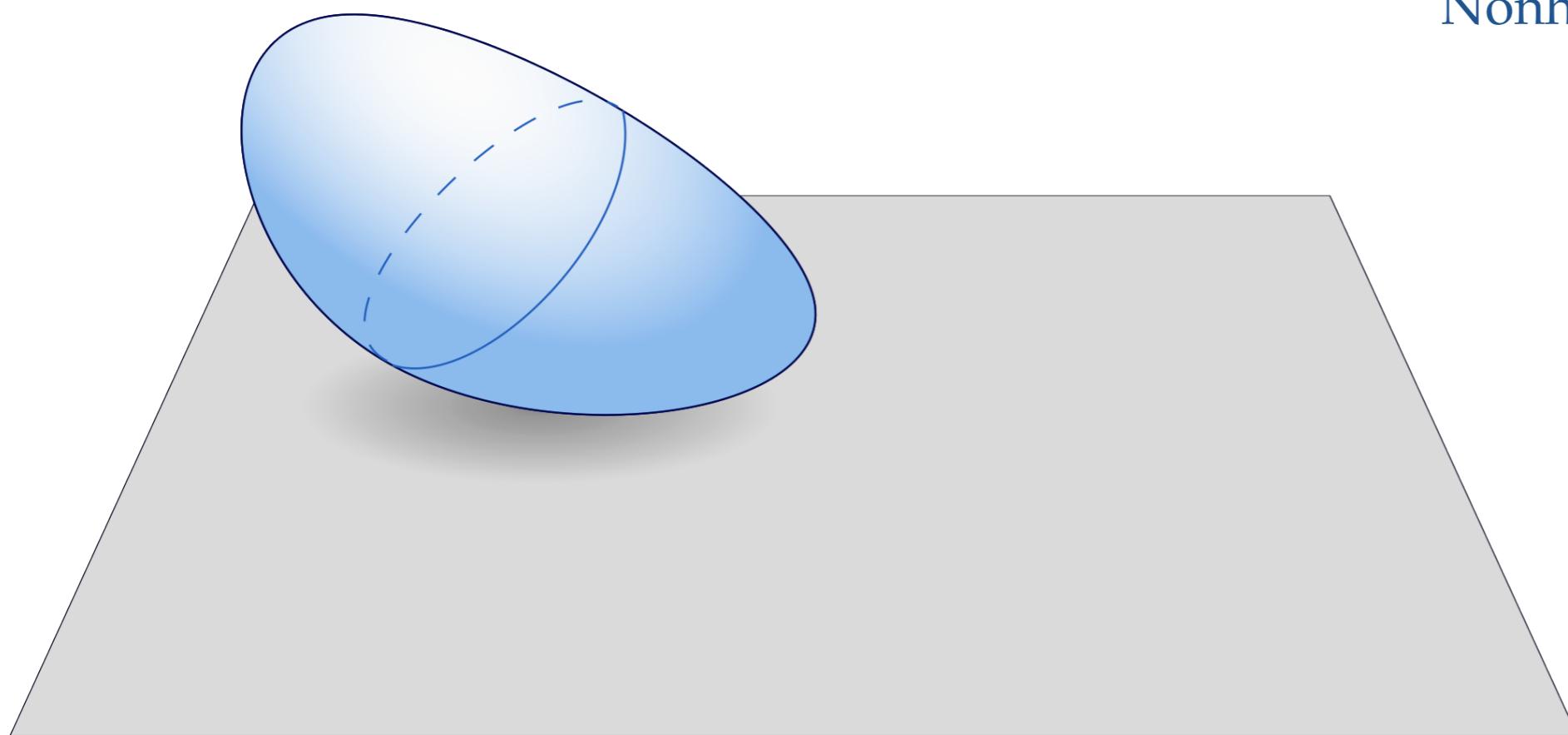
Nonholonomic system

Configuration manifold:  $Q = SO(3) \times \mathbb{R}^2$   
 $(B, \underline{x}) \in Q$

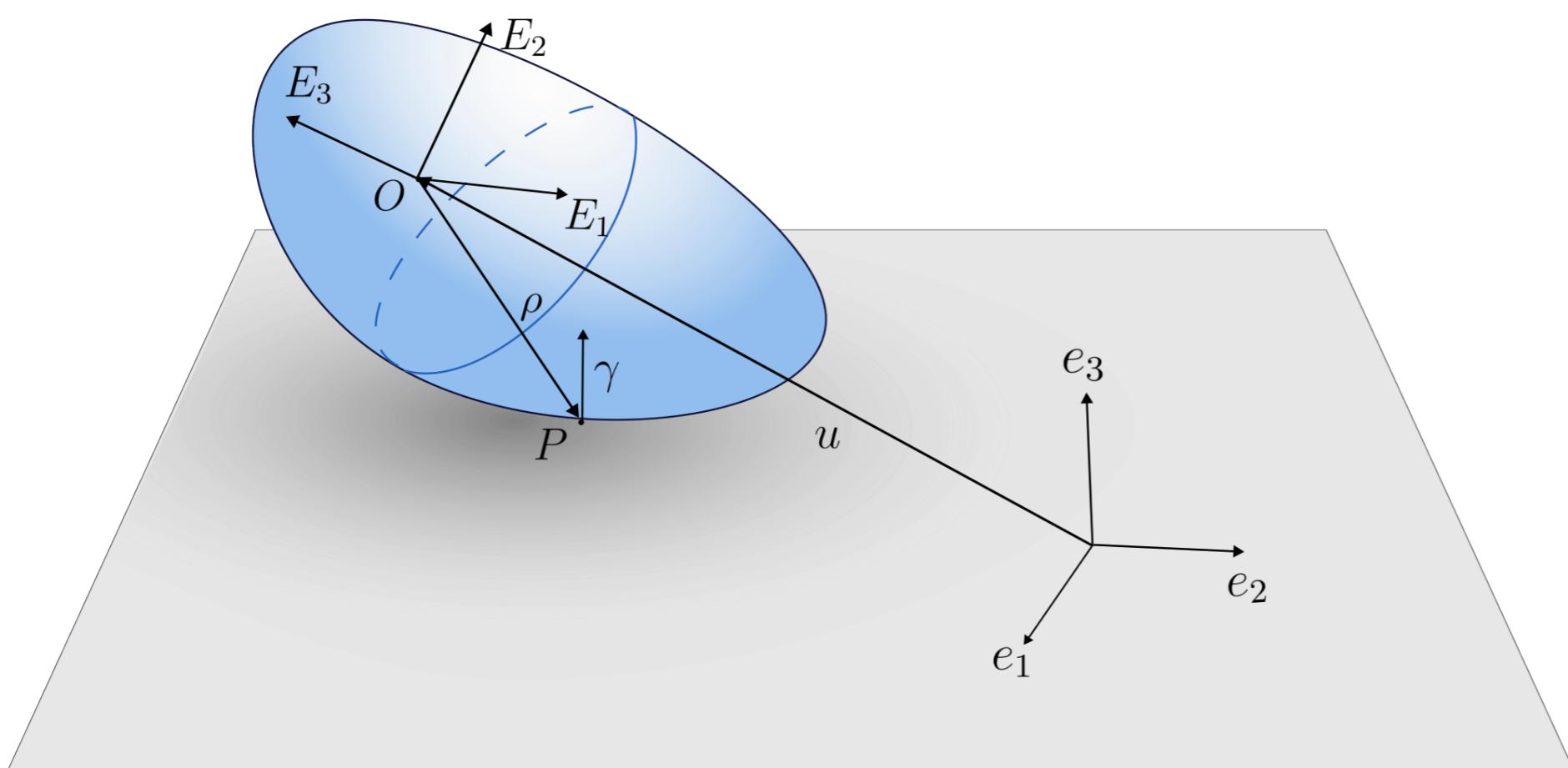
Nonholonomic linear constraint

Phase space: dim 8

$SE(2)$ - symmetry on plane



# Classic convex body rolling without slipping on a plane



$$\gamma = B^{-1}e_3$$

$$\rho = \rho(\gamma)$$

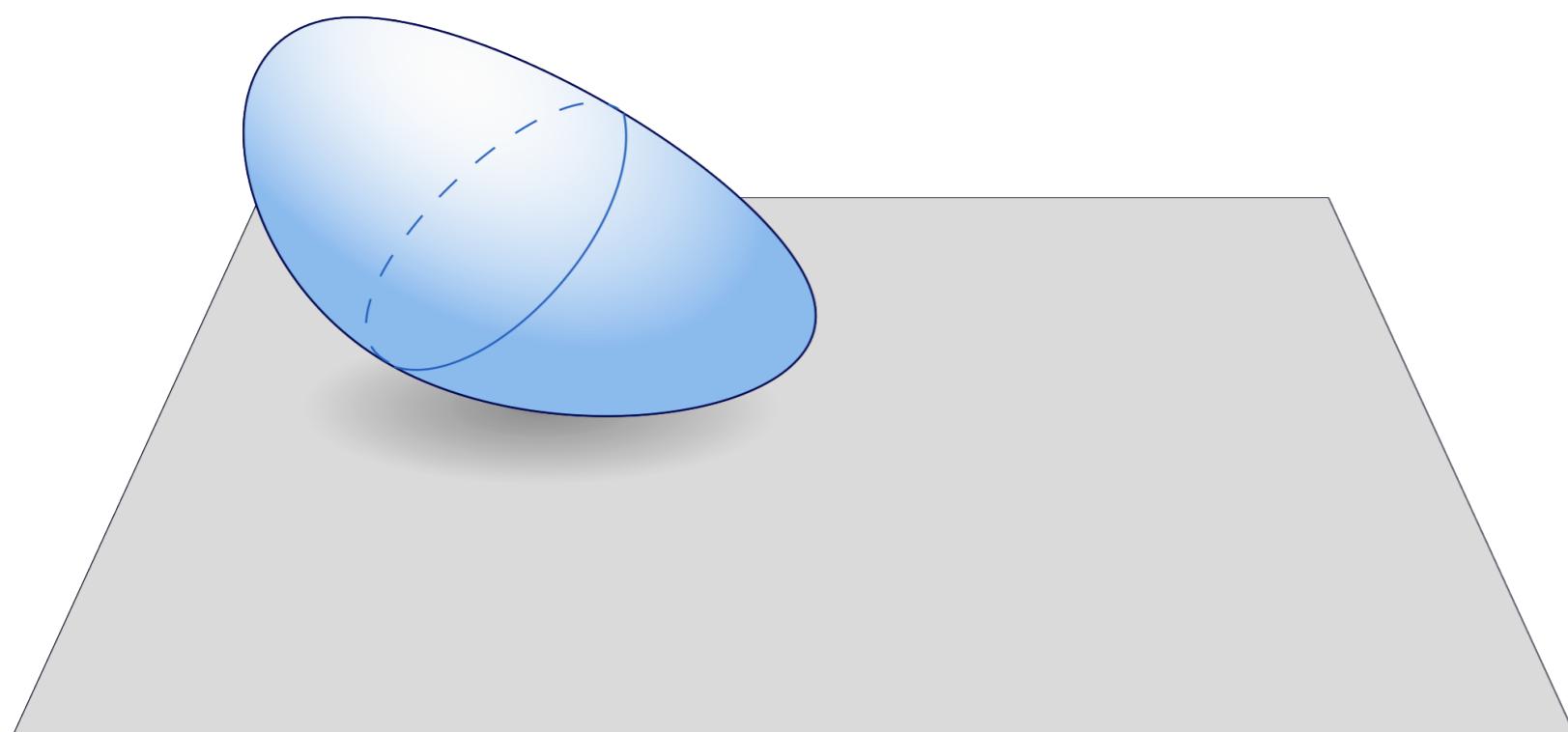
$$\underline{x} = u + B\rho$$

Equations of motion:

$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$

$$\dot{\gamma} = \Omega \times \gamma$$

$$M = I\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$



Equations of motion:

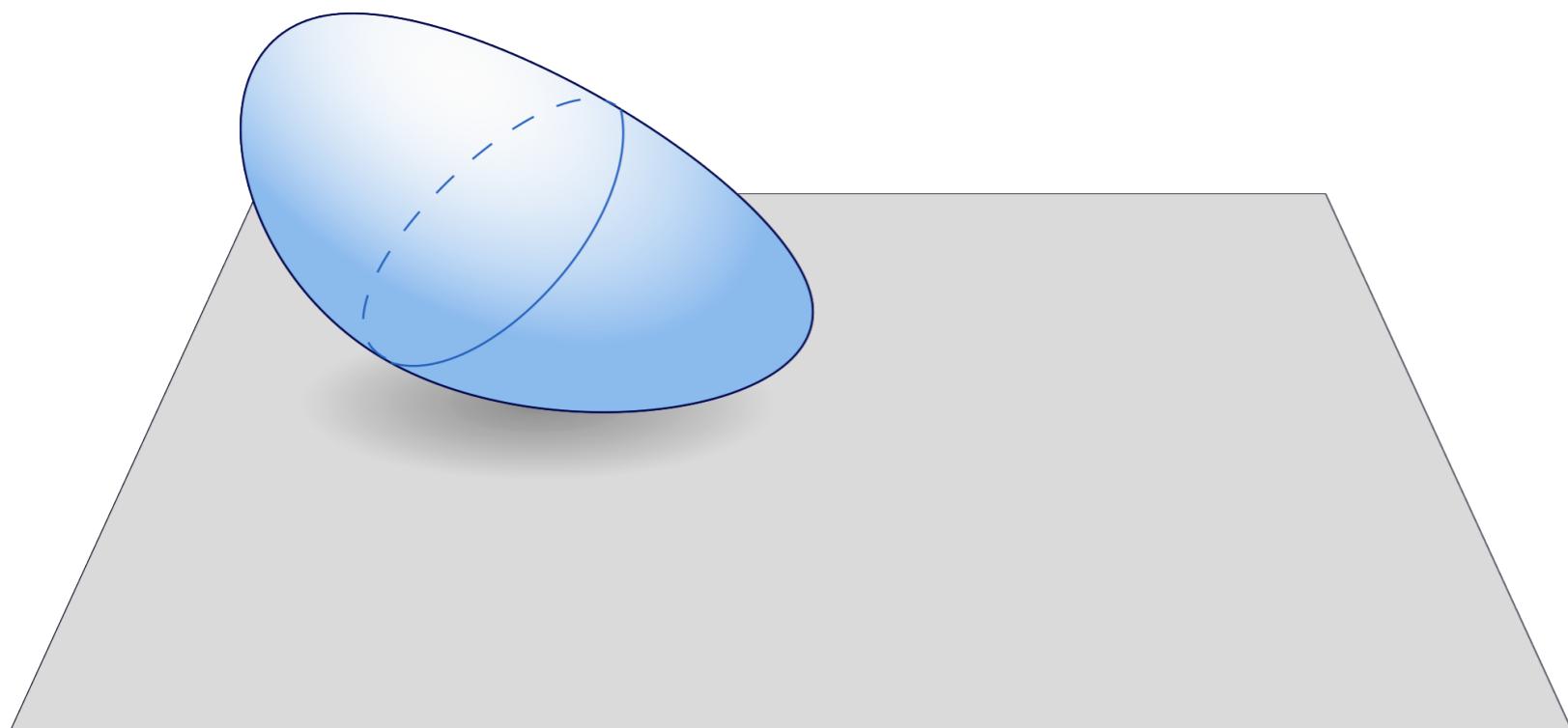
$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$

$$\dot{\gamma} = \Omega \times \gamma$$

$$M = I\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$

First integrals:

$$H = \frac{1}{2}\langle M, \Omega \rangle - mg\langle \rho, \gamma \rangle, \quad \|\gamma\|^2 = 1$$



Equations of motion:

$$\dot{M} = M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma$$

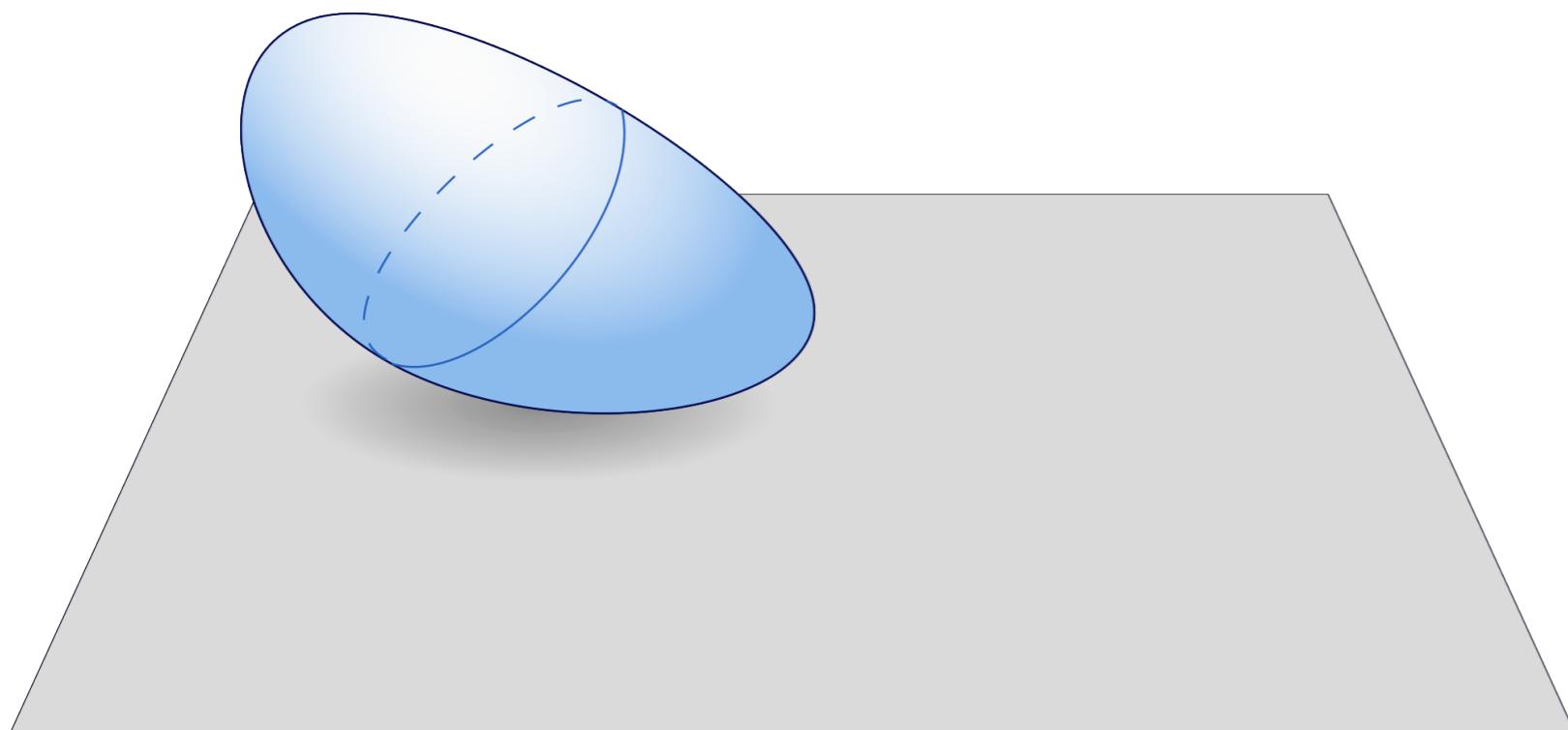
$$\dot{\gamma} = \Omega \times \gamma$$

$$M = I\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$

First integrals:

$$H = \frac{1}{2}\langle M, \Omega \rangle - mg\langle \rho, \gamma \rangle, \quad \|\gamma\|^2 = 1$$

- No invariant measure
- Generally chaotic



Equations of motion:

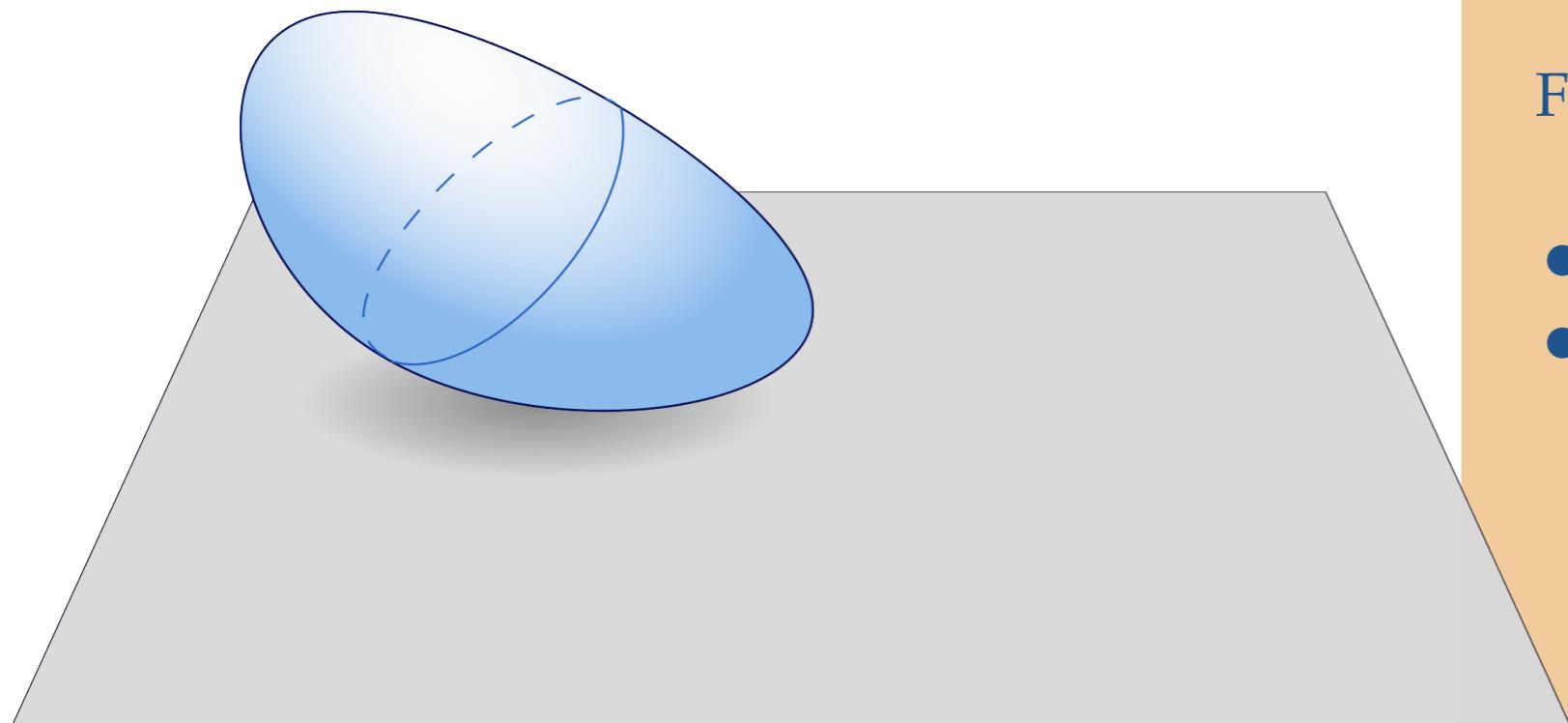
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$$M = I\Omega + m\rho \times (\Omega \times \rho), \quad \rho = \rho(\gamma)$$

First integrals:

$$H = \frac{1}{2}\langle M, \Omega \rangle - mg\langle \rho, \gamma \rangle, \quad \|\gamma\|^2 = 1$$



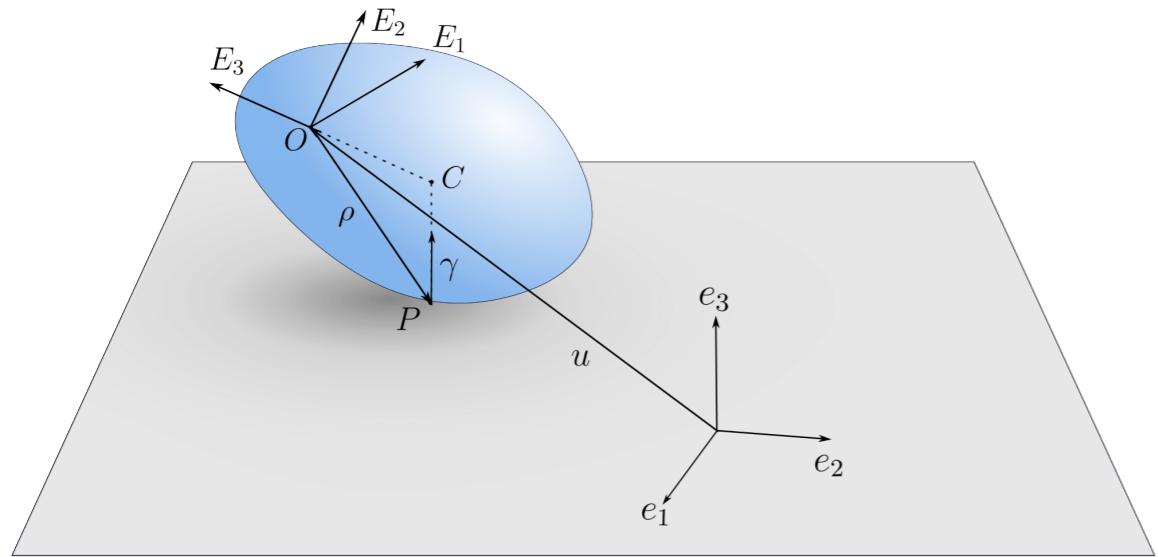
- No invariant measure
- Generally chaotic

For it to be integrable we need:

- 2 additional first integrals
- invariant measure.

# Integrable cases

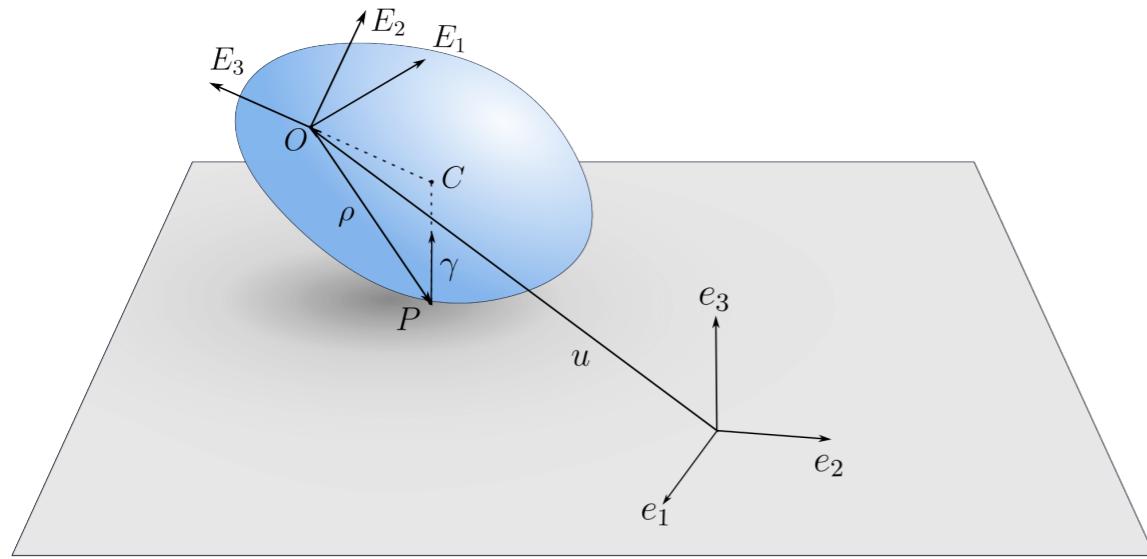
# Integrable cases



Body of revolution (Chaplygin, 1897)

$$I_1 = I_2 \neq I_3, \quad O \neq C$$

# Integrable cases



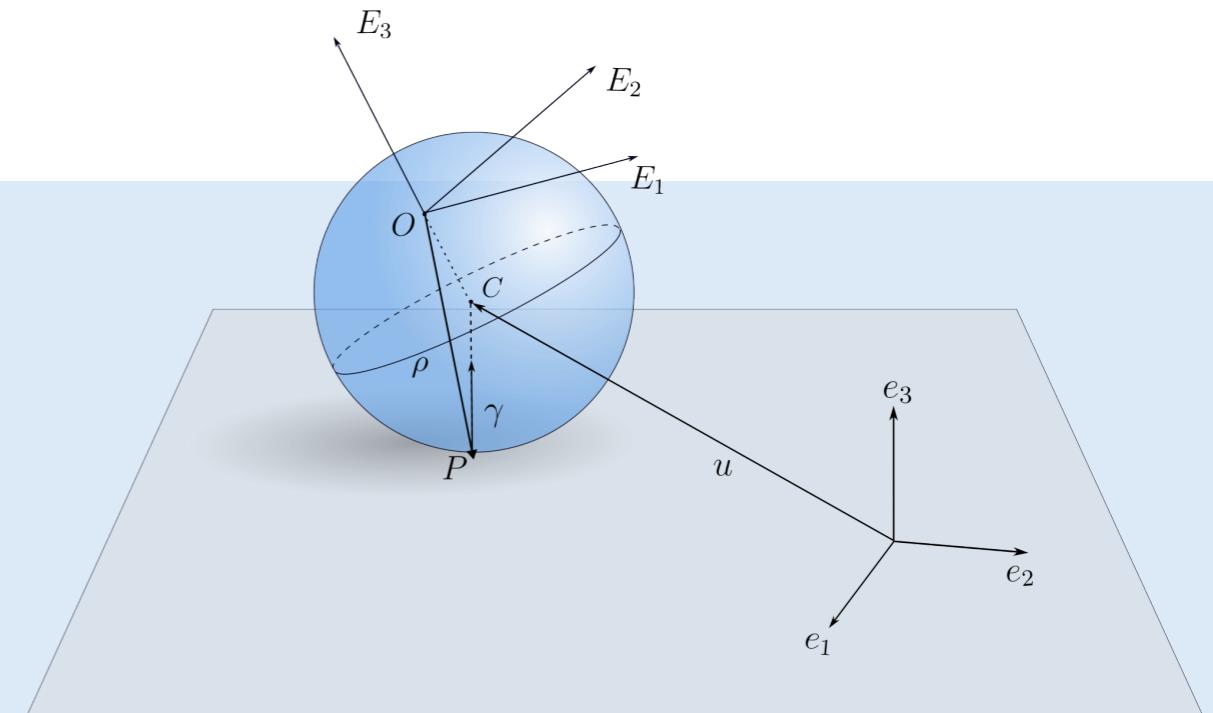
Routh's sphere (Routh, 1884)

$$I_1 = I_2 \neq I_3, \quad O \neq C$$

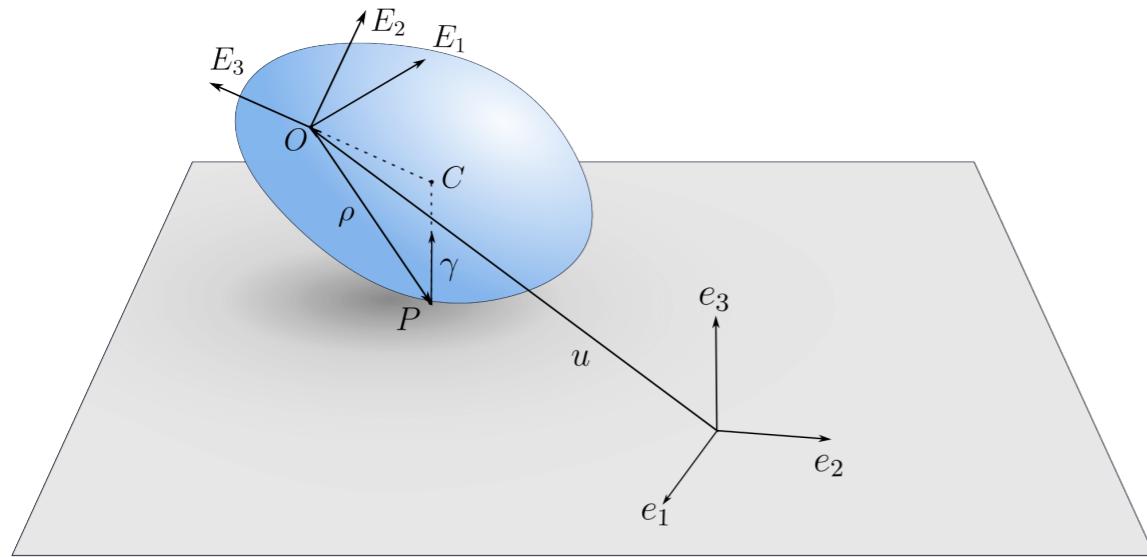
$$\rho(\gamma) = -r\gamma - lE_3$$

Body of revolution (Chaplygin, 1897)

$$I_1 = I_2 \neq I_3, \quad O \neq C$$



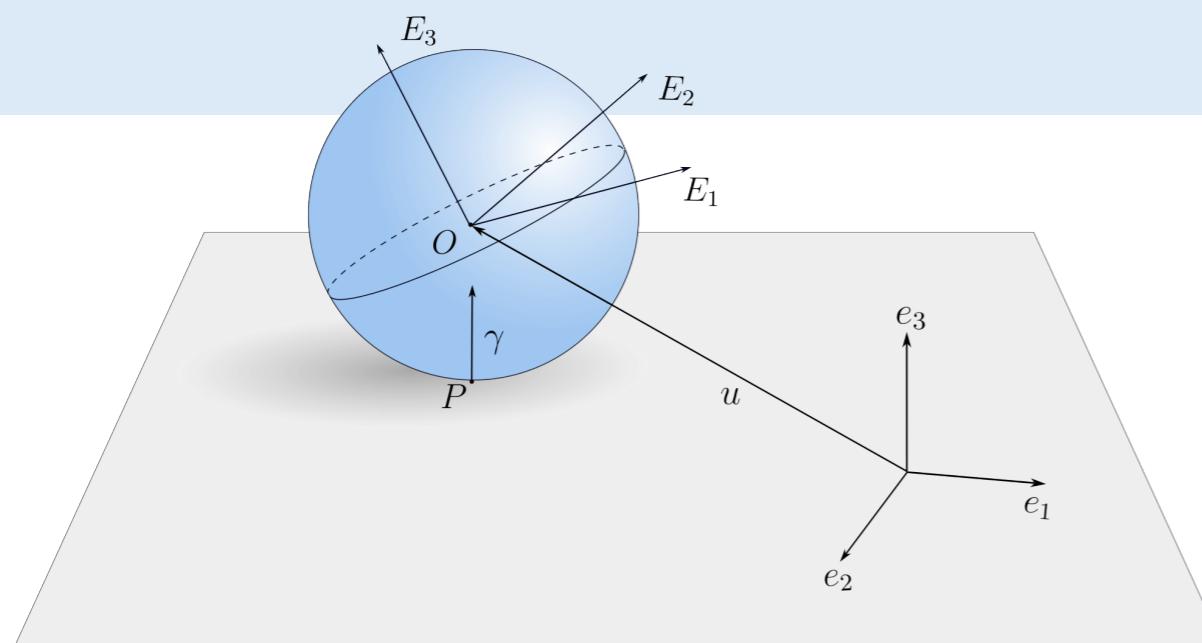
# Integrable cases



Routh's sphere (Routh, 1884)

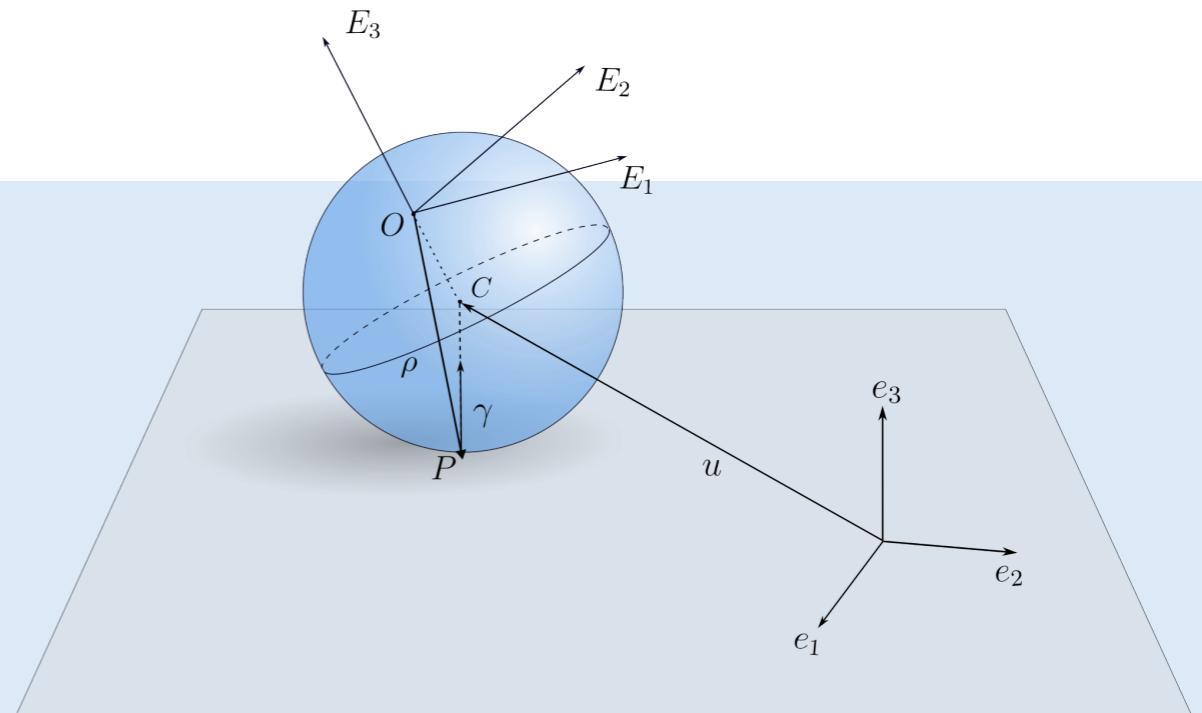
$$I_1 = I_2 \neq I_3, O \neq C$$

$$\rho(\gamma) = -r\gamma - lE_3$$



Body of revolution (Chaplygin, 1897)

$$I_1 = I_2 \neq I_3, O \neq C$$



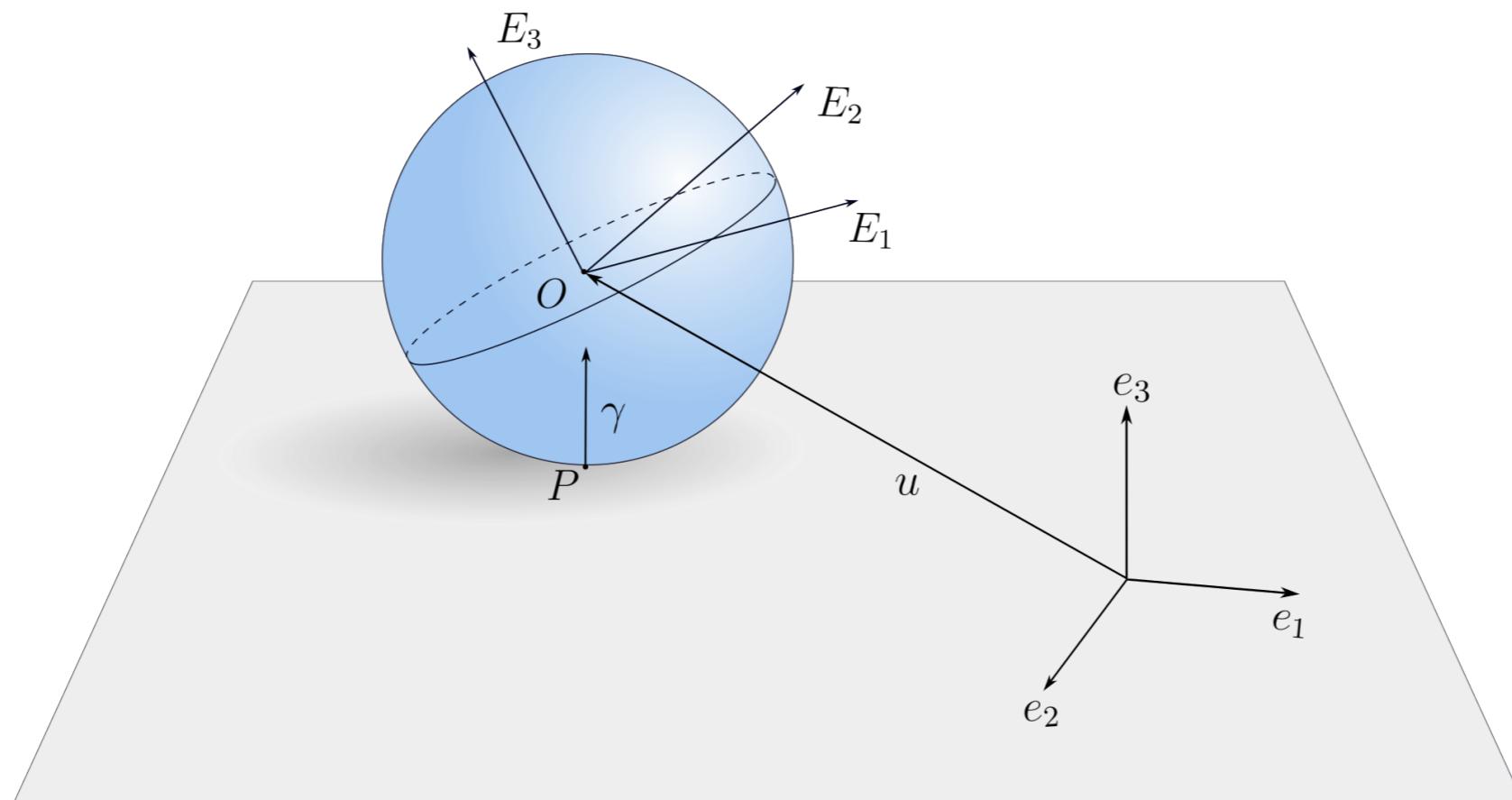
Chaplygin sphere (Chaplygin, 1903)

$$I_1 \neq I_2 \neq I_3, O = C$$

# Integrable cases

## Homogeneous sphere

$$I_1 = I_2 = I_3, \quad O = C$$

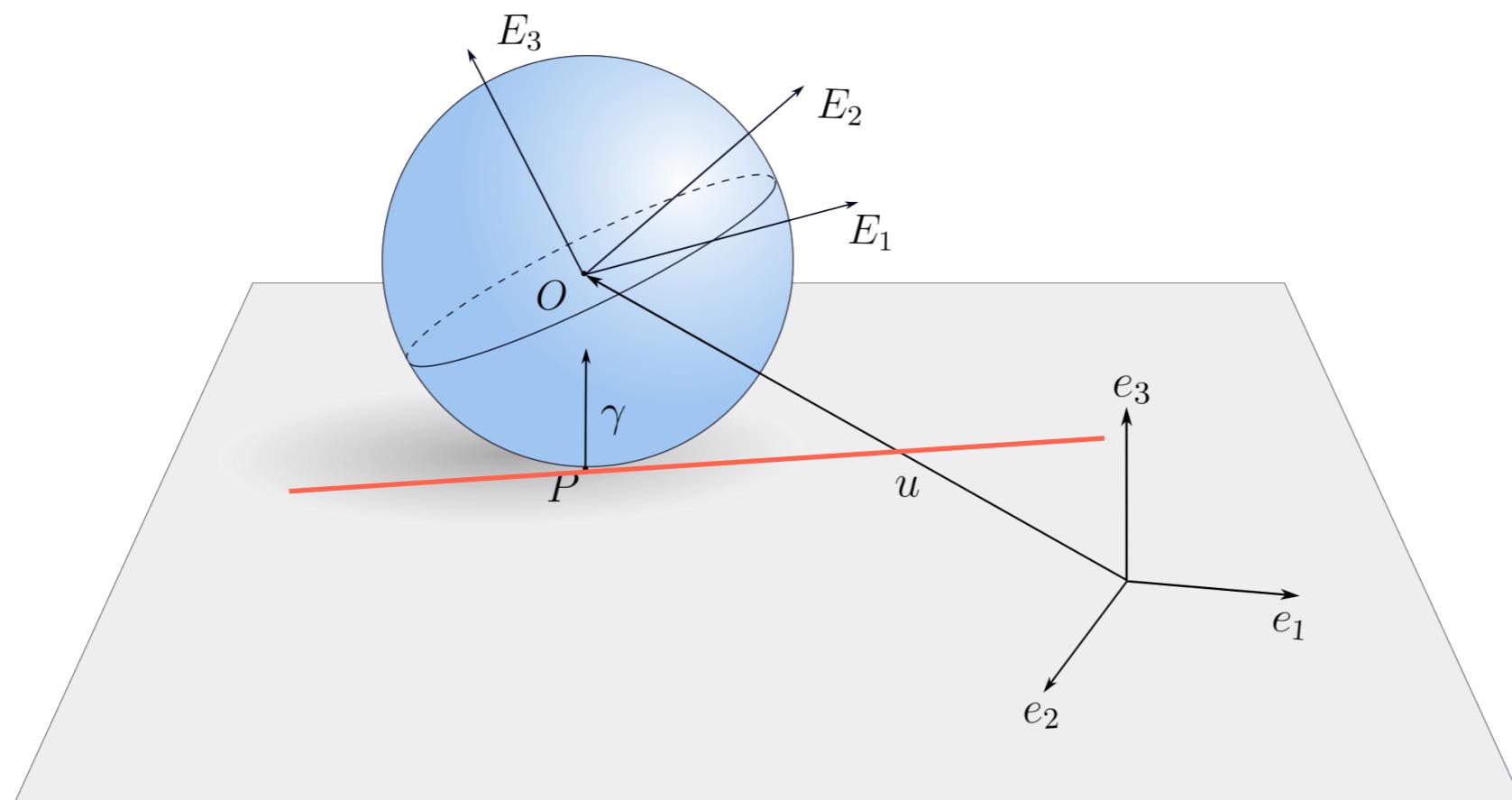


# Integrable cases

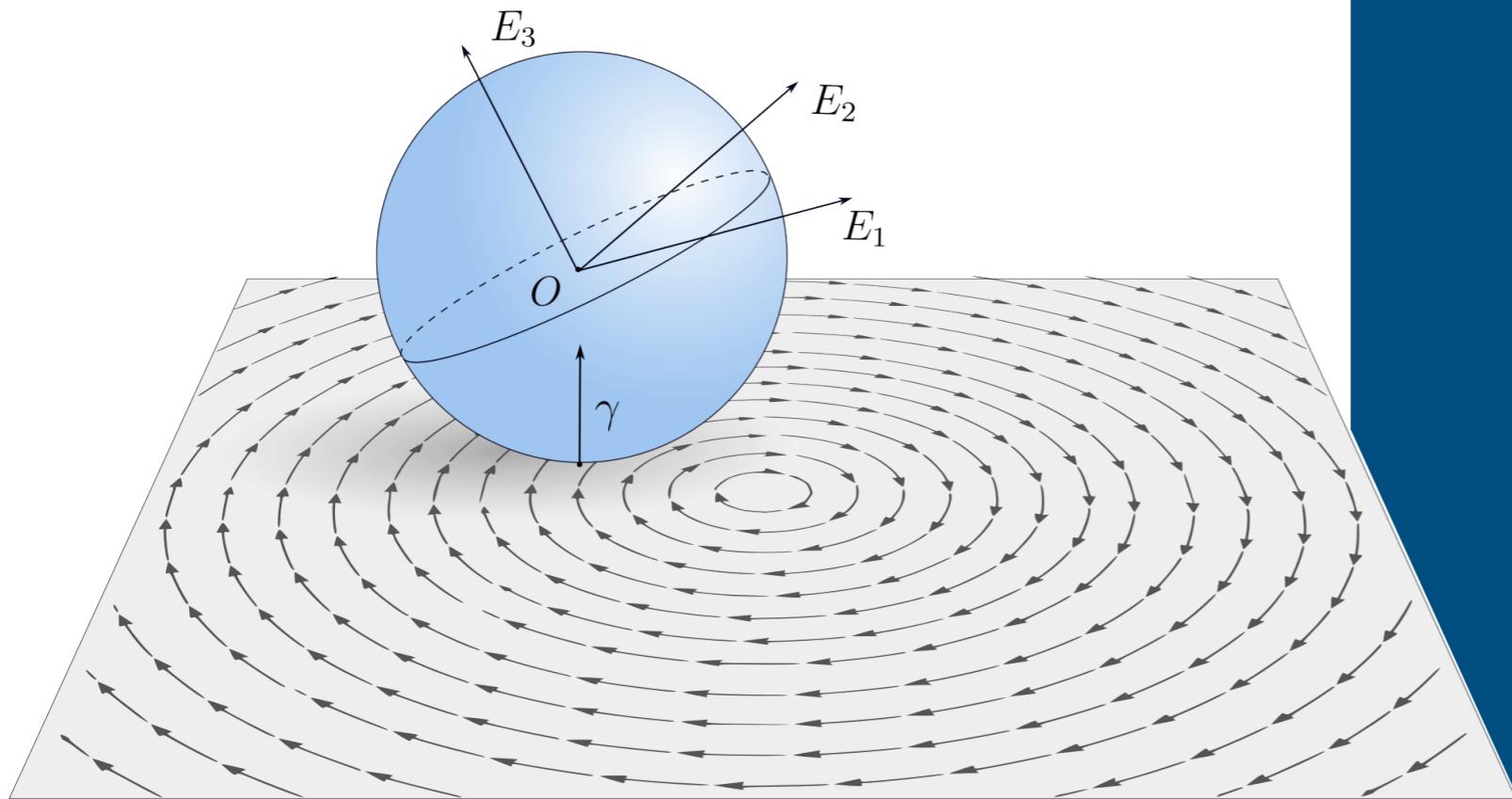
Homogeneous sphere

$$I_1 = I_2 = I_3, \quad O = C$$

Moves in straight line

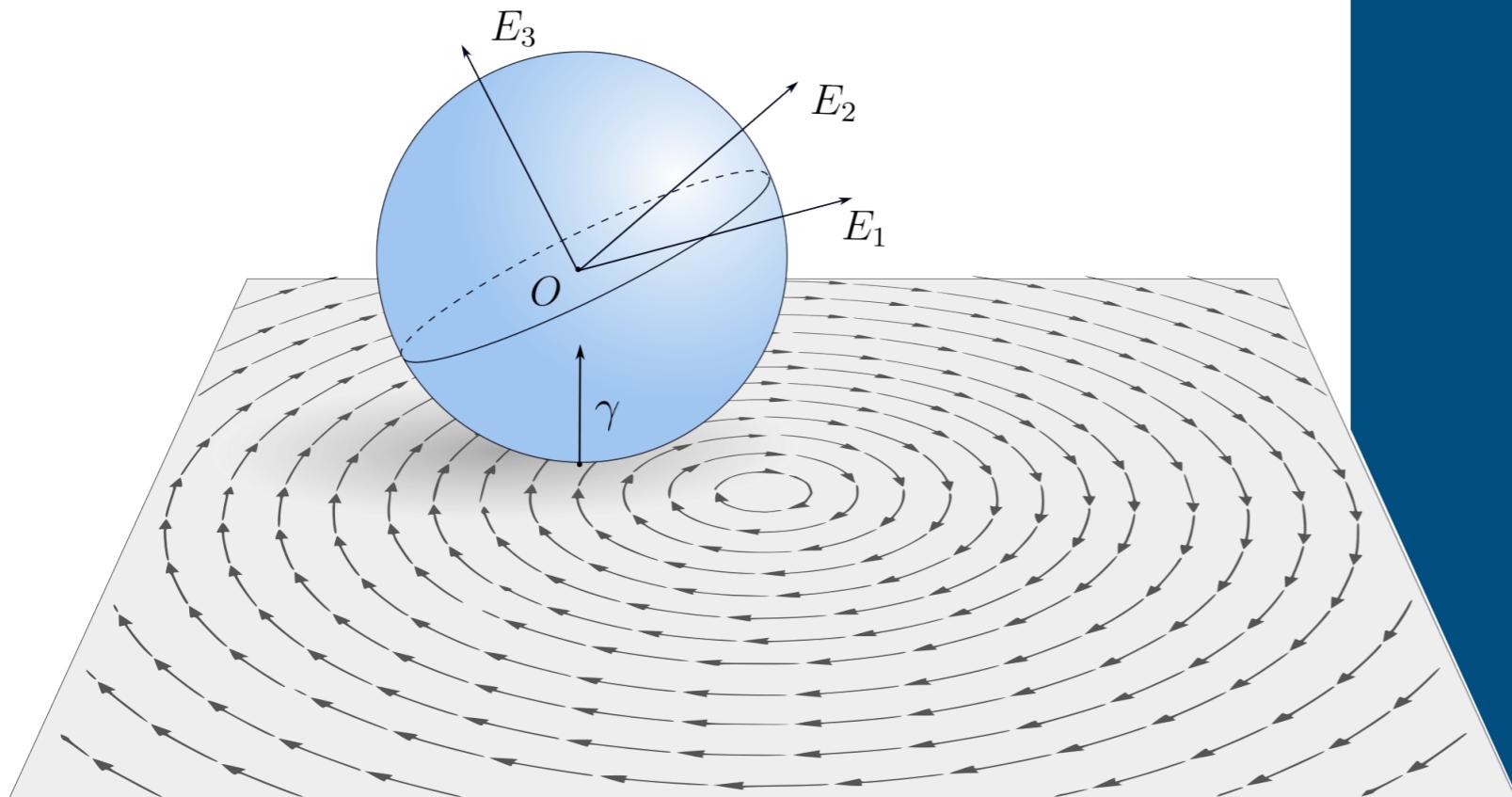


What happens if the plane  
is moving?

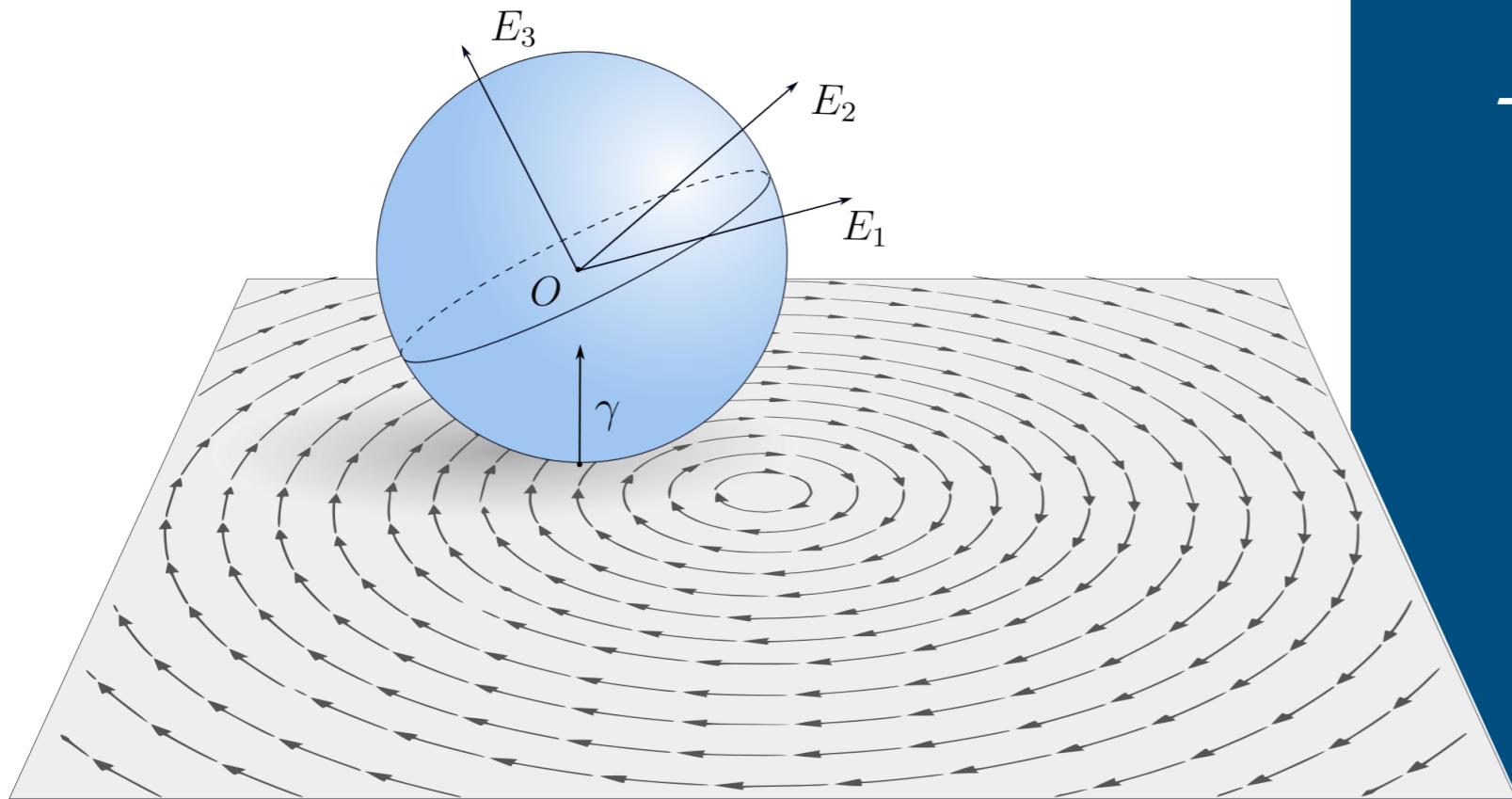


## What happens if the plane is moving?

- Lose symmetry on the plane
- Lose energy first integral



## What happens if the plane is moving?



- Lose symmetry on the plane
- Lose energy first integral
- Moving energy  
(Fassò, Sansonetto, 2015)

# Homogeneous sphere on rotating plane

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - \eta\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + \eta\gamma \times (\gamma \times U)$$

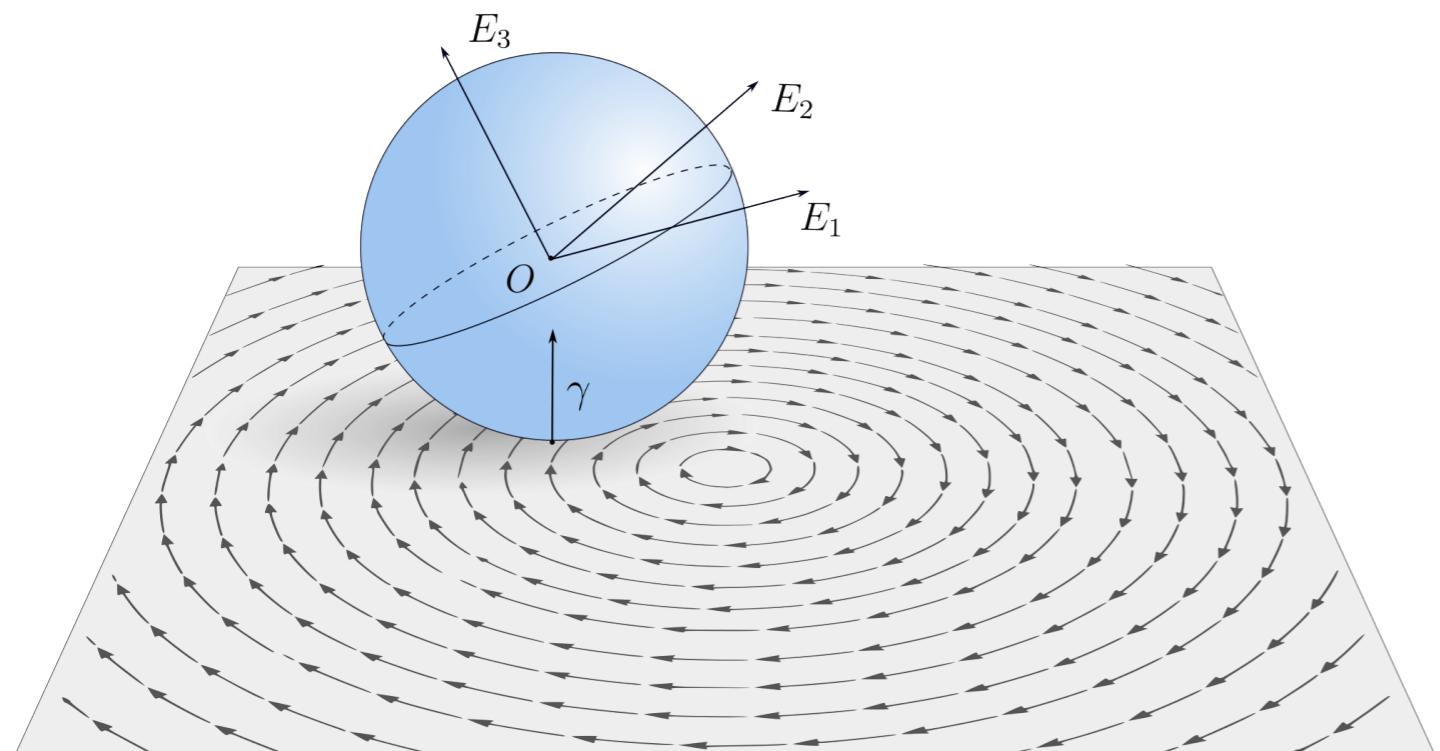
First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0,$$

$$E_{mov}, \quad \|\Omega + \eta U\|^2,$$

Invariant measure

$$dM dx dy d\gamma$$



# Homogeneous sphere on rotating plane

Integrable

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - \eta\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + \eta\gamma \times (\gamma \times U)$$

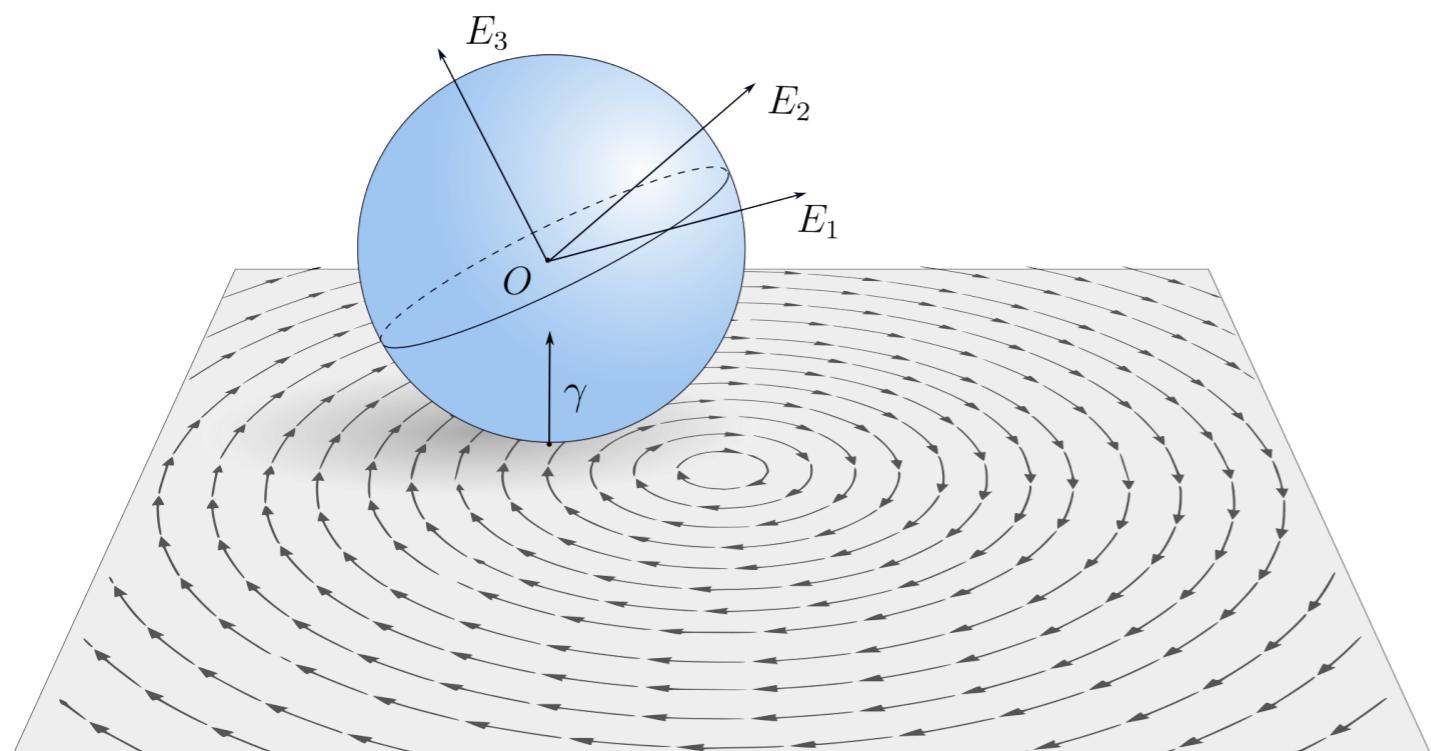
First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0,$$

$$E_{mov}, \quad \|\Omega + \eta U\|^2,$$

Invariant measure

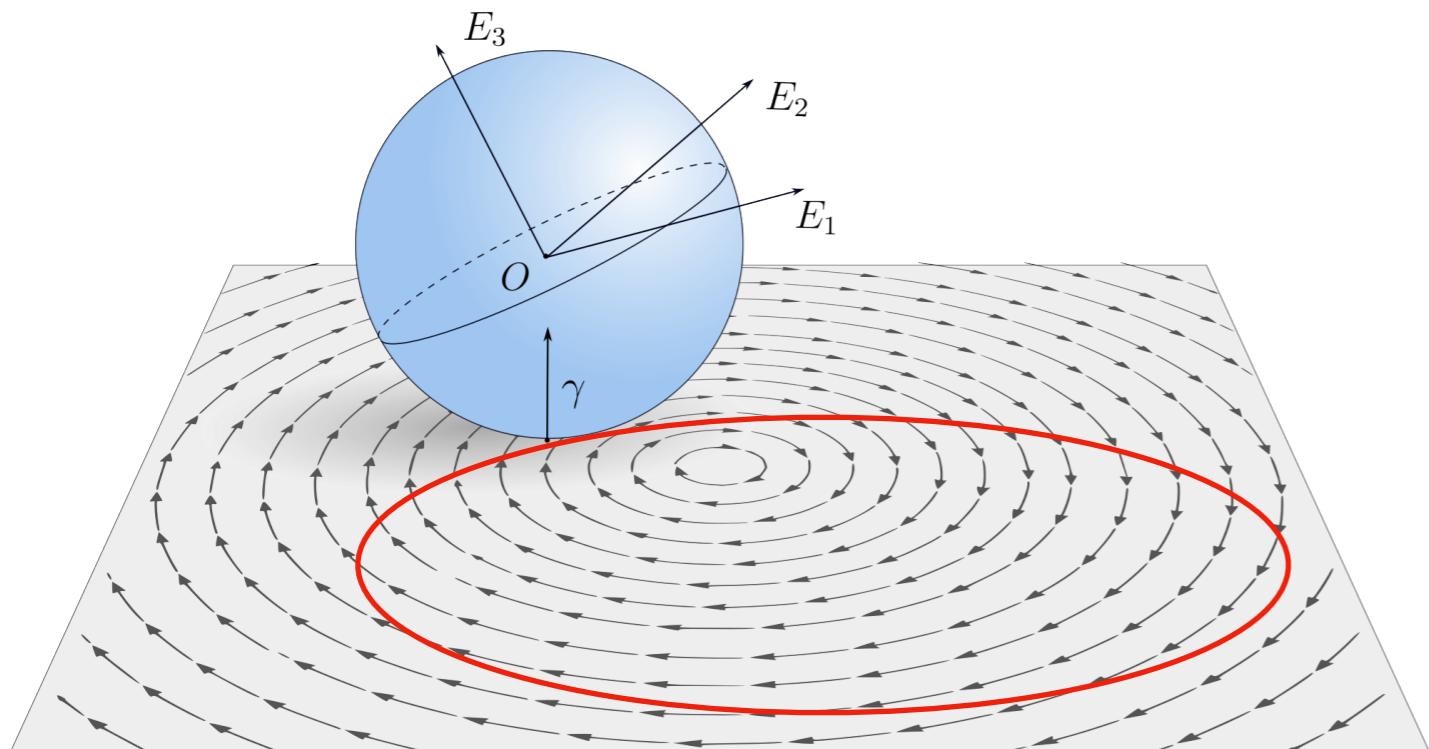
$$dM dx dy d\gamma$$



# Homogeneous sphere on rotating plane

Integrable

Moves in circle



Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - \eta\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + \eta\gamma \times (\gamma \times U)$$

First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0,$$

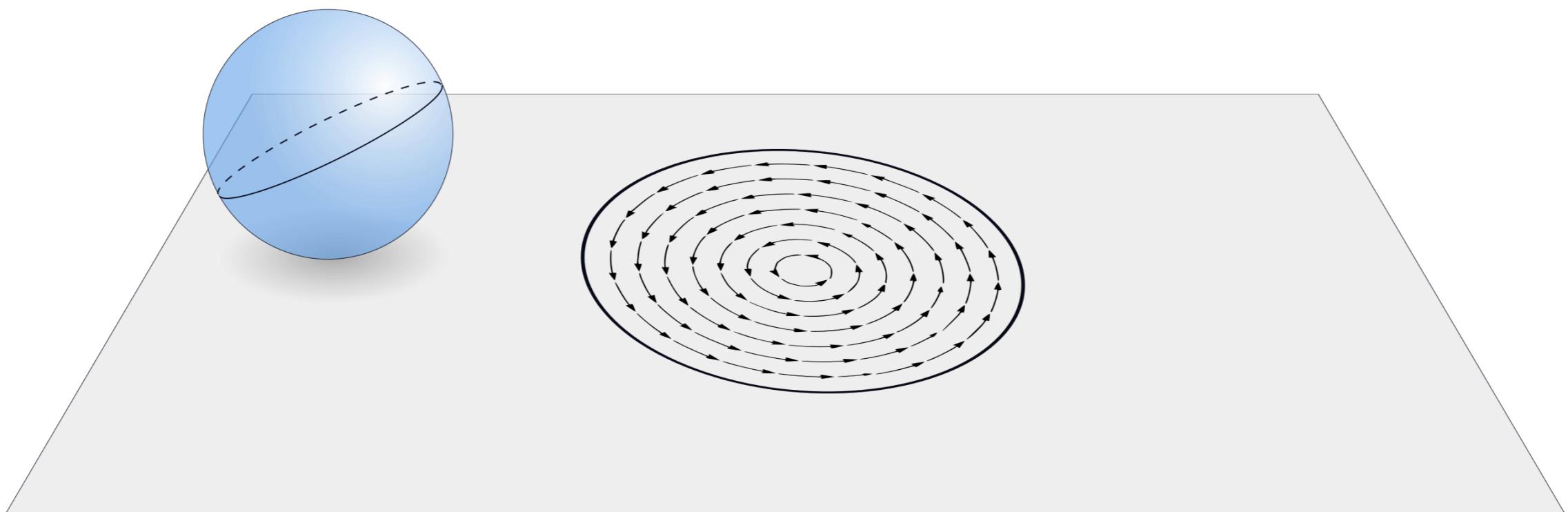
$$E_{mov}, \quad \|\Omega + \eta U\|^2,$$

Invariant measure

$$dM dx dy d\gamma$$

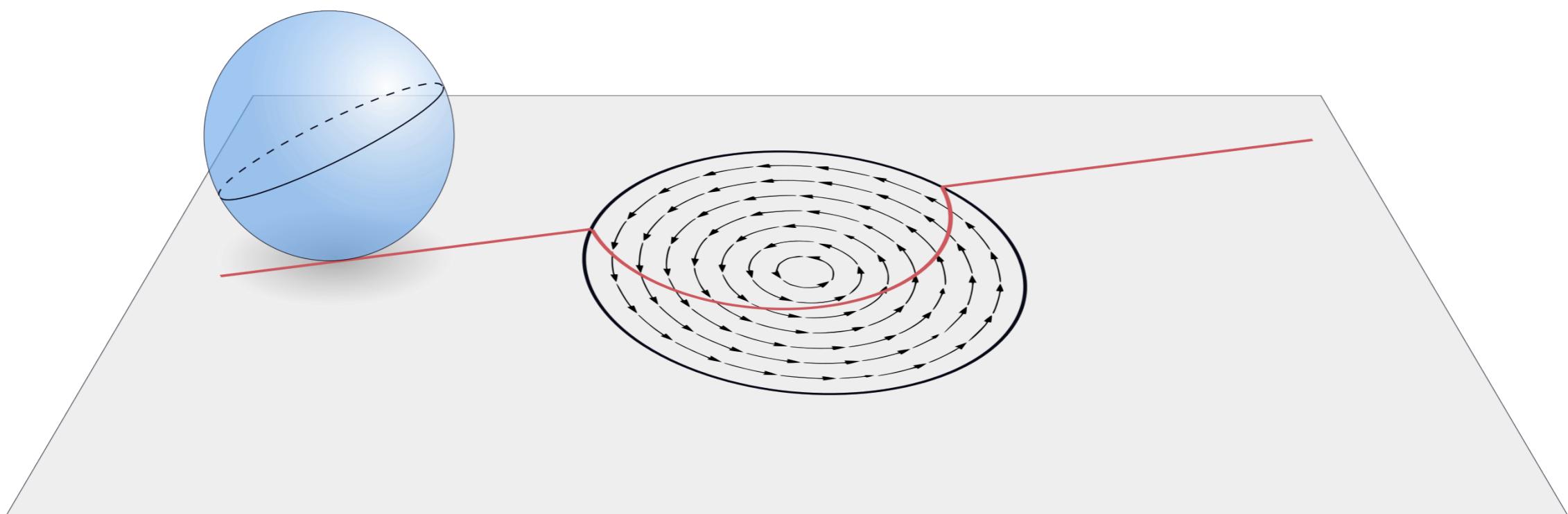
# Homogeneous sphere on rotating plane

Anais-billiard phenomenon  
(Lévy- Leblond, 1986)



# Homogeneous sphere on rotating plane

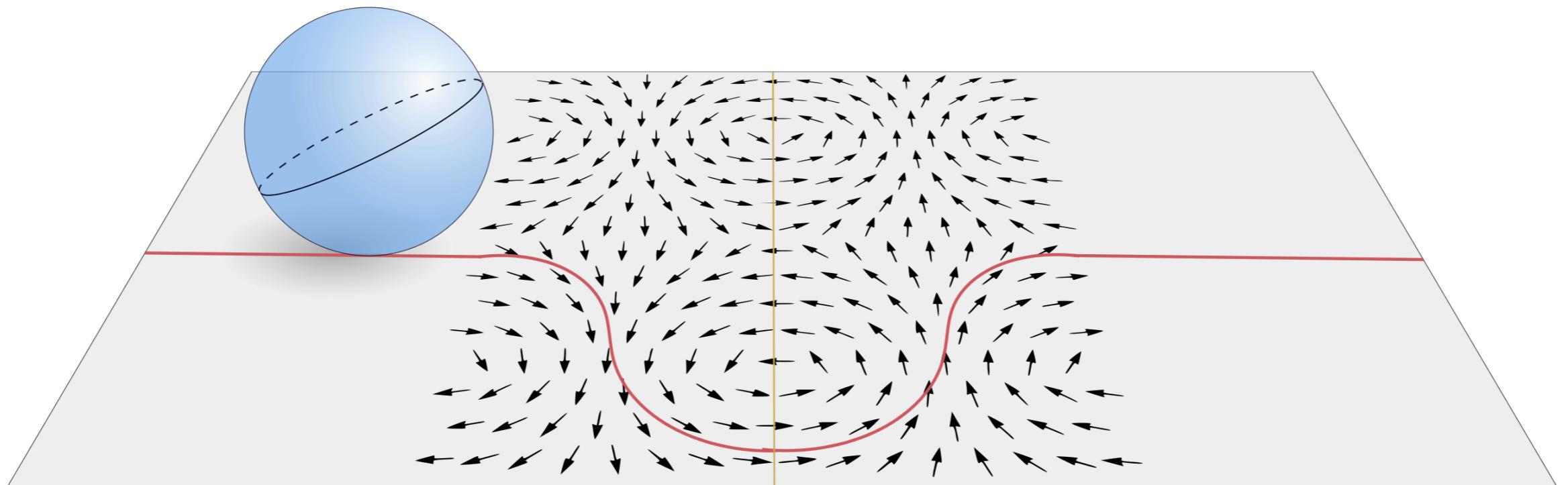
Anais-billiard phenomenon  
(Lévy- Leblond, 1986)



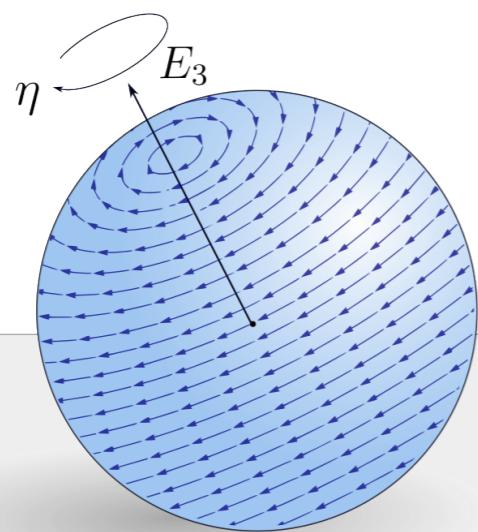
# Homogeneous sphere in plane with symmetric vector field

Anais-billiard phenomenon  
(Lévy- Leblond, 1986)

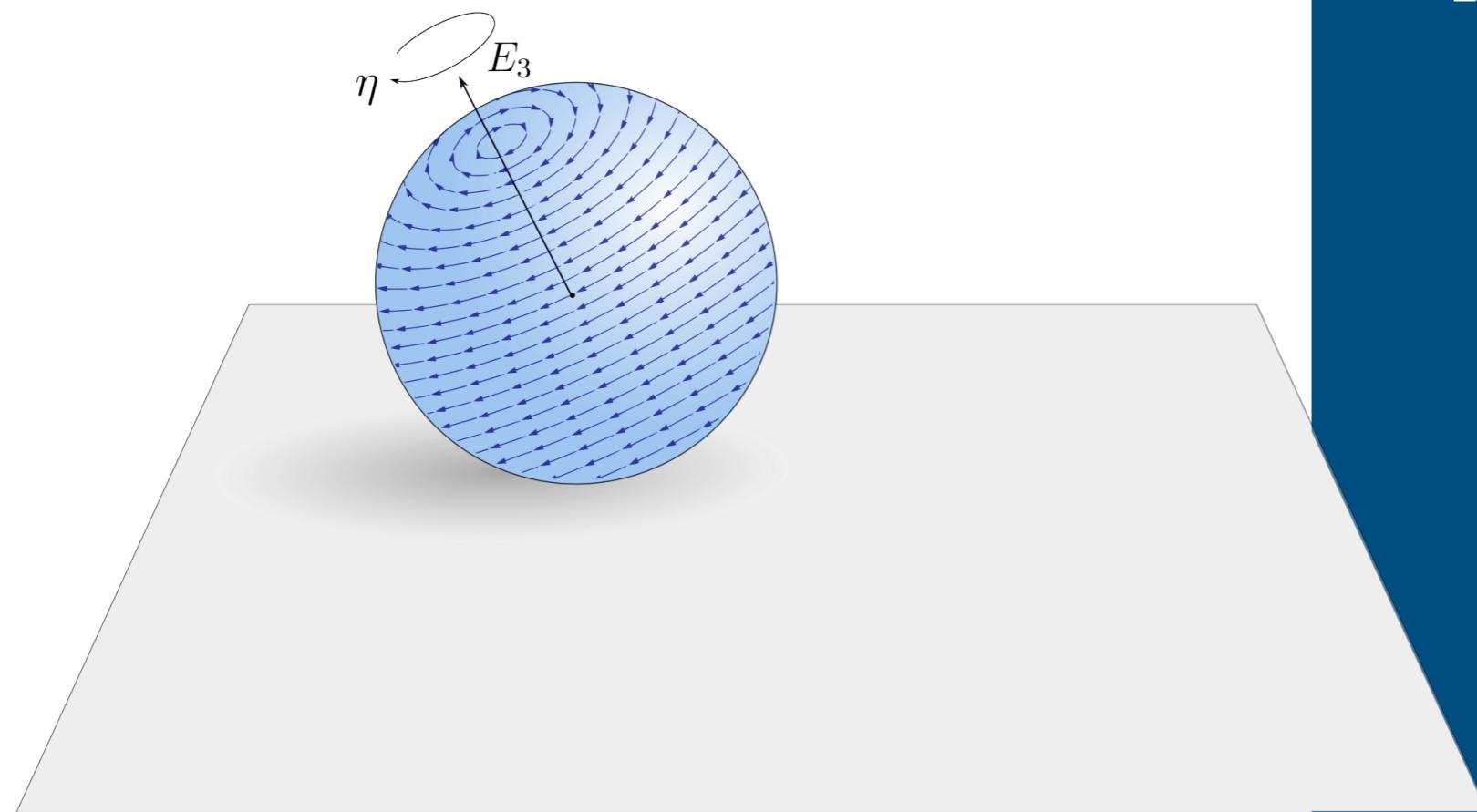
$$\begin{pmatrix} V_1(x, y) \\ V_2(x, y) \end{pmatrix} = \begin{pmatrix} V_1(-x, y) \\ -V_2(-x, y) \end{pmatrix}$$



What happens if the shell of the sphere is moving?

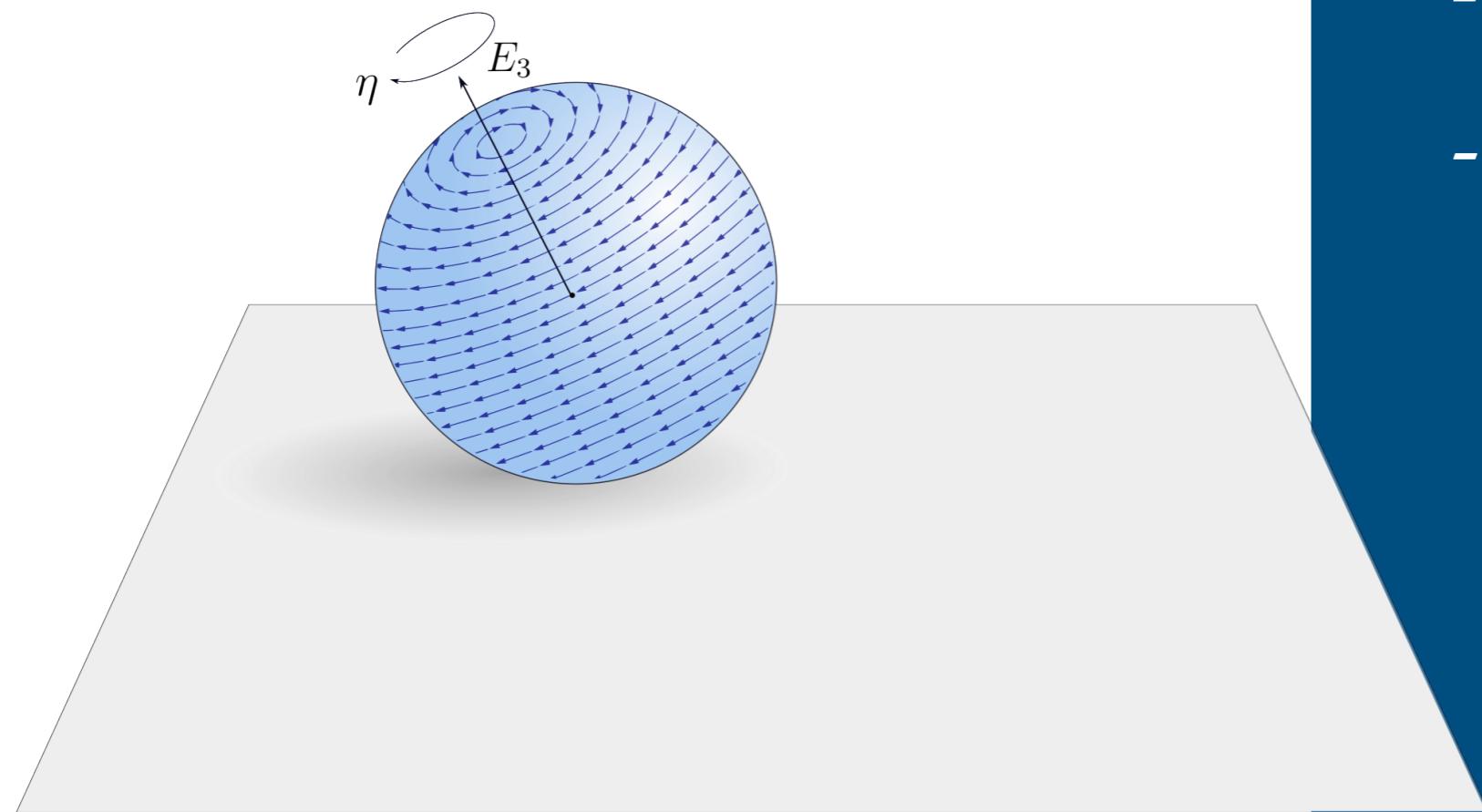


## What happens if the shell of the sphere is moving?



- Lose symmetry on the sphere
- Lose energy first integral

## What happens if the shell of the sphere is moving?



- Lose symmetry on the sphere
- Lose energy first integral
- Moving energy

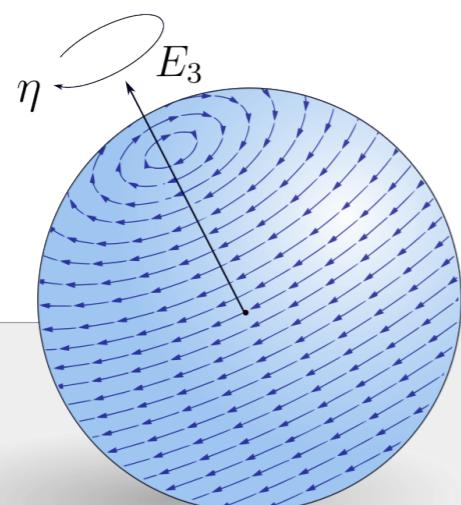
# Homogeneous sphere with rotating shell

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3)$$



First integrals:

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad E_{mov}$$

Invariant measure:

$$dM dx dy dd\gamma$$

# Homogeneous sphere with rotating shell

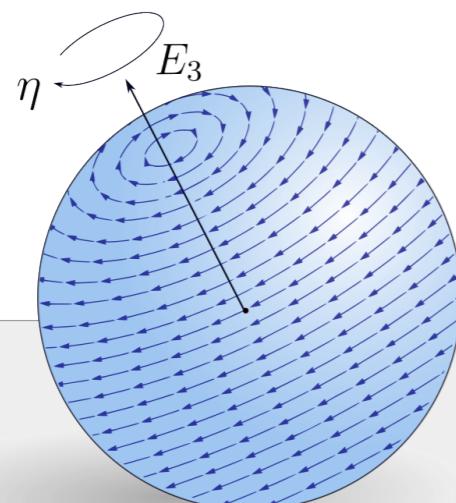
Integrable

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3)$$



First integrals:

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad E_{mov}$$

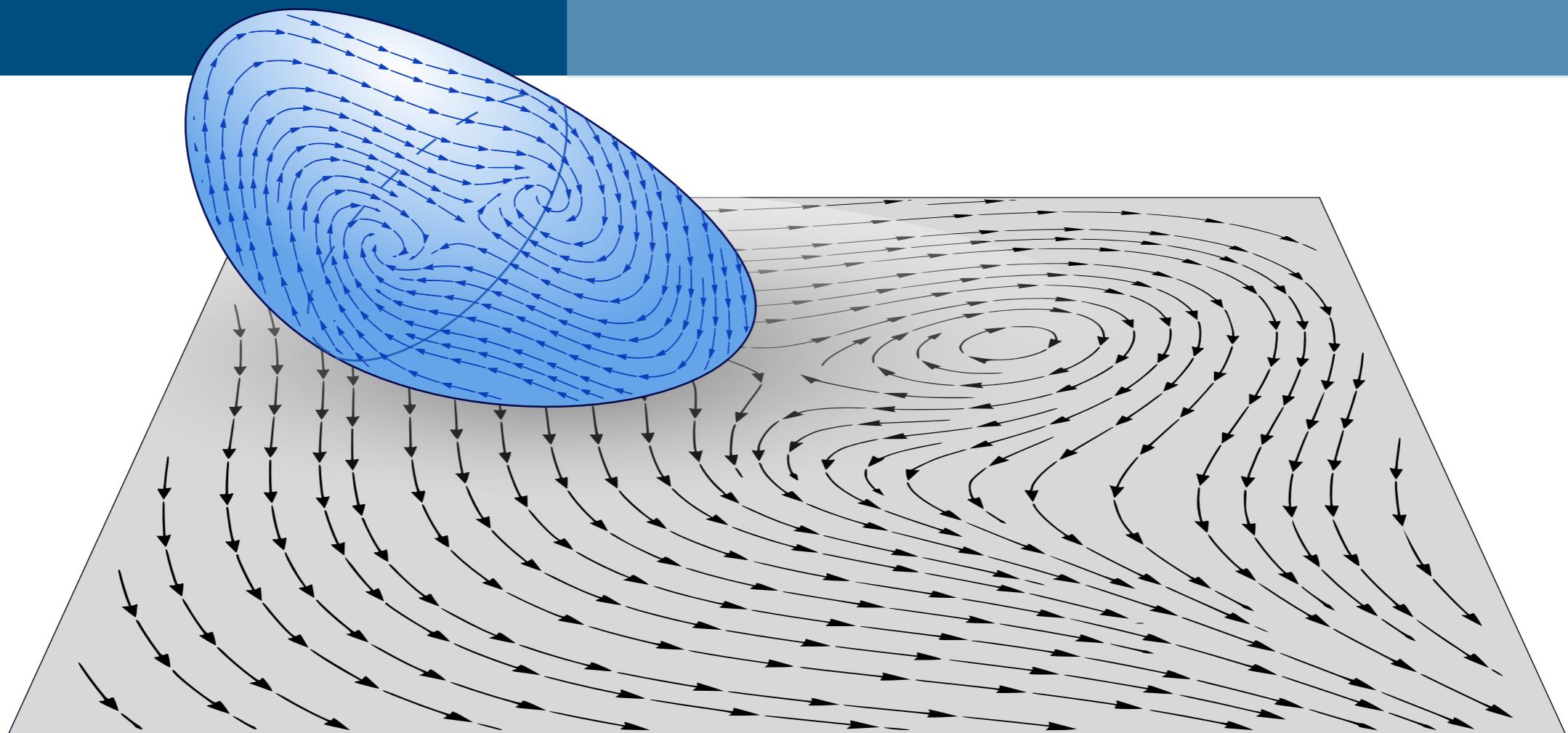
Invariant measure:

$$dM dx dy dd\gamma$$

# Generalization

Convex body rolling on a plane

Add general vector fields to plane and body shell

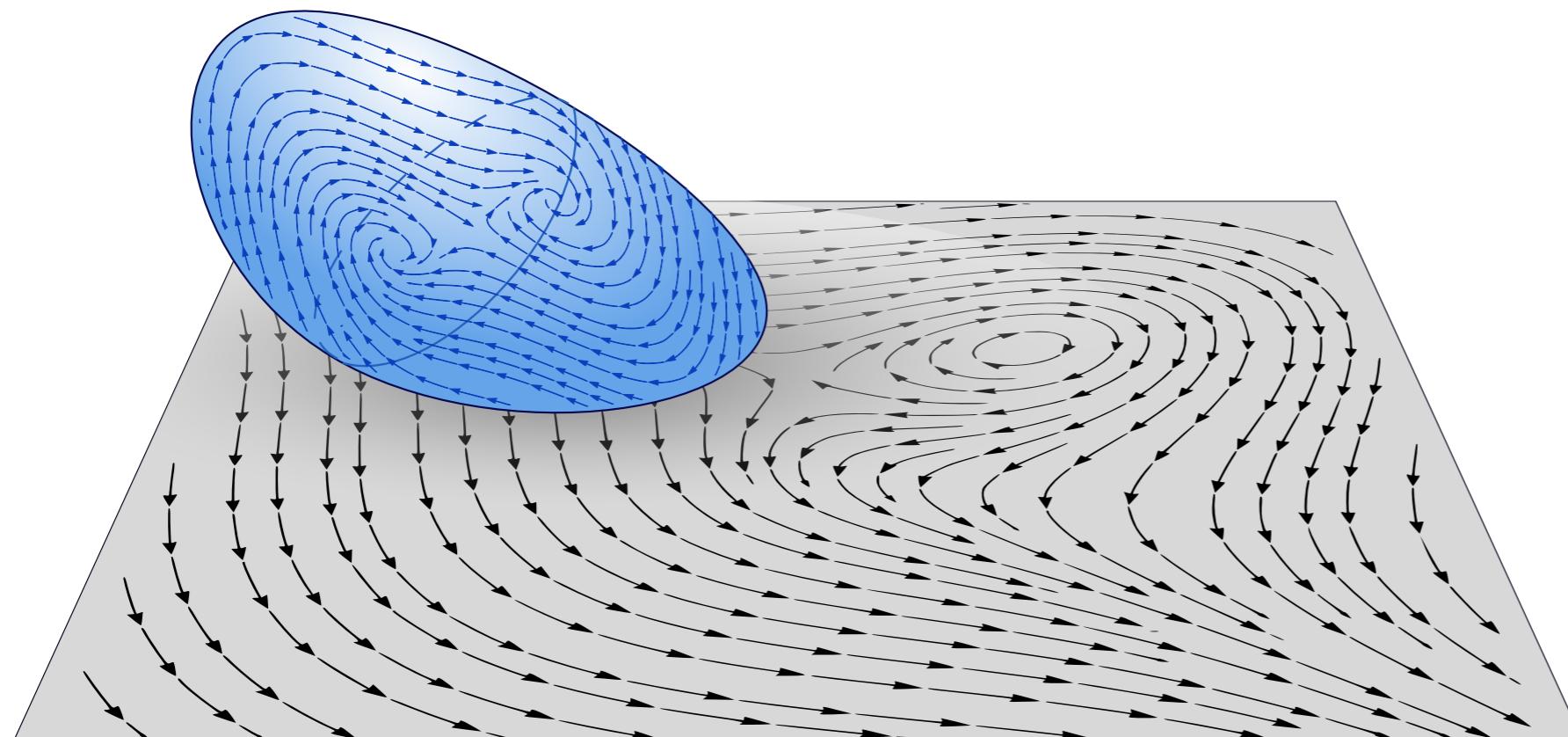


# Generalization

## Goal:

Identify general principles leading to

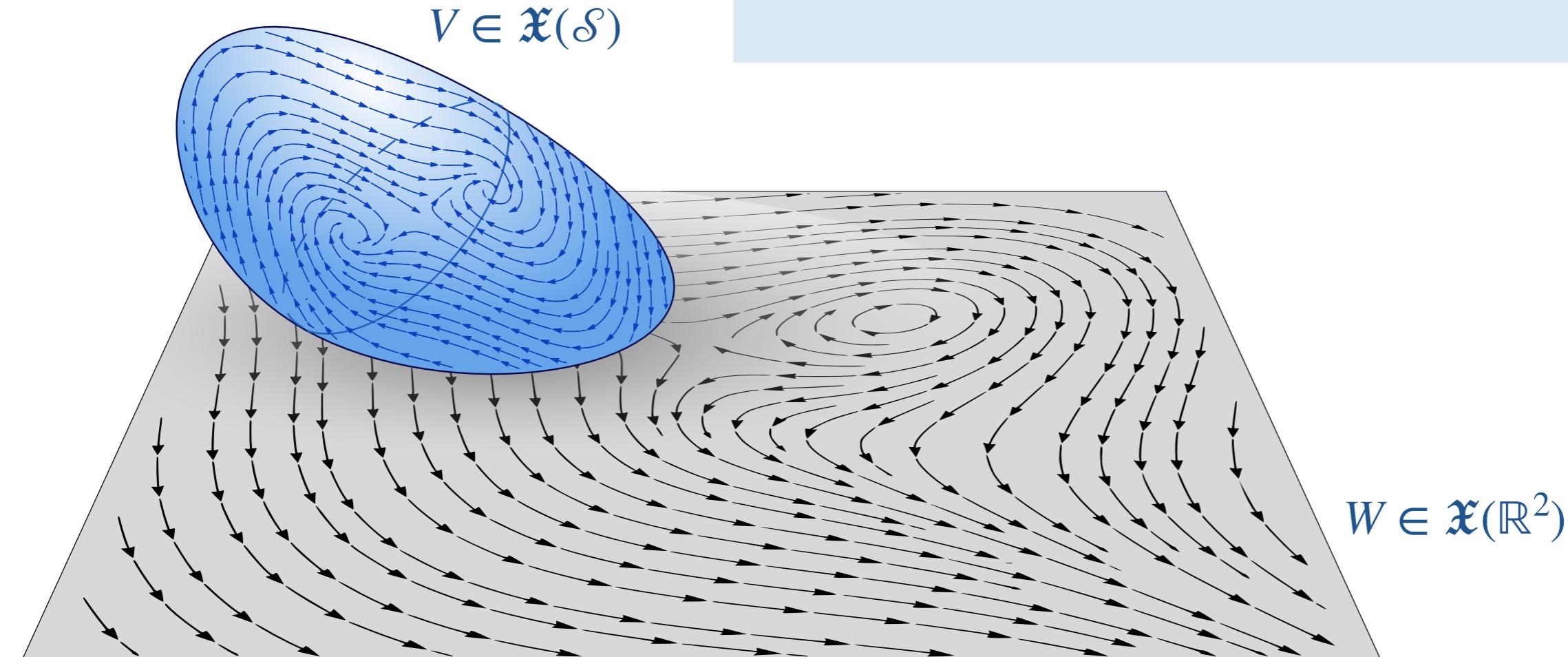
- First integrals
- Moving energy
- Invariant measure



# Generalization

Configuration manifold:  $Q = SO(3) \times \mathbb{R}^2$   
 $(B, \underline{x}) \in Q$

Affine nonholonomic constraint:  
 $\dot{u} = B(\rho \times \Omega) + BV + W$

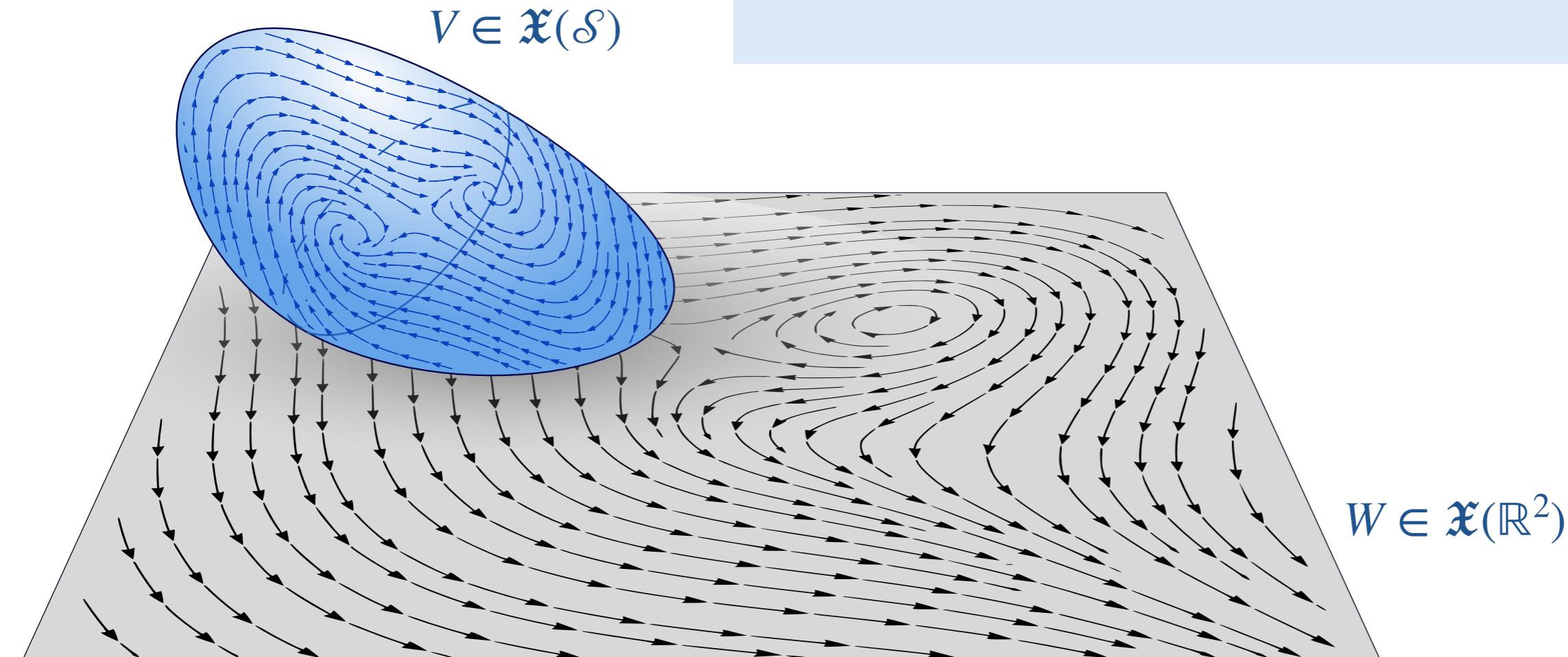


# Generalization

No symmetries on  
plane or body

Configuration manifold:  $Q = SO(3) \times \mathbb{R}^2$   
 $(B, \underline{x}) \in Q$

Affine nonholonomic constraint:  
 $\dot{u} = B(\rho \times \Omega) + BV + W$



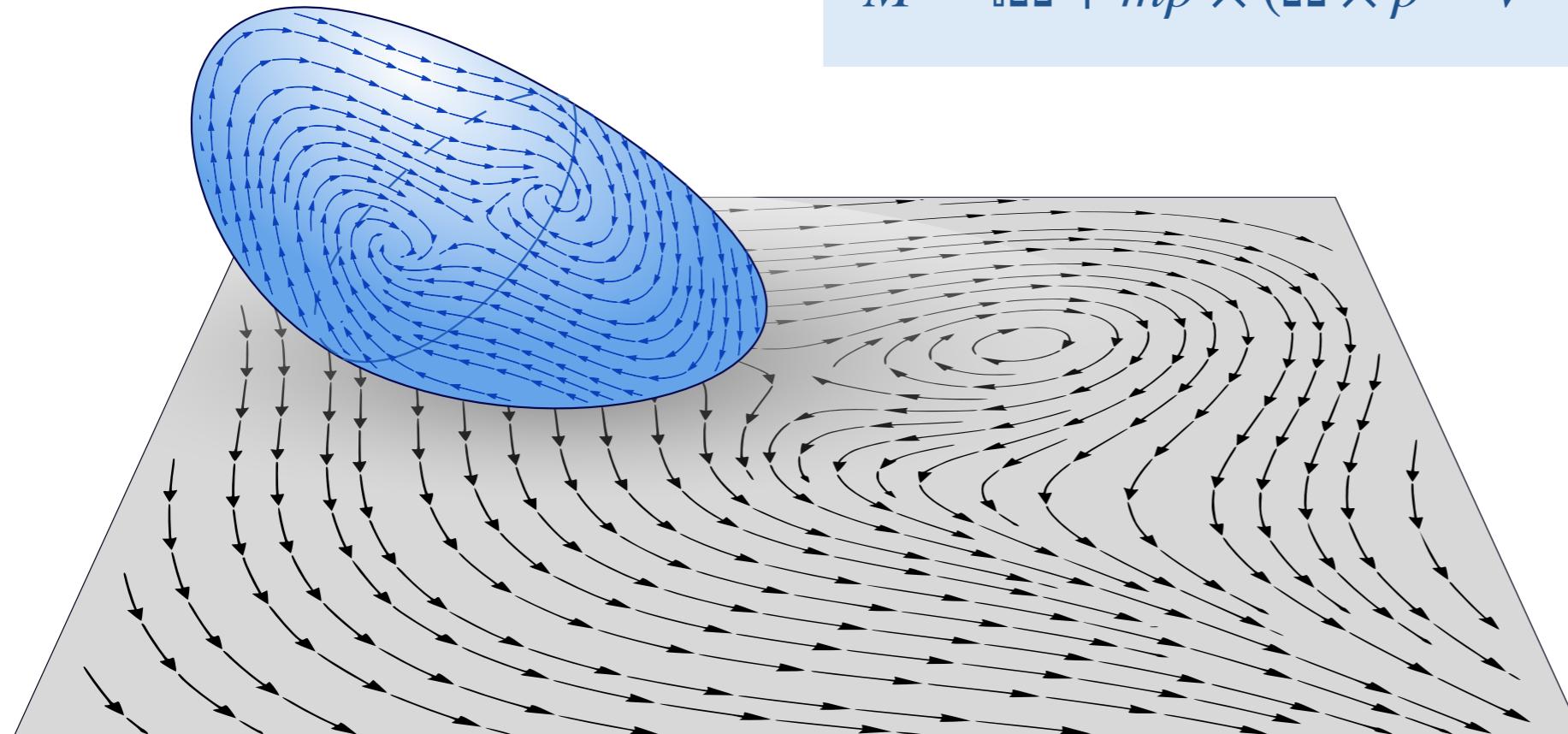
# General system

Theorem (C., García Naranjo, 2023):

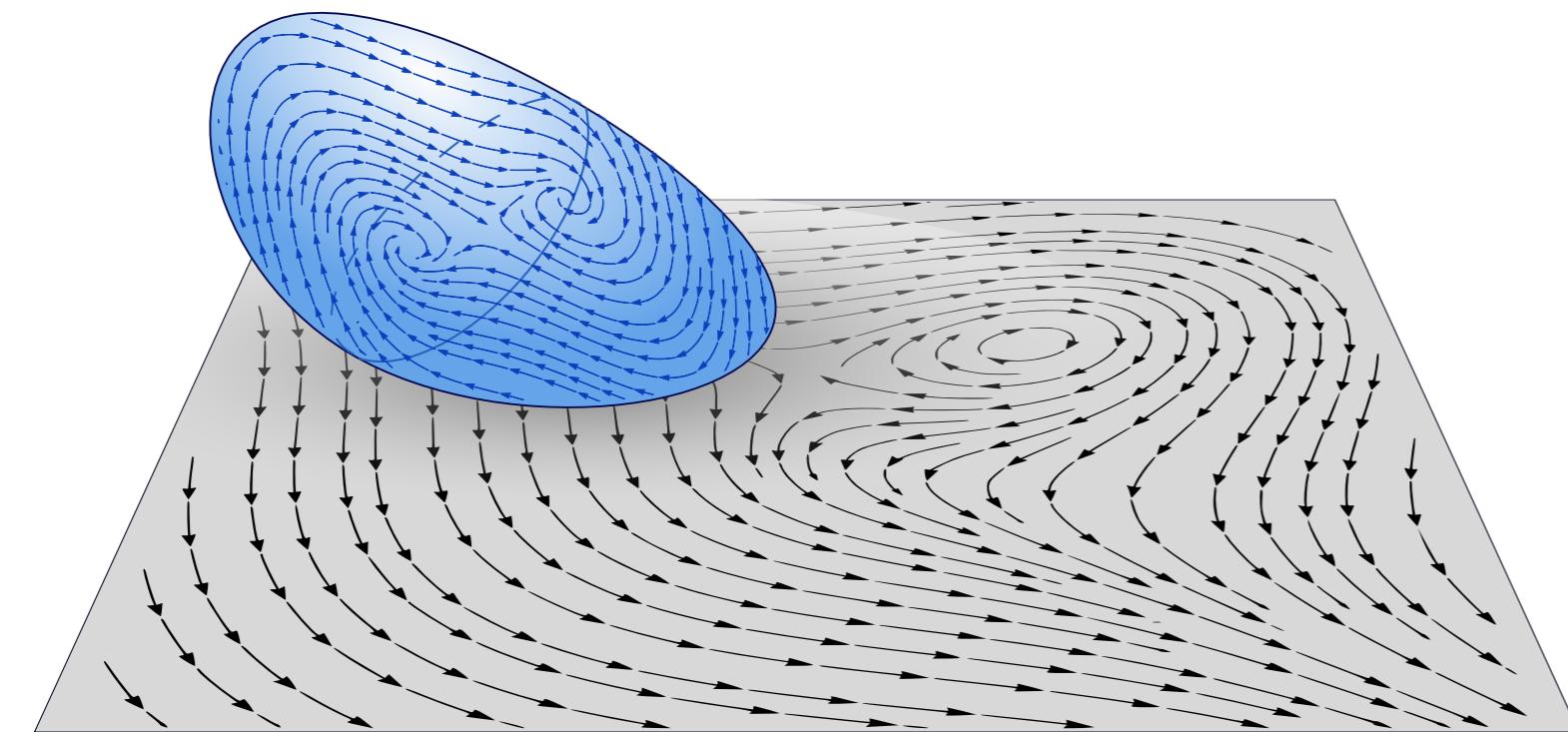
Equations of motion :

$$\begin{aligned}\dot{M} &= M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma \\ &\quad + m(V + B^{-1}W) \times (\dot{\rho} + \Omega \times \rho) \\ \dot{u} &= B(\rho \times \Omega) + BV + W \\ \dot{B} &= B\hat{\Omega}\end{aligned}$$

$$M = \mathbb{I}\Omega + m\rho \times (\Omega \times \rho - V - B^{-1}W)$$

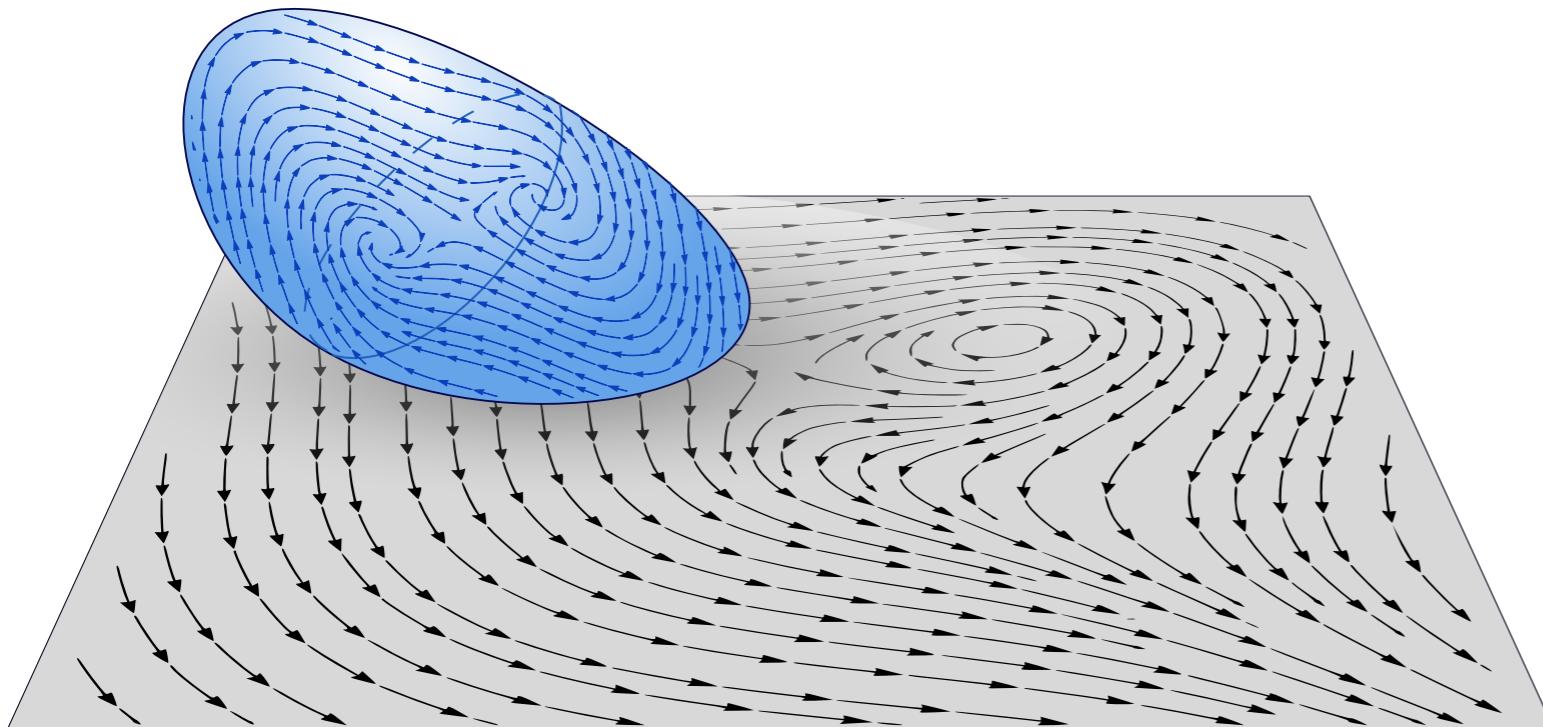


# General system



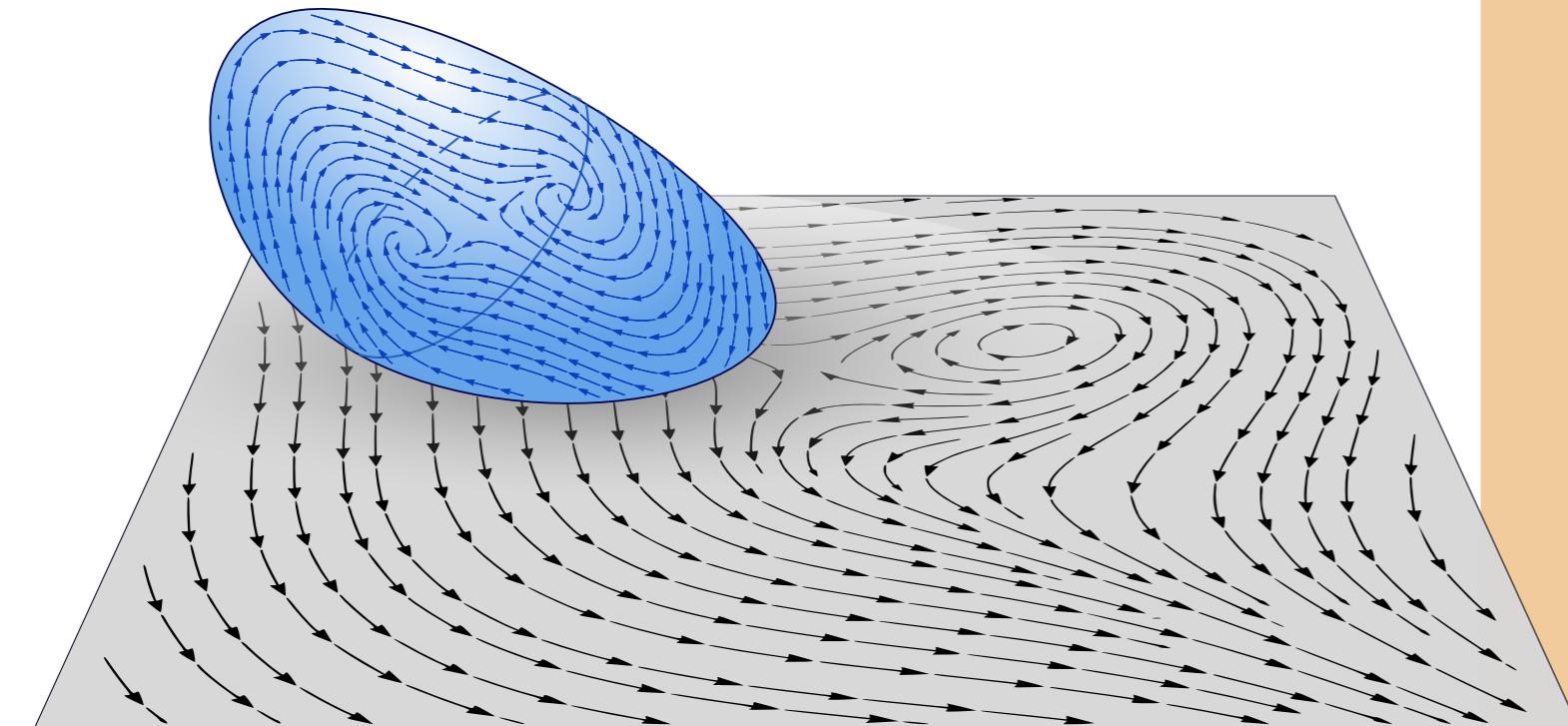
# General system

- No invariant measure
- Generally chaotic



# General system

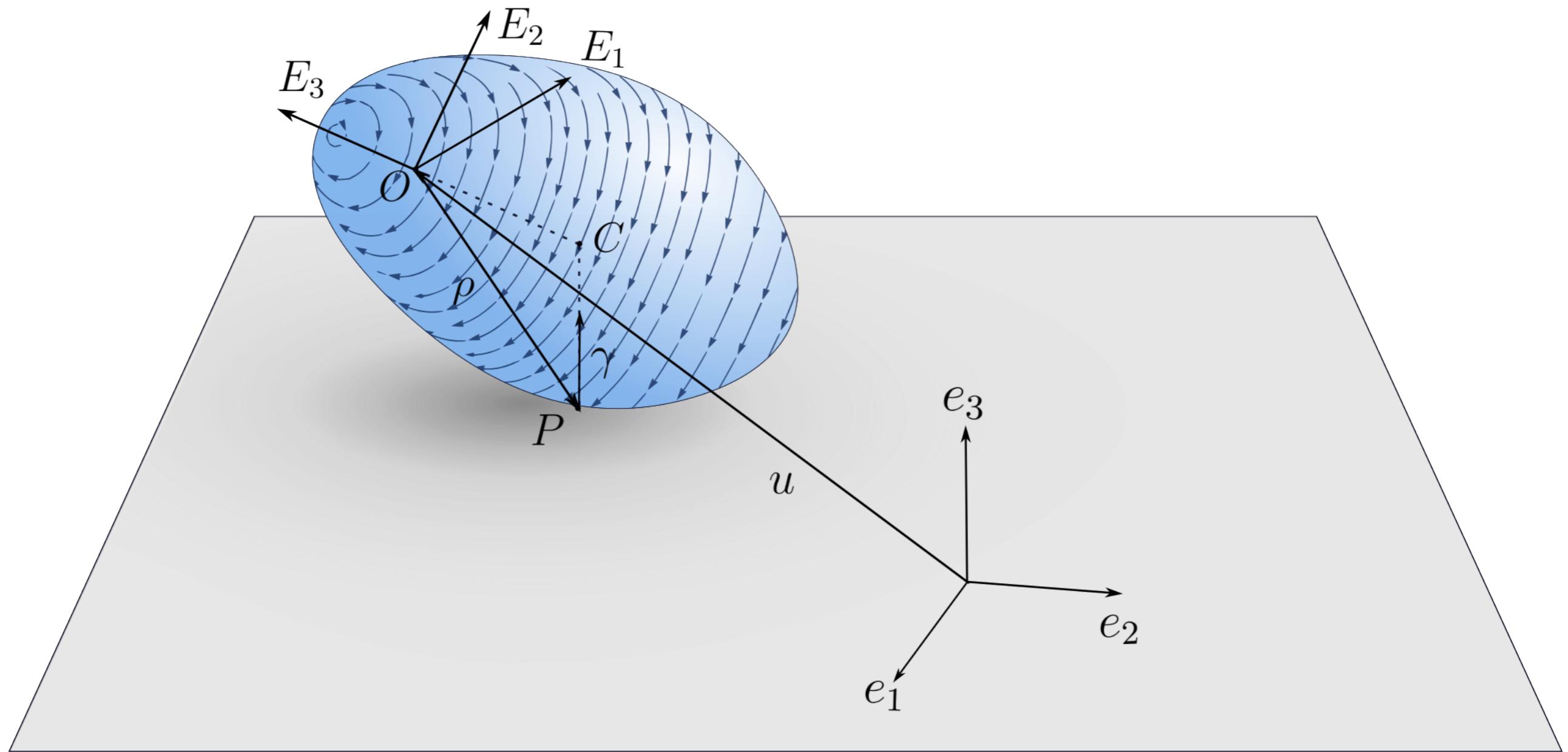
- No invariant measure
- Generally chaotic



For it to be integrable we need extra symmetries that lead to a sufficient number of:

- First integrals
- Invariant measure

# Body of revolution with rotating shell



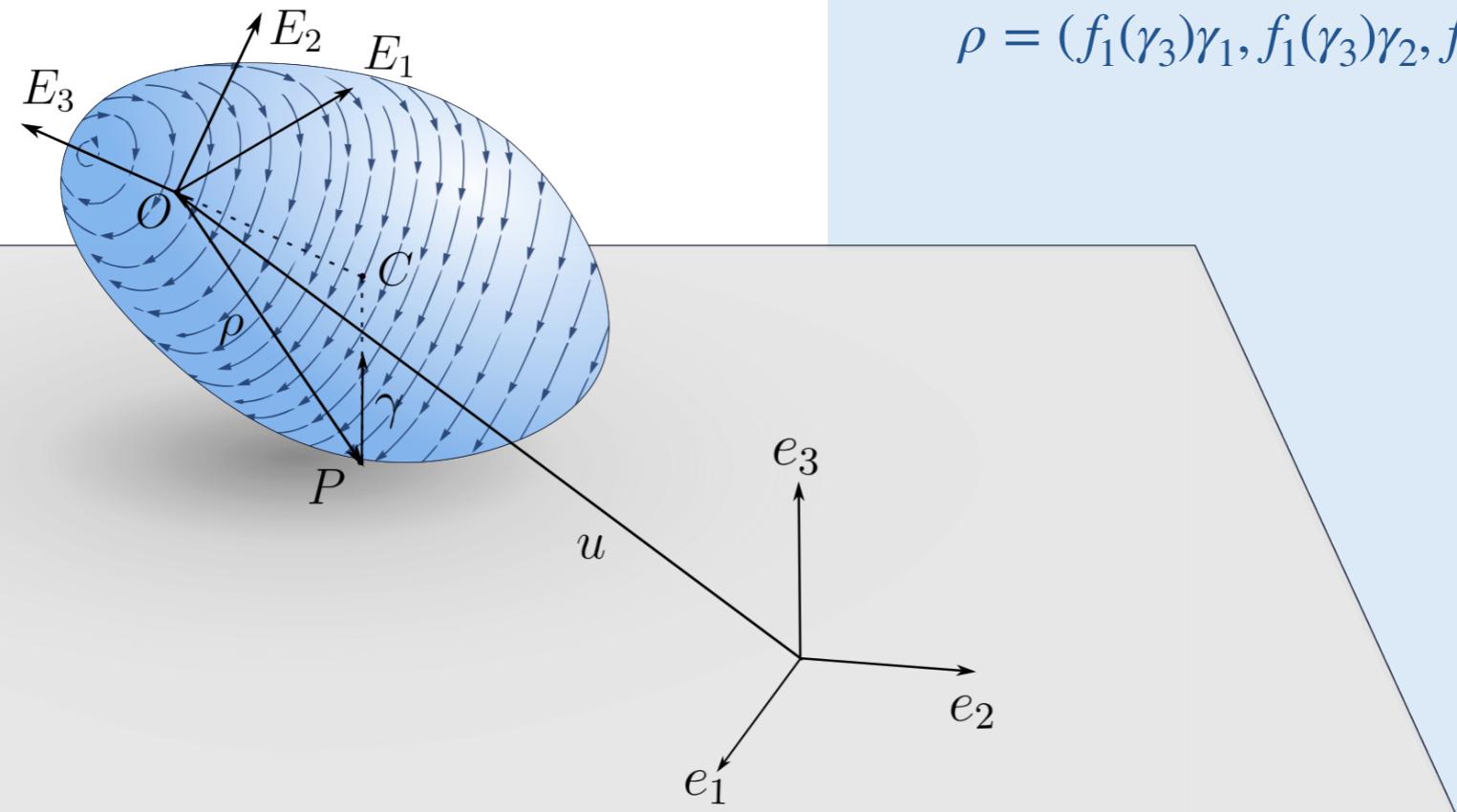
# Body of revolution with rotating shell

Equations of motion:

$$\begin{aligned}\dot{M} &= M \times \Omega + m\dot{\rho} \times (\Omega \times \rho) + mg\rho \times \gamma \\ &\quad + m\eta(\rho \times E_3) \times (\dot{\rho} + \Omega \times \rho) \\ \dot{\gamma} &= \gamma \times \Omega\end{aligned}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\eta\gamma \times (\rho \times E_3)$$

$$\rho = (f_1(\gamma_3)\gamma_1, f_1(\gamma_3)\gamma_2, f_2(\gamma_3))$$



# Body of revolution with rotating shell

First integrals:

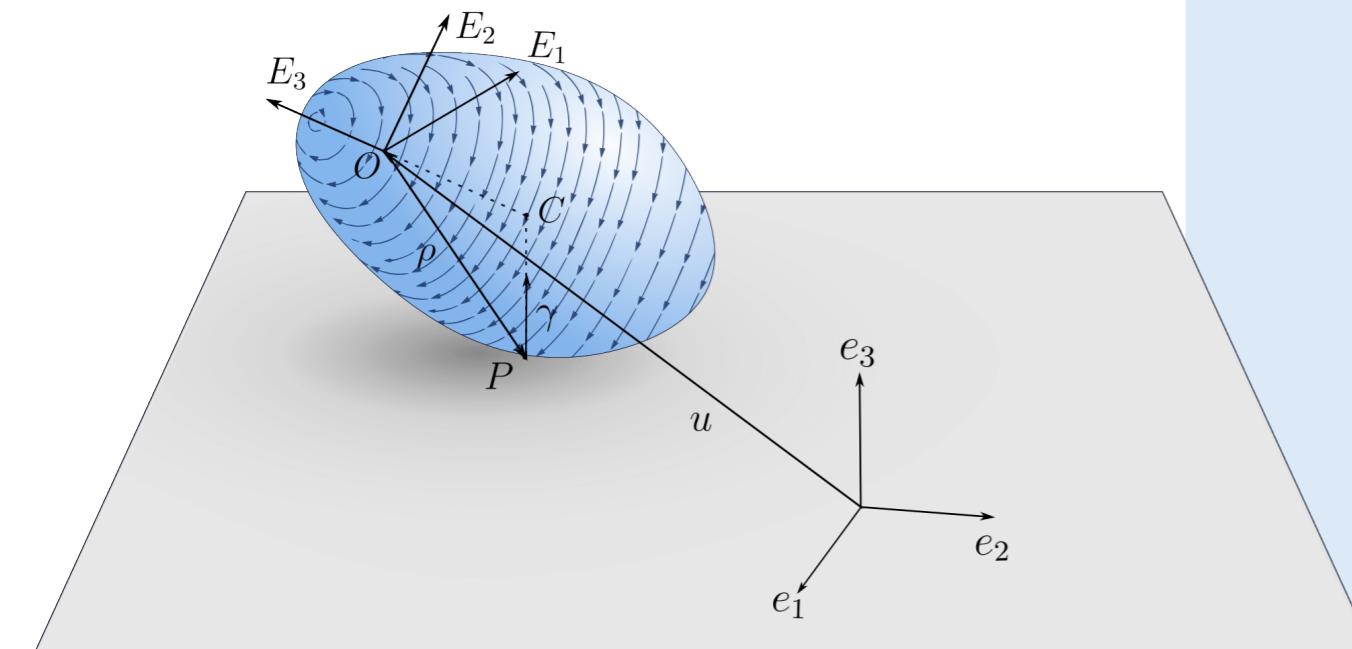
$$\begin{pmatrix} J_1 \\ J_2 \end{pmatrix} = U^{-1} \left( \begin{pmatrix} K_1 \\ K_2 \end{pmatrix} - \eta u \right),$$

$$K_1 = \frac{\langle M, \rho \rangle}{f_1}, \quad K_2 = \frac{\Omega_3}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}},$$

$U$  solution matrix of  $U' = G(\gamma_3)U, \quad U(0) = Id$

$u$  solution of  $u' = Gu + b.$

$$G = G(\gamma_3), b = b(\gamma_3)$$



# Body of revolution with rotating shell

Theorem (C., García Naranjo, 2023):

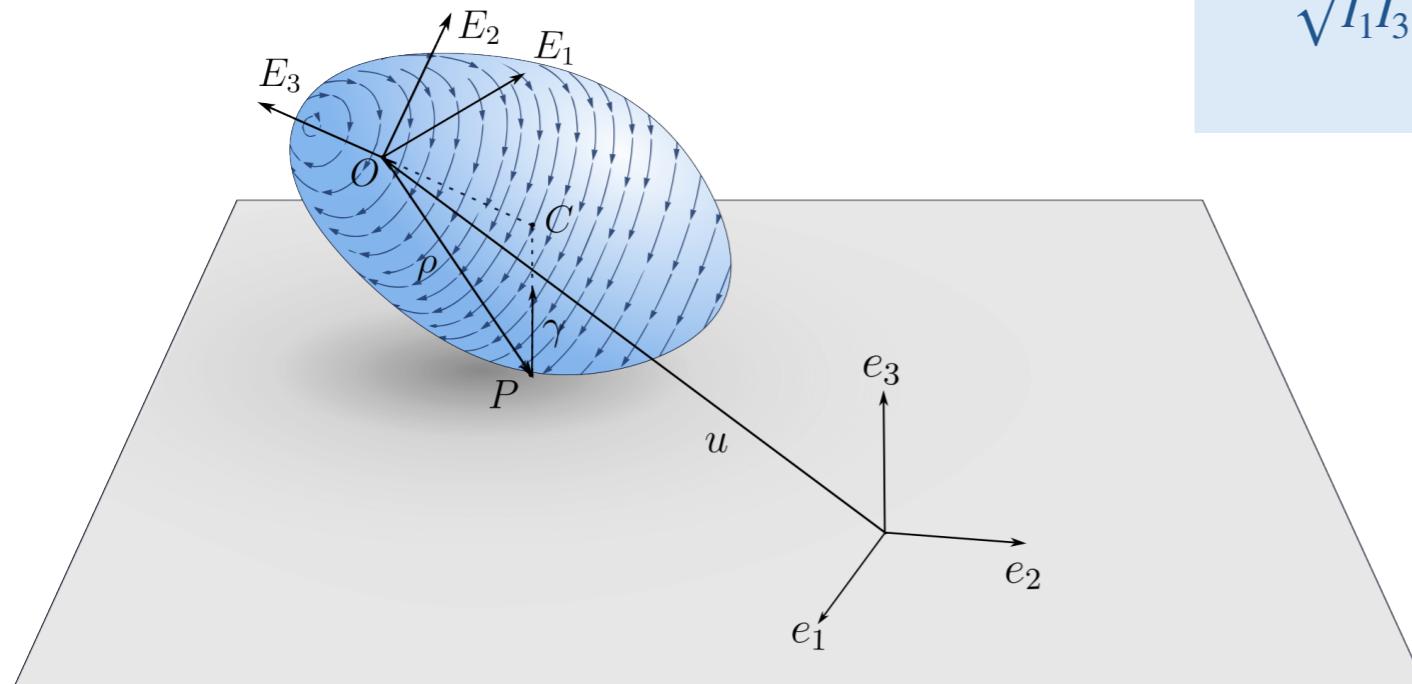
First integrals:

$$||\gamma||^2, J_1, J_2$$

$$\begin{aligned} E_{mov} = & \frac{1}{2} \langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle \\ & + \frac{m}{2} \|\rho \times (\Omega + \eta E_3)\|^2 - mr \langle \rho, \gamma \rangle \end{aligned}$$

Invariant measure:

$$\frac{1}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}} dM d\gamma$$



# Body of revolution with rotating shell

Integrable

Theorem (C., García Naranjo, 2023):

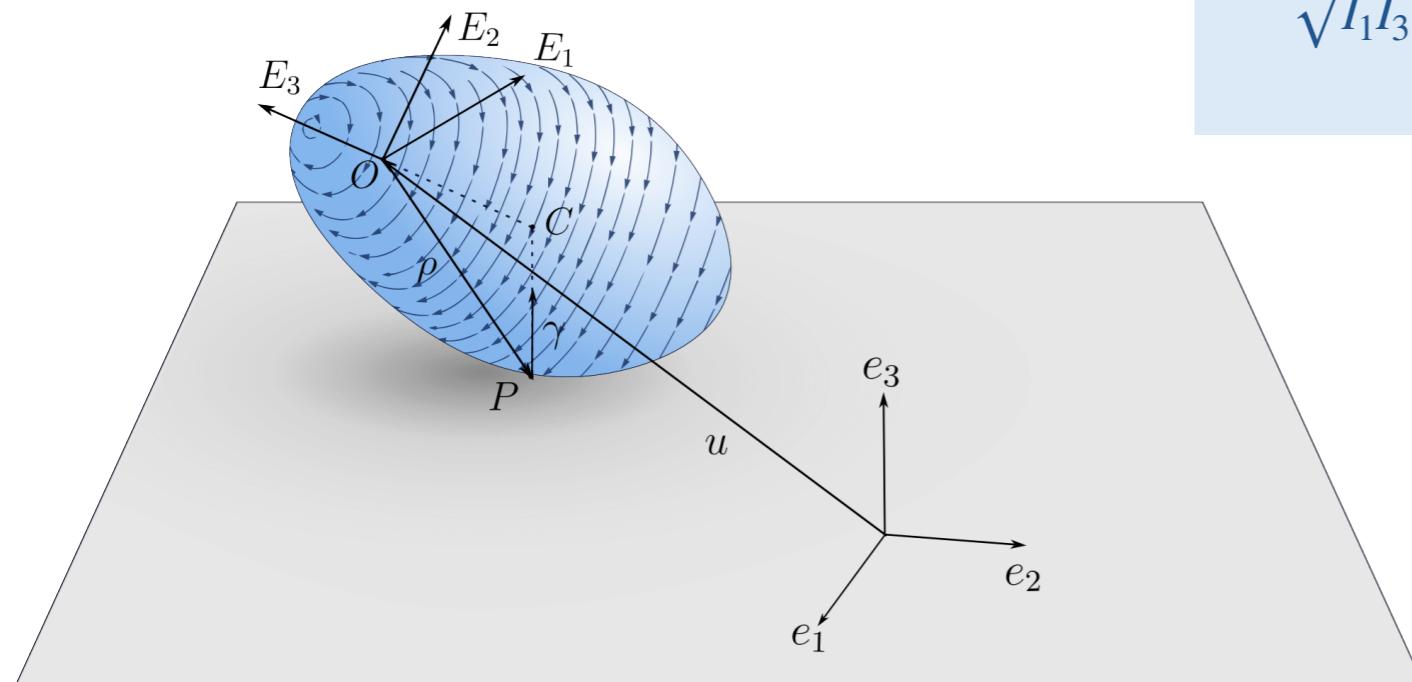
First integrals:

$$||\gamma||^2, J_1, J_2$$

$$\begin{aligned} E_{mov} = & \frac{1}{2} \langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle \\ & + \frac{m}{2} \|\rho \times (\Omega + \eta E_3)\|^2 - mr \langle \rho, \gamma \rangle \end{aligned}$$

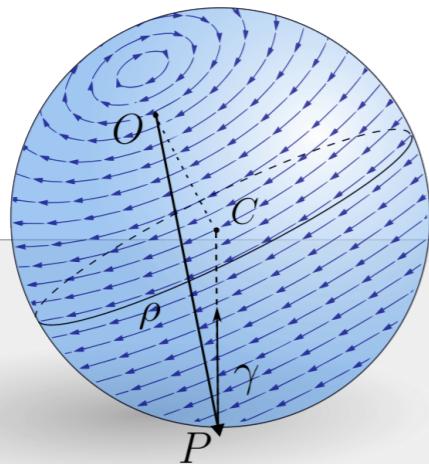
Invariant measure:

$$\frac{1}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}} dM d\gamma$$



# Routh's sphere with rotating shell

$$\rho(\gamma) = -r\gamma - lE_3$$



First integrals

$$||\gamma||^2,$$

$$J_1 = \frac{1}{r} \langle M, \rho \rangle,$$

$$J_2 = \Omega_3 + \eta \left( \frac{I_1 I_3 \sqrt{m}}{(I_1 - I_3)^{3/2}} \arctan \left( \frac{\mu'}{r \sqrt{m(I_1 - I_3)}} \right) - \frac{I_1}{I_1 - I_3} \mu \right),$$

$$\mu = \sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}$$

$$E_{mov} = \frac{1}{2} \langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle + \frac{m}{2} ||\rho \times (\Omega + \eta E_3)||^2 - mr \langle \rho, \gamma \rangle$$

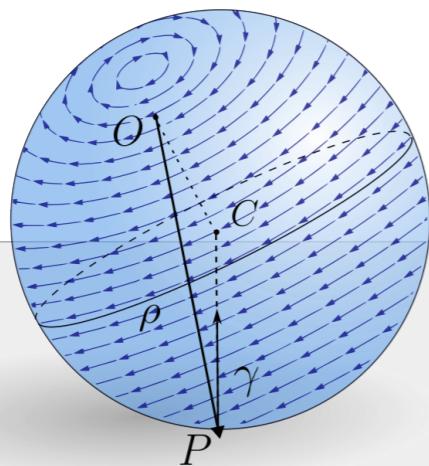
Invariant measure

$$\frac{1}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}} dM d\gamma$$

# Routh's sphere with rotating shell

Integrable

$$\rho(\gamma) = -r\gamma - lE_3$$



First integrals

$$||\gamma||^2,$$

$$J_1 = \frac{1}{r} \langle M, \rho \rangle,$$

$$J_2 = \Omega_3 + \eta \left( \frac{I_1 I_3 \sqrt{m}}{(I_1 - I_3)^{3/2}} \arctan \left( \frac{\mu'}{r \sqrt{m(I_1 - I_3)}} \right) - \frac{I_1}{I_1 - I_3} \mu \right),$$
$$\mu = \sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}$$

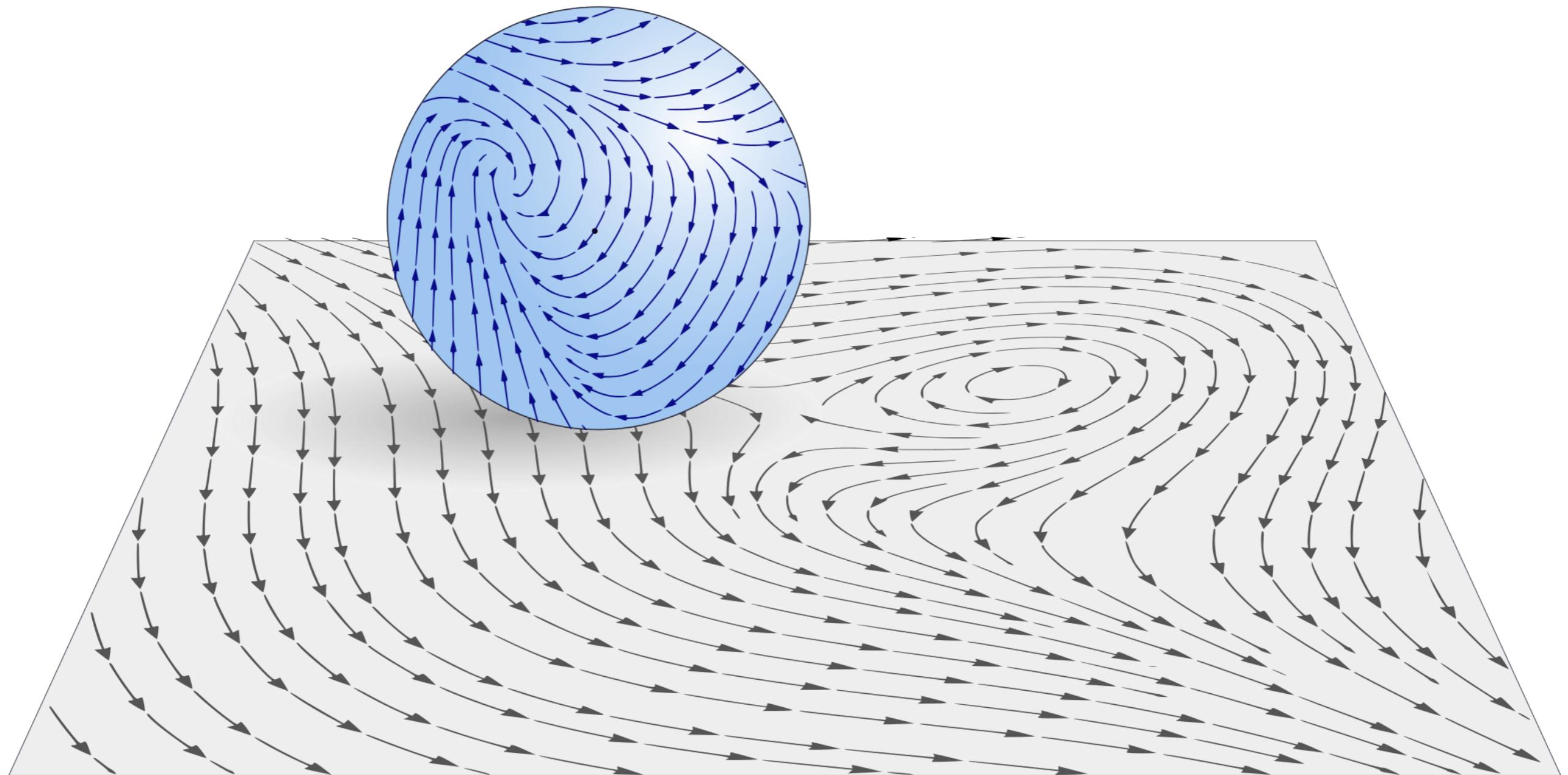
$$E_{mov} = \frac{1}{2} \langle \mathbb{I}(\Omega + \eta E_3), \Omega + \eta E_3 \rangle + \frac{m}{2} ||\rho \times (\Omega + \eta E_3)||^2 - mr \langle \rho, \gamma \rangle$$

Invariant measure

$$\frac{1}{\sqrt{I_1 I_3 + m \langle \rho, \mathbb{I} \rho \rangle}} dM d\gamma$$

# Chaplygin sphere

$$I_1 \neq I_2 \neq I_3$$



# Chaplygin sphere

$$I_1 \neq I_2 \neq I_3$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

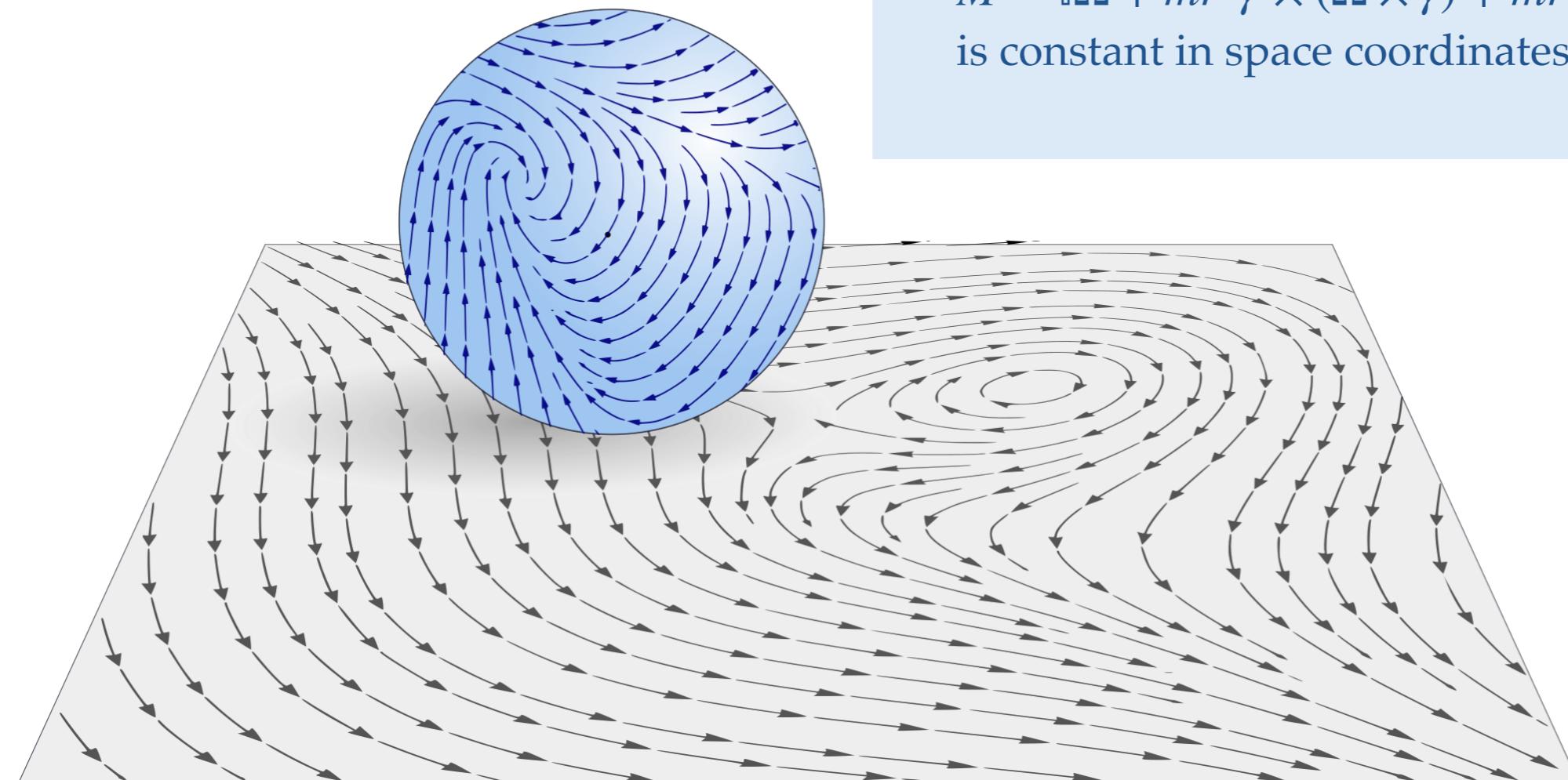
$$\dot{u} = -rB(\gamma \times \Omega) + BV + W$$

$$\dot{B} = B\hat{\Omega}$$

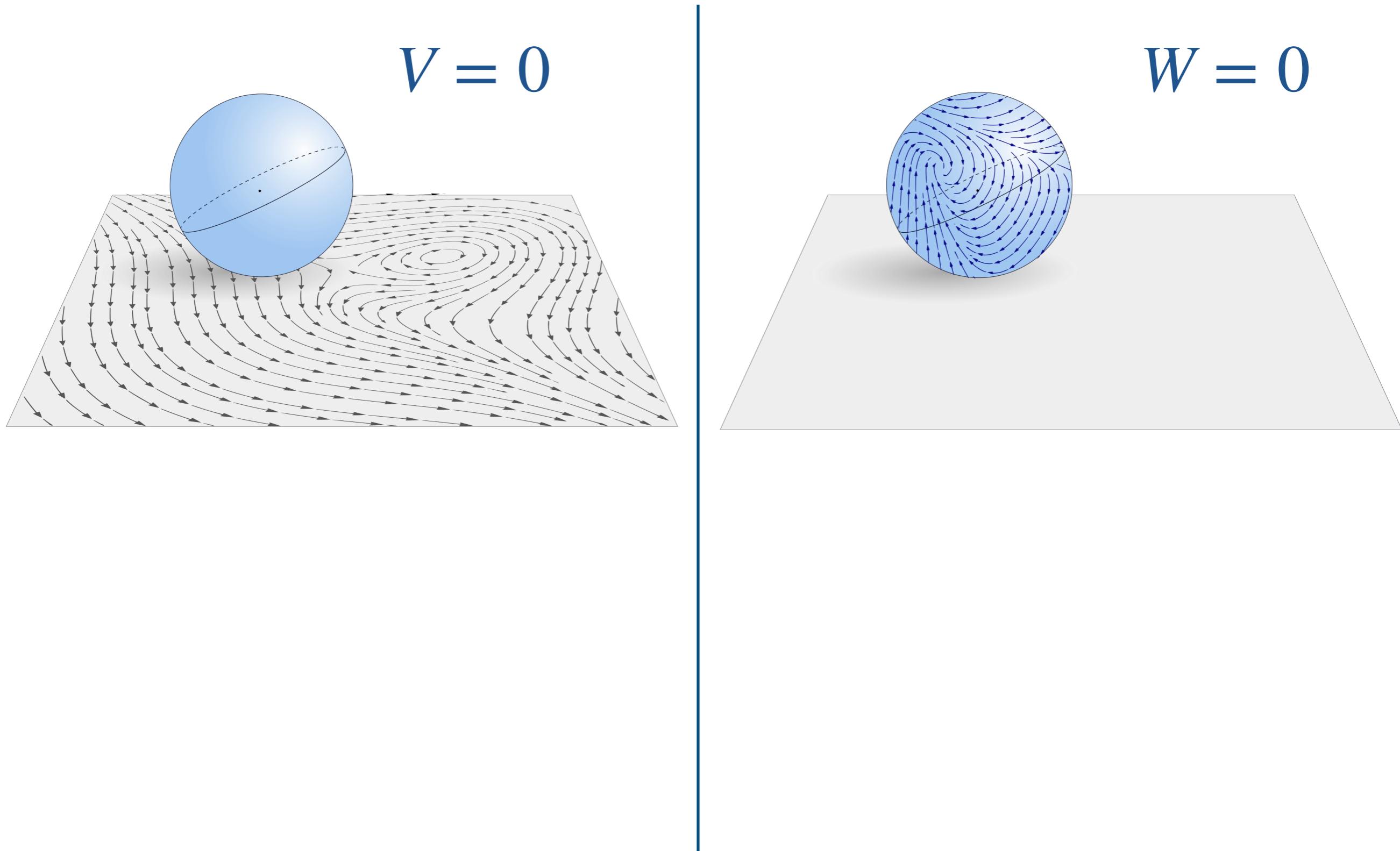
Theorem (C. García Naranjo, 2023):

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (V + B^{-1}W)$$

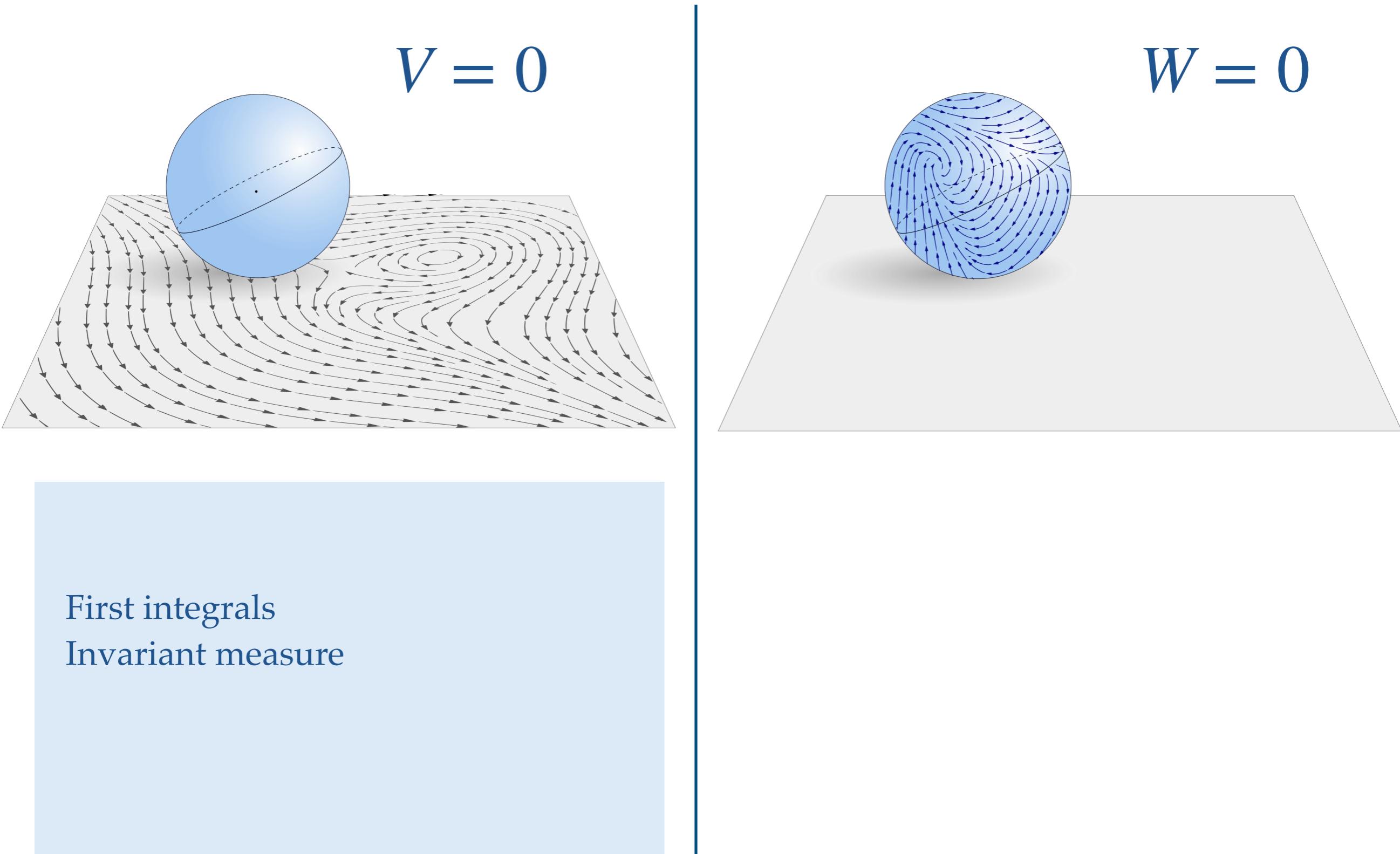
is constant in space coordinates



# Chaplygin sphere

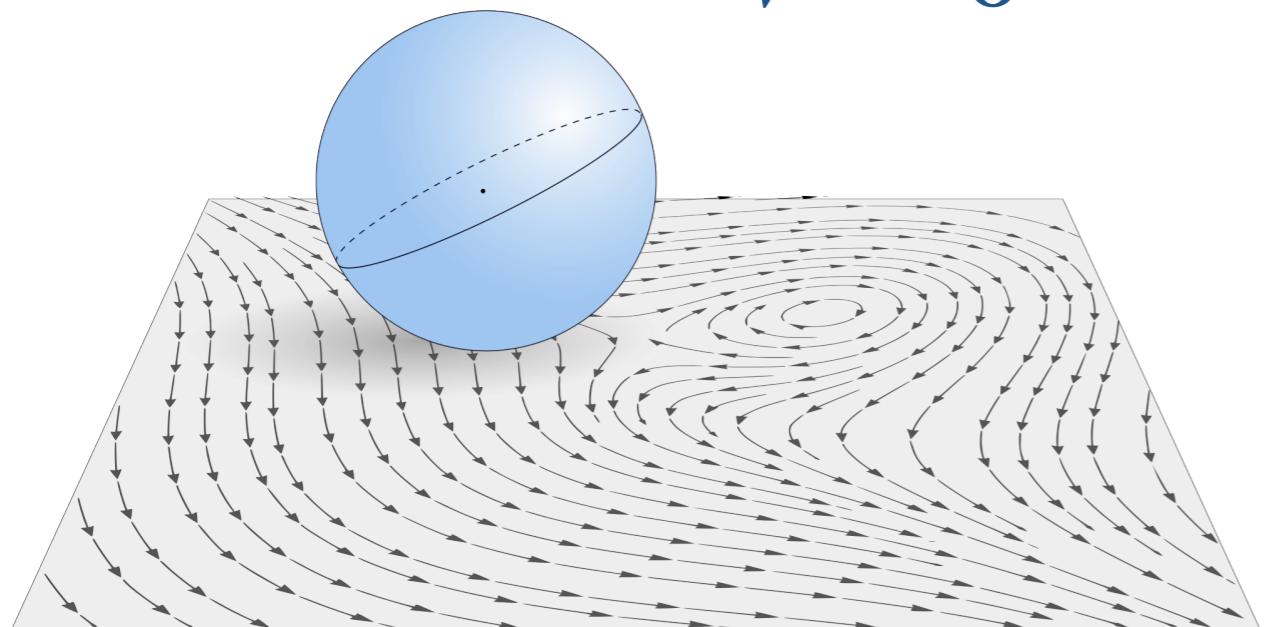


# Chaplygin sphere

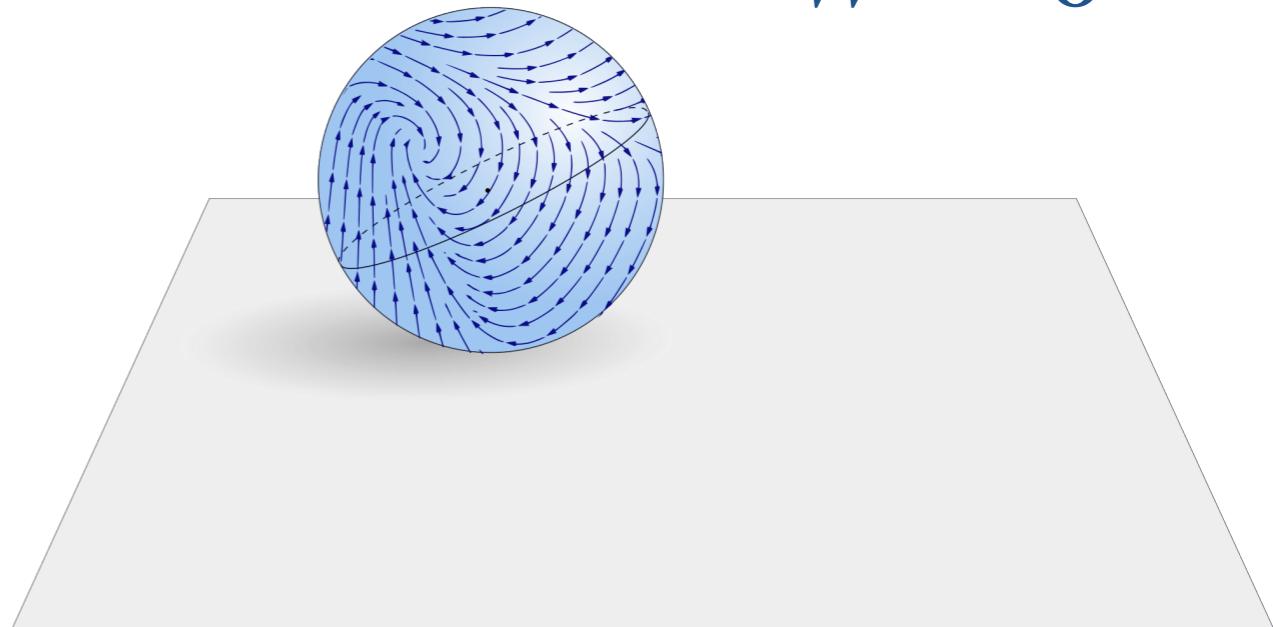


# Chaplygin sphere

$$V = 0$$



$$W = 0$$



First integrals  
Invariant measure

SE(2)-symmetry  
First integrals

# Chaplygin sphere

$$V = 0$$

Examples:

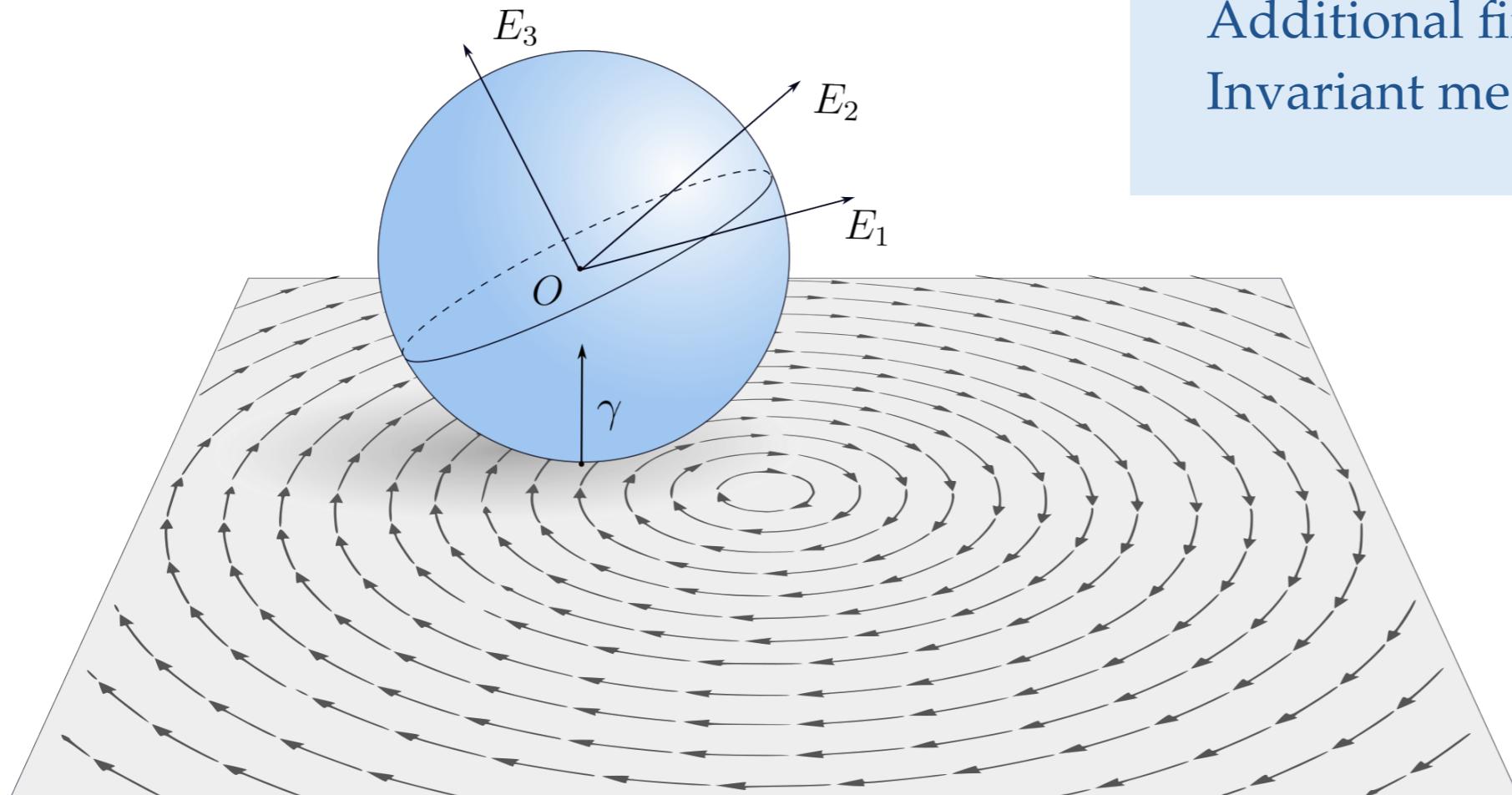
Rotating plane:

Bizyaev, Borisov, Mamaev, 2018

Vibrating plane:

Kilin, Pivovarova, 2021

Additional first integral- moving energy  
Invariant measure



# Chaplygin sphere

$$V = 0$$

Generally chaotic

Examples:

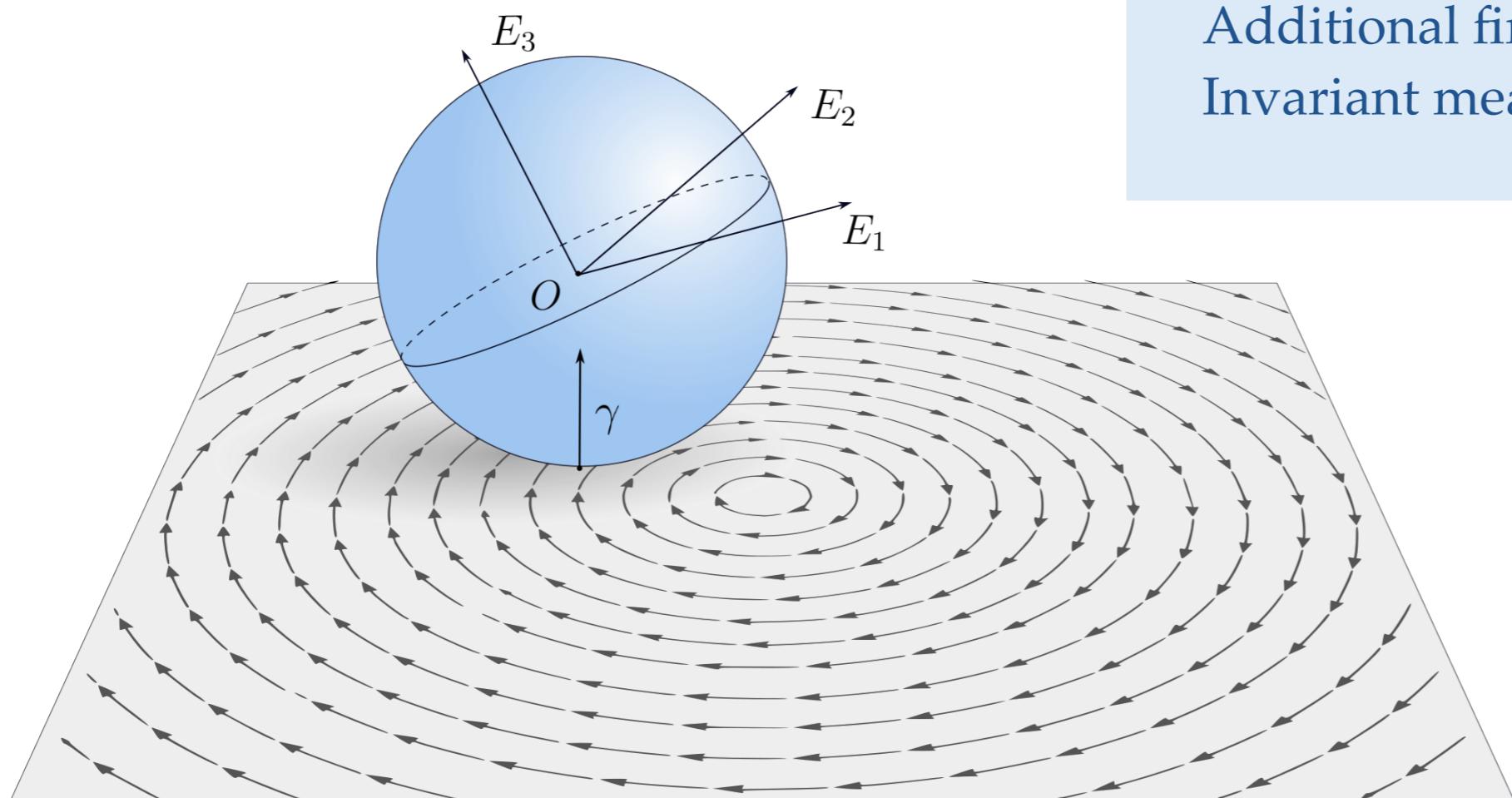
Rotating plane:

Bizyaev, Borisov, Mamaev, 2018

Vibrating plane:

Kilin, Pivovarova, 2021

Additional first integral- moving energy  
Invariant measure



# Chaplygin sphere

$$V = 0$$

$$W \in \mathfrak{X}(\mathbb{R}^2)$$

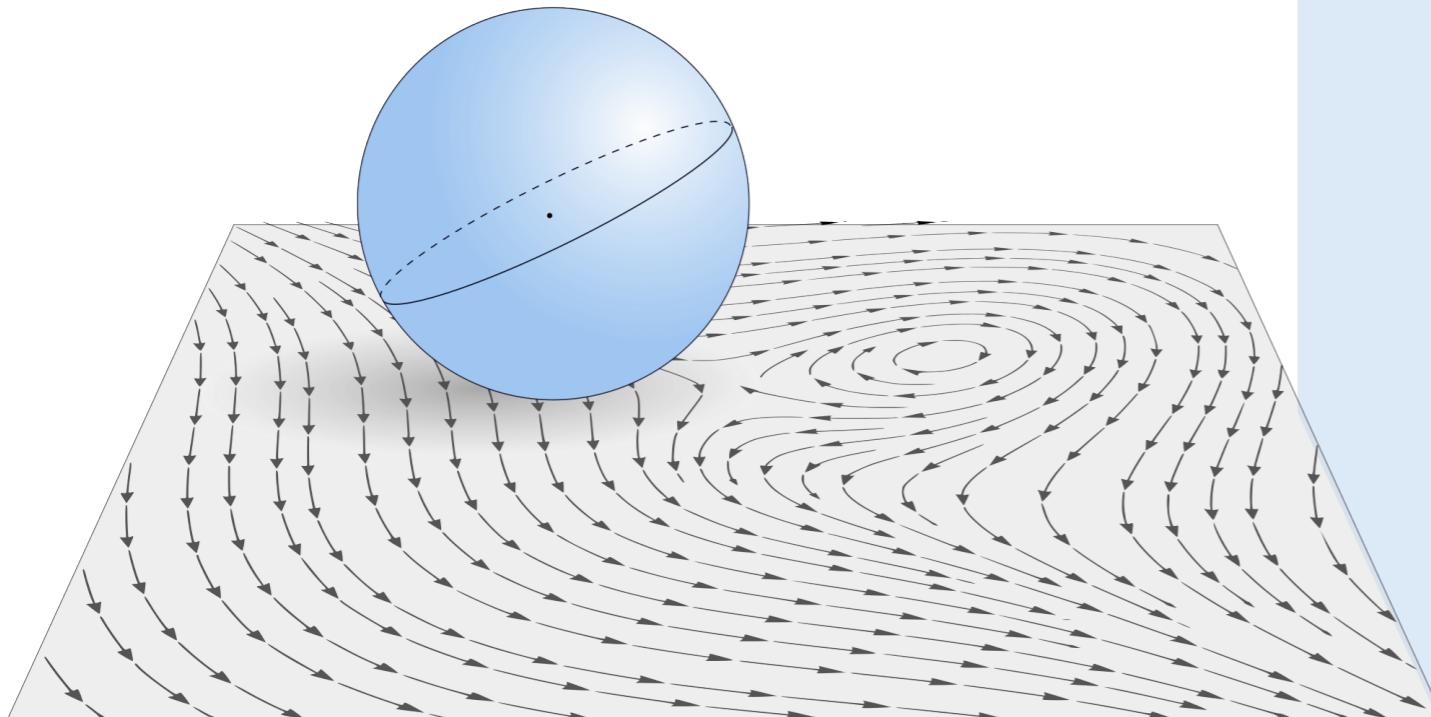
Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{u} = -rB(\gamma \times \Omega) + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (B^{-1}W)$$



# Chaplygin sphere

$$V = 0$$

$$W \in \mathfrak{X}(\mathbb{R}^2)$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

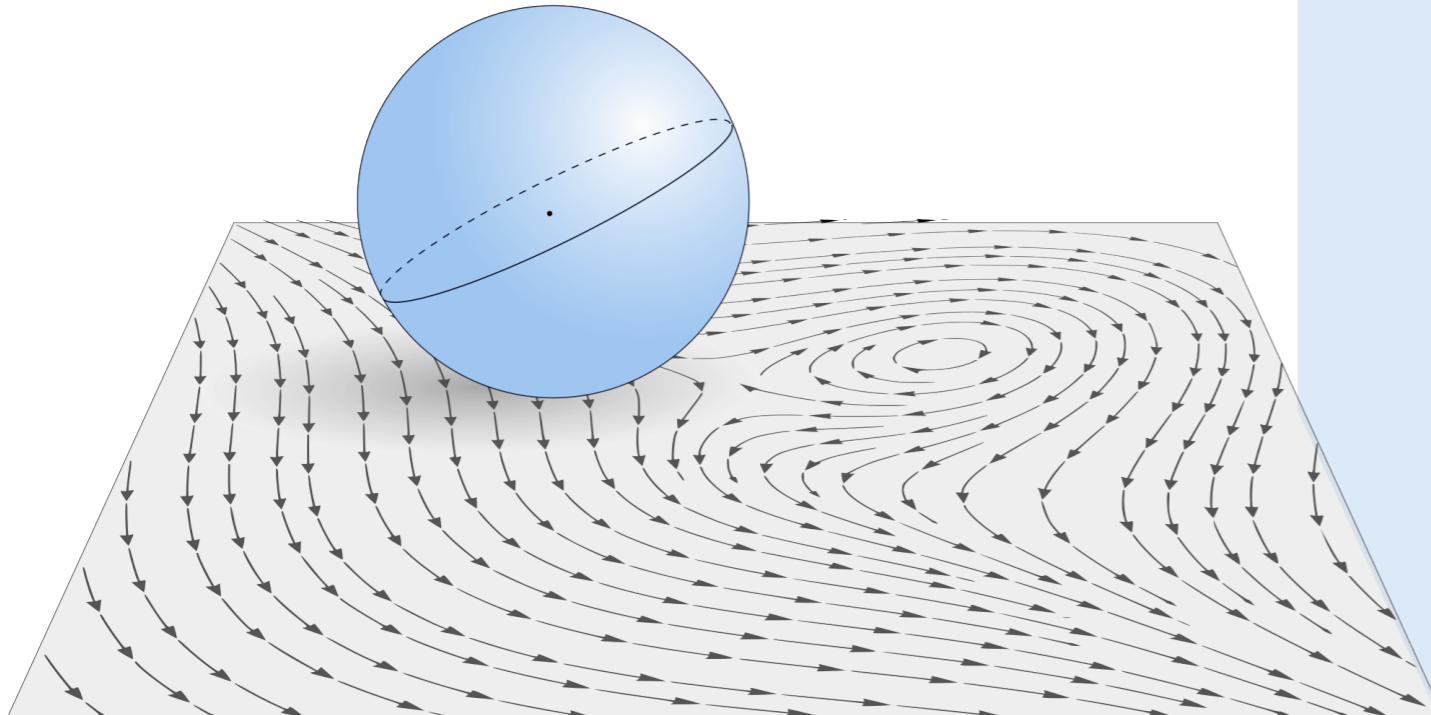
$$\dot{u} = -rB(\gamma \times \Omega) + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (B^{-1}W)$$

First integrals

$$\langle M, \alpha \rangle, \langle M, \beta \rangle, \langle M, \gamma \rangle$$



# Chaplygin sphere

$$V = 0$$

$$W \in \mathfrak{X}(\mathbb{R}^2)$$

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{u} = -rB(\gamma \times \Omega) + W$$

$$\dot{B} = B\hat{\Omega}$$

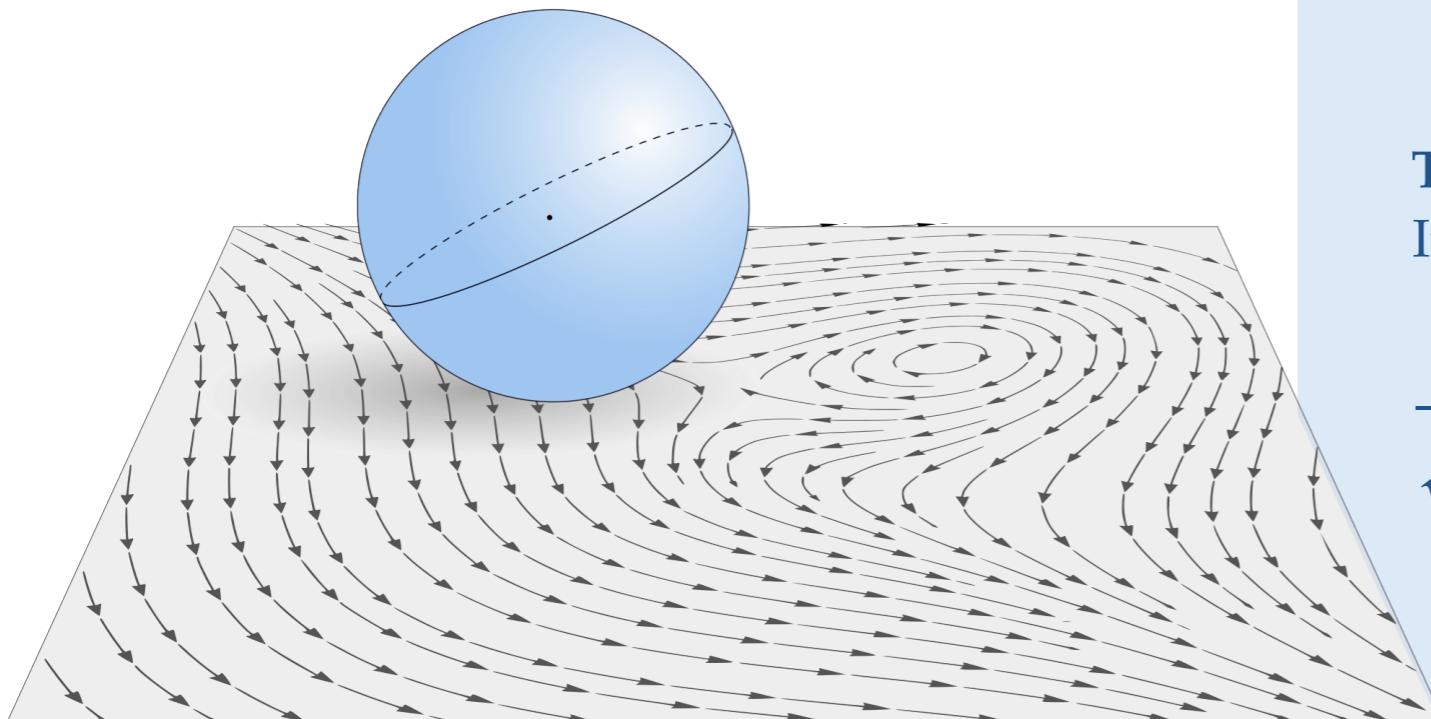
$$M = \mathbb{I}\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (B^{-1}W)$$

First integrals

$$\langle M, \alpha \rangle, \langle M, \beta \rangle, \langle M, \gamma \rangle$$

**Theorem (C. García Naranjo, 2023):**  
Invariant measure if  $\operatorname{div}_{\mathbb{R}^2} W = 0$

$$\frac{1}{\sqrt{1 - mr^2 \langle \gamma, (\mathbb{I} + mr^2)^{-1} \gamma \rangle}} dM dx dy da d\beta dy$$



# Chaplygin sphere

$$W = 0$$

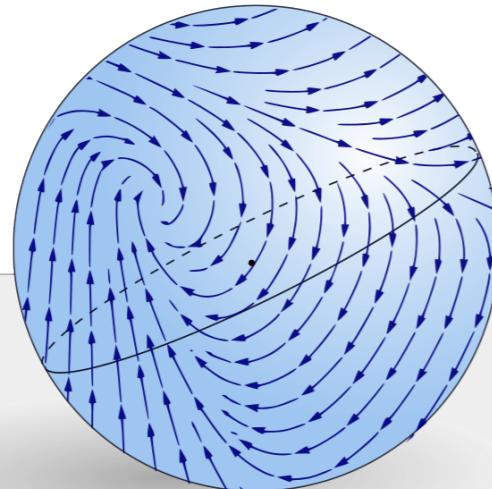
Equations of motion:

$$\begin{aligned}\dot{M} &= M \times \Omega \\ \dot{\gamma} &= \gamma \times \Omega\end{aligned}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times V$$

First integrals

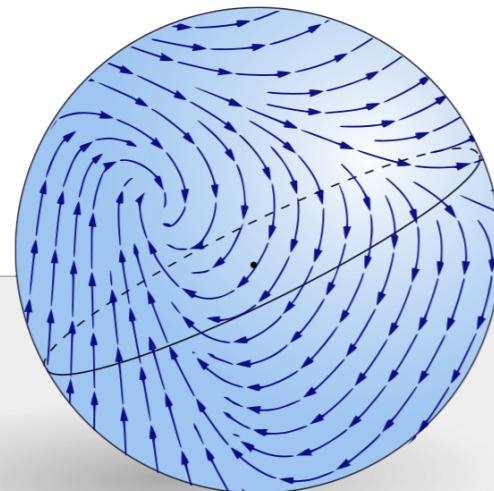
$$\|\gamma\|^2, \quad \|M\|^2, \quad \langle M, \gamma \rangle$$



# Chaplygin sphere

$$W = 0$$

No invariant measure  
No moving energy



Equations of motion:

$$\begin{aligned}\dot{M} &= M \times \Omega \\ \dot{\gamma} &= \gamma \times \Omega\end{aligned}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times V$$

First integrals

$$\|\gamma\|^2, \quad \|M\|^2, \quad \langle M, \gamma \rangle$$

# Chaplygin sphere with rotating shell

$$V = -r\eta\gamma \times E_3$$

Equations of motion:

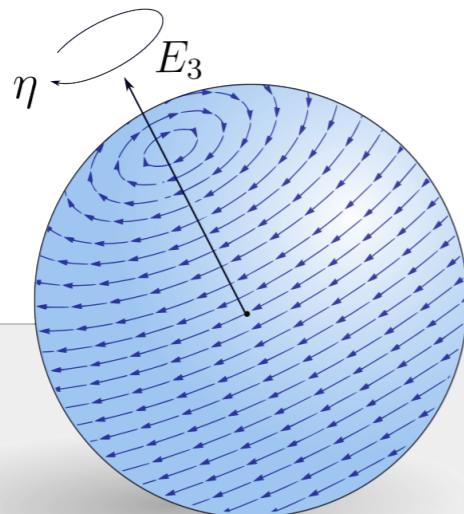
$$\begin{aligned}\dot{M} &= M \times \Omega \\ \dot{\gamma} &= \gamma \times \Omega\end{aligned}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - m^2\eta\gamma \times (\gamma \times E_3)$$

First integrals:

$$\|\gamma\|^2, \quad \|M\|^2, \quad \langle M, \gamma \rangle$$

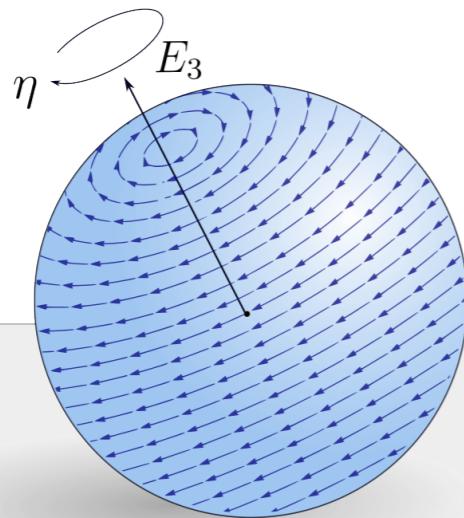
3 dimensional level sets



# Chaplygin sphere with rotating shell

$$V = -r\eta\gamma \times E_3$$

Generally chaotic



Equations of motion:

$$\begin{aligned}\dot{M} &= M \times \Omega \\ \dot{\gamma} &= \gamma \times \Omega\end{aligned}$$

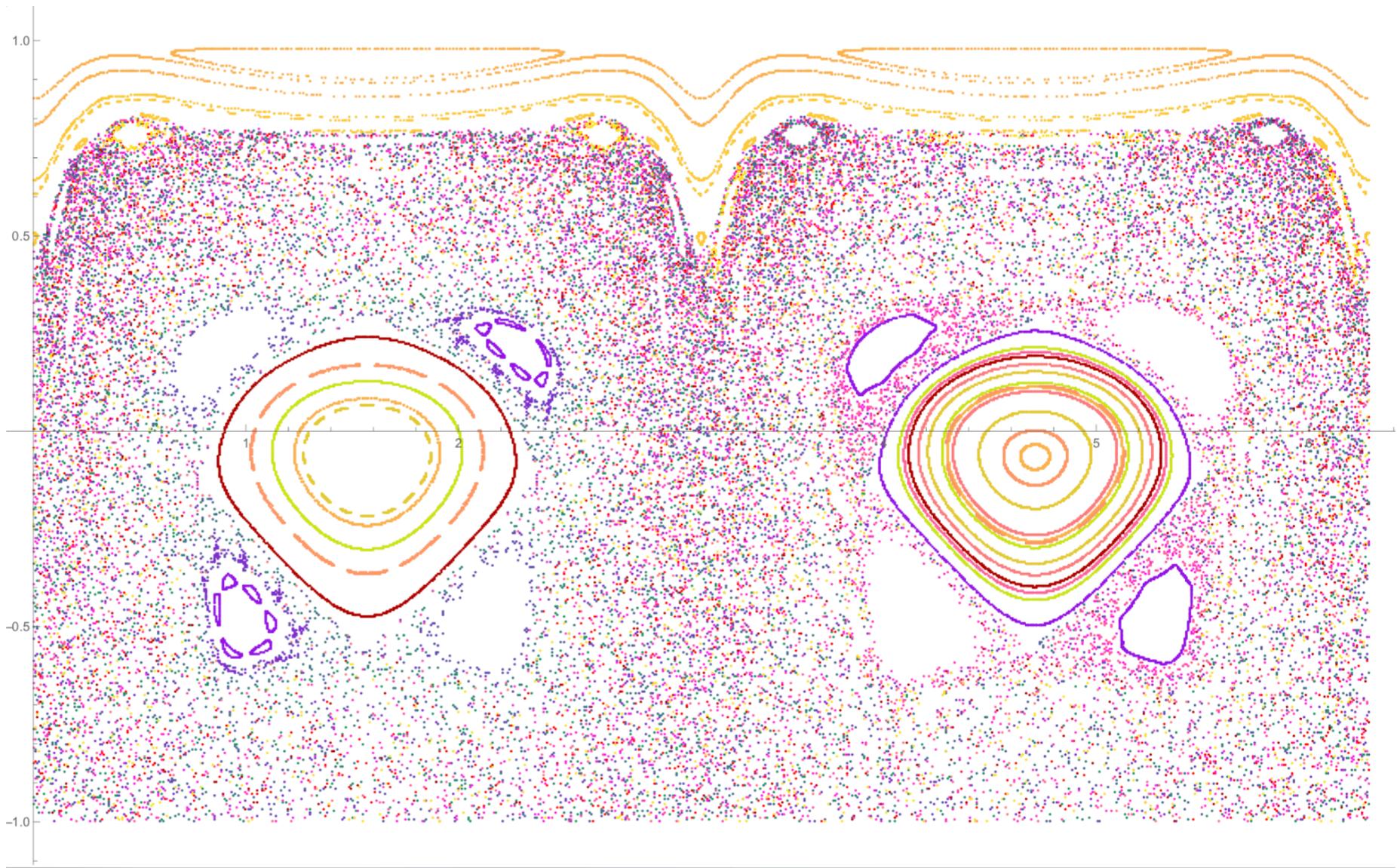
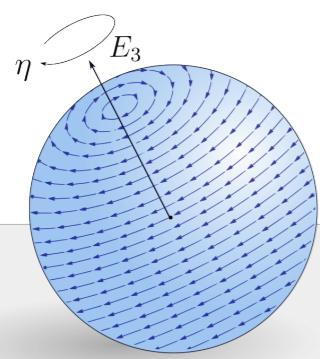
$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - m^2\eta\gamma \times (\gamma \times E_3)$$

First integrals:

$$\|\gamma\|^2, \quad \|M\|^2, \quad \langle M, \gamma \rangle$$

3 dimensional level sets

# Poincaré map



# Poincaré map

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

# Poincaré map

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

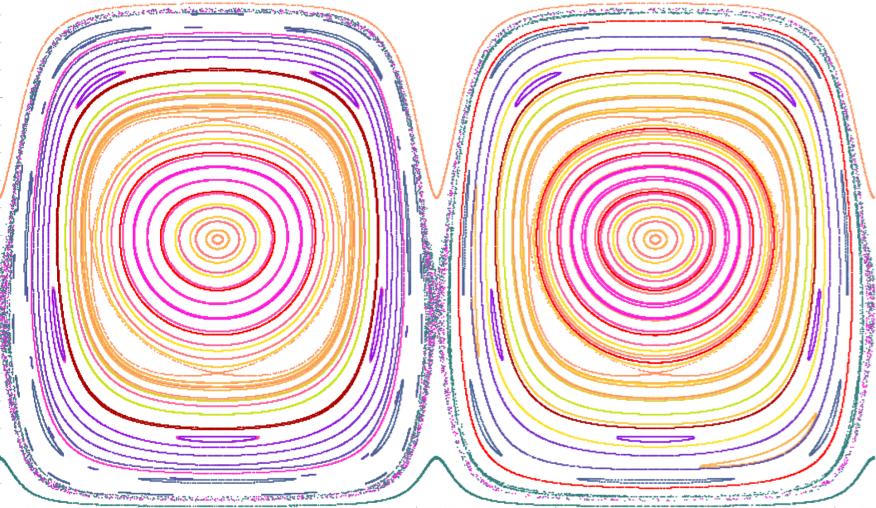
Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

# Poincaré map

$\varepsilon \gg 1$



Equations of motion:

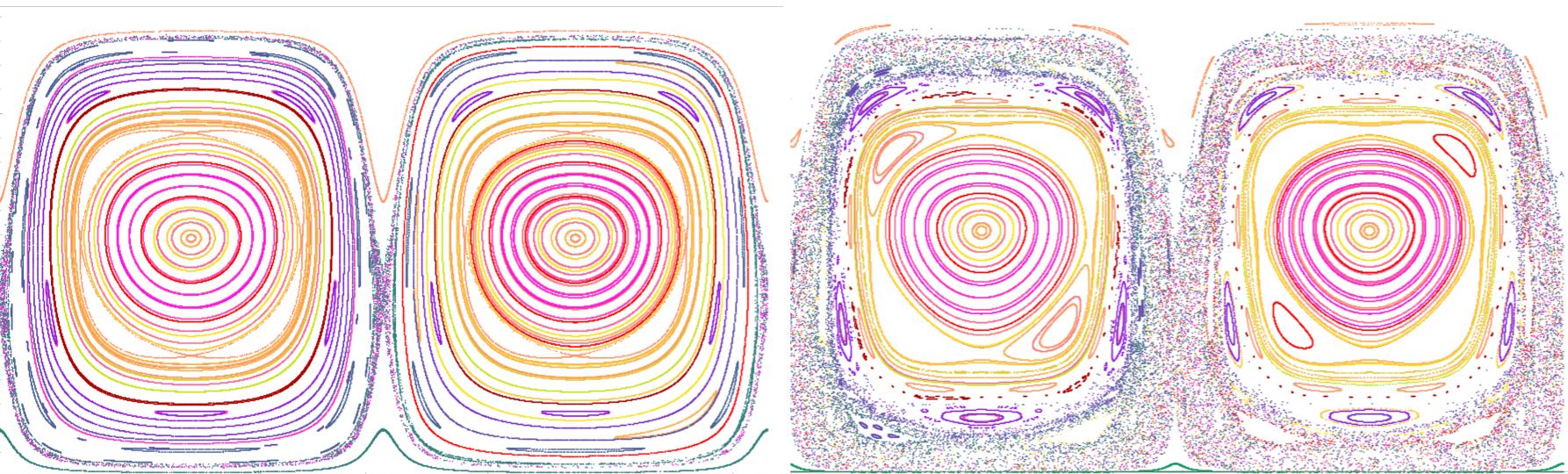
$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

# Poincaré map

$\varepsilon \gg 1$



Equations of motion:

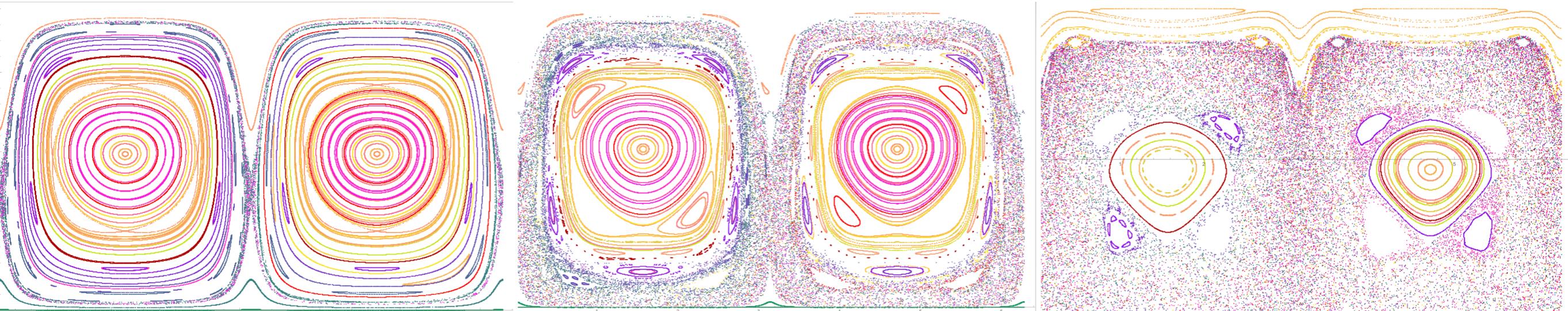
$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

# Poincaré map

$\varepsilon \gg 1$



Equations of motion:

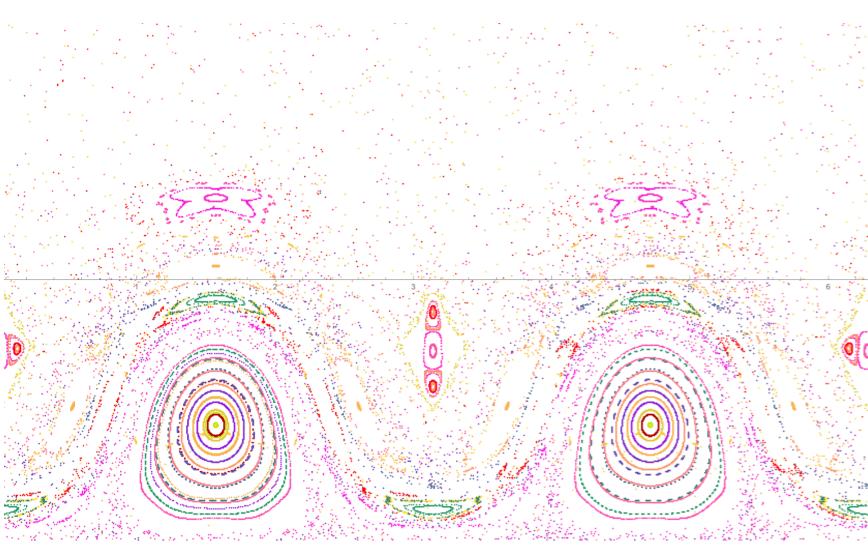
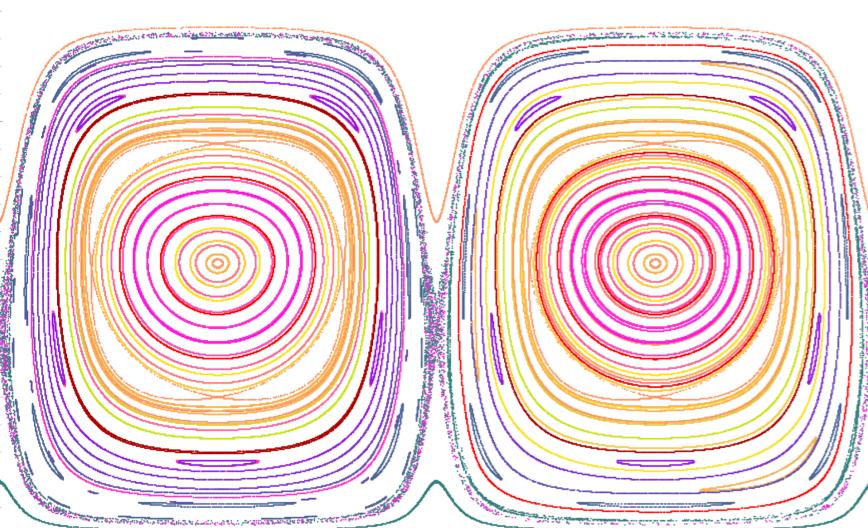
$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

# Poincaré map

$\varepsilon \gg 1$

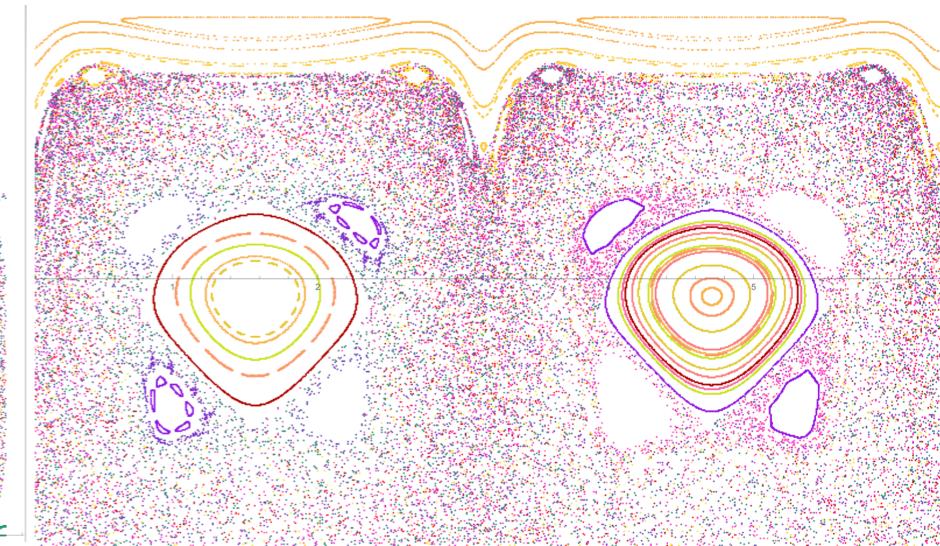
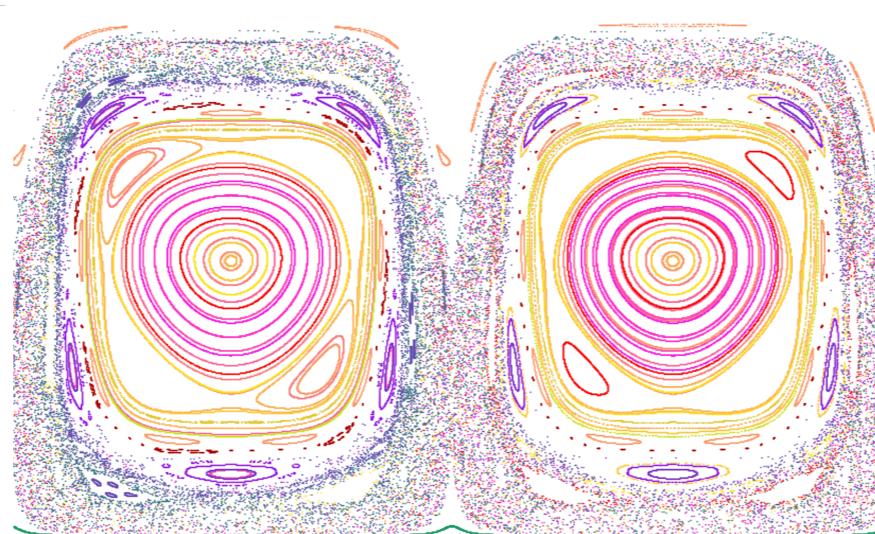


$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$



# Poincaré map

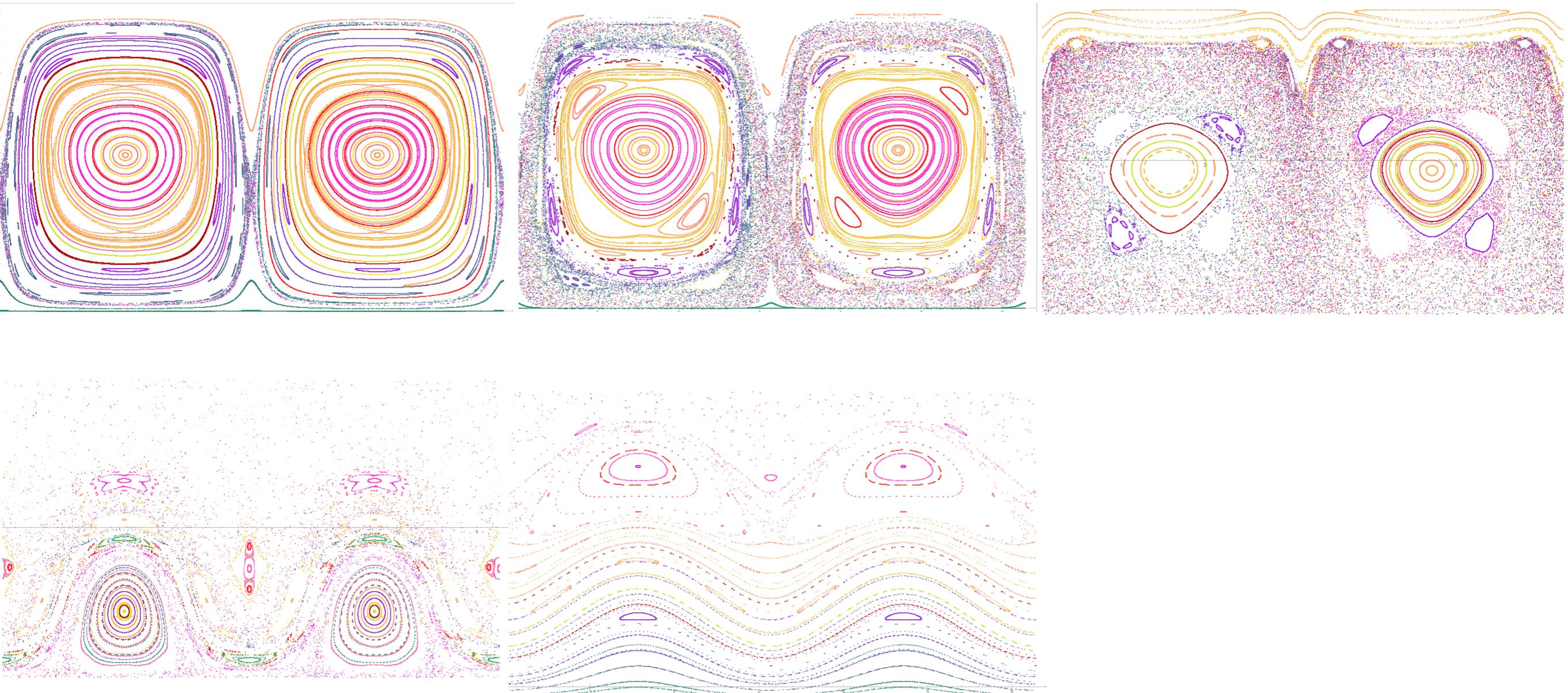
$\varepsilon \gg 1$

Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

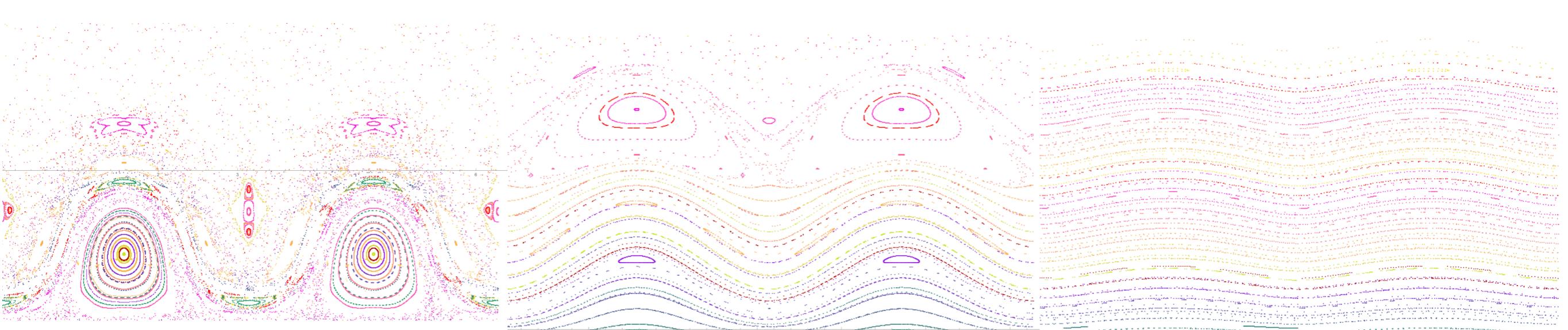
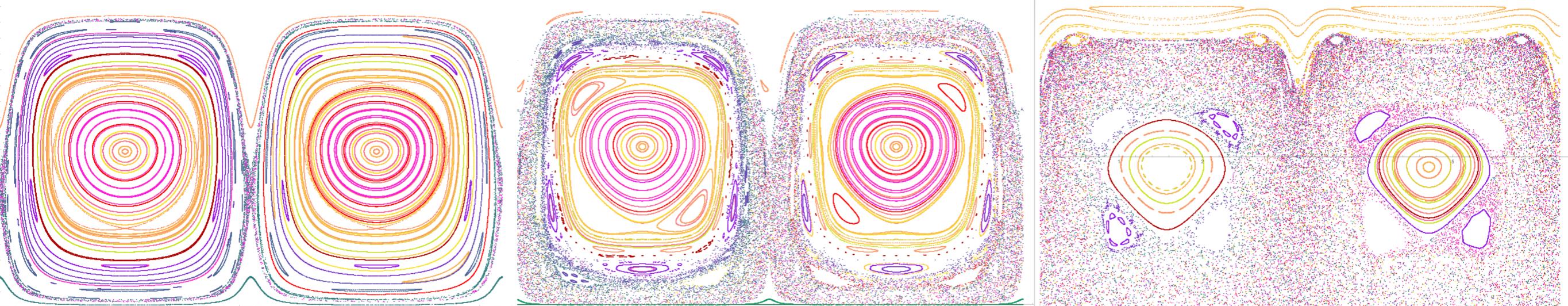
$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$



# Poincaré map

$\varepsilon \gg 1$



Equations of motion:

$$\dot{M} = \varepsilon (kM \times \Omega_c) + M \times \Omega_a$$

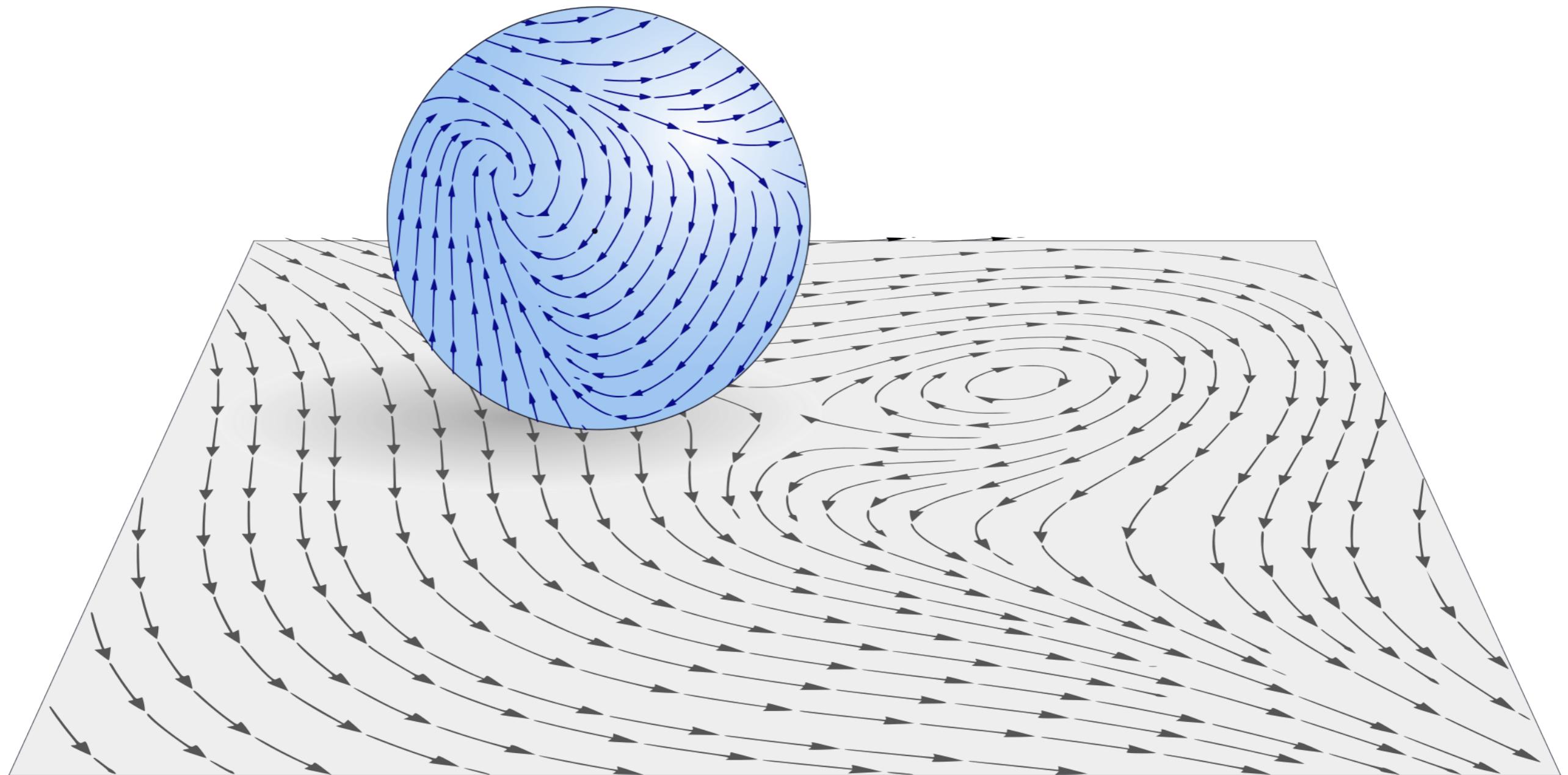
$$\dot{\gamma} = \varepsilon (k\gamma \times \Omega_c) + \gamma \times \Omega_a$$

$\varepsilon \ll 1$

$$\varepsilon = \frac{\|M\|}{mr^2\eta}$$

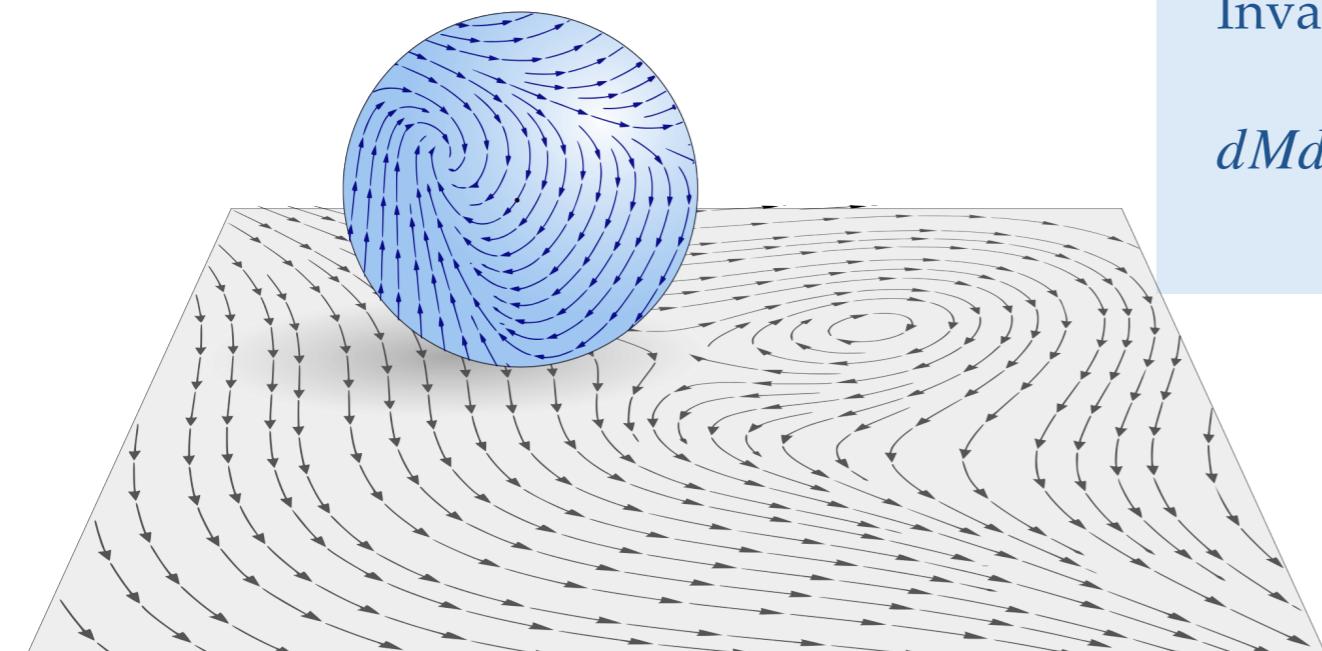
# Homogeneous sphere

$$I_1 = I_2 = I_3$$



# Homogeneous sphere

$$I_1 = I_2 = I_3$$



Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{u} = -rB(\gamma \times \Omega) + BV + W$$

$$\dot{B} = B\hat{\Omega}$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) + mr\gamma \times (V + B^{-1}W)$$

First integrals:

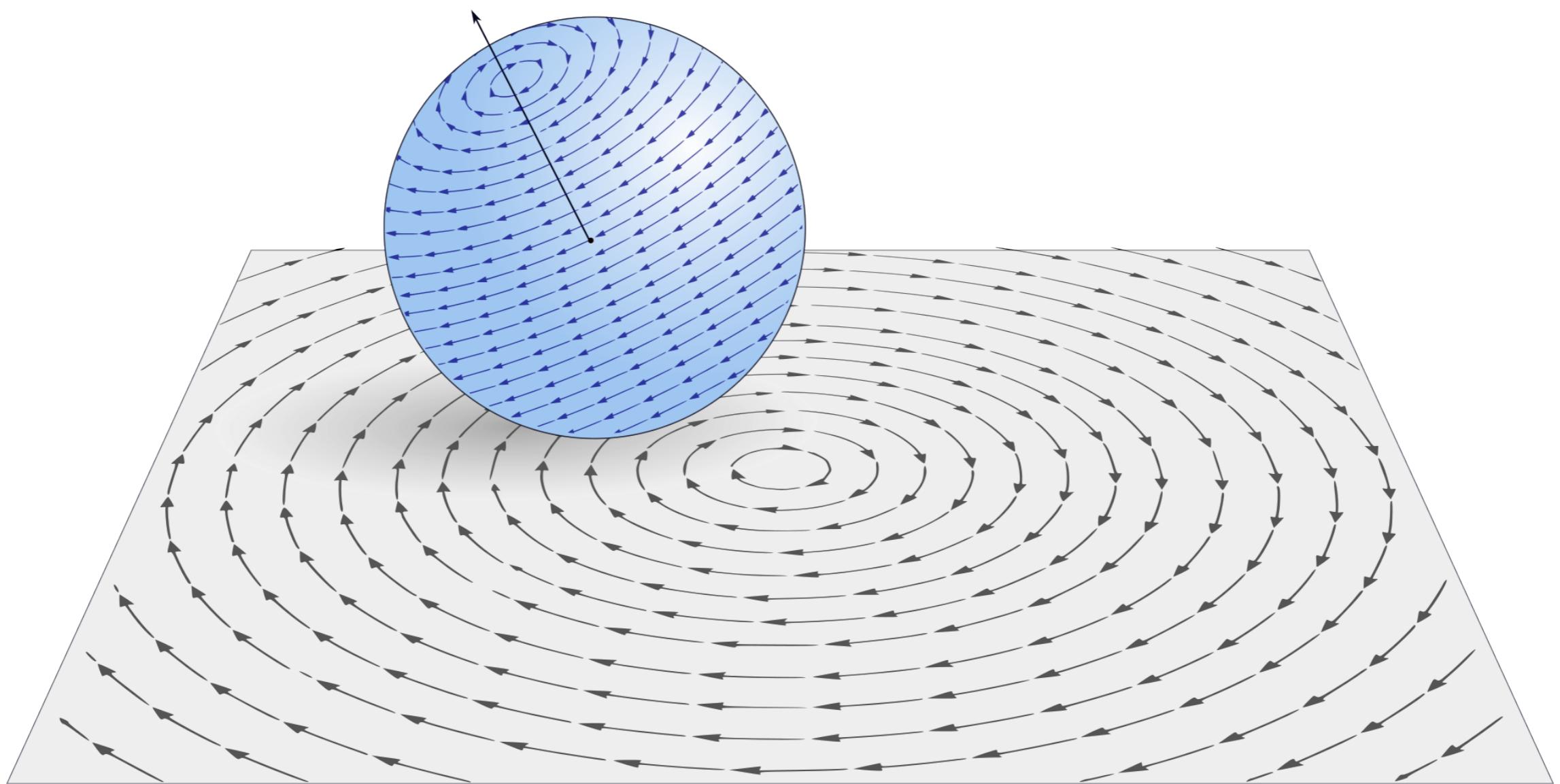
$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2$$

**Theorem (C., García Naranjo, 2023)**

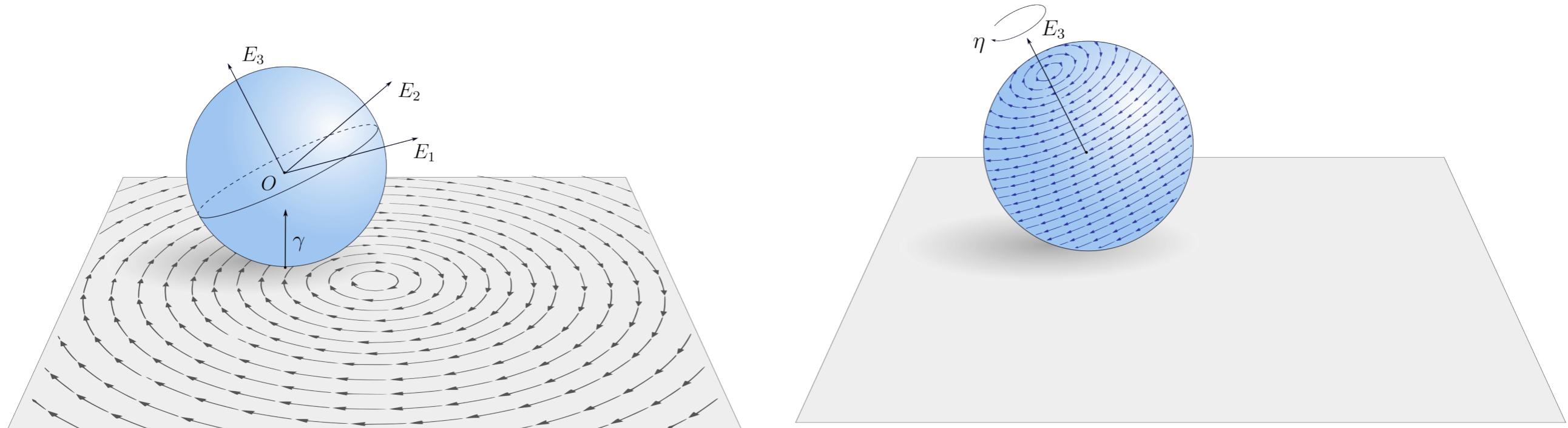
Invariant measure:

$$dM dx dy d\alpha d\beta d\gamma \text{ if } \operatorname{div}_{\mathbb{R}^2} W = 0 \text{ and } \operatorname{div}_{S^2} V = 0$$

# Homogeneous sphere

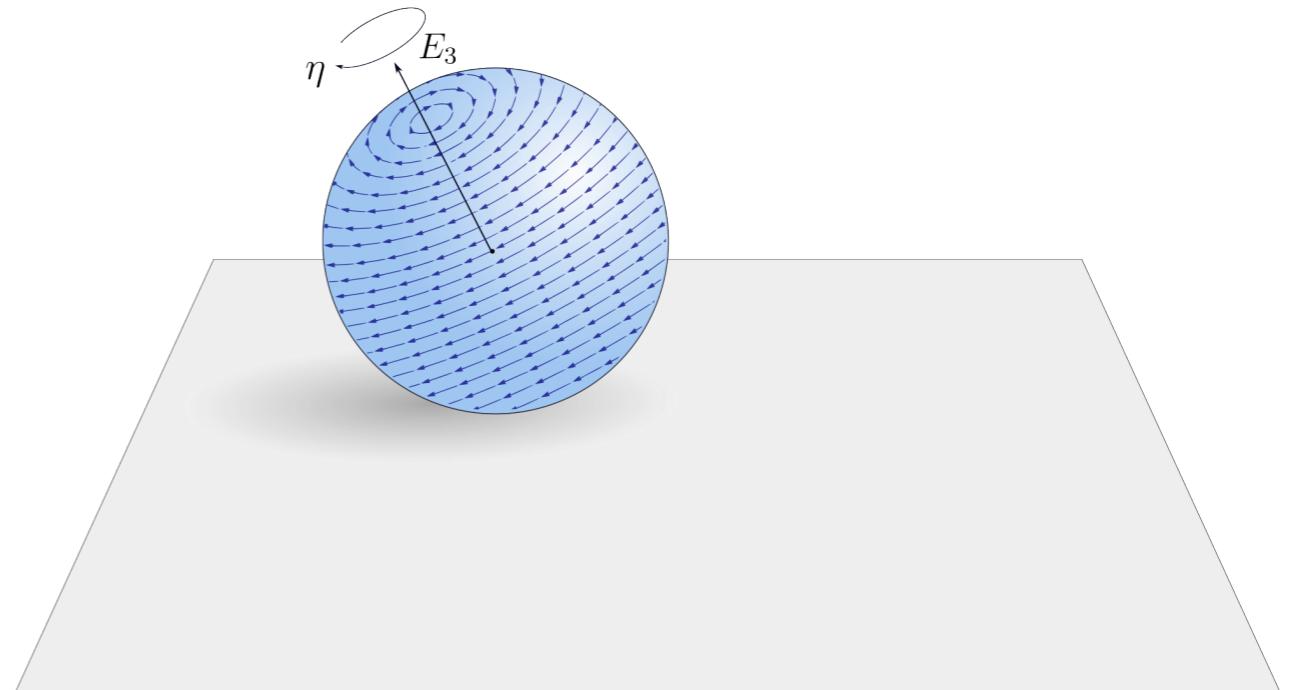
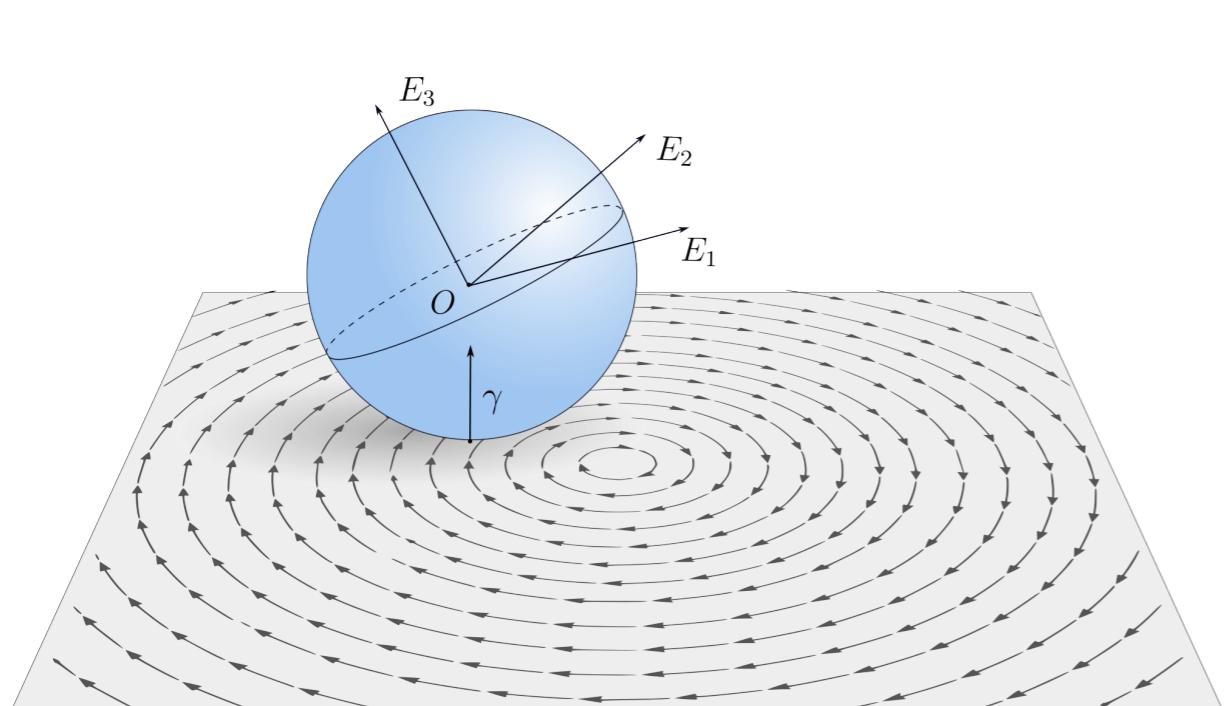


# Homogeneous sphere



# Homogeneous sphere

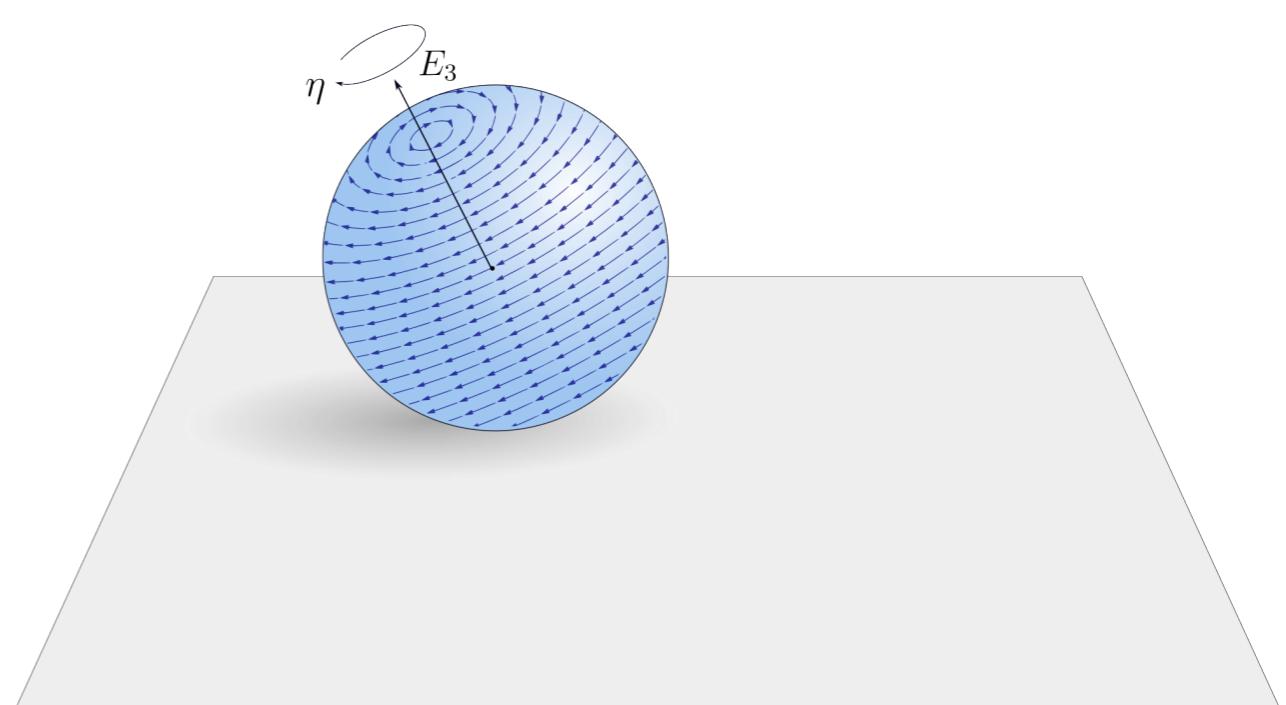
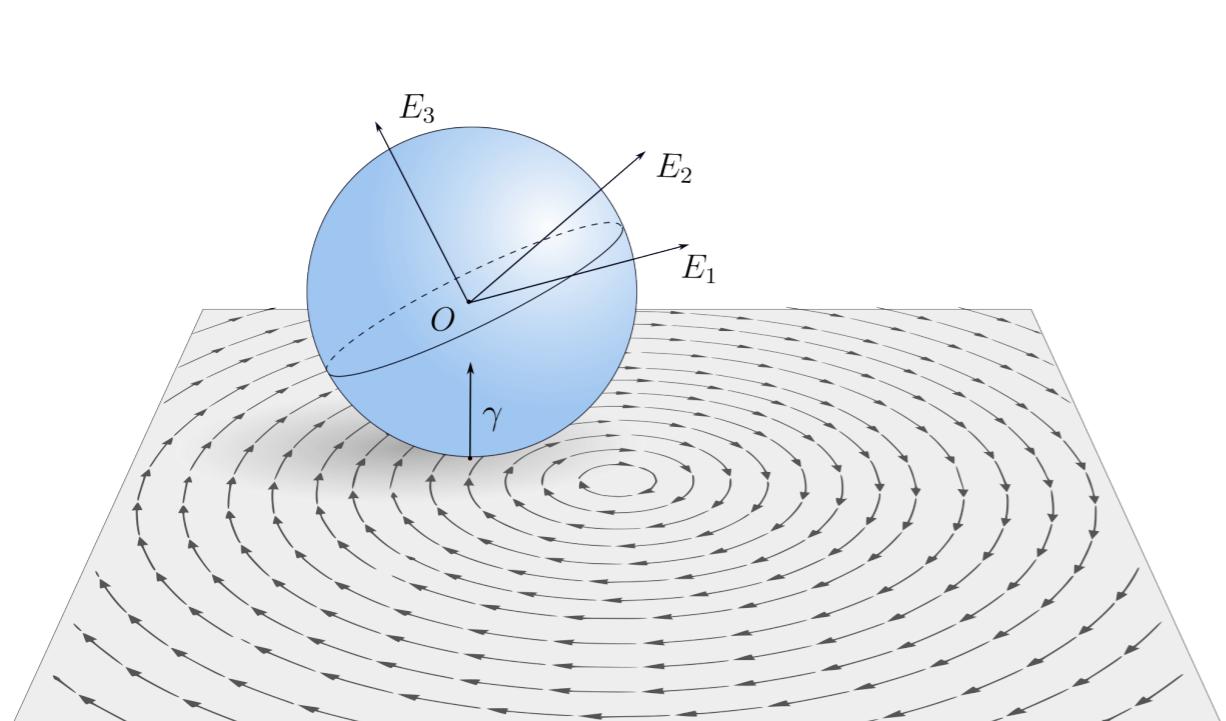
Integrable



# Homogeneous sphere

Integrable

Integrable



# Homogeneous sphere

Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - r\eta_1(\gamma \times E_3) - \eta_2\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

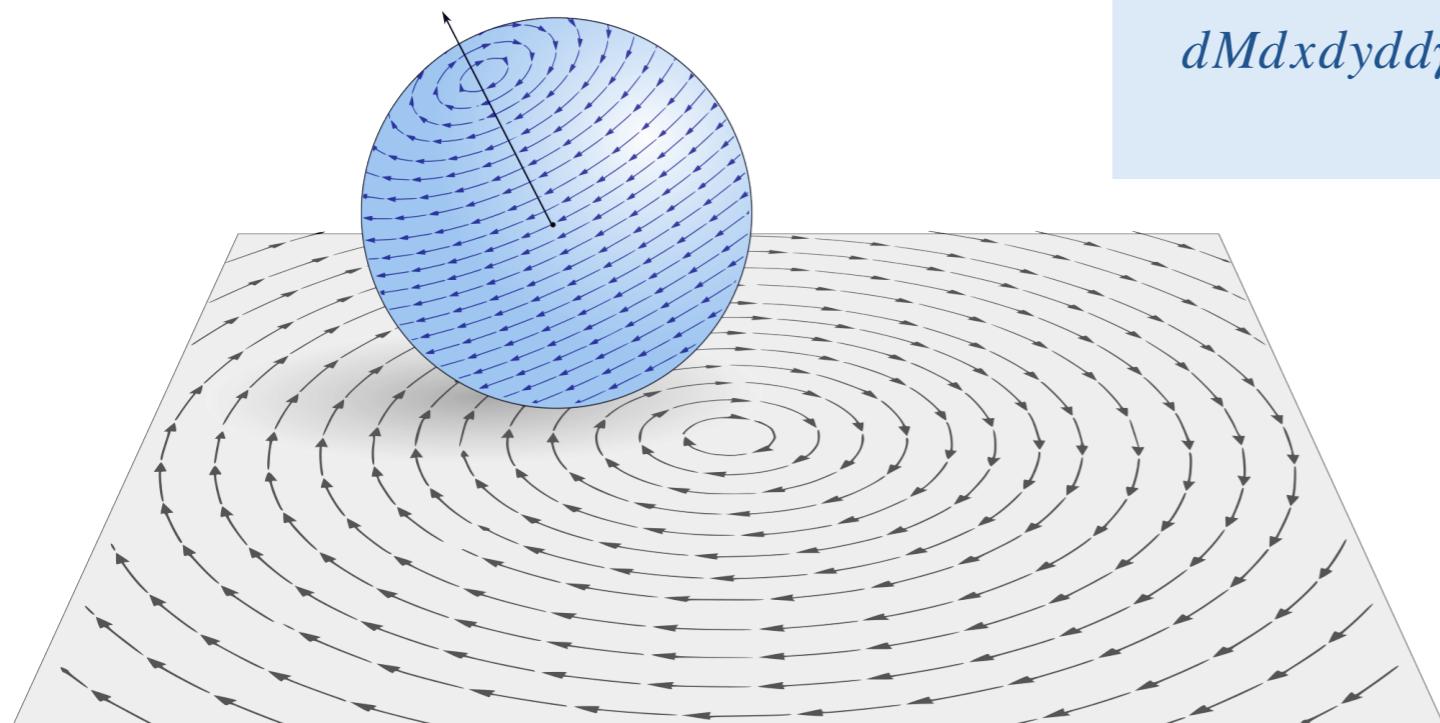
$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3) + \eta_2\gamma \times (\gamma \times U)$$

First integrals

$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0, \quad E_{mov}$$

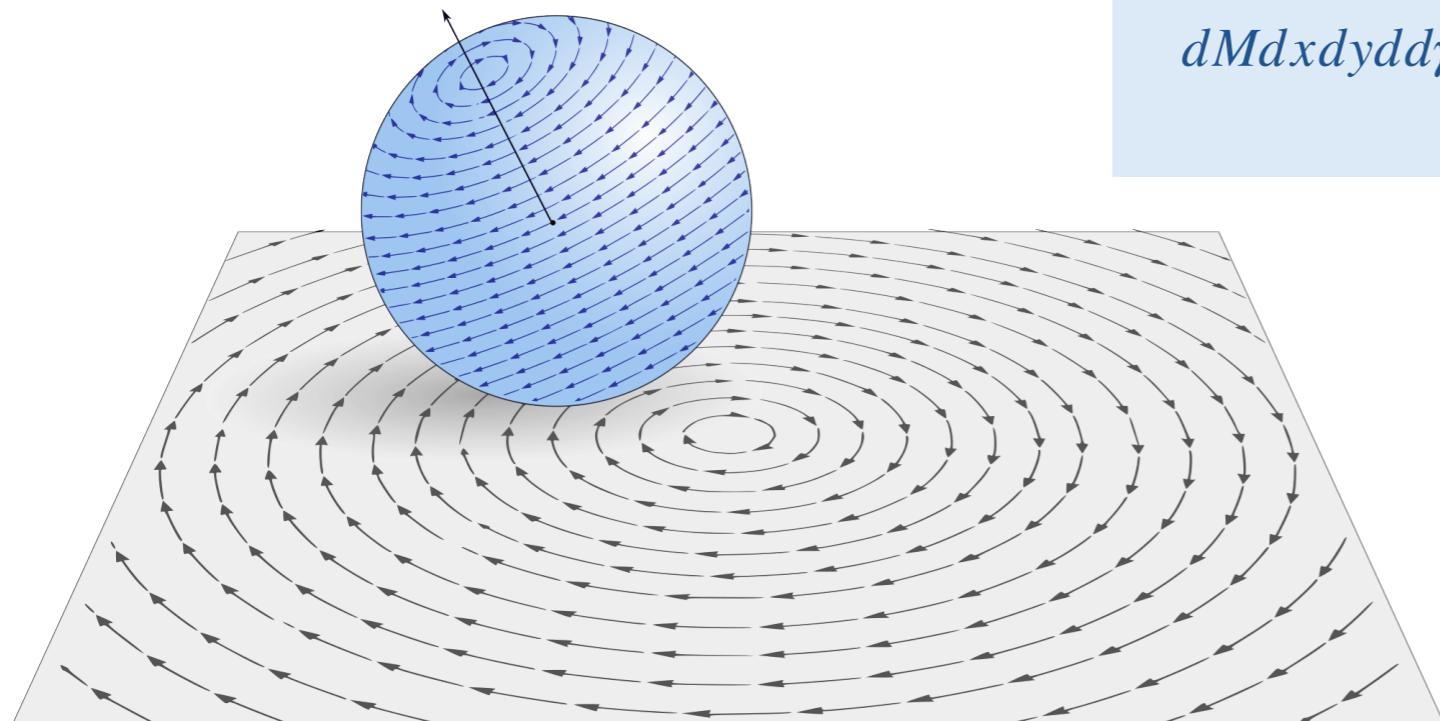
Invariant measure

$$dM dx dy dd\gamma$$



# Homogeneous sphere

Current work



Equations of motion:

$$\dot{M} = M \times \Omega$$

$$\dot{U} = -r(\gamma \times \Omega) + U \times \Omega - r\eta_1(\gamma \times E_3) - \eta_2\gamma \times U$$

$$\dot{\gamma} = \gamma \times \Omega$$

$$M = I\Omega + mr^2\gamma \times (\Omega \times \gamma) - mr^2\eta_1\gamma \times (\gamma \times E_3) + \eta_2\gamma \times (\gamma \times U)$$

First integrals

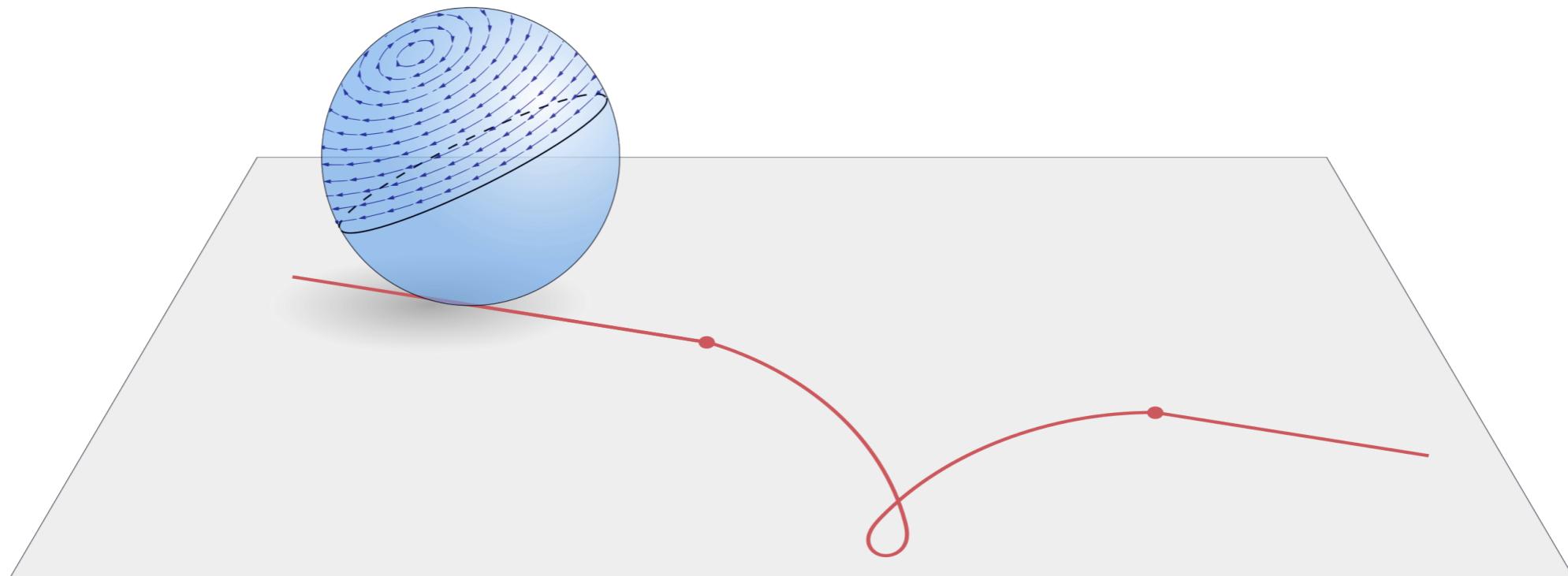
$$\|M\|^2, \quad \langle M, \gamma \rangle, \quad \|\gamma\|^2, \quad \langle U, \gamma \rangle = 0, \quad E_{mov}$$

Invariant measure

$$dM dx dy dd\gamma$$

# Homogeneous sphere with rotating shell

We proved an analogous to the  
Anais-billiard phenomenon



# Conclusions

# Conclusions

- Identified a rich class of examples to analyze dynamical properties of affine nonholonomic systems
- Recognized a robust mechanism leading to existence of first integrals (momenta)
- Found new instances of moving energy
- Observed integrable and chaotic behavior and transition as function of problem parameters

# Future work

# Future work

- Role of reversibility in affine nonholonomic systems
- Momentum first integrals for affine nonholonomic systems
- Develop a theory to frame these examples:  
Integrability, perturbation theory and chaos of nonholonomic systems

Thank you!