# A Shock Tube Study of $OH + H_2O_2 \rightarrow H_2O + HO_2$ and $H_2O_2 + M \rightarrow 2OH + M$ using Laser Absorption of $H_2O$ and OH

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The rate constants of the reactions

$$H_2O_2 + M \rightarrow 2OH + M \tag{1}$$

$$OH + H2O2 \rightarrow H2O + HO2$$
 (2)

were measured in shock-heated  $H_2O_2/Ar$  mixtures using laser absorption diagnostics for  $H_2O$  and OH. Time-histories of  $H_2O$  were monitored using tunable diode laser absorption at 2550.96 nm, and time-histories of OH were achieved using ring dye laser absorption at 306 nm. Initial  $H_2O_2$  concentrations were also determined utilizing the  $H_2O$  diagnostic. On the basis of simultaneous time-history measurements of OH and  $H_2O$ ,  $k_2$  was found to be  $4.6 \times 10^{13} \exp(-2630 \text{ K/T}) \text{ [cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$  over the temperature range 1020-1460 K at 1.8 atm; additional measurements of  $k_2$  near 1 atm showed no significant pressure dependence. Similarly,  $k_1$  was found to be  $9.5 \times 10^{15} \exp(-21\ 250\ \text{K/T}) \text{ [cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$  over the same temperature and pressure range.

### 1. Introduction

Hydrocarbon ignition at intermediate temperatures (850–1200 K) is controlled by hydrogen peroxide ( $H_2O_2$ ) and hydroperoxyl radical ( $HO_2$ ) reactions. These reactions are regarded as "the central kinetic feature in engine knock in spark ignition engines, in ignition in liquid-fueled diesel engines, and in the operation of homogeneous charge, compression ignition (HCCI) engines." The  $H_2O_2$  thermal decomposition system provides a unique opportunity to isolate and study  $H_2O_2$  and  $HO_2$  reactions at combustion temperatures.

For example, a reliable, clean HO<sub>2</sub> precursor is desirable to study HO<sub>2</sub> reactions with alkyl radicals. The H<sub>2</sub>O<sub>2</sub> decomposition system is a possible way to fulfill this need, since HO<sub>2</sub> is produced by a simple two-step process:

$$H_2O_2 + M \rightarrow OH + OH + M$$
 (1)

$$OH + H2O2 \rightarrow H2O + HO2$$
 (2)

Reaction 2 has received considerable attention near room temperature<sup>2–12</sup> because OH and HO<sub>2</sub> catalyze the destruction of ozone (O<sub>3</sub>), and H<sub>2</sub>O<sub>2</sub> is an important reservoir for both radicals. The low temperature experimental data for  $k_2$  show good consistency, ranging from  $9.3 \times 10^{11}$  to  $1.20 \times 10^{12}$  [cm<sup>3</sup> mol<sup>-1</sup> s<sup>-1</sup>] at 298 K. However, some researchers reported a small positive temperature dependence of the rate constant<sup>2–6,8</sup> whereas others found a negative temperature dependence<sup>7,9</sup> over their measured temperature ranges. Atkinson et al. <sup>12</sup> reviewed the available low temperature data and proposed an evaluated rate

As opposed to the extensive studies at low temperatures, only two studies of reaction  $2^{13-15}$  were conducted at combustion temperatures, and these were both performed in the same laboratory. The limited availability of experimental data for  $k_2$  is due to the unstable nature of  $H_2O_2$  at high temperatures and the lack of sensitive diagnostics for  $H_2O_2$ . The earlier study  $^{13}$  essentially inferred the ratio of the rates for the two reactions,  $OH + H_2O_2 \rightarrow H_2O + HO_2$  and  $OH + HO_2 \rightarrow H_2O + O_2$ . The expression for  $k_2$  from the most recent Hippler et al. study  $^{14}$  is widely used in chemical kinetics models for combustion  $^{16-19}$  and accepted in the latest evaluation by Baulch et al.  $^{20}$  An unusual up-turn in the reaction rate at temperatures higher than 800 K was reported by Hippler et al.  $^{14}$ 

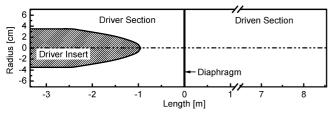
Recently, a new systematic approach for studying the thermal decomposition of  $H_2O_2$  at combustion temperatures was reported. In light of this newly developed method, better-controlled experiments can be designed to study other major reactions within the  $H_2O_2$  decomposition system. Here we report our latest work on the reaction  $OH + H_2O_2 \rightarrow H_2O + HO_2$  and extend our previous measurements of  $H_2O_2 + M \rightarrow 2OH + M$  to higher temperatures, both with additional measurements of OH using ring dye laser absorption in the ultraviolet.

#### 2. Experimental Setup

**2.1.** Shock Tube with Modified Driver Section. The experiments were carried out in a high-purity, 304 stainless steel shock tube with inner diameter of 14.1 cm. The driven section of the shock tube is 8.5 m long, and the driver section is 3.3 m long. Conditions behind reflected shocks were determined using normal-shock relations. Incident-shock velocity measurements

expression of  $k_2 = 1.75 \times 10^{12} \exp(-160 \text{ K/T}) [\text{cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$  with a value of  $1.02 \times 10^{12}$  at room temperature.

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**Figure 1.** Scale drawing of the driver-insert configuration. After the diaphragm rupture, the incident shock wave propagates from left to right, and an expansion wave propagates from right to left.

were made using five piezoelectric pressure transducers (PCB) over the last 1.5 m of the shock tube, allowing accurate extrapolation to infer the incident shock velocity at the endwall. Uncertainty in the initial temperature behind the reflected shock wave,  $T_5$ , was  $\pm 0.8\%$ .<sup>22</sup> In addition to the five PCB pressure transducers for incident-shock velocity measurements, another sidewall piezoelectric pressure transducer (Kistler model 603B) located 2 cm from the endwall was used to monitor pressure time-histories.

Accurate rate measurements in reflected shock wave experiments rely on the near-constant-volume performance of largediameter shock tubes at short test times (less than 1-2 ms). At longer test times, small increases in temperature and pressure in the reflected shock conditions can occur due to facility-related nonideal effects, such as boundary layer growth and shock attenuation. These nonideal effects can be compensated for by using driver inserts. The design of the driver insert was based on a recently developed semianalytical model.<sup>23</sup> A scale drawing of the primary insert configuration used in the current work is shown in Figure 1. We have demonstrated that by using the driver insert method, we can successfully maintain a highly uniform pressure and temperature profile for up to 8 ms behind the reflected shock. These small variations in temperature have been confirmed by in situ temperature measurements. <sup>23,24</sup> In the current study, postshock pressure fluctuations for all the shocks are less than  $\pm 2\%$ , and the corresponding long-time temperature uncertainty is evaluated to be less than  $\pm 0.8\%$ .

**2.2. Tunable Diode Laser Absorption of H<sub>2</sub>O near 2.55**  $\mu$ m. The H<sub>2</sub>O concentration time-histories were measured 2 cm from the end wall of the shock tube using laser absorption of H<sub>2</sub>O at 2550.96 nm (3920.09 cm<sup>-1</sup>) within the  $v_3$  fundamental vibrational band. Almost all spectral parameters for this transition were taken from HITRAN.<sup>25</sup> The collisional-broadening coefficient (half-width at half-maximum per atm) for argon (not available in HITRAN) was measured in the shock tube over a temperature range between 1000 and 1600 K by matching the peak of the fitted Voigt profile to the measured line center absorbance of a known amount of H<sub>2</sub>O and was found to be  $\gamma_{Ar}(T) = 0.0277(296/T)^{0.50}$  [cm<sup>-1</sup> atm<sup>-1</sup>]. Details of this H<sub>2</sub>O absorption diagnostic can be found in the literature.<sup>21</sup>

A distributed feedback (DFB) diode laser near 2.55  $\mu$ m from Nanoplus GmbH was used. Shortly before taking data, the laser was scanned over the water line to ensure it was located at the line center. The laser beam was collimated by a lens, transmitted through the shock tube and a JDSU narrow-bandpass filter (JDSU N02548-8, center wavelength: 2546 nm, half width: 46 nm), focused by a short-focal-length convex lens into an integrating sphere (SphereOptics SPH-1G-3) to reduce beamsteering effects, and detected by a liquid nitrogen-cooled InSb detector (IR Associates IS-2.0, 1 MHz bandwidth). The horizontal laser beam crosses the axis of the shock tube so that the path length of the measurement is the inner diameter of the shock tube. The beam path outside the shock tube was purged with pure N<sub>2</sub> to minimize laser attenuation from ambient H<sub>2</sub>O.

# 2.3. Ring-Dye Laser Absorption of OH near 306.7 nm.

The laser beam path for the OH absorption diagnostic was in the same axial plane (but rotated by  $45^{\circ}$ ) as the  $H_2O$  diagnostic (2 cm from the shock tube endwall), allowing for simultaneous measurement of OH and  $H_2O$  time-histories. A pair of fused-silica windows was installed on the shock tube sidewalls to allow the transmission of the OH laser.

The  $R_1(5)$  line of the OH  $A^2\Sigma^+-X^2\Pi$  (0,0) band near 306.7 nm, chosen for the current study, has been well-characterized and extensively utilized in chemical kinetics studies in the past.<sup>22,26,27</sup> The absorption coefficient of the OH radical is well-established, and measured OH concentrations are accurate to better than  $\pm 5\%$ .<sup>22</sup>

To generate 306.7 nm light, a 5 W, 532 nm, cw beam produced by a Coherent Verdi laser pumps a Spectra-Physics 380 ring-dye laser operating with Rhodamine R6G dye was used to generate visible light at 613.4 nm. UV light (1–2 mW) at 306.7 nm was generated by intracavity frequency doubling using a temperature-tuned AD\*A crystal. Wavelength was determined using a Burleigh WA-1000 wavemeter, and mode quality of the visible light was examined by a scanning interferometer.

A part of the UV beam was split off as the reference beam; the remaining part was transmitted through the shock tube for OH absorption. Both the reference and the transmitted beams were monitored by Thorlabs PDA 36A detectors modified with large area Hamamatsu S1722-02 photo diodes (with an effective active area of 13.2 mm²). Both detectors were shielded by two identical spectral filters (Newport FSR-UG11, center wavelength: 340 nm, half width: 70 nm) to reduce broadband interference. Further details of the OH ring-dye laser absorption diagnostic may be found in Herbon et al.<sup>22</sup>

**2.4.** H<sub>2</sub>O<sub>2</sub> Source. Urea-H<sub>2</sub>O<sub>2</sub> (carbamide peroxide, formula: (NH<sub>2</sub>)<sub>2</sub>CO•H<sub>2</sub>O<sub>2</sub>) has proven to be a reliable laboratory source of H<sub>2</sub>O<sub>2</sub> vapor.<sup>21,28</sup> Approximately 10 g of urea-hydrogen peroxide powder (Sigma-Aldrich, 15–17% active oxygen basis) was mixed with roughly an equivalent amount of sand (SiO<sub>2</sub>, Sigma-Aldrich) in a polycarbonate flask to prevent agglomeration. The flask was sealed with a platinum-cured silicone stopper and placed in a water bath. Upon gentle heating (48 °C), urea-H<sub>2</sub>O<sub>2</sub> powder decomposes to yield H<sub>2</sub>O<sub>2</sub> vapor and urea solid. The H<sub>2</sub>O<sub>2</sub> formed was carried into the shock tube test section by a stream of research grade argon (99.999%) that was passed through the flask at a flow rate of 0.4-0.5 SLPM (standard liters per minute) to obtain a typical H<sub>2</sub>O<sub>2</sub> concentration of about 1500 ppm. A stable supply of H<sub>2</sub>O<sub>2</sub> at this rate can be generated for approximately 3-4 h. The H<sub>2</sub>O<sub>2</sub>/carrier gas mixture was then directed into the driven section of the shock tube from a filling port near the endwall.

Previous work<sup>21</sup> has found that  $H_2O_2$  decomposes inside the shock tube on time scales of the order of 15 min, and this decomposition has been attributed to the heterogeneous reaction of  $H_2O_2$  on the shock tube wall surfaces. To minimize the contact time between  $H_2O_2$  vapor and the shock tube walls, a stratified filling technique has been developed.<sup>21</sup> The shock tube was first filled with  $H_2O_2$ /argon mixture to about 10-14 Torr (3-4 min); then the shock tube was raised to the target preshock pressure  $P_1$  with pure argon using a filling port near the diaphragm location, thereby compressing the initial  $H_2O_2$  mixture into a smaller volume in the test section adjacent to the endwall.

#### 3. Analysis

**3.1. Overview of Approach.** Four reactions dominate the thermal decomposition of H<sub>2</sub>O<sub>2</sub>. The two reactions mentioned

in the introduction section yield OH and HO<sub>2</sub> radicals, and two others consume the radicals:

$$OH + HO_2 \rightarrow H_2O + O_2 \tag{3}$$

$$HO_2 + HO_2 \rightarrow H_2O_2 + O_2$$
 (4)

OH radicals are primarily formed by reaction 1 and removed by reactions 2 and 3. Therefore, the time evolution of OH can be described by:

$$\frac{d[OH]}{dt} = 2k_1[H_2O_2][M] - k_2[OH][H_2O_2] - k_3[OH][HO_2]$$
(1)

OH initially builds up following the decomposition of  $H_2O_2$  and is gradually depleted due to reactions 2 and 3. When OH concentration peaks, the above equation can be rewritten as:

$$k_2 = 2k_1 \frac{[M]}{[OH]_{\text{peak}}} - k_3 \frac{[HO_2]}{[H_2O_2]}$$
 (3)

Considering that reaction 2 is the major  $HO_2$  formation channel, only a small amount of  $HO_2$  is present when the OH concentration is at its peak. Therefore, the second term on the right can be dropped, and to first order, the rate constant of the title reaction can be related to  $k_1$  and  $[OH]_{peak}$  by:

$$k_2 \approx 2k_1 \frac{[M]}{[OH]_{\text{peak}}}$$
 (4a)

**3.2. Analysis Using Detailed Kinetics Model.** Data analysis and simulations were performed using an updated version of the chemical kinetics mechanism, GRI-Mech 3.0.<sup>16</sup> The mechanism has been updated to incorporate the latest data for the OH heat of formation,  $^{22,30}$  the HO<sub>2</sub> heat of formation,  $^{31}$  the rate constant for the reaction OH + HO<sub>2</sub>  $\rightarrow$  H<sub>2</sub>O + O<sub>2</sub> ( $k_3$  = 4.28  $\times$  10<sup>13</sup> (T/300)<sup>-0.21</sup> exp(113/T) [cm³ mol<sup>-1</sup> s<sup>-1</sup>]),  $^{32}$  and the thermal decomposition rate of H<sub>2</sub>O<sub>2</sub>.<sup>21</sup> Using this mechanism, a sensitivity analysis<sup>21</sup> for OH was performed with the Senkin<sup>29</sup> kinetics code package. The results of the sensitivity analysis coincide with our expectation that two reactions control the time evolution of OH: H<sub>2</sub>O<sub>2</sub> + M  $\rightarrow$  2OH + M, which shows predominantly positive sensitivity, and OH + H<sub>2</sub>O<sub>2</sub>  $\rightarrow$  H<sub>2</sub>O + HO<sub>2</sub>, which shows predominantly negative sensitivity. Sensitivity analysis plots are presented later in the paper.

As will become evident from these sensitivity analyses, the kinetics that control OH and  $H_2O$  formation in the  $H_2O_2$  thermal decomposition system can be well-described by a small set of reactions, which consists of two major reactions, 1 and 2, and three minor reactions, 3–5. Reactions 1–4 have been mentioned earlier in the paper, and reaction 5 is:

$$OH + OH \rightarrow H_2O + O \tag{5}$$

The limited number of reaction rate constants that do show a significant effect on the modeling of the experimental measurements are called active parameters.<sup>33,34</sup> Of the five identified active parameters, the rate constants of reactions 1, 2, and 5 will be determined experimentally in the current study.

The coefficient for reaction 3 is updated from our recent study,<sup>32</sup> and the rate constant expression for reaction 4 is that used in almost all other recent mechanisms<sup>17–19</sup> and is based on the Kappel et al. study.<sup>15</sup>

Because of the limited number of important reaction rate constants, the results of this study are nearly immune to the choice of the base kinetic mechanism. With this expectation, we tested two representative mechanisms, the modified GRI-Mech 3.0 as previously described and the Li et al. model. If identical  $k_1$  and  $k_2$  expressions are used, the two mechanisms predict effectively indistinguishable  $H_2O$  time-histories and nearidentical OH time-histories. The small difference in the predicted OH concentrations stems from the different rate expressions used for  $k_5$ , as GRI-Mech 3.0 uses Wooldridge et al. State and the Li et al. model uses results from the work by Sutherland et al. However, this small difference does not affect our determination of  $k_1$  and  $k_2$ , as will become evident in Section 4. In this study, we use the modified GRI-Mech 3.0 as the base kinetic model.

**3.3. Test Mixture Nonuniformity Considerations.** It is possible that the stratified filling procedure used may introduce some axial nonuniformity in the test mixtures, but this should not affect incident shock speed. As mentioned above, we have used the in situ H<sub>2</sub>O diagnostic described in ref 21 to accurately determine the mixture composition at the test location. Because extremely diluted test mixtures were used in this study, small axial variations in mixture composition will have very small influence on the incident shock speed and reflected shock temperatures. Comparisons between the 4000 ppm H<sub>2</sub>O<sub>2</sub>/H<sub>2</sub>O mixtures used in this study and pure Ar shock wave experiments indicate that at most, a maximum of only 6 K difference may be expected in reflected shock temperatures.

It may be possible that radial  $H_2O_2$  concentration gradients could exist if  $H_2O_2$  undergoes significant surface decomposition on the sidewall of the shock tube. However, although some reaction of  $H_2O_2$  may occur at the shock tube wall, such effects are limited by the slow  $H_2O_2$  diffusion rate to the wall. A rearrangement of eq 3 shows that the measured  $[OH]_{peak}$  is not a strong function of  $H_2O_2$  concentration, thus the  $k_2$  rate coefficient determination is also not a strong function of  $H_2O_2$  concentration or its varying radial distribution. In fact, simulations using variable amounts of  $H_2O_2$  have indicated that the entire OH profile is not a strong function of the initial  $H_2O_2$  concentration.

#### 4. Results and Discussion

**4.1. Lower Temperature (1000** < T < **1200** K) Example Case. It has been demonstrated in the study by Hong et al.<sup>21</sup> that in the temperature range from 1000 to 1200 K, H<sub>2</sub>O formation during the decomposition of H<sub>2</sub>O<sub>2</sub> is predominantly controlled by  $k_1$ . No other reaction shows significant sensitivity at temperatures ranging from 1000 to 1200 K.<sup>21</sup>

From the analysis in Section 3, it is natural to select the OH diagnostic for the determination of  $k_2$ . We chose to also implement the simultaneous  $H_2O$  diagnostic for two reasons: (1) with initial  $H_2O_2$  concentrations measured using the  $H_2O$  diagnostic,<sup>21</sup> comprehensive kinetics analyses can be performed for OH time-histories; and (2) with real-time  $k_1$  determinations using the  $H_2O$  diagnostic, the temperature uncertainty in  $k_1$  does not propagate to  $k_2$ .

Considering that the overall reaction of  $H_2O_2$  decomposition is described by

$$H_2O_2 = H_2O + \frac{1}{2}O_2$$

and noting the mole-for-mole conversion between H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O, initial H<sub>2</sub>O<sub>2</sub> loadings can be inferred from H<sub>2</sub>O concentration profiles by taking the difference between the initial and final H<sub>2</sub>O concentrations.<sup>21</sup>

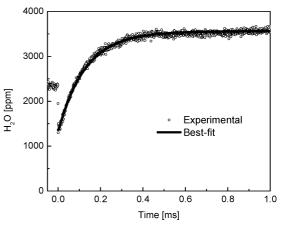
Figure 2 shows an example H<sub>2</sub>O time-history obtained at 1192 K, 1.95 atm, in an argon bath gas. H<sub>2</sub>O concentrations at time zero and in the plateau region at later times are 1364 ppm and 3580 ppm, respectively, corresponding to an initial H<sub>2</sub>O<sub>2</sub> loading of 2216 ppm. Using GRI-Mech 3.0<sup>16</sup> with the modifications described in Section 3, water sensitivity coefficients were calculated with the Senkin<sup>29</sup> kinetics code package. Figure 3 is the H<sub>2</sub>O sensitivity plot for conditions of Figure 2, with the sensitivity coefficient  $\alpha$  defined the same as in the ref 21.

The rate constant of the reaction  $H_2O_2 + M \rightarrow 2OH + M$  $(k_1)$  was determined by changing  $k_1$  in the chemical kinetics model to best-fit the experimental H<sub>2</sub>O time-histories, as described in ref 21. The value of  $k_1$  for the example case is 1.8  $\times$  10<sup>8</sup> [cm<sup>3</sup> mol<sup>-1</sup> s<sup>-1</sup>] with an estimated fitting error of  $\pm$ 10%.<sup>21</sup> The relatively high activation energy of the reaction H<sub>2</sub>O<sub>2</sub> +  $M \rightarrow 2OH + M (42 \text{ kcal/mol}^{21}) \text{ renders the rate } k_1 \text{ very}$ temperature sensitive. The uncertainty in  $k_1$  associated with the temperature uncertainty was estimated to be 21%.<sup>21</sup> However, when the chemical kinetics model was updated with the experimental  $k_1$  values for the purpose of determining  $k_2$ , the temperature uncertainty in  $k_1$  does not propagate to  $k_2$  since both  $k_1$  and  $k_2$  were evaluated simultaneously and were at exactly the same temperature. Only a 10% fitting uncertainty in  $k_1$ propagates to  $k_2$ .

With the initial  $H_2O_2$  concentration and  $k_1$  determined from the  $H_2O$  diagnostic, the rate of the reaction  $OH + H_2O_2 \rightarrow H_2O$ + HO<sub>2</sub> ( $k_2$ ) can be derived from the OH time-history, as OH yield is predominantly controlled by the competition between  $k_1$  and  $k_2$ . Figure 4 shows the OH concentration time-history from the same test as Figures 2 and 3. The corresponding OH sensitivity plot is presented in Figure 5. The OH sensitivity coefficients were calculated using GRI-Mech  $3.0^{16}$  with  $k_1$  and  $k_2$  values updated to the experimentally determined values.

By changing  $k_2$  in the chemical kinetics model, the calculated OH profile matches measurement in the peak OH concentration, as shown in Figure 4. The associated fitting uncertainty is estimated to be  $\pm 3\%$ . The other two reactions that the OH yield shows sensitivity to at early times are  $OH + HO_2 \rightarrow H_2O + O_2$  $(k_3)$  and  $HO_2 + HO_2 \rightarrow H_2O_2 + O_2$   $(k_4)$ .  $HO_2$  concentration is still low when OH concentration peaks. Therefore, the calculated [OH] maximum is not significantly affected by the choice of a  $k_3$  value. Figure 6 illustrates that even a factor of 2 uncertainty in  $k_3$  results in only a  $\pm 3\%$  uncertainty in the inferred  $k_2$ , and this is a much larger uncertainty in  $k_3$  than assigned ( $\pm 27\%$ ) in our recent study.<sup>32</sup> The rate of the reaction  $HO_2 + HO_2 \rightarrow H_2O_2$ + O<sub>2</sub> does not affect the peak OH concentration, as can be seen in Figure 6.

Uncertainties in  $k_2$  come from various error sources, noticeably uncertainties in  $k_1$ , OH concentration, temperature, secondary reactions, and fitting procedures. As discussed previously, the uncertainty in  $k_1$  that propagates to  $k_2$  is estimated to be  $\pm 10\%$ . Secondary reactions and fitting procedures introduce uncertainties of  $\pm 3$  and  $\pm 5\%$ , respectively. The OH line strength has an estimated maximum uncertainty of  $\pm 5\%$ .<sup>22</sup> At the same time, a 0.02 cm<sup>-1</sup> uncertainty in UV light wavelength introduces additional  $\pm 1.5\%$  uncertainty. The combined temperature uncertainty was evaluated to be 11 K,21 and the associated



**Figure 2.** The dissociation rate of  $H_2O_2(k_1)$  is fitted to be  $1.8 \times 10^8$ [cm<sup>3</sup> mol<sup>-1</sup> s<sup>-1</sup>] with an estimated fitting error of  $\pm 10\%$ . Initial test mixture: 2216 ppm H<sub>2</sub>O<sub>2</sub>/1364 ppm H<sub>2</sub>O/682 ppm O<sub>2</sub>/Ar; initial reflected shock conditions: 1.95 atm, 1192 K.

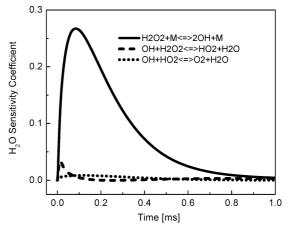
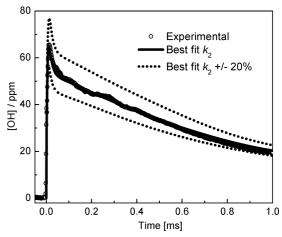


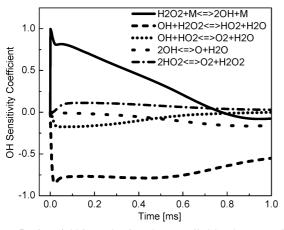
Figure 3. The formation of H<sub>2</sub>O is predominantly controlled by the dissociation rate of H<sub>2</sub>O<sub>2</sub>. Conditions are those of Figure 2.



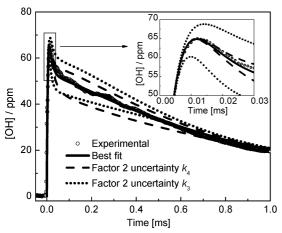
**Figure 4.** The rate of the reaction OH +  $H_2O_2 \rightarrow H_2O + HO_2$  ( $k_2$ ) is fitted to be  $5.1 \times 10^{12} \, [\text{cm}^3 \, \text{mol}^{-1} \, \text{s}^{-1}]$  with an estimated fitting error of  $\pm 3\%$ . Special attention was paid to match the peak in OH concentration. Conditions are those of Figure 2.

uncertainty in  $k_2$  is estimated to be  $\pm 2\%$ . Therefore, the overall uncertainty of the example case at 1192 K can be estimated to be  $\pm 13\%$ .

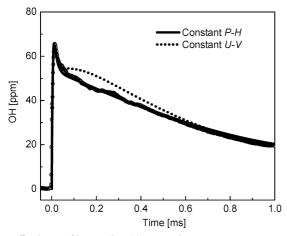
It should be pointed out that OH time-histories were calculated assuming test mixtures were at constant pressure and enthalpy (constant P-H) behind reflected shock waves in the current study, as opposed to a usually assumed constant volume



**Figure 5.** OH yield is predominantly controlled by the competition between the OH formation reaction  $H_2O_2 + M \rightarrow 2OH + M$  ( $k_1$ ) and the OH removal reaction  $OH + H_2O_2 \rightarrow H_2O + HO_2$  ( $k_2$ ). Conditions are those of Figure 2.

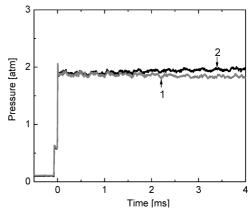


**Figure 6.** [OH]<sub>peak</sub> is controlled by  $k_1$  and  $k_2$ . Other reactions show very minor or no impact on [OH]<sub>peak</sub>. Conditions are those of Figure 2.



**Figure 7.** OH profiles predicted by assuming a constant P-H process or a constant U-V process behind the reflected shock wave. Identical  $k_1$  and  $k_2$  were used in those calculations. The two models agree on the peak OH concentration, although the constant U-V model predicts a slight hump behavior after the OH spike. Conditions are those of Figure 2.

and internal energy (constant U-V). Shown in Figure 7 is a comparison of OH profiles calculated assuming constant P-H and constant U-V reactors with the identical  $k_1$  and  $k_2$  values. The  $k_1$  value was determined from the  $H_2O$  time-history (Figure



**Figure 8.** Comparison of pressure profiles: (1) 1134 K, 2900 ppm  $H_2O_2/816$  ppm  $H_2O/Ar$  driven gas, with driver insert, dP/dt = 0%/ms; (2) 1100 K, 2444 ppm  $H_2O_2/1538$  ppm  $H_2O/Ar$  driven gas, without driver insert, dP/dt = 1.3%/ms. All pressure profiles have been rescaled to match the initial  $P_5$ .

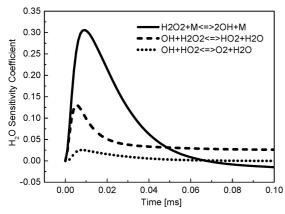
2), where the two reactor models result in almost identical  $\rm H_2O$  profiles. The  $k_2$  value was inferred from the OH peak concentration (Figure 4), where the two models agree very well, as can been seen from the comparison in Figure 7. However, the two models start to diverge after the early time OH spike. The constant U-V reactor model predicts a hump behavior, which was not observed in the experiment. No discussion on this issue was found in the work by Hippler et al. <sup>14</sup> as their experimental data were only compared to model calculations at early times before this discrepancy could appear.

We initially attempted to resolve the discrepancy from a chemical kinetics perspective. The rates of the other two reactions that also show OH sensitivities, OH + HO<sub>2</sub>  $\rightarrow$  H<sub>2</sub>O + O<sub>2</sub> ( $k_3$ ) and HO<sub>2</sub> + HO<sub>2</sub>  $\rightarrow$  H<sub>2</sub>O<sub>2</sub> + O<sub>2</sub> ( $k_4$ ), were both adjusted by a factor of 2 while retaining the constant *P-H* assumption (Figure 6). None of those adjusted OH profiles show a similar hump behavior, suggesting that the discrepancy could not be explained by chemical kinetics parameters.

It is significant that  $k_1$  is a strong function of temperature, with an activation energy of 42.1 kcal/mol.<sup>21</sup> In contrast,  $k_2$  only weakly depends on temperature with an activation energy of only 5.2 kcal/mol, as will be discussed later in the paper. Therefore, a temperature rise favors  $k_1$  and shifts the balance toward higher OH concentrations. For an exothermic system, such as the decomposition of  $H_2O_2$ , temperature rises higher in a constant U-V reactor than at constant P-H, which apparently explains the discrepancy between the two models. At late times (>0.6 ms in Figure 7),  $H_2O_2$  is close to depletion and temperature effects on  $k_1$  diminish; as a result, the OH profiles calculated with both the models converge again.

The mixtures used in this study were very dilute so that increases in temperature due to the exothermicity of  $\rm H_2O_2$  decomposition were minimized. For the example case discussed here, GRI-Mech  $\rm 3.0^{16}$  predicts final temperature rises of 20 and 11 K by the constant U-V reactor model and the constant P-H reactor model, respectively. It may be possible in the future to distinguish these two reactor models by taking advantage of a highly sensitive diagnostic recently developed for shock tube temperature measurements. <sup>24</sup>

To ensure that the observed departure in the experimental OH profile from the commonly adopted constant U-V reactor model is not a result of using the driver insert, a test was conducted without the driver insert. Pressure profiles for these experiments are shown in Figure 8.



**Figure 9.** At higher temperatures (T > 1200 K), the formation of H<sub>2</sub>O shows significant sensitivity to  $OH + H_2O_2 \rightarrow H_2O + HO_2$ , whereas  $H_2O_2 + M \rightarrow 2OH + M$  remains the most important one. Initial test mixture: 2540 ppm H<sub>2</sub>O<sub>2</sub>/1234 ppm H<sub>2</sub>O/617 ppm O<sub>2</sub>/Ar; initial reflected shock conditions: 1.91 atm, 1398 K.

Neither pressure profile in Figure 8 shows any evidence of reaction-induced pressure change, as the energy release due to chemical reaction is so small that it cannot be resolved by the pressure measurements. (The linear rise in pressure seen in profile 2 is due to facility effects such as boundary layer growth.) The measured OH time-history for the experiment without driver insert (not shown here) was very similar to the profile observed with driver insert (i.e., similar to the measured OH profile in Figure 7) and does not show the hump feature that appears in the constant U-V model in Figure 7. Thus, the departure from the constant U-V reactor model was not caused by using the driver insert.

4.2. Higher Temperature (1200 K < T < 1460 K) Example Case. As temperature rises above 1200 K, H<sub>2</sub>O formation starts to show increasingly pronounced sensitivity to the reaction OH  $+ H_2O_2 \rightarrow H_2O + HO_2$ , as can be seen from Figure 9 at 1398 K. The apparent reason is that H<sub>2</sub>O<sub>2</sub> is mainly consumed by its thermal decomposition at lower temperatures (1000 < T < 1200K). However, at elevated temperatures, OH concentration increases dramatically such that the reaction OH +  $H_2O_2 \rightarrow$  $H_2O + HO_2$  becomes an important channel in the removal of H<sub>2</sub>O<sub>2</sub>. Radical-radical reactions that lead to final products (H<sub>2</sub>O and O<sub>2</sub>) are typically very fast, and the consumption of H<sub>2</sub>O<sub>2</sub> is the rate-limiting step, either through  $H_2O_2 + M \rightarrow 2OH + M$ or OH +  $H_2O_2 \rightarrow H_2O$  +  $HO_2$ . Therefore, at elevated temperatures,  $k_1$  could not be solely determined from  $H_2O$ concentration time-histories.

Fortunately, the OH sensitivity coefficient plot in Figure 10 shows that both the reaction rates ( $k_1$  and  $k_2$ ) that control H<sub>2</sub>O formation also control OH yields. With two constraints (H<sub>2</sub>O and OH), it is possible to determine two independent unknowns  $(k_1 \text{ and } k_2)$ . In addition, the problem is well posed because  $k_1$ and k<sub>2</sub> sensitivities to H<sub>2</sub>O are similar, whereas the OH sensitivities have opposite signs. In the current study, H<sub>2</sub>O and OH time-histories of all the tests performed at higher temperatures (1200 < T < 1460 K) were analyzed at the same time to infer both  $k_1$  and  $k_2$  with small uncertainties.

Example H<sub>2</sub>O and OH profiles of the test conditions of Figure 9 are presented in Figures 11 and 12, respectively. We used a search algorithm to evaluate  $k_1$  and  $k_2$ . Specifically, the best-fit  $k_1$  and  $k_2$  were determined iteratively between the profiles of the two species, beginning by adjusting  $k_1$  and  $k_2$  together to match  $H_2O$  time-histories with  $k_1/k_2$  ratios fixed at the values estimated from  $[OH]_{peak}$ .  $k_2$  values were then fine-tuned to match OH time-histories while keeping  $k_1$  unchanged to get updated

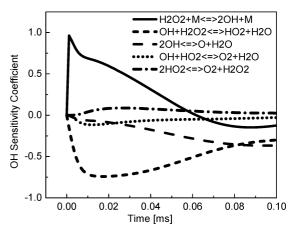
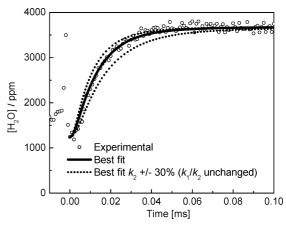


Figure 10. Similar to the example case at a lower temperature (Figure 5), OH yield at an elevated temperature is controlled by the OH formation reaction  $H_2O_2 + M \rightarrow 2OH + M$  and the OH removal reaction  $OH + H_2O_2 \rightarrow H_2O + HO_2$ . Conditions are those of Figure 9.



**Figure 11.** H<sub>2</sub>O time-history recorded at the conditions of Figure 9. The best-fit H<sub>2</sub>O profile was achieved by setting  $k_1 = 2.4 \times 10^9$  [cm<sup>3</sup>  $\text{mol}^{-1} \text{ s}^{-1}$ ] and  $k_2 = 6.8 \times 10^{12} \text{ [cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$ . In comparison, the calculated curves with both  $k_1$  and  $k_2$  changed  $\pm 30\%$  while keeping  $k_1/k_2$  as determined from the corresponding OH profile.

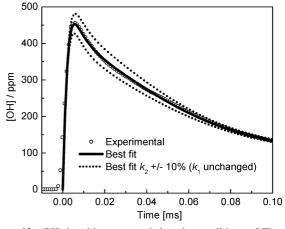


Figure 12. OH time-history recorded at the conditions of Figure 9. The best-fit OH profile was achieved by the same pair of  $k_1$  and  $k_2$ values that best-fits the water profile in Figure 11. The calculated curves with  $k_2$  changed  $\pm 10\%$  while keeping  $k_1$  at its overall optimal value.

 $k_1/k_2$  ratios. The process continued until good fits had been achieved for both species time-histories. The sensitivity coefficient plots in Figures 9 and 10 were generated using GRI-Mech<sup>16</sup> with experimentally evaluated  $k_1$  and  $k_2$ .

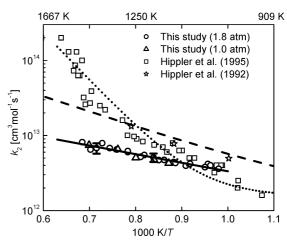
Besides the search algorithm, other optimization methods<sup>34,37,38</sup> have been developed to simultaneously analyze a variety of experimental data, including ignition delay times, flame speeds, and species time-histories. These experimental data are the target constraints of combustion mechanisms, although the data usually contain similar and redundant information about the systems. Typically, the optimal solution for one set of data is not the optimal solution for another set of data, due to experimental errors. Fortunately, the current search algorithm is suitable for a low-dimensionality system.<sup>34</sup> Of the H<sub>2</sub>O<sub>2</sub> thermal decomposition system under investigation, there are only two major active parameters ( $k_1$  and  $k_2$ ). Furthermore, the two active parameters have well-behaved and different functional influence on the time-histories of OH and H<sub>2</sub>O obtained from individual experiments, thus permitting a unique determination of both reaction rate constants.

Recall that at lower temperatures (1000 < T < 1200 K), the initial  $H_2O_2$  loading was determined by taking the difference between the initial and the final plateau  $H_2O$  levels without corrections. The accuracy of the approach is satisfactory at lower temperatures since near-complete (>99%) conversions to final products ( $H_2O$  and  $O_2$ ) following the decomposition of  $H_2O_2$  can be assumed. Chemical kinetics calculations show that  $H_2O_2$  and  $HO_2$  are depleted when a  $H_2O$  plateau is reached. A small fraction of the initial  $H_2O_2$  is converted to and remains as OH, but this represents typically less than 0.5% of the initial water concentration.

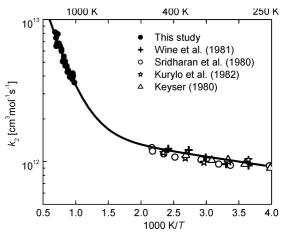
At high temperatures, initial  $H_2O_2$  concentration must be related to both the final  $H_2O$  and OH concentrations. For the example case at 1398 K, there is a residual OH concentration of approximately 130 ppm at 0.1 ms (Figure 12) when the apparent  $H_2O$  plateau was achieved (Figure 11). The initial  $H_2O_2$  concentration must thus be increased by 65 ppm to correct for the residual OH. Similar adjustments were made for all the tests conducted at temperatures higher than 1200 K. Good agreement was found between chemical kinetics calculations using corrected  $H_2O_2$  initial concentrations and experimental observations.

 $k_1$  and  $k_2$  errors stem from uncertainties in fitting, uncertainty in temperature, and uncertainties in interfering reactions. The fitting uncertainties of  $k_1$  and  $k_2$  were estimated jointly. By setting a constant  $k_1/k_2$  ratio,  $k_1$  and  $k_2$  can both be determined within  $\pm 15\%$  accuracy by fitting the H<sub>2</sub>O profile, as illustrated in Figure 11. In addition, the  $k_1/k_2$  ratio has an estimated fitting error of  $\pm 4\%$  as derived from the OH profile (Figure 12). The overall fitting uncertainty for both  $k_1$  and  $k_2$  are bounded by  $\pm 16\%$ , which is the root-sum-square (RSS) of the two separate fitting uncertainties. The uncertainty in  $k_1$  associated with temperature uncertainty was estimated to be  $\pm 21\%$ .<sup>21</sup> By contrast, the uncertainties in  $k_2$  resulting from temperature uncertainty and OH absorption cross-section uncertainty were evaluated to be  $\pm 2$  and  $\pm 5\%$  as discussed in Section 4.1 of the current study. Uncertainty from interfering reactions was estimated to be less than  $\pm 3\%$  for both  $k_1$  and  $k_2$ . Therefore, the overall uncertainties for  $k_1$  and  $k_2$  are less than  $\pm 27\%$  and  $\pm 17\%$ , respectively, for the example case at T = 1398 K.

We also found that the difference between OH profiles calculated using the constant U-V and the constant P-H reactor models diminishes as temperature increases. The trend coincides with our expectation, because the decomposition of  $H_2O_2$  is greatly enhanced at elevated temperatures and the OH formation rate via  $H_2O_2$  thermal decomposition becomes comparable to the OH removal rate via the reaction OH +  $H_2O_2 \rightarrow H_2O + HO_2$ . As the temperature-sensitive  $H_2O_2$  decomposition reaction losses its dominance in the  $H_2O_2$  decomposition system, OH



**Figure 13.** Arrhenius plot of the rate of the reaction  $OH + H_2O_2 \rightarrow H_2O + HO_2$  ( $k_2$ ). The solid line is the linear-fit to all the experimental data of the current study. The dashed line is the rate expression used by Ó Conaire et al.<sup>17</sup> and Li et al.,<sup>18</sup> whereas the dotted line is the one adopted by GRI-Mech.<sup>16</sup>

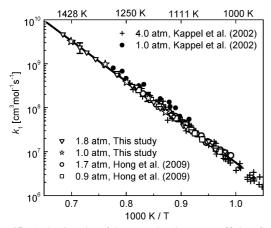


**Figure 14.** The reaction rate of  $OH + H_2O_2 \rightarrow H_2O + HO_2$  ( $k_2$ ), displays a non-Arrhenius behavior over a wide temperature range. Experimental data are well-represented by a sum of two Arrhenius expressions (solid line).

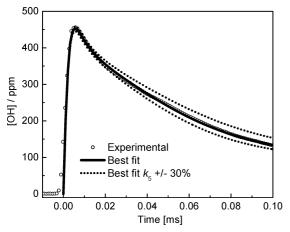
profiles show less sensitivity to temperatures. This observation may in turn justify the constant  $P\!-\!H$  reactor model we adopted to address the discrepancies between experimentally observed OH profiles and calculations based on the constant  $U\!-\!V$  reactor model.

**4.3.** Arrhenius Plot of the Reaction OH +  $H_2O_2 \rightarrow H_2O$  +  $HO_2$ . Using the approaches described in Sections 4.1 and 4.2, the rate constant of reaction 2, OH +  $H_2O_2 \rightarrow H_2O + HO_2$ , was determined at 1.8 atm over the entire temperature range from 1000 to 1460 K. Experimental data were plotted in an Arrhenius plot in Figure 13 ,and are compared to the results reported by Hippler et al. <sup>13,14</sup> The current study shows reasonable agreement with the previous studies at lower temperatures (T < 1300 K). However, at temperatures higher than 1300 K, the study by Hippler et al. <sup>14</sup> reported a dramatic increase in the reaction rate. Since the measurements made by Hippler et al. <sup>13,14</sup> were the only available experimental data, their values were widely used in combustion mechanisms, <sup>16–19</sup> as represented by the dotted line and the dashed line in Figure 13.

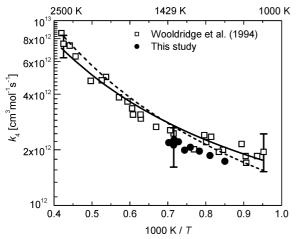
We also carried out rate measurements at 1 atm. As the total pressure decreased from 1.8 to 1 atm, the OH formation rate is reduced almost by half, since the major OH formation channel  $\rm H_2O_2 + M \rightarrow 2OH + M$  is pressure dependent, whereas the



**Figure 15.** Arrhenius plot of the second-order rate coefficient for  $H_2O_2$  thermal decomposition  $(k_1)$ . Excellent agreement was found between the current study and previous studies. <sup>15,21</sup> The solid line is the linear-fit to the experimental data of this study at 1.8 atm (bath gas: argon).



**Figure 16.** OH time-history recorded at the conditions of Figure 9. The optimal rate of the reaction  $2OH \rightarrow H_2O + O(k_5)$  was determined by best-fitting the OH time-history. The calculated curves with  $k_5$  changed  $\pm 30\%$  while retaining the values of all other parameters.



**Figure 17.** Arrhenius plot of the secondary reaction  $2OH \rightarrow H_2O + O(k_5)$ . The results obtained in the current study are within the estimated error bars of Wooldridge et al.<sup>35</sup> The solid line is the rate expression used by Wooldridge et al.<sup>35</sup> The dashed line is the rate expression reported by Michael<sup>42</sup> corrected for the updated OH heat of formation.

OH removal channel OH +  $H_2O_2 \rightarrow H_2O + HO_2$  is not.<sup>14,39</sup> Measured OH yields at 1 atm were expected to be approximately half of those at 1.8 atm to recover a pressure-independent  $k_2$ .

As demonstrated in Figure 13, the data obtained from 1 and 1.8 atm experiments show remarkable consistency, suggesting a low likelihood for systematic error. A linear fit to all the data (1 and 1.8 atm) yields an expression

$$k_2 = 10^{13.66} \exp(-2630 \text{ K/T}) [\text{cm}^3 \text{mol}^{-1} \text{ s}^{-1}] \quad (1020 < T < 1460 \text{ K})$$

We performed a simple Arrhenius fit to all of the Hippler et al. data  $^{14}$  to yield an A-coefficient of 4.6  $\times$   $10^{16}$  [cm  $^3$  mol  $^{-1}$  s  $^{-1}$ ] and an activation energy of 19.7 kcal/mol. The ab initio calculations by Ginovska et al.  $^{39}$  reveal that the barrier heights on the ground state pathways is only 7.3 or 7.8 kcal/mol  $^{39}$  above the precursor complex (H<sub>3</sub>O<sub>3</sub>), whereas the precursor complex is 6.2 kcal/mol lower in energy than the separated H<sub>2</sub>O<sub>2</sub> and OH species. The activation energy from this study is 5.2 kcal/mol, which is in better agreement with the theoretical calculations. As well, the effective A-coefficient from Hippler et al.  $^{14}$  is unusually large. By comparison, the A-coefficient resulting from the current study,  $4.6 \times 10^{13}$  [cm  $^3$  mol  $^{-1}$  s  $^{-1}$ ], is a more reasonable value.

To account for the discrepancy between the current study and that of Hippler et al.,  $^{14}$  experimental procedures of the previous study were carefully reviewed. The thermal decomposition rate of  $H_2O_2$  ( $k_1$ ) utilized in their study was approximately 40% larger than determined in the current study, which translated to an approximate 40% difference in  $k_2$ . Their first example (Figure 1 of ref 14, T=1152 K) yields a new  $k_2=4.2$  [cm³ mol $^{-1}$  s $^{-1}$ ], in good agreement with our current study, if their data are reevaluated with our  $H_2O_2$  thermal decomposition rate.

However, at elevated temperatures, the large discrepancy in  $k_2$  (more than a factor of 3) cannot be explained by utilizing different values for  $k_1$ . Unfortunately, absolute OH yields for a high temperature example case (Figure 3 of ref 14, T = 1566K) were not provided. However, by examining the experimental conditions included in the paper (Table 3 of ref 14), we found a dramatic decrease in [H<sub>2</sub>O<sub>2</sub>] of the test mixtures, dropping from 1220 ppm at 931 K to 54 ppm at 1678 K. Given that a sensitive H<sub>2</sub>O<sub>2</sub> diagnostic was not available when the work<sup>14</sup> was accomplished, they had to derive [H<sub>2</sub>O<sub>2</sub>] from the early rises in OH time-histories and the thermal decomposition rate of H<sub>2</sub>O<sub>2</sub> determined in separate studies. 40,41 Without providing an explanation for such dramatic reduction in [H<sub>2</sub>O<sub>2</sub>], we conjecture that the measured OH yields were only fractions of the true values. This may in turn explain their rate of the OH removal reaction OH +  $H_2O_2 \rightarrow H_2O + HO_2$  (k<sub>2</sub>) being too large at high temperatures. We also noticed that Hippler et al. 14 conducted their experiments at pressures near 0.4 atm. The OH absorption feature is narrow at those low pressures as the pressure broadening<sup>22</sup> is less pronounced, resulting in a higher uncertainty in OH absorption cross-section.

Most previous studies of  $k_2$  were carried out near room temperature. The studies by Wine,<sup>2</sup> Sridharan,<sup>8</sup> Kurylo,<sup>4</sup> and Keyser<sup>5</sup> show remarkable consistency and are plotted in Figure 14 along with the results of this study (both at 1 and 1.8 atm). A strong non-Arrhenius behavior is apparent. All the experimental data can be well represented by a sum of two Arrhenius expressions:

$$k_2 = 10^{12.24} \exp(-160 \text{ K/T}) + 10^{13.88} \exp(-3660 \text{ K/T}) [\text{cm}^3 \text{mol}^{-1} \text{s}^{-1}] \quad (280 < T < 1640 \text{ K})$$

 $k_5$ [H<sub>2</sub>O] $[H_2O_2]$  b  $k_1 (M = Ar)$  $k_2$  $(cm^3 mol^{-1} s^{-1})$  $(cm^3 mol^{-1} s^{-1})$ T(K)P (atm) (ppm) (ppm)  $(cm^3 mol^{-1} s^{-1})$  $5.5\times10^{12}$ 1221 1.921 1793  $3.0 \times 10^{8}$ 2827 1192 1.949 2216  $1.8 \times 10^{8}$  $5.1 \times 10^{12}$ 1364  $4.7 \times 10^{12}$ 1160 2.008 480 1500  $9.8 \times 10^{7}$  $4.4\times10^{12}$ 1133 2.070 842 1291  $6.4 \times 10^{7}$  $4.0\times10^{12}$ 1106 2.126 1043 378  $3.9 \times 10^{7}$ 1235 1670  $4.2 \times 10^{12}$ 1045 2.066  $1.6 \times 10^{7}$  $4.0\times10^{12}$ 1039 2.192 2196 3632  $1.1 \times 10^{7}$ 1218 2.192  $3.6 \times 10^{12}$ 1020 3360  $1.0 \times 10^{7}$  $2.9 \times 10^{7}$  $4.2 \times 10^{12}$ 1089 2.239 725 1678  $6.0\times10^{12}$ 1277 2.031 324 2200  $6.2 \times 10^{8}$  $3.8\times10^{12}$ 1.966 2204  $1.7 \times 10^{7}$ 1052 826 1100 1.947 1538 2444  $3.6 \times 10^{7}$  $4.0\times10^{12}$  $4.3 \times 10^{12}$ 1128 1.884 1042 3210  $6.2 \times 10^{7}$  $1.3 \times 10^{8}$  $4.7 \times 10^{12}$  $1.7 \times 10^{12}$ 1.810 988 3256 1176 1.746 514 2228  $3.1 \times 10^{8}$  $5.2 \times 10^{12}$  $1.9 \times 10^{12}$ 1233  $6.8\times10^{12}$  $2.0\times10^{12}$ 1344 1.974 608 2836  $1.3 \times 10^{9}$  $6.2 \times 10^{12}$  $2.0 \times 10^{12}$ 1277 2.020 1038 3011  $5.6 \times 10^{8}$  $6.5 \times 10^{12}$  $2.1 \times 10^{12}$ 1317 2.011 1166 2729  $9.3 \times 10^{8}$ 1374 1.945 1062 2684  $1.8 \times 10^{9}$  $7.9 \times 10^{12}$  $2.2\times10^{12}$  $6.8\times10^{12}$  $2.3\times10^{12}$ 1.909 1234 2540  $2.4 \times 10^{9}$ 1398 1424 1.868 1357 3148  $3.0 \times 10^{9}$  $6.5\times10^{12}$  $2.2\times10^{12}$  $8.2\times10^{12}$ 1461 1.834 1346 2837  $4.7 \times 10^{9}$  $6.9 \times 10^{7}$  $4.5\times10^{12}$ 1.874 2900 1134 816 1430 1.039 1154  $3.2 \times 10^{9}$  $7.4 \times 10^{12}$ 3367  $6.5\times10^{12}$ 3250  $9.7 \times 10^{8}$ 1313 1.057 1170  $5.1\times10^{12}$ 1249 1.113 1218 3806  $3.8 \times 10^{8}$ 

 $1.7 \times 10^{8}$ 

 $8.0 \times 10^{7}$ 

TABLE 1: Test Conditions and Results of H<sub>2</sub>O<sub>2</sub> Decomposition Experiments in Argon Bath Gas<sup>a</sup>

1.152

1.224

1190

1147

**4.4.** Arrhenius Plot of the Reaction  $H_2O_2 + M \rightarrow 2OH + M$ . By combining the  $H_2O$  and OH diagnostics, the measurements of  $H_2O_2$  decomposition rate have been extended beyond the higher limit of the temperature range of our previous study (1200 K) to 1460 K.<sup>21</sup> Results of this study and previous studies<sup>15,21</sup> are shown in Figure 15.

786

775

3650

3157

Excellent agreement was found between the current study and the previous study  $^{21}$  over the entire overlapped temperature range from 1000 to 1200 K. It was also confirmed that the pressure dependence of  $k_1$  is negligible at pressures lower than 2 atm, as reported in the previous study. The linear-fit to the data obtained at 1.8 atm of this study thus yields a determination of the low-pressure limit of  $H_2O_2$  decomposition rate in argon:

$$k_1 = 10^{15.98} \exp(-21250 \text{ K/T}) [\text{cm}^3 \text{ mol}^{-1} \text{ s}^{-1}] \quad (1020 < T < 1460 \text{ K})$$

By comparison, the low-pressure limit in argon obtained in the previous study<sup>21</sup> can be represented by the expression  $k_1 = 10^{15.97} \exp(-21\ 220\ \text{K/T}) \ [\text{cm}^3\ \text{mol}^{-1}\ \text{s}^{-1}]$  for temperatures between 1000 and 1200 K, showing extraordinary consistency.

**4.5. Secondary Reaction 2OH**  $\rightarrow$  **H**<sub>2</sub>**O** + **O.** The OH sensitivity coefficient plot for the example case at the elevated temperature (Figure 10) reveals that the reaction 2OH  $\rightarrow$  H<sub>2</sub>O + O becomes an important OH removal channel after the OH peak. The rates of the two other reactions that show substantial sensitivities during the OH decay, H<sub>2</sub>O<sub>2</sub> + M  $\rightarrow$  2OH + M and OH + H<sub>2</sub>O<sub>2</sub>  $\rightarrow$  H<sub>2</sub>O + HO<sub>2</sub>, have been determined relatively accurately in the present study. Therefore, a good estimation can be obtained for the rate of the secondary reaction 2OH  $\rightarrow$  H<sub>2</sub>O + O ( $k_5$ ) by varying its value in the chemical kinetics mechanism to capture the trend of OH decay. For the example case at 1398 K of this study, the best-fit OH profile is presented

in Figure 16, in comparison with the ones calculated with  $k_5$  values 30% above or below its optimal value.

 $4.7 \times 10^{12}$ 

 $4.3 \times 10^{12}$ 

Recalling the discussion earlier in the paper (Section 4.1) regarding the difference between using a constant U-V and a constant P-H model, the question arises, are the  $k_5$  values inferred from the decay of OH susceptible to the choice of these two models. The issue is found to be minor by a careful examination of Figures 5 and 7, because the predictions by the two models converge when the reaction  $2OH \rightarrow H_2O + O$  becomes an important OH removal channel.

Major error sources of  $k_5$  include uncertainty in fitting, uncertainties in  $k_1$  and  $k_2$ , and uncertainty in temperature. The fitting uncertainty was estimated to be  $\pm 15\%$ . Following the discussion in Section 4.1, the temperature uncertainty in  $k_1$  and  $k_2$  does not propagate to  $k_5$ . The uncertainty inherited from  $k_1$  and  $k_2$  is mainly from their fitting errors ( $\pm 16\%$ , see Section 4.2) and was evaluated to be  $\pm 17\%$  by varying  $k_5$  to best-fit the OH time-history while  $k_1$  and  $k_2$  were at the limits of their fitting uncertainties. The uncertainty in  $k_5$  associated with temperature uncertainty was estimated to be  $\pm 5\%$ . Therefore, the overall uncertainty of  $k_5$  is  $\pm 24\%$ .

Similar best-fit procedures were carried out for the tests at temperatures higher than 1176 K. For the tests conducted at temperatures lower than 1176 K, OH yields were so low that 2OH  $\rightarrow$  H<sub>2</sub>O + O is a minor channel for OH decay. All of the inferred  $k_5$  values are plotted in Figure 17.  $k_5$  has been measured by Wooldridge et al.<sup>35</sup> with estimated uncertainties of -16% to +11% at T > 2100 K and -22% to +25% at T = 1050 K. The  $k_5$  values determined in the present study are in good agreement with those of the Wooldridge et al.<sup>35</sup>

In addition, the reverse reaction  $O + H_2O \rightarrow 2OH$  has been experimentally investigated<sup>36,43</sup> and an expression for  $k_5$  has been derived from these two studies using equilibrium constants.<sup>42</sup> Correcting for the updated OH heat of formation,<sup>22,30</sup>  $k_5$  is well

<sup>&</sup>lt;sup>a</sup> The H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O concentrations are initial values.

represented by  $4.34 \times 10^3 \ T^{2.7} \exp(951/T) \ [\text{cm}^3 \ \text{mol}^{-1} \ \text{s}^{-1}]$ , in good agreement with the expression  $k_5 = 3.57 \times 10^4 \ T^{2.40} \exp(1063/T) \ [\text{cm}^3 \ \text{mol}^{-1} \ \text{s}^{-1}]$  recommended by Wooldridge et al.<sup>35</sup>

#### 5. Conclusions

The  $H_2O_2$  thermal decomposition system was studied behind reflected shock waves over the temperature range between 1020 and 1460 K using laser absorption diagnostics for both  $H_2O$  and OH. Good detectivity for  $H_2O$  was achieved using tunable diode laser absorption of water at 2550.96 nm within its  $\upsilon_3$  fundamental band. The initial compositions of the reactant mixtures were determined using the sensitive  $H_2O$  diagnostic. OH absorption was measured using the well-characterized  $R_1(5)$  line of the OH  $A^2\Sigma^+-X^2\Pi$  (0,0) band near 306.7 nm.

The determinations of the thermal decomposition rate of  $H_2O_2$  ( $k_1$ ) were extended to 1460 K by jointly interpreting  $H_2O$  and OH time-histories. The low-pressure reaction rate was found to be  $k_1 = 10^{15.98} \exp(-21\ 250\ \text{K/T}) \ [\text{cm}^3\ \text{mol}^{-1}\ \text{s}^{-1}]$ , in excellent agreement with the previous study. The estimated uncertainties of  $k_1$  were  $\pm 27\%$  for temperatures higher than 1200 K and  $\pm 23\%$  for temperatures between 1000 and 1200 K.

The rate of the reaction OH +  $H_2O_2 \rightarrow H_2O + HO_2$  ( $k_2$ ) was evaluated at pressures of 1 and 2 atm, over the entire temperature range between 1020 and 1460 K by analyzing the time-histories of  $H_2O$  and OH. No pressure dependence of the reaction rate was found. The measurements can be represented by the Arrhenius expression  $k_2 = 10^{13.66} \exp(-2630 \text{ K/T}) \text{ [cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$  over the measured temperature range with an overall uncertainty ranging from  $\pm 13\%$  at low temperatures to  $\pm 17\%$  at high temperatures. Non-Arrhenius behavior was found by comparing the results of the current study to the previous measurements near room temperature.  $^{2,4,5,8}$  Over a wide temperature range from 280 to 1640 K, the behavior can be described by a sum of two Arrhenius expressions:  $k_2 = 10^{12.24} \exp(-160 \text{ K/T}) + 10^{13.88} \exp(-3660 \text{ K/T}) \text{ [cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$ .

The rate of a secondary reaction  $2OH \rightarrow H_2O + O(k_5)$  was evaluated by examining the decay of OH at elevated temperatures (T > 1176 K). The results were well within the estimated uncertainty limits of the work of Wooldridge et al. <sup>35</sup> Their rate expression  $k_5 = 10^{4.55} T^{2.40} \exp(1063 \text{ K/T}) \text{ [cm}^3 \text{ mol}^{-1} \text{ s}^{-1}]$  (298 < 2380 K) is supported by the present study.

## **Appendix**

Acknowledgment. This work was primarily supported by the National Science Foundation under award No. 0649936 with partial support from the Department of Energy, Office of Basic Energy Sciences, with Dr. Wade Sisk as technical monitor, and the Department of Energy [National Nuclear Security Administration] under Award No. NA28614. The authors are grateful to their associates at Stanford University, Professor David M. Golden for helpful discussions, and Shengkai Wang for technical support.

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