

The derivation of the Lindblad equation for the electron-electron scattering

I. LINDBLAD DYNAMICS

According to Eq. 4 and 5 of PHYSICAL REVIEW B 90, 125140 (2014) or Eq. 3.20 and 3.21 of Rosati's thesis, the scattering term has the form

$$\frac{d\hat{\rho}}{dt} = \hat{A}\hat{\rho}\hat{A}^\dagger - \frac{1}{2} \left\{ \hat{A}^\dagger \hat{A}, \hat{\rho} \right\},$$

where

$$\hat{A} = \sqrt{\frac{1}{\sqrt{2\pi t}}} \int dt' e^{-\frac{H_0 t'}{i\hbar}} \hat{H} e^{\frac{H_0 t'}{i\hbar}} e^{-\frac{1}{4} \left(\frac{t'}{t} \right)^2}. \quad (1)$$

Therefore,

$$\begin{aligned} \frac{d\rho_{12}}{dt} &= \text{Tr} \left(\hat{c}_2^\dagger \hat{c}_1 \frac{d\hat{\rho}}{dt} \right) \\ &= \frac{1}{2} \text{Tr} \left(\hat{c}_2^\dagger \hat{c}_1 \hat{A} \hat{\rho} \hat{A}^\dagger + \hat{A} \hat{\rho} \hat{A}^\dagger \hat{c}_2^\dagger \hat{c}_1 - \hat{c}_2^\dagger \hat{c}_1 \hat{A}^\dagger \hat{A} \hat{\rho} - \hat{\rho} \hat{A}^\dagger \hat{A} \hat{c}_2^\dagger \hat{c}_1 \right) \\ &= \frac{1}{2} \text{Tr} \left(\left[\hat{A}^\dagger, \hat{c}_2^\dagger \hat{c}_1 \right] \hat{A} \hat{\rho} + \hat{\rho} \hat{A}^\dagger \left[\hat{c}_2^\dagger \hat{c}_1, \hat{A} \right] \right) \\ &= \frac{1}{2} \text{Tr} \left(\left[\hat{A}^\dagger, \hat{c}_2^\dagger \hat{c}_1 \right] \hat{A} \hat{\rho} \right) + H.C.. \end{aligned} \quad (2)$$

II. ELECTRON-ELECTRON

The electron-electron Hamiltonian reads (see Eq. 4 in PHYSICAL REVIEW B 92, 235423 (2015))

$$\begin{aligned} \hat{H} &= \frac{1}{2} \sum_{1234} g_{1234} c_1^\dagger c_2^\dagger c_4 c_3, \\ g_{1234} &= \langle 1(r) | \langle 2(r') | V(r - r') | 3(r) \rangle | 4(r') \rangle \\ &= V(q_{13}, \omega_{13}) \delta_{k_1+k_2, k_3+k_4} \langle 1|3 \rangle \langle 2|4 \rangle, \end{aligned}$$

where $q_{13} = k_1 - k_3$ and $\omega_{13} = \epsilon_1 - \epsilon_3$.

Insert it in 1,

$$\begin{aligned} \hat{A} &= \frac{1}{2} \sum_{1234} A_{1234} c_1^\dagger c_2^\dagger c_4 c_3, \\ A_{1234} &= \sqrt{\frac{1}{\sqrt{2\pi t}}} \int dt' e^{-\frac{(\epsilon_1 + \epsilon_2 - \epsilon_3 - \epsilon_4)t'}{i\hbar}} g_{1234} e^{-\frac{1}{4} \left(\frac{t'}{t} \right)^2} \\ &= \sqrt{\frac{2\pi}{\hbar}} g_{1234} \delta^{G,1/2}(\epsilon_1 + \epsilon_2 - \epsilon_3 - \epsilon_4). \end{aligned} \quad (3)$$

For later derivation, we define (we will use the symmetries of the matrix later)

$$\mathcal{A}_{1234} = \frac{1}{4} (A_{1234} - A_{2134} - A_{1243} - A_{2143}).$$

Since $V(q, \omega) = V(-q, -\omega)$, we have

$$\mathcal{A}_{1234} = \frac{1}{2} (A_{1234} - A_{1243}).$$

It satisfies (see Eq. 3.22 in Rosati's thesis)

$$\hat{A} = \frac{1}{2} \sum_{1234} \mathcal{A}_{1234} c_1^\dagger c_2^\dagger c_4 c_3.$$

Since $\hat{A}^\dagger = \frac{1}{2} \sum_{3456} \mathcal{A}_{3456}^* c_5^\dagger c_6^\dagger c_4 c_3$,

$$\begin{aligned} [\hat{A}^\dagger, \hat{c}_2^\dagger \hat{c}_1] &= \left(\frac{1}{2} \sum_{356} \mathcal{A}_{3256}^* c_5^\dagger c_6^\dagger c_1 c_3 + \frac{1}{2} \sum_{456} \mathcal{A}_{1456}^* c_5^\dagger c_6^\dagger c_4 c_2 \right. \\ &\quad \left. - \frac{1}{2} \sum_{346} \mathcal{A}_{3416}^* c_2^\dagger c_6^\dagger c_4 c_3 - \frac{1}{2} \sum_{345} \mathcal{A}_{3451}^* c_5^\dagger c_1^\dagger c_4 c_3 \right) \\ &= \sum_{345} \mathcal{A}_{3245}^* c_4^\dagger c_5^\dagger c_1 c_3 - \sum_{345} \mathcal{A}_{3415}^* c_2^\dagger c_5^\dagger c_4 c_3. \end{aligned} \quad (4)$$

The above euqaiton is the same as Eq. A5 in PHYSICAL REVIEW B 90, 125140 (2014). Insert Eq. 3 and 4 into Eq. 2,

$$\begin{aligned} \frac{d\rho_{12}}{dt} &= \frac{1}{2} \text{Tr} \left([\hat{A}^\dagger, \hat{c}_2^\dagger \hat{c}_1] \hat{A} \hat{\rho} \right) + H.C. \\ &= \frac{1}{4} \text{Tr} \left(\left(\sum_{345} \mathcal{A}_{3245}^* c_4^\dagger c_5^\dagger c_1 c_3 - \sum_{345} \mathcal{A}_{3415}^* c_2^\dagger c_5^\dagger c_4 c_3 \right) \sum_{6789} \mathcal{A}_{6789} c_6^\dagger c_7^\dagger c_9 c_8 \hat{\rho} \right) + H.C. \\ &= \frac{1}{4} \sum_{3456789} (\mathcal{A}_{3245}^* k_{45136798} - \mathcal{A}_{3415}^* k_{25436798}) \mathcal{A}_{6789} + H.C., \\ k_{12345678} &= \text{Tr} \left(c_1^\dagger c_2^\dagger c_3 c_4 c_5^\dagger c_6^\dagger c_7 c_8 \hat{\rho} \right). \end{aligned} \quad (5)$$

To close the equation, we first approximate $k_{12345678}$,

$$k_{12345678} \approx \text{Tr} \left(c_7 c_8 \hat{\rho} c_1^\dagger c_2^\dagger \right) \text{Tr} \left(c_5^\dagger c_6^\dagger \hat{\rho} c_3 c_4 \right).$$

In the mean-field approximation (or Hartree-Fock approximation, see Eq. 74 in Reviews of Modern Physics 74, 895 (2002)),

$$\begin{aligned} \text{Tr} \left(c_7 c_8 \hat{\rho} c_1^\dagger c_2^\dagger \right) &\approx \text{Tr} \left(\hat{\rho} c_1^\dagger c_8 \right) \text{Tr} \left(\hat{\rho} c_2^\dagger c_7 \right) - \text{Tr} \left(\hat{\rho} c_2^\dagger c_8 \right) \text{Tr} \left(\hat{\rho} c_1^\dagger c_7 \right) \\ &= \rho_{81} \rho_{72} - \rho_{82} \rho_{71}, \\ \text{Tr} \left(c_5^\dagger c_6^\dagger \hat{\rho} c_3 c_4 \right) &\approx \text{Tr} \left(\hat{\rho} c_3 c_6^\dagger \right) \text{Tr} \left(\hat{\rho} c_4 c_5^\dagger \right) - \text{Tr} \left(\hat{\rho} c_4 c_6^\dagger \right) \text{Tr} \left(\hat{\rho} c_3 c_5^\dagger \right) \\ &= (I - \rho)_{36} (I - \rho)_{45} - (I - \rho)_{46} (I - \rho)_{35}, \end{aligned}$$

using the relation $\hat{c}_1 \hat{c}_2^\dagger = \delta_{12} - \hat{c}_2^\dagger \hat{c}_1$. Insert the above three equations in Eq. 5,

$$\frac{d\rho_{12}}{dt} = \frac{1}{4} \sum_{3456789} \left(\begin{aligned} &\mathcal{A}_{3245}^* [(I - \rho)_{17} (I - \rho)_{36} - (I - \rho)_{37} (I - \rho)_{16}] (\rho_{84} \rho_{95} - \rho_{85} \rho_{94}) \\ &- \mathcal{A}_{3415}^* [(I - \rho)_{47} (I - \rho)_{36} - (I - \rho)_{37} (I - \rho)_{46}] (\rho_{82} \rho_{95} - \rho_{85} \rho_{92}) \end{aligned} \right) \mathcal{A}_{6789} + H.C..$$

Use the relations $\mathcal{A}_{1234} = \mathcal{A}_{2143} = -\mathcal{A}_{2134} = -\mathcal{A}_{1243}$,

$$\frac{d\rho_{12}}{dt} = \frac{1}{4} \sum_{3456789} \left\{ \begin{aligned} &(I - \rho)_{17} (I - \rho)_{36} \mathcal{A}_{7698} \mathcal{A}_{2354}^* \rho_{84} \rho_{95} + (I - \rho)_{17} (I - \rho)_{36} \mathcal{A}_{7698} \mathcal{A}_{2345}^* \rho_{85} \rho_{94} \\ &+ (I - \rho)_{16} (I - \rho)_{37} \mathcal{A}_{6789} \mathcal{A}_{2345}^* \rho_{95} \rho_{84} + (I - \rho)_{16} (I - \rho)_{37} \mathcal{A}_{6789} \mathcal{A}_{2354}^* \rho_{85} \rho_{94} \\ &- (I - \rho)_{36} [(I - \rho)_{74} \mathcal{A}_{3415} \mathcal{A}_{6789}^* \rho_{59}]^* \rho_{82} - (I - \rho)_{36} [(I - \rho)_{74} \mathcal{A}_{3415} \mathcal{A}_{6798}^* \rho_{58}]^* \rho_{92} \\ &- (I - \rho)_{37} [(I - \rho)_{64} \mathcal{A}_{3415} \mathcal{A}_{7689}^* \rho_{59}]^* \rho_{82} - (I - \rho)_{37} [(I - \rho)_{64} \mathcal{A}_{3415} \mathcal{A}_{7698}^* \rho_{58}]^* \rho_{92} \end{aligned} \right\} + H.C..$$

Notice that in the bracket “{” of the above equation, the first 4 terms are the same after sum and the last 4 terms are the same after sum, so that

$$\frac{d\rho_{12}}{dt} = \frac{1}{2} \sum_{3456789} \left\{ \begin{array}{l} (I - \rho)_{17} [2 (I - \rho)_{36} \mathcal{A}_{7698} \mathcal{A}_{2354}^* \rho_{84}] \rho_{95} \\ - (I - \rho)_{36} [2 (I - \rho)_{74} \mathcal{A}_{3415} \mathcal{A}_{6789}^* \rho_{59}]^* \rho_{82} \end{array} \right\} + H.C..$$

Define

$$P_{1234} = 2 \sum_{5678} (I - \rho)_{65} \mathcal{A}_{1537} \mathcal{A}_{2648}^* \rho_{78}.$$

We will have

$$\frac{d\rho_{12}}{dt} = \frac{1}{2} \sum_{795} (I - \rho)_{17} P_{7295} \rho_{95} - \sum_{368} (I - \rho)_{36} P_{3618}^* \rho_{82} + H.C..$$

This is the same as Eq. 21, 23, 24 in PHYSICAL REVIEW B 92, 235423 (2015).