

# Link

## Knot soliton in models with B-L and Peccei-Quinn symmetries

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Yu Hamada (KEK)

arXiv: 2301.XXXXX (work in progress)

w/ M. Eto (Yamagata U.) and M. Nitta (Keio U.)



# Introduction

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# Topological Soliton

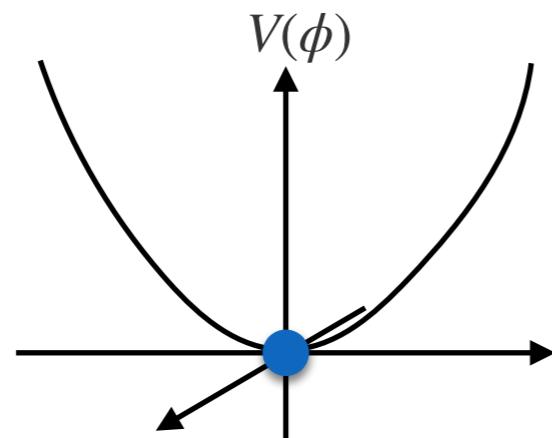
- **Non-perturbative object in field theories**
  - monopole, vortex string, skyrmion, instanton, etc..

# Topological Soliton

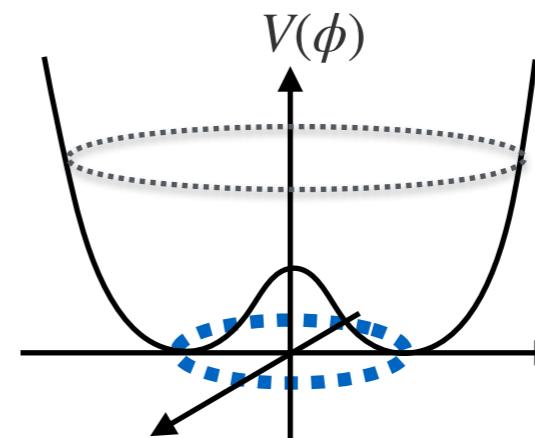
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# Topological Soliton

- Non-perturbative object in field theories
  - monopole, vortex string, skyrmion, instanton, etc..
- It appears if **vacuum has non-trivial topology.**



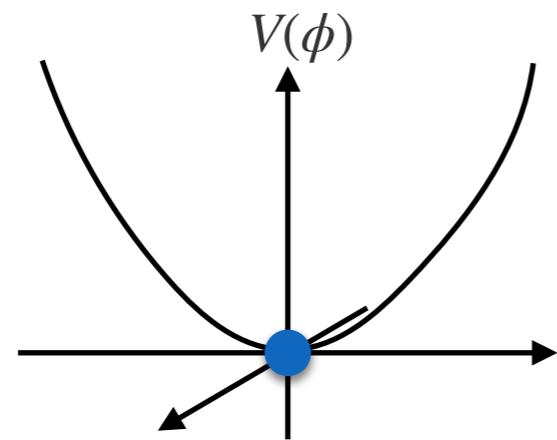
**trivial**  
(point)



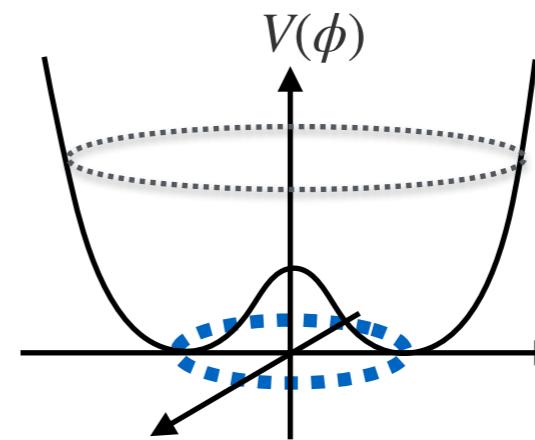
**non-trivial**  
(circle  $S^1$ )

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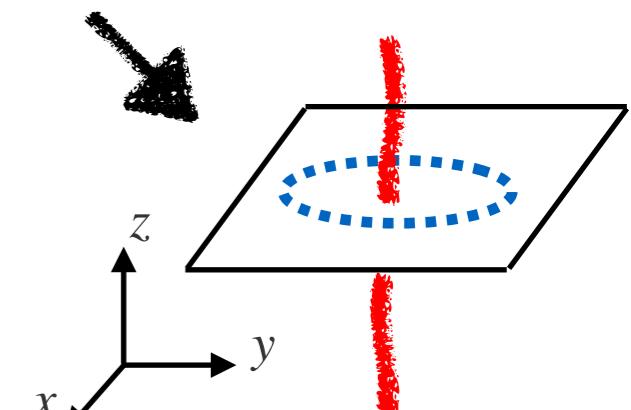
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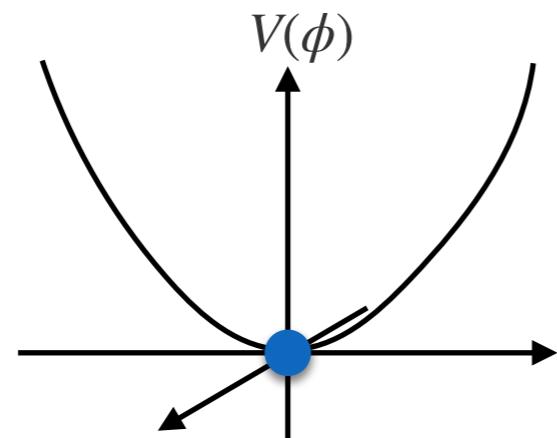
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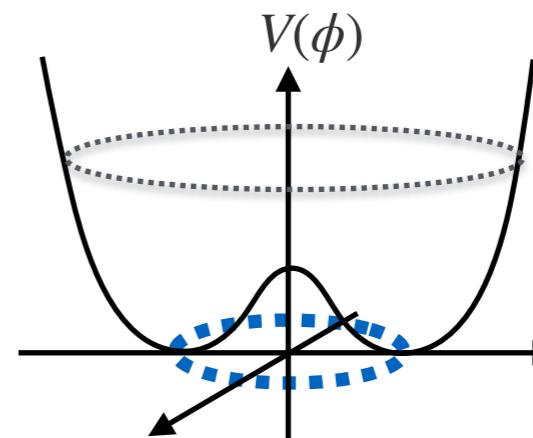
vortex string

# Topological Soliton

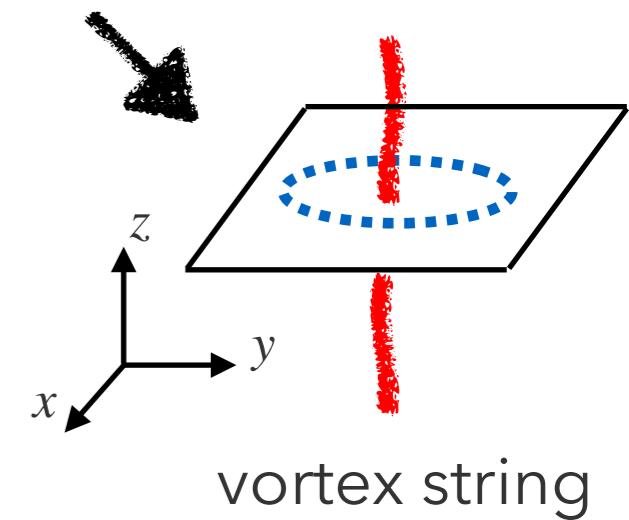
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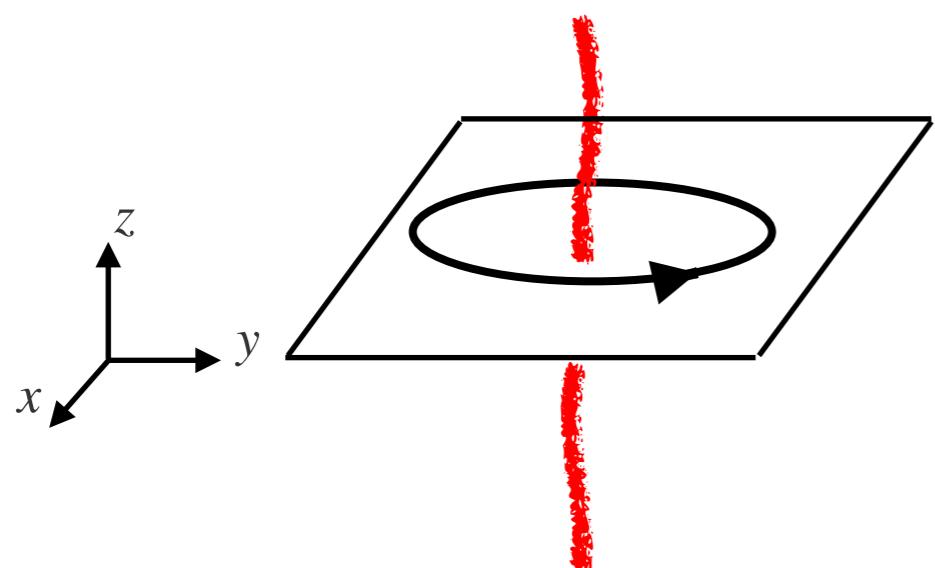
- Vortex string appears in many systems:
  - cosmic string, superconductor, neutron star, etc.

# Global vs Local strings

- SSB of **gauged**  $U(1)$  sym  $\rightarrow$  **local** vortex string
  - $\rightarrow$  magnetic flux is squeezed w/ finite width

- SSB of **global**  $U(1)$  sym  $\rightarrow$  **global** vortex string

$\rightarrow$  w/o magnetic flux



NG boson phase changes from 0 to  $2\pi$

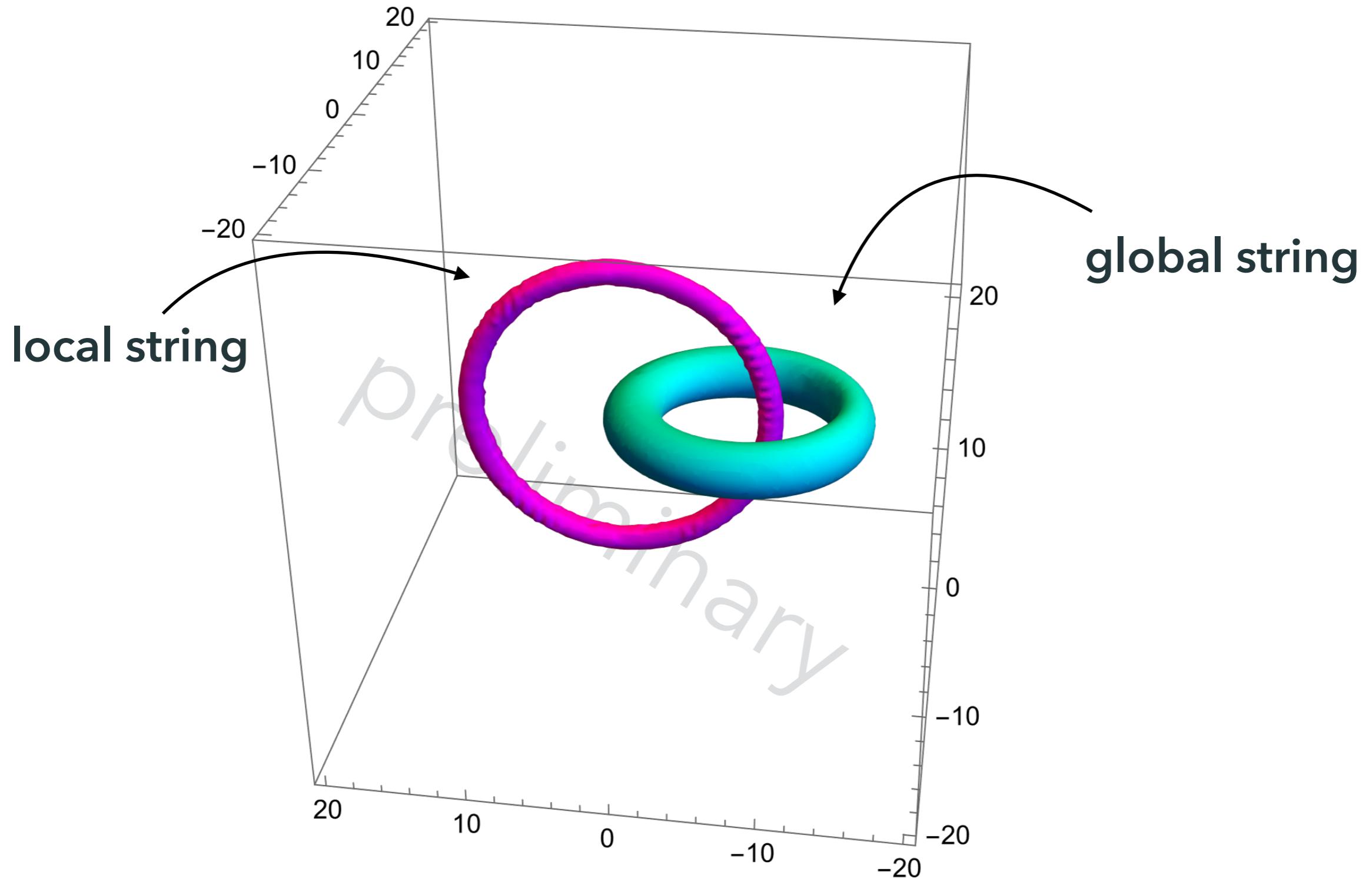
# Global vs Local strings



# Global vs Local strings



# Link soliton

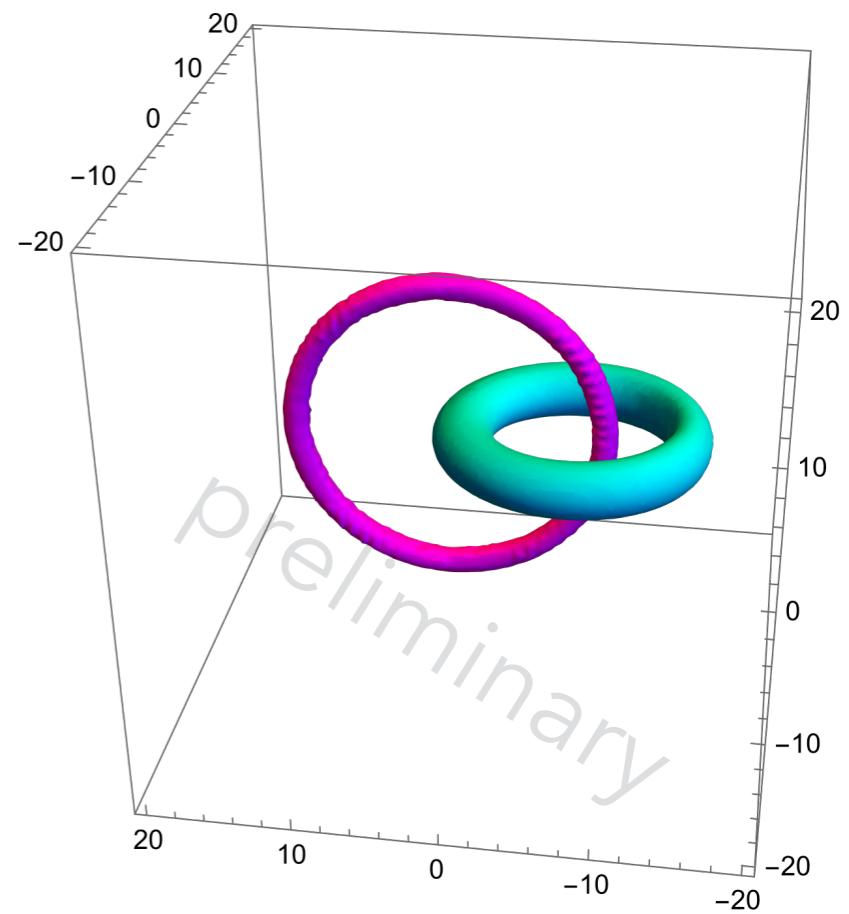


link soliton made of local & global strings!

# Link soliton

- stable solution of EOM (!)
- Key: Chern-Simons coupling like  $\frac{c}{16\pi^2} \int d^4x aF\tilde{F}$
- 1st example of link soliton in gauge theory!
- Natural choice:
$$\begin{cases} U(1)_{global} \rightarrow U(1)_{PQ} \\ U(1)_{gauge} \rightarrow U(1)_{B-L} \end{cases}$$
- generally, global string can contain small flux  
→ linking flux, applicable for baryogenesis

**link = origin of matter**



# Plan of talk

- Introduction
- Vortices w/ CS coupling (review)
- Stability of link soliton
- Baryogenesis
- Summary

# Vortices w/ CS coupling (review)

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[e.g., Horvathy-Zhang '08]

# Abelian-Higgs w/ Chern-Simons

Let's start from 2+1D Abelian-Higgs w/ CS term:

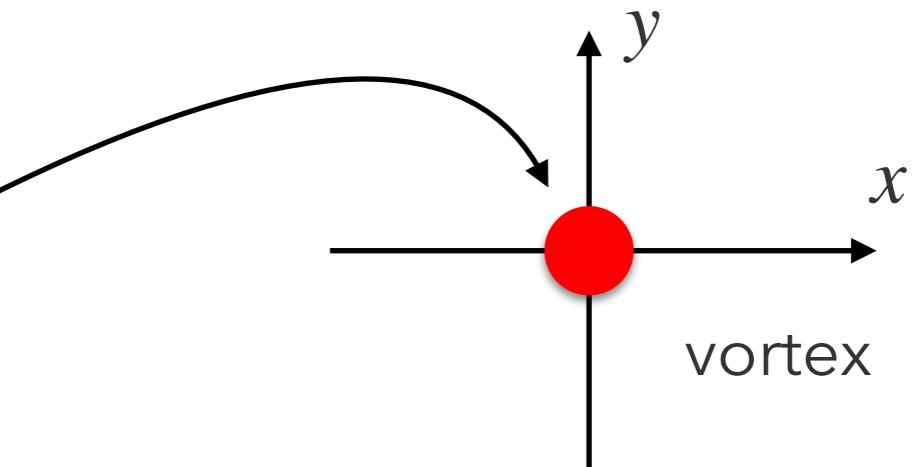
$$\mathcal{L} = |D_\mu \phi|^2 - \frac{1}{4} F_{\mu\nu}^2 - V(\phi) + g^2 c \epsilon^{\mu\nu\lambda} A_\mu \partial_\nu A_\lambda$$

$$V(\phi) = -m^2 |\phi|^2 + \lambda |\phi|^4$$

- For  $c = 0$ ,  $A_0$  is decoupled for static configurations.

→ static solution:  $\begin{cases} \phi = v f(r) e^{i\theta} \\ A_\theta = a(r)/g \\ A_0 = A_r = 0 \end{cases} \quad \begin{cases} f(0) = 0, f(\infty) = 1 \\ a(0) = 0, a(\infty) = 1 \end{cases}$

→ quantized magnetic flux  $\int d^2x B = 2\pi/g$



# Chern-Simons vortex

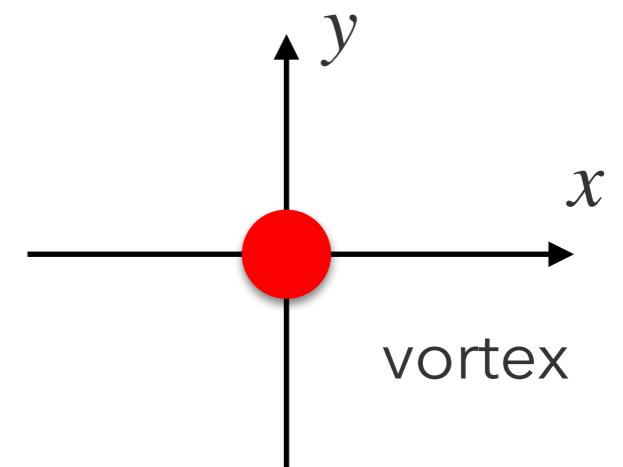
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- For  $c \neq 0$ ,  $A_0$  is NOT, due to Gauss law constraint:

$$\frac{\delta \mathcal{L}}{\delta A_0} = \partial_i E_i - g^2 J^0 + g^2 c B = 0 \quad \begin{aligned} E_i &= \partial_i A_0 \\ J^0 &\equiv \phi^\dagger i D^0 \phi + (h.c.) \end{aligned}$$

→ magnetic flux sauces electric field!



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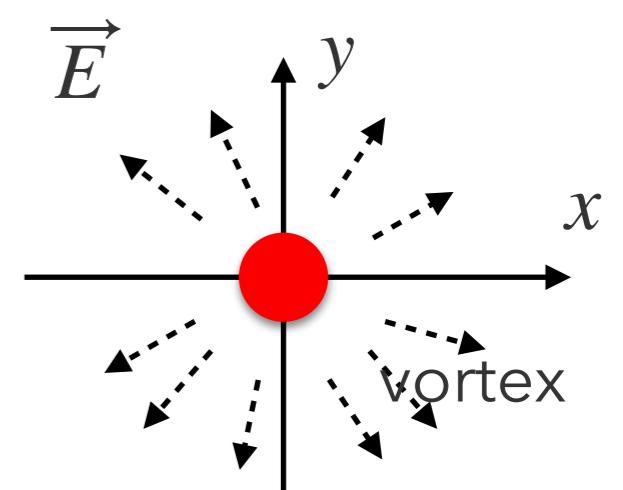
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$$E_i = \partial_i A_0$$
$$J^0 \equiv \phi^\dagger i D^0 \phi + (h.c.)$$

→ magnetic flux sauces electric field!

- quantized magnetic flux & electric charge

$$\int d^2x B = 2\pi/g \quad \int d^2x J^0 = 2\pi c/g$$



# Chern-Simons vortex

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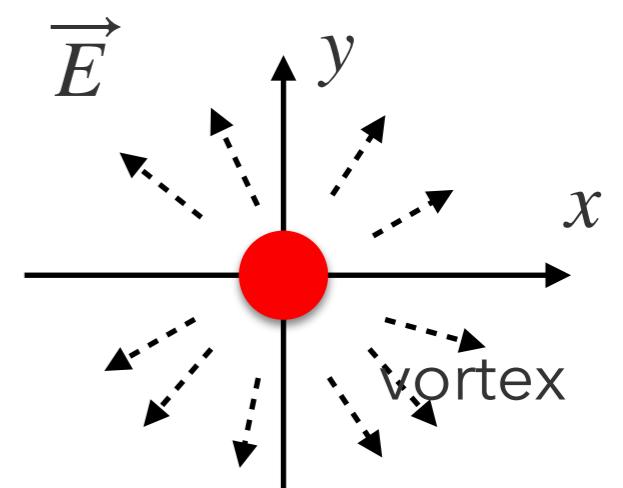
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$$\int d^2x B = 2\pi/g$$

$$\int d^2x J^0 = 2\pi c/g$$

called Chern-Simons vortex



# Interaction of Chern-Simons vortices

- Single static solution:  $\begin{cases} \phi = v f(r) e^{i\theta} \\ A_\theta = a(r)/g \\ A_0 = b(r)/g \end{cases}$   $\begin{cases} f(0) = 0, f(\infty) = 1 \\ a(0) = 0, a(\infty) = 1 \\ b(0) = 0, b(\infty) = 0 \end{cases}$
- Asymptotic behavior at  $r \rightarrow \infty$ 
$$1 - f(r) \sim e^{-M_\phi r} \quad b(r) \sim 1 - a(r) \sim e^{-M_c r}$$
w/  $M_c \equiv gv \left( \frac{1}{2}\sqrt{4+c^2} - \frac{c}{2} \right) \simeq \mathcal{O}(gv/c)$  for  $c \rightarrow \infty$

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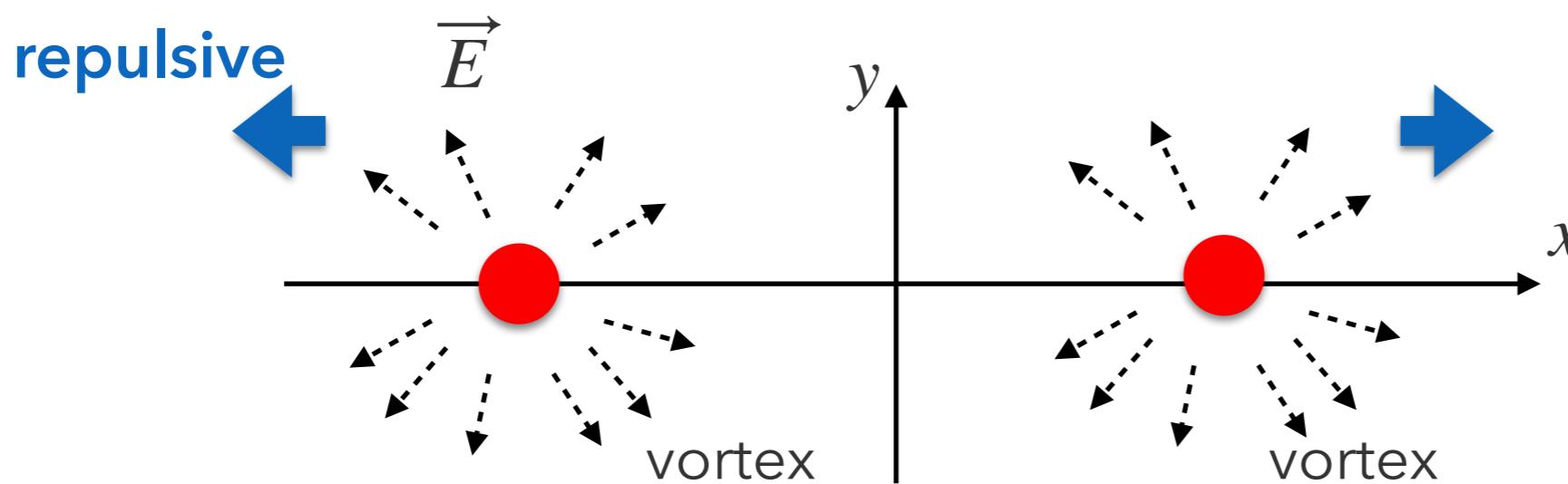
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- For large  $c$ , **long-range repulsive force!**



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# Stability of link soliton

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# The model

3+1D theory:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |\partial_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2)$$

$$D_\mu \phi_1 = (\partial_\mu - igA_\mu) \phi_1$$

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- Symmetries:

$$U(1)_{gauge} : \phi_1 \rightarrow e^{i\theta_1} \phi_1 \quad U(1)_{global} : \phi_2 \rightarrow e^{i\theta_2} \phi_2$$

- For  $\kappa > 0$  &  $\lambda > 0$ , both symmetries are broken at the vacuum:

$$\langle \phi_1 \rangle = v_1, \langle \phi_2 \rangle = v_2$$

→ global & local strings

# The model

3+1D theory:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |\partial_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2)$$

CS coupling

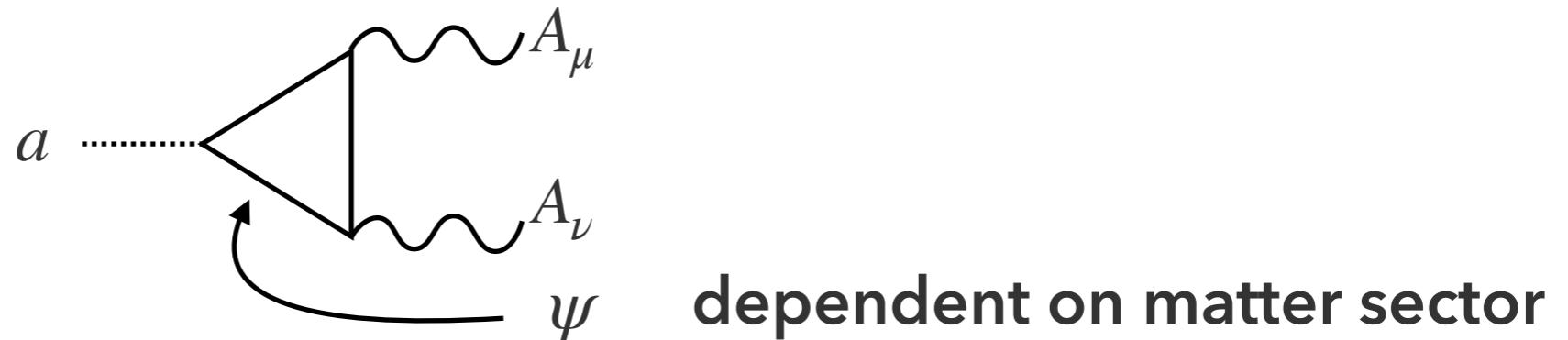
$$+\frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

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$$a \equiv -i \arg(\phi_2)$$

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- At the broken phase, CS coupling is induced by triangle anomaly.



$$\Rightarrow c = \sum_f Q_{global}^f (Q_{gauge}^f)^2$$

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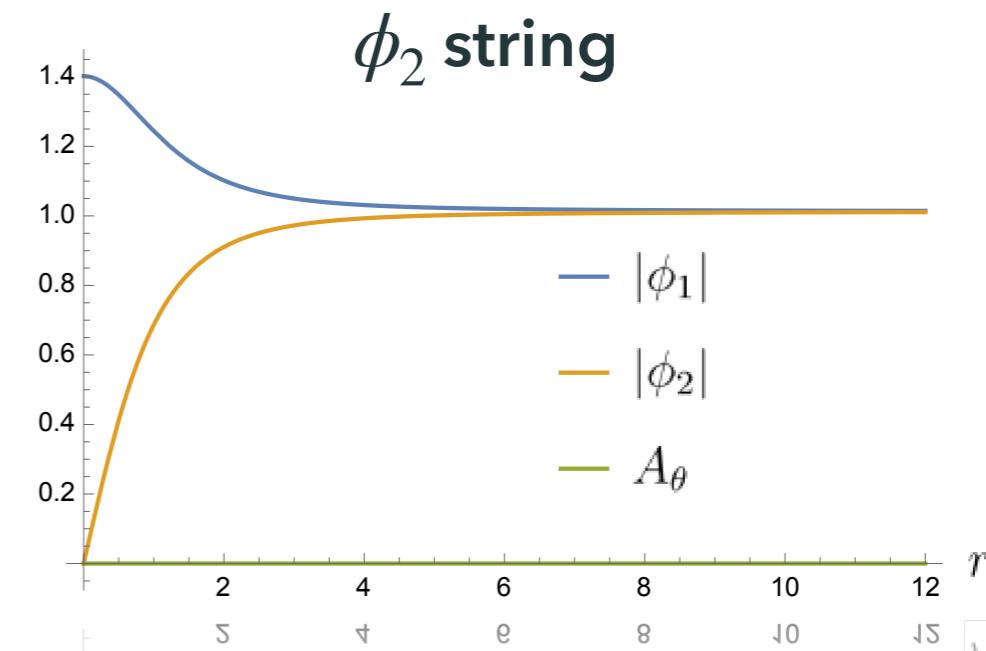
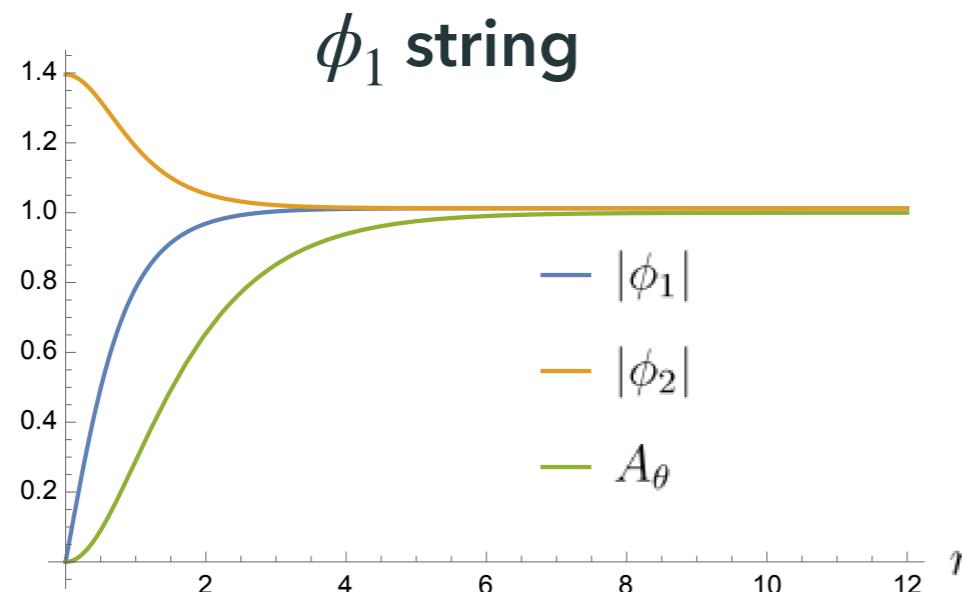
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- CS coupling does not affect single strings.



# Charged string

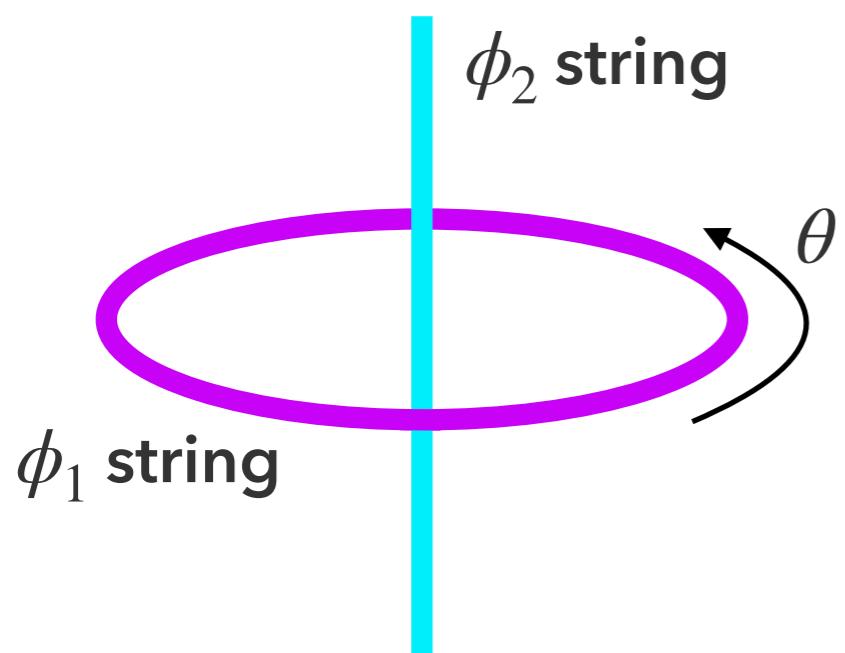
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$$a = \theta \Rightarrow (\partial_i a) A_0 B^i = \frac{1}{R} A_0 |\vec{B}|$$



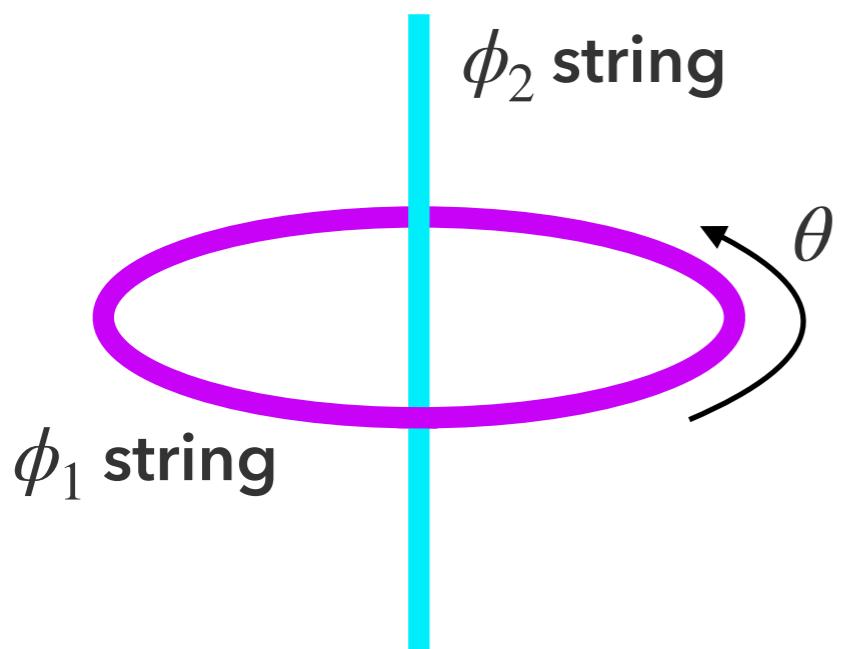
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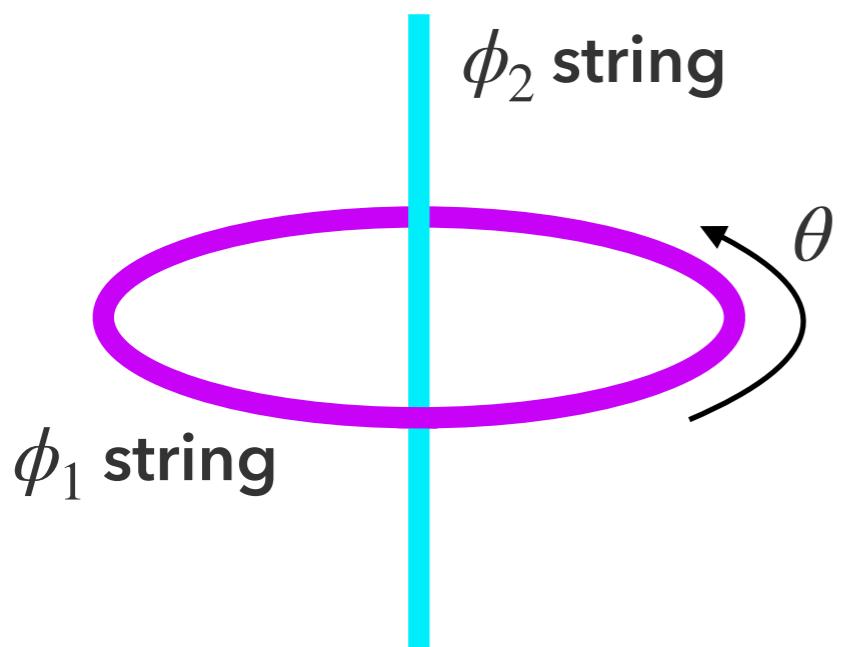
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$$\Rightarrow \int d^3x J^0 = 2\pi R \int d^2x \frac{c}{16\pi^2 R} |\vec{B}| = \frac{c}{4g}$$

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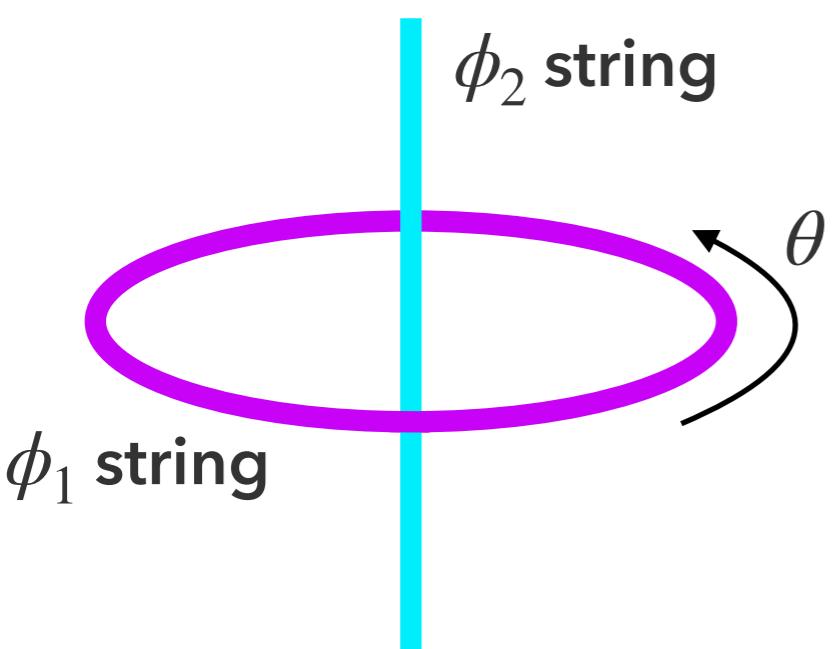
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ϕ₁ string is electrically charged! → doesn't shrink

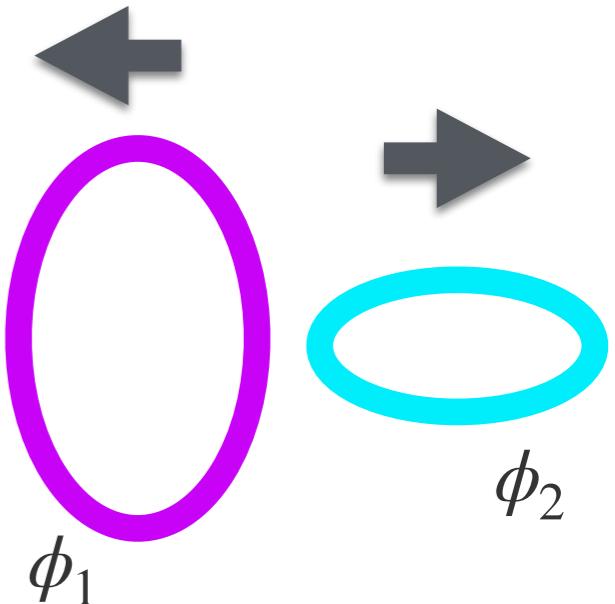
# Stability

- Delinking by passing through each other?

→ prevented by taking  $\lambda \gg g^2, \kappa, \chi$

**Overlap of strings ( $\phi_1 = \phi_2 = 0$ ) cost large energy**

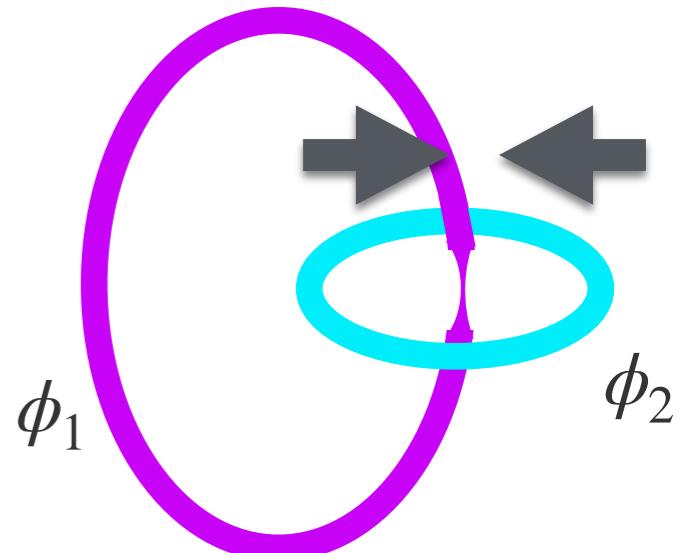
$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$



- $\phi_2$  string is not charged and thus can shrink ?

→ prevented by taking  $v_2/v_1 \ll 1$

**$\phi_2$  string is too light to pinch  $\phi_1$  string**



# Numerical calculation

Energy:

$$\begin{aligned}\mathcal{E} = & |D_i \phi_1|^2 + |\partial_i \phi_2|^2 + V(\phi_1, \phi_2) + \frac{1}{2g^2} (\partial_i A_j)^2 \\ & - g^2 |\phi_1|^2 A_0^2 - \frac{1}{2g^2} (\partial_i A_0)^2 - \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}\end{aligned}$$

- Not positive definite  $\rightarrow$  remove  $A_0$  by solving Gauss law:

$$\frac{\delta \mathcal{L}}{\delta A_0} = \partial^2 A_0 - 2g^2 |\phi_1|^2 A_0 + \frac{g^2 c}{16\pi^2} (\vec{\nabla} a) \cdot \vec{B} = 0$$

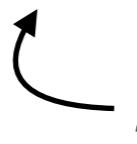
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 $\sim M_c^2 A_0 \sim \mathcal{O}(g^2 v_1^2 / c^2) A_0 \text{ for large } c$

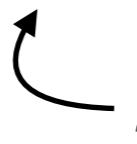
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$$\therefore A_0 \approx \frac{g^2 c}{16\pi^2} \frac{(\vec{\nabla} a) \cdot \vec{B}}{2g^2 |\phi_1|^2}$$

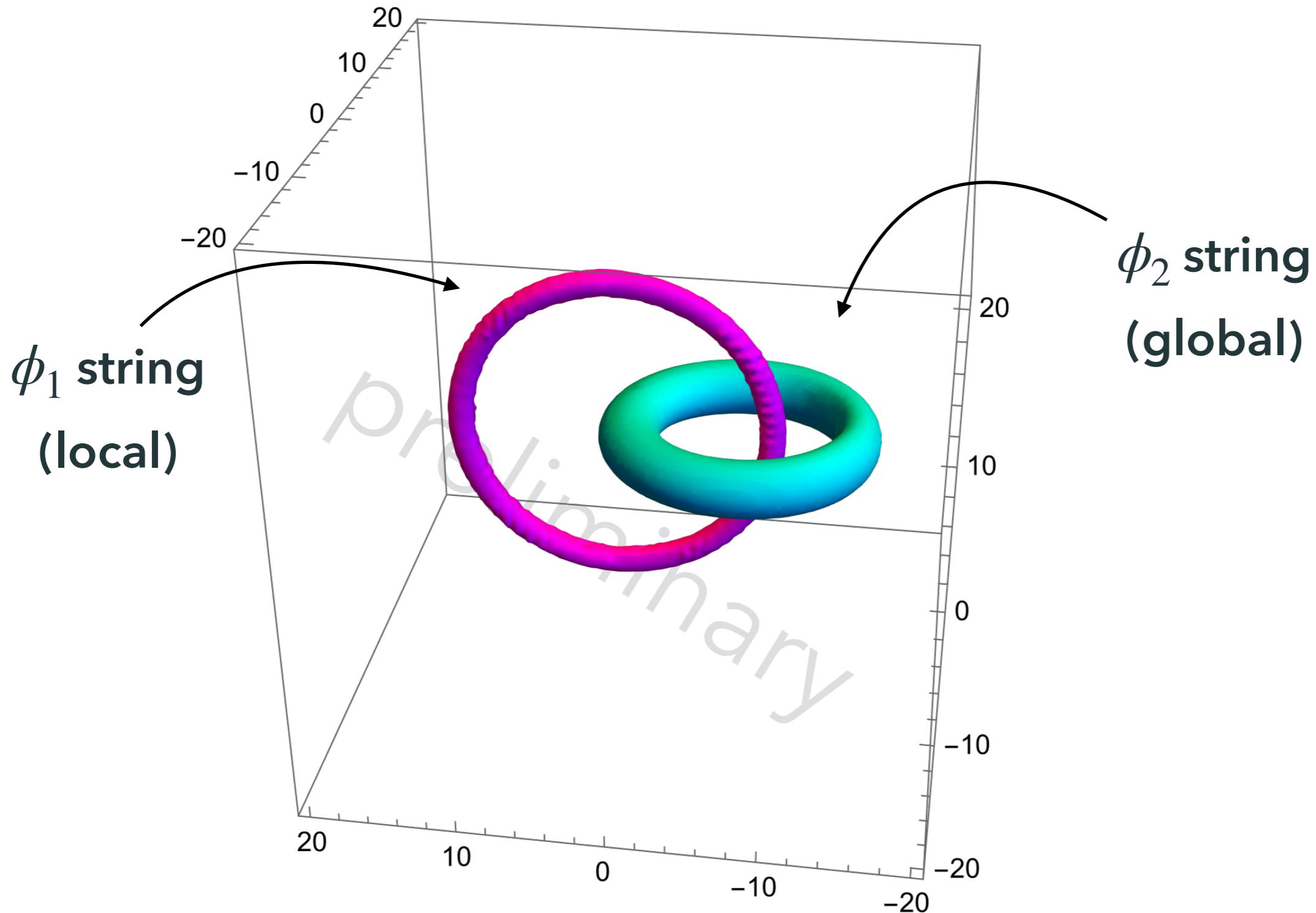
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- positive definite -> no obstacle
- Minimizing energy via non-linear conjugate gradient method
- CPU 400-cores parallelizing on YITP computer cluster
- lattice spacing =  $0.4/gv_1$ ,  $N = 100^3$ , converged w/ O(1) hours

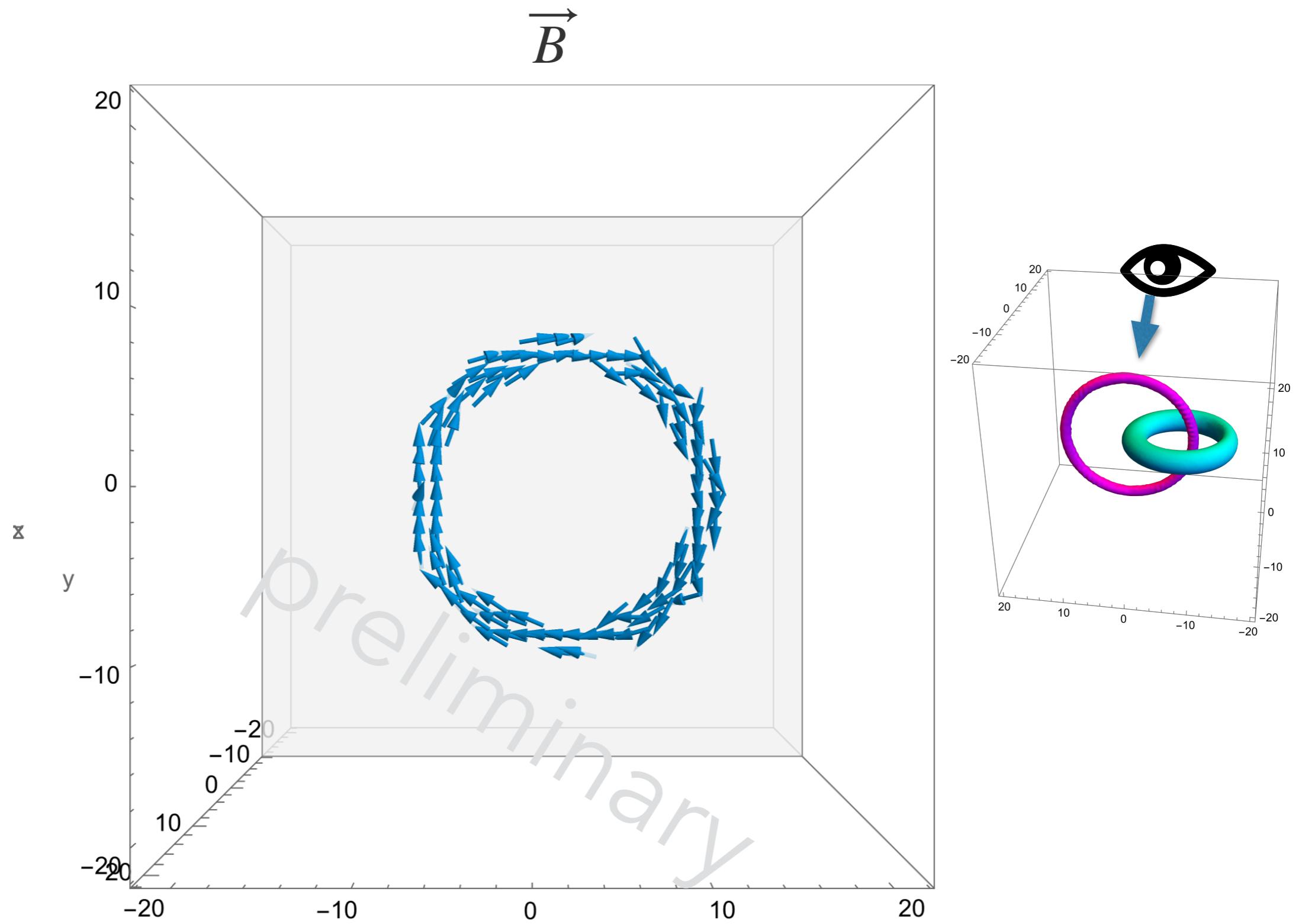
# Numerical solution



$$\lambda/g^2 = 100, \chi/g^2 = 19.5, \kappa/g^2 = 0.1, v_2/v_1 = 0.05$$

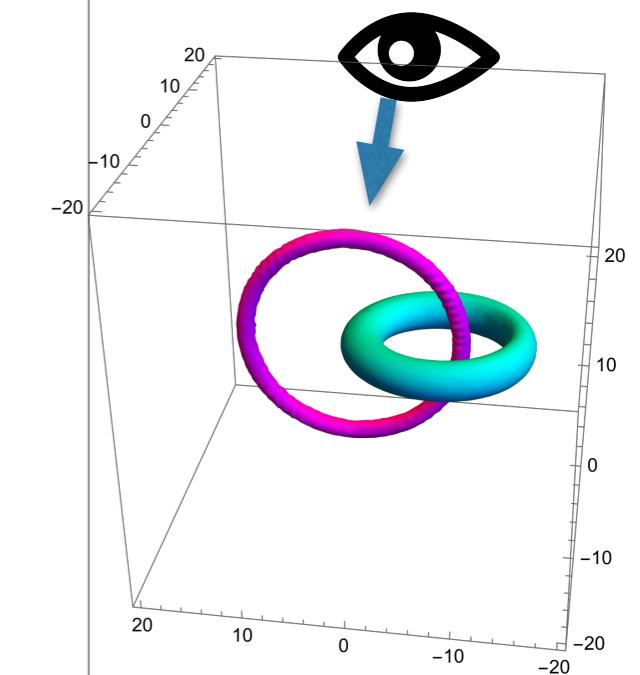
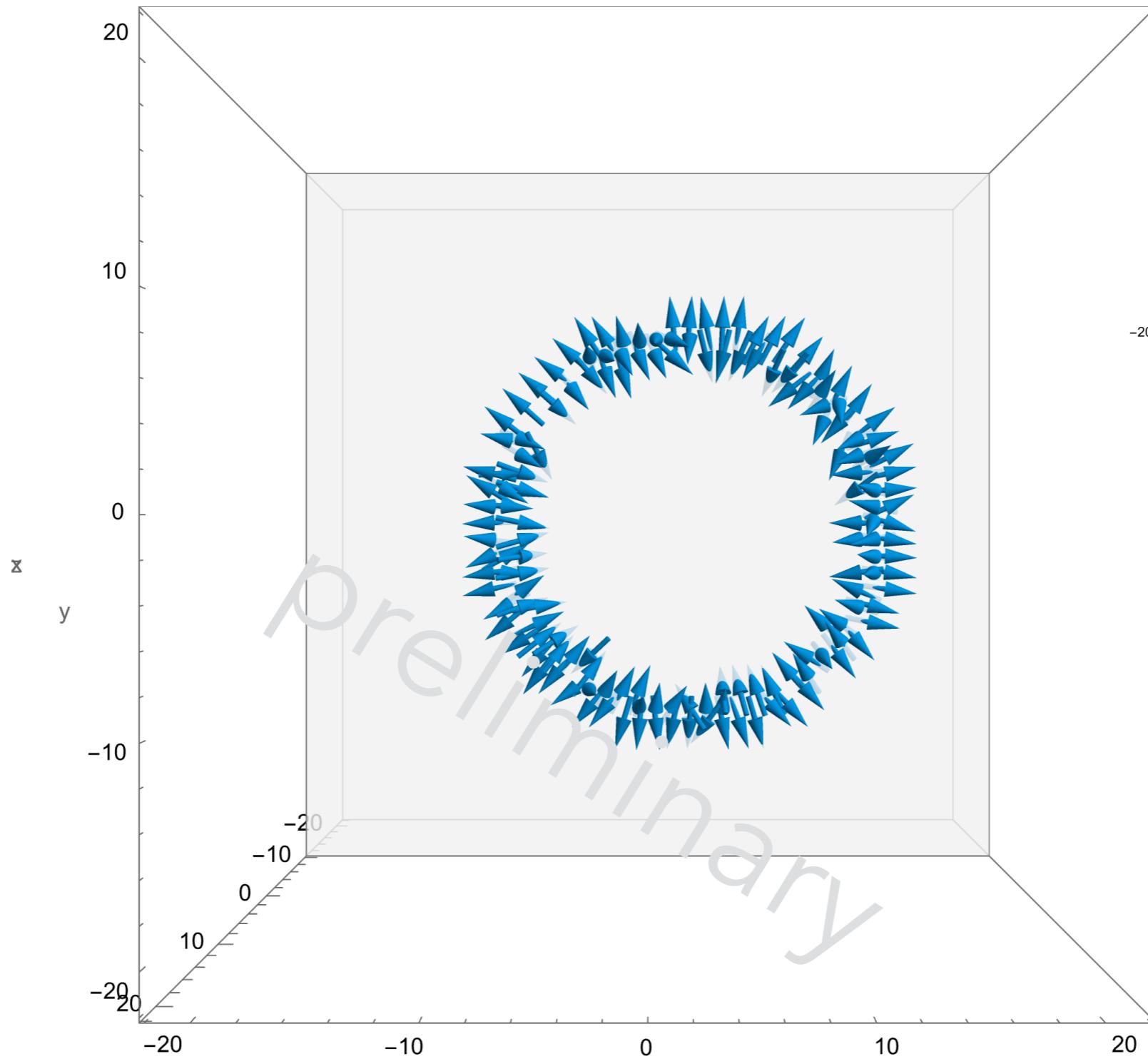
$$g^2 c/(16\pi^2) = 16$$

# Magnetic field



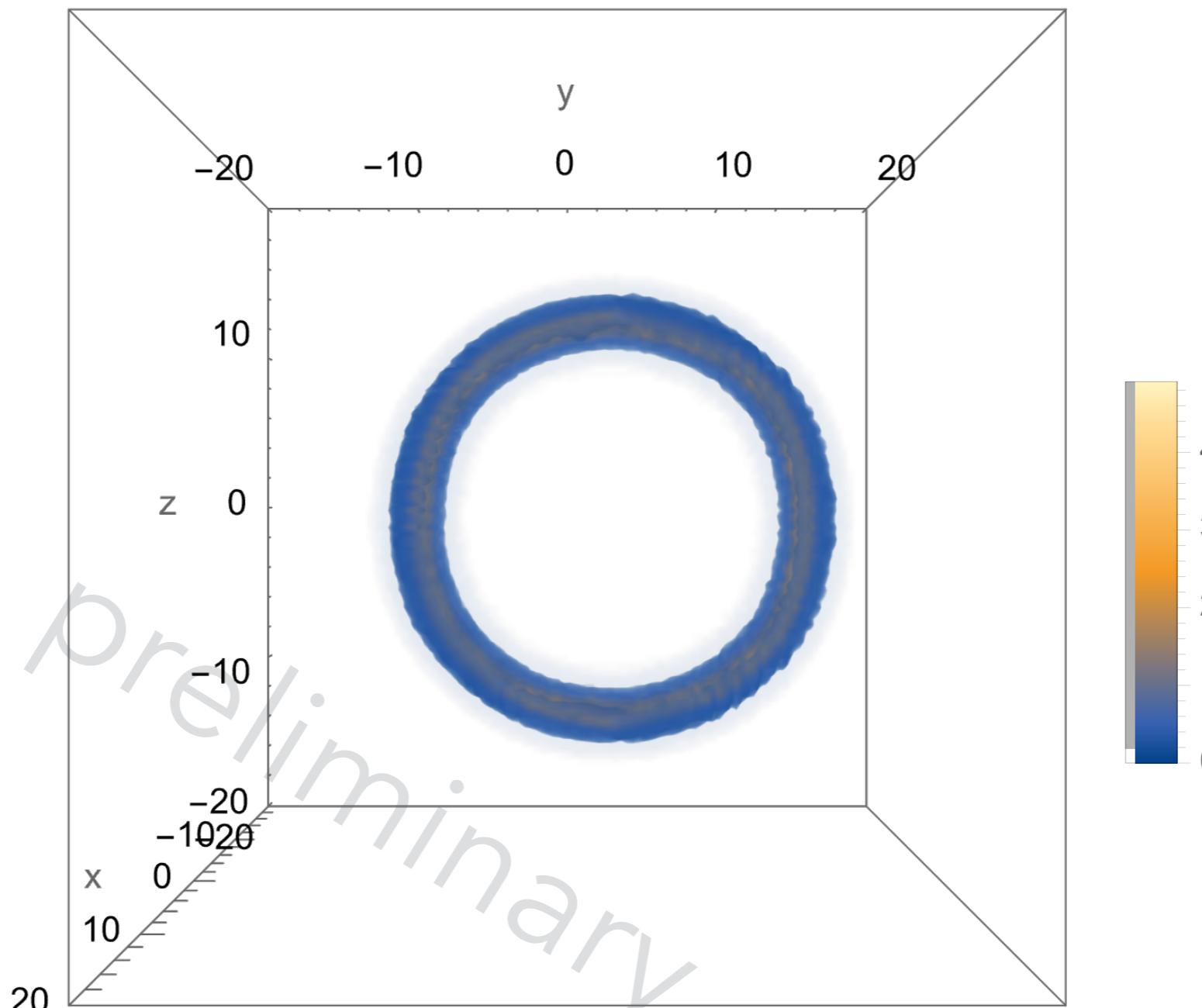
# Electric field

$$\vec{E} = \vec{\nabla} A_0$$

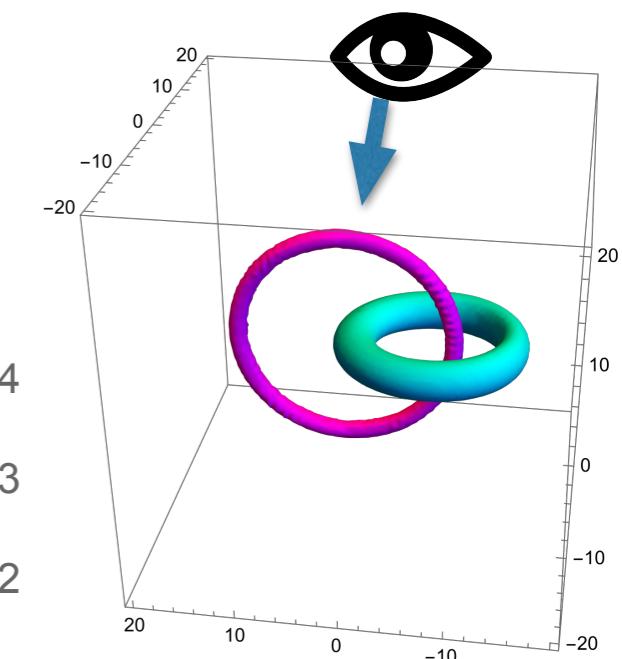


# Energy

Energy is dominated by  $\phi_1$  string



total energy:  $E \sim 401 g v_1$



# General setup

More general charge assignment:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2)$$

$$D_\mu \phi_1 = (\partial_\mu - igq_1 A_\mu) \phi_1 \quad D_\mu \phi_2 = (\partial_\mu - igq_2 A_\mu) \phi_2 \quad q_2/q_1 \ll 1$$

$$V(\phi) = \lambda \left( |\phi_1|^2 + |\phi_2|^2 - \mu^2 \right)^2 - \kappa |\phi_1|^2 |\phi_2|^2 + \chi |\phi_2|^4$$

- Symmetries:

$$U(1)_{gauge} : \begin{cases} \phi_1 \rightarrow e^{iq_1\theta_1} \phi_1 \\ \phi_2 \rightarrow e^{iq_2\theta_1} \phi_2 \end{cases}$$

$$U(1)_{global} : \begin{cases} \phi_1 \rightarrow e^{-iq_2\theta_1} \phi_1 \\ \phi_2 \rightarrow e^{iq_1\theta_1} \phi_2 \end{cases}$$

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→ Also  $\phi_2$  string contains magnetic flux, but the solution is almost the same.

# General setup

More general charge assignment:

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2) + \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

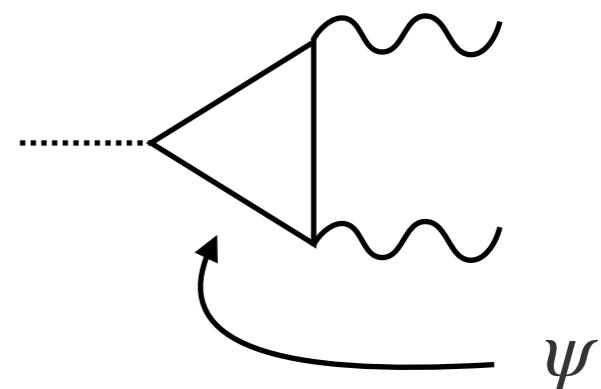
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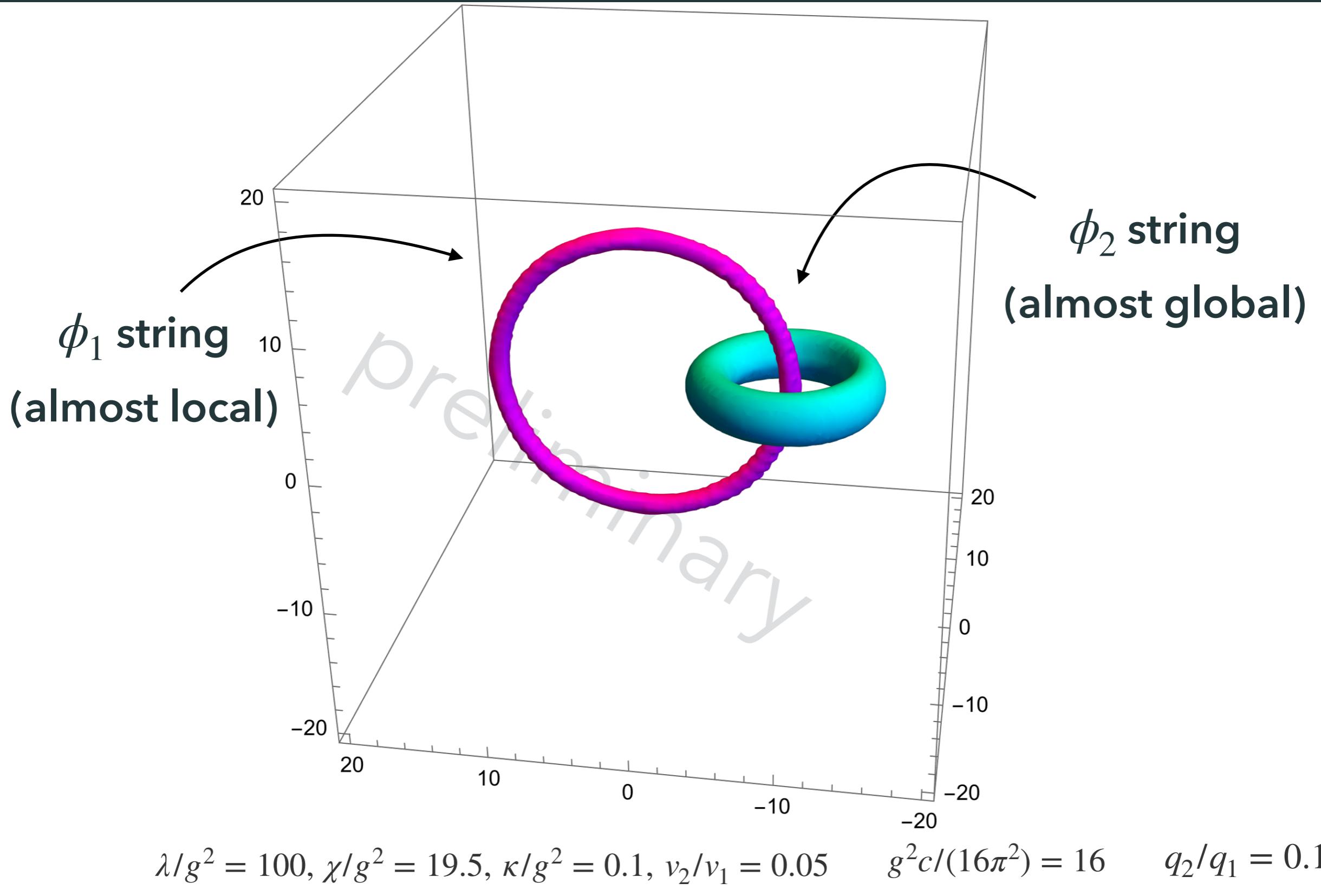
- Definition of "axion" is more complicated

$$a \equiv \frac{1}{i\sqrt{q_2^2 v_1^2 + q_1^2 v_2^2}} (-q_2 v_1 \arg(\phi_1) + q_1 v_2 \arg(\phi_2))$$

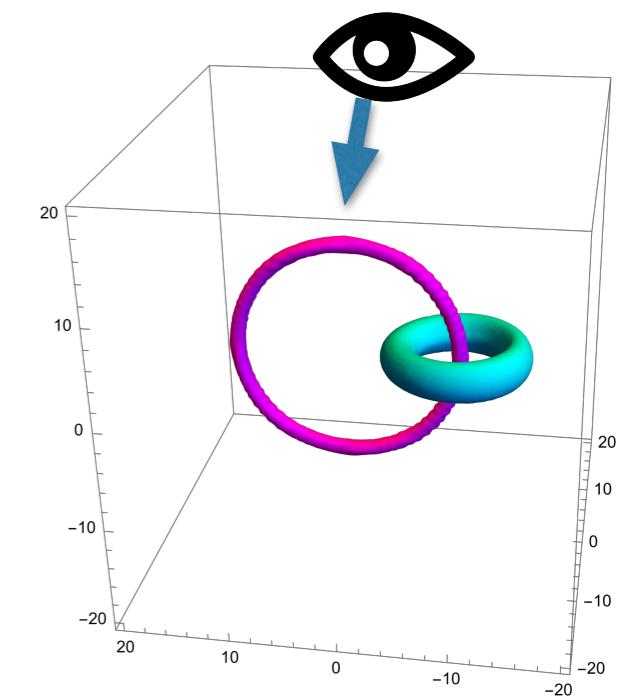
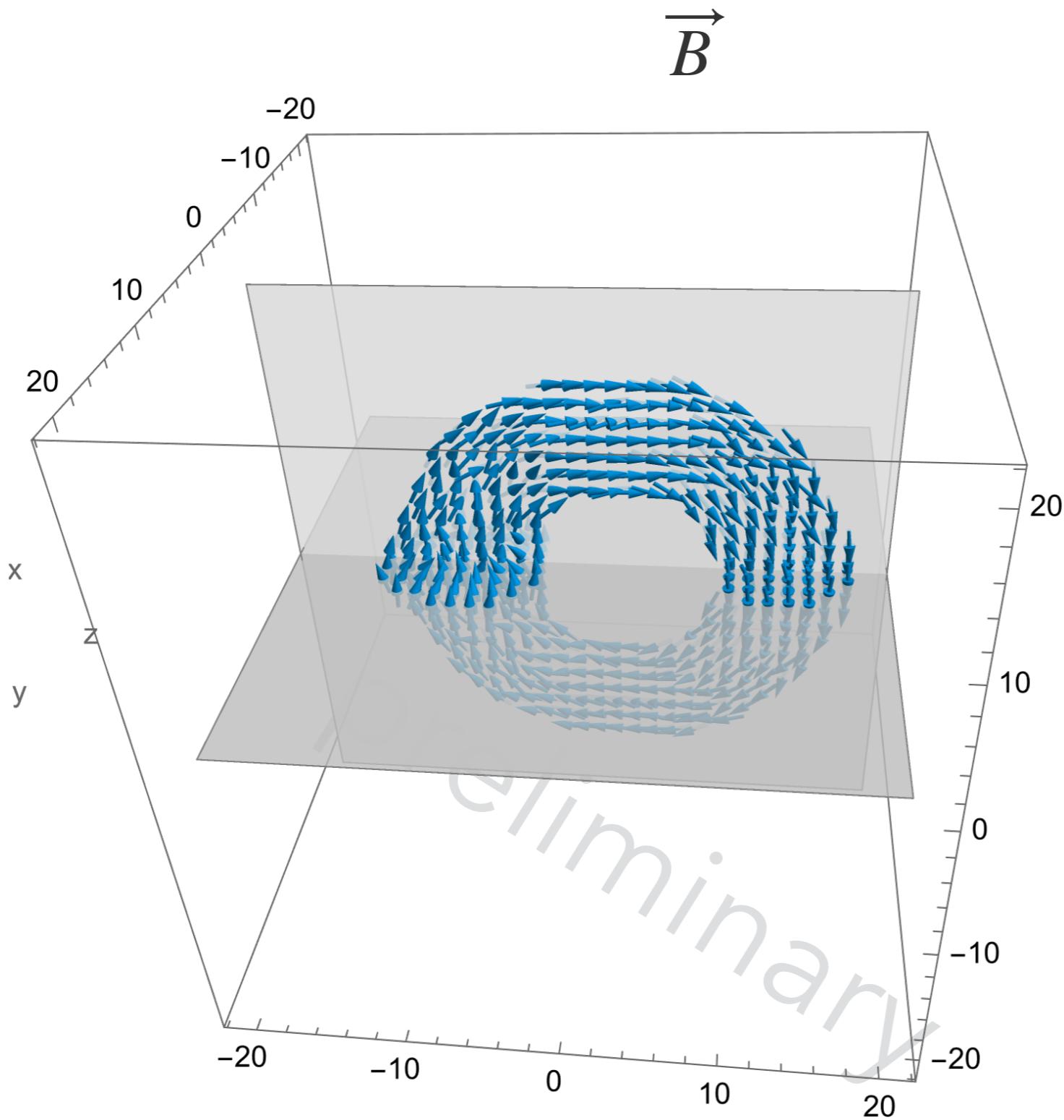
- Triangle anomaly is also complicated  
→  $c$  will be taken as free parameter.



# Numerical solution



# Magnetic field



# Helical magnetic field

- Since the magnetic fluxes are linked, this soliton has finite helicity (Chern-Simons number):

$$N_{CS}[A] \equiv \frac{1}{16\pi^2} \int d^3x A dA = \frac{1}{16\pi^2} \int d^3x \vec{A} \cdot \vec{B}$$

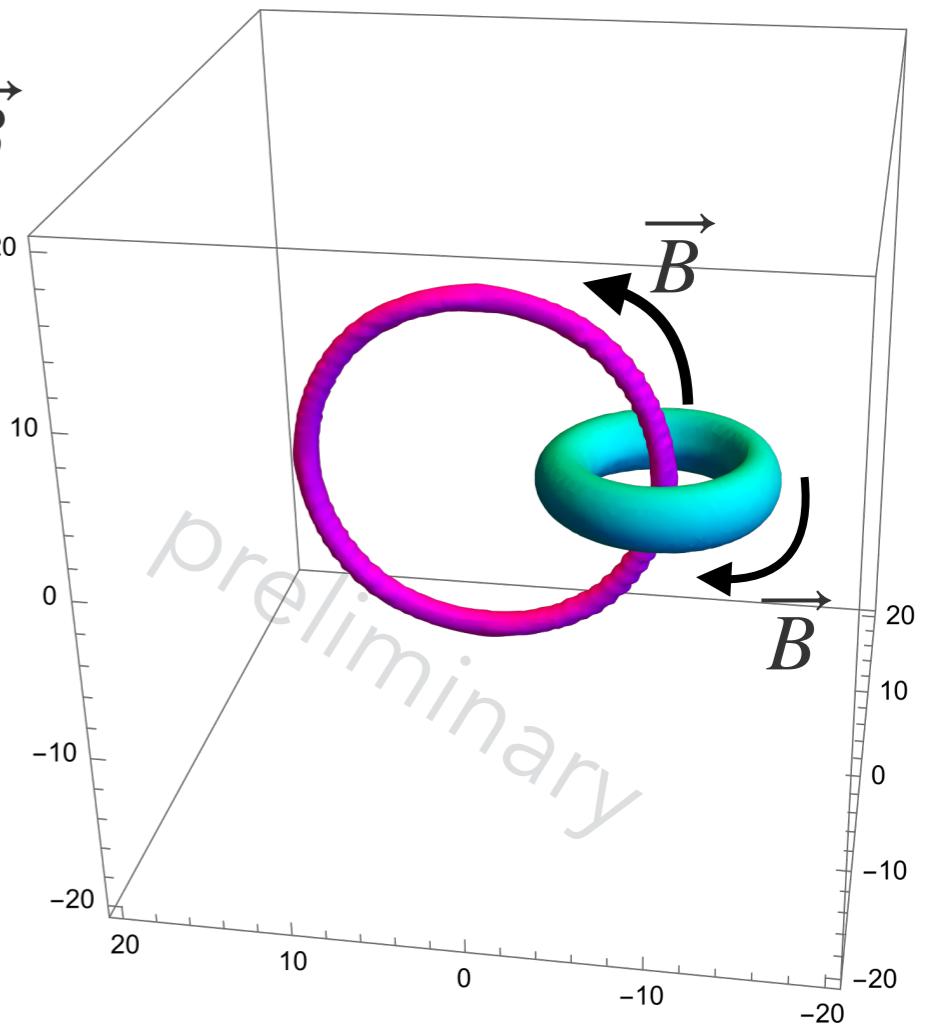
For  $q_1 = 1, q_2 = 0.1$ ,  $N_{CS}[A] \simeq 0.28$

( $N_{CS} \simeq 0$  for  $q_1 = 1, q_2 = 0$ )

**can be used for baryogenesis!**

(cf: baryogenesis by helical  $U(1)_Y$  field)

[Kamada-Long '16]



# Plan of talk

- Introduction
- Vortices w/ CS coupling (review)
- Stability of link soliton
- Baryogenesis
- Summary

# Baryogenesis (work in progress)

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# The model

$$\mathcal{L} = |D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 - \frac{1}{4g^2} F_{\mu\nu}^2 - V(\phi_1, \phi_2) + \frac{c}{16\pi^2} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- Natural setup:  $U(1)_{gauge} = U(1)_{B-L}$  &  $U(1)_{global} = U(1)_{PQ}$

Type-I seesaw  $\rightarrow \nu$ -mass

QCD axion  $\rightarrow$  strong CP & Dark matter

$$\Rightarrow v_1 \sim v_2 \sim 10^{9-12} \text{ GeV}$$

- Axion quality problem can be avoided by gauged PQ mechanism.

$$\Rightarrow q_1 = 1, q_2 = 0.1$$

[Fukuda-Ibe-Suzuki-Yanagida '17]

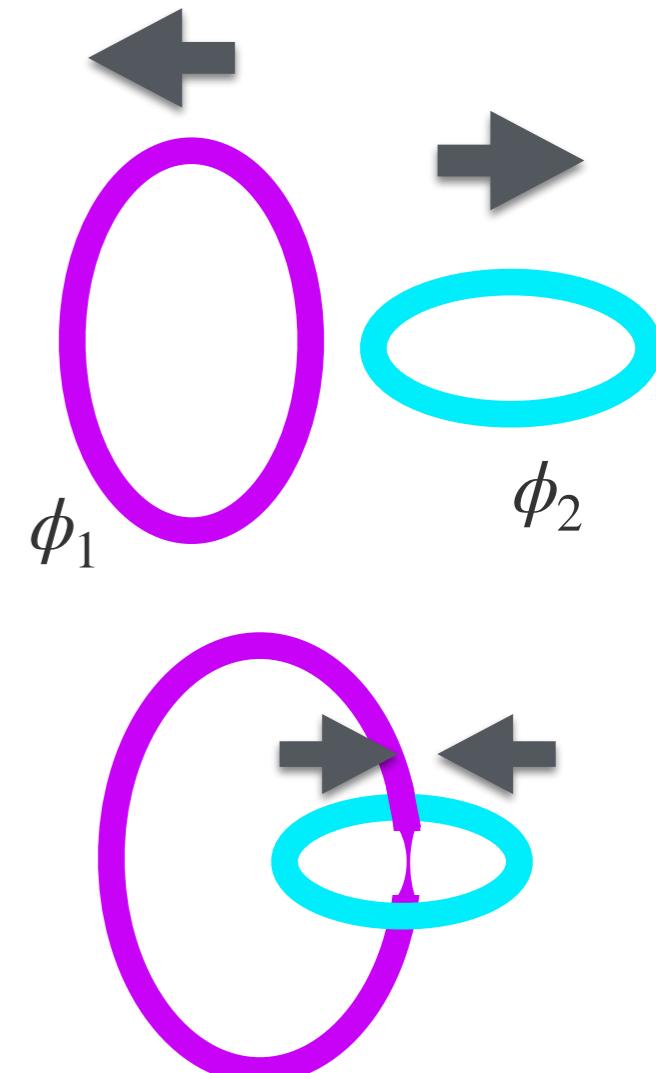
- Assume kinetic mixing with  $U(1)_Y$  in SM:  $\mathcal{L} \supset \frac{\epsilon}{2} Y_{\mu\nu} F^{\mu\nu}$   
 $\rightarrow$  contains  $U(1)_Y$  helicity:  $N_{CS}[Y] \simeq \epsilon^2 N_{CS}[A]$

# Fate of link soliton

- produced by Kibble mechanism or thermal fluctuation at  $T \sim v_1, v_2$
- classically stable but can decay by quantum tunneling
- decay after electroweak phase transition
  - change of helicity:  $\Delta N_{CS}[Y] \simeq \epsilon^2 N_{CS}[A]$
  - baryon # is generated through chiral anomaly:

$$\boxed{\Delta B = \epsilon^2 N_{CS}[A] \text{ per link}}$$

$$(\partial_\mu J_B^\mu \sim Y\tilde{Y} + \text{tr } W\tilde{W})$$



# Baryon # from link

- Naively, anti-link is also produced  $\rightarrow n_{link} - n_{\overline{link}} = 0$  ?
- need "chemical potential"  $\mu$  (discussed later) when produced:

$$\frac{n_{link} - n_{\overline{link}}}{s} \simeq \frac{\mu}{T} \frac{n_{link}}{s} \simeq \frac{\mu}{T} 10^{-6}$$

We have used  $n_{link} \sim 10^{-4} T^3$

[Vachaspati '84]

$\rightarrow$  generated total baryon # due to decay:

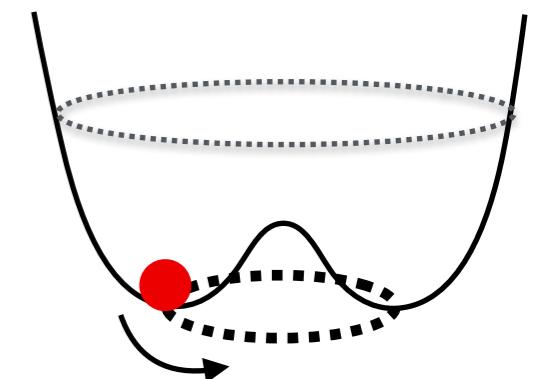
$$\eta_B = \frac{n_B}{s} \sim 10^{-10} \left( \frac{\epsilon}{0.1} \right)^2 \left( \frac{\mu/v_1}{0.1} \right)$$

# Origin of chemical potential?

- Simplest choice: rotating pseudo scalar (axion-like particle)

$$\Delta \mathcal{L} = \frac{c}{16\pi^2} a' F_{\mu\nu} \tilde{F}^{\mu\nu} = -\frac{c}{16\pi^2} (\partial_0 a') A dA \equiv \mu_{eff}$$

(cf: Affleck-Dine mechanism, axiogenesis [Co-Harigaya '19])



$$\eta_B = \frac{n_B}{s} \sim 10^{-10} \left( \frac{\epsilon}{0.1} \right)^2 \left( \frac{\dot{a}'|_{T \sim v_1}}{0.1 v_1} \right)$$

Once the link asymmetry is produced, it remains until decay.

→ later & earlier dynamics of ALP are irrelevant.

# Testability

- Before the links decay, they dominate the energy density of universe.

$$\left. \frac{\rho_{link}}{\rho_\gamma} \right|_{T \sim T_{EW}} = \left. \frac{M_{link} n_{link}}{T^4} \right|_{T \sim T_{EW}} \simeq 10^{-2} \frac{\nu_1}{\nu_{EW}} \gg 1$$

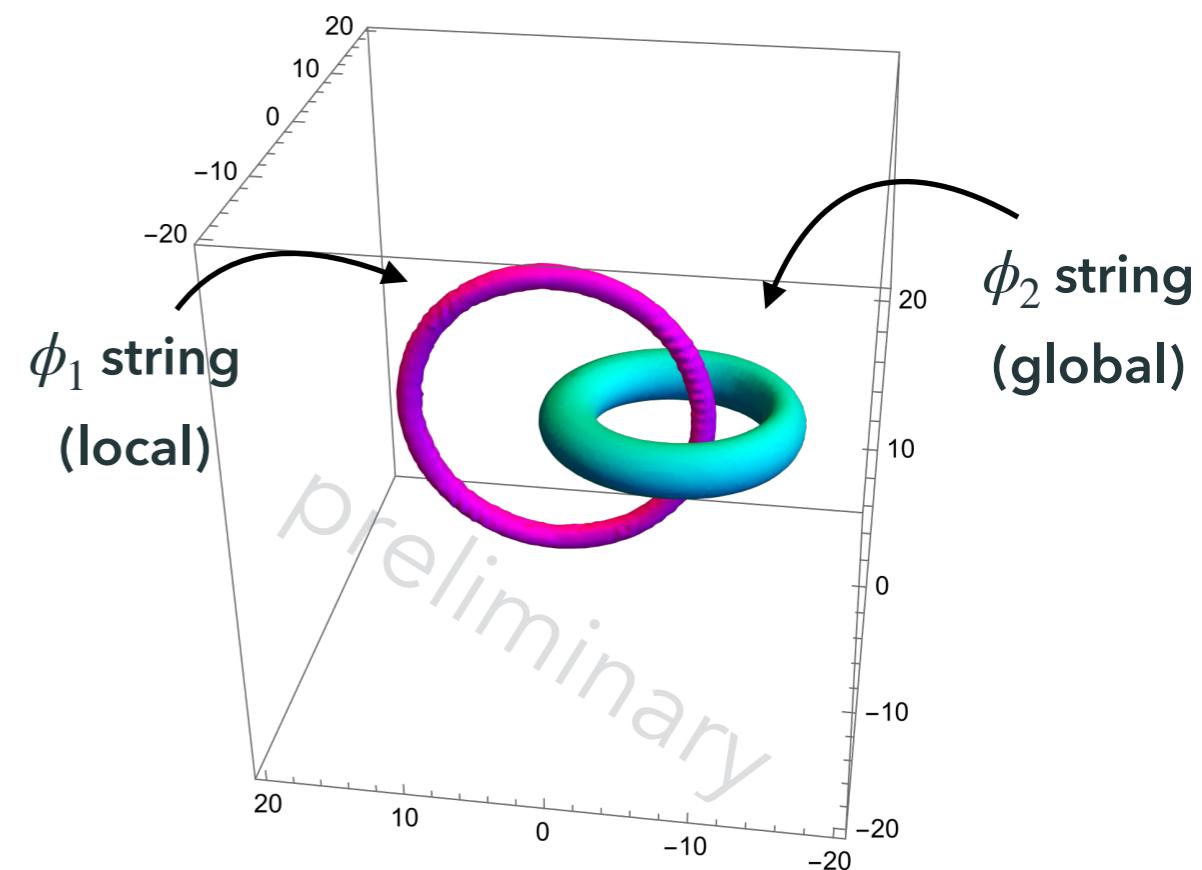
- The entropy production due to decay cannot be ignored.
  - distorts spectrum of primordial gravitational wave
  - **probed by primordial gravitational wave?**

# Summary

- 1st example of link soliton in realistic models
- Key: CS coupling  $\frac{c}{16\pi^2} \int d^4x aF\tilde{F}$
- Natural setup

$$\begin{cases} U(1)_{global} \rightarrow U(1)_{PQ} \\ U(1)_{gauge} \rightarrow U(1)_{B-L} \end{cases}$$

→ linking flux, applicable for baryogenesis



**link = origin of matter**