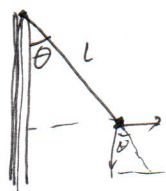


# 习题课5

1.



$$L_{iz} = m \mathbf{r} \times \mathbf{v} = m \cdot l \cdot \sin \theta \cdot v = L_0 \quad (\text{z方向无力矩})$$

$$\tan \theta = \frac{m v^2}{L \sin \theta} = \frac{v^2}{g l \sin \theta} = \frac{\sin \theta}{\cos \theta} \quad \text{对于 } \theta, l, \frac{v}{\theta} \text{ 均有上限}$$

$$\begin{cases} v \cdot \sin \theta = \frac{L_0}{m l} \\ v^2 \cos \theta = g l \sin^2 \theta \end{cases}$$

解得

$$\text{令 } \frac{L_0}{m l} = C \quad v^2 = x \quad \begin{cases} \sin \theta = \frac{C}{x} \\ x \cdot \sqrt{1 - \frac{C^2}{x^2}} = g l \frac{C^2}{x} \end{cases}$$

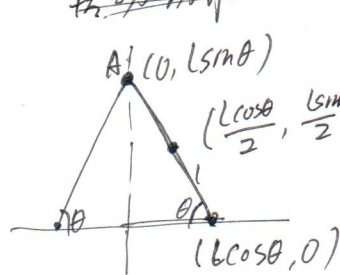
好吧, 解丑无比, 其实答案也没解, 只是说变化趋势.

2.

设碰地速度  $V_f$

$$\text{平动动能: } V_{cm} = \frac{V_f}{2} \quad E_{K1} = \frac{1}{2} \cdot 2m \cdot \left(\frac{V_f}{2}\right)^2 = \frac{1}{4} m V_f^2$$

转动动能 只看一个杆.



$$V_{cm} = \left(-\frac{l}{2} \sin \theta \cdot \dot{\theta}, \frac{l}{2} \cos \theta \cdot \dot{\theta}\right)$$

$$E_K = \frac{1}{2} m \left[ \left(-\frac{l}{2} \sin \theta \cdot \dot{\theta}\right)^2 + \left(\frac{l}{2} \cos \theta \cdot \dot{\theta}\right)^2 \right] + \frac{1}{2} I \cdot \dot{\theta}^2$$

$$\text{其中 } I = \frac{m l^2}{12} \quad \text{代入}$$

$$E_K = \frac{m l^2}{8} \dot{\theta}^2 + \frac{m l^2}{24} \dot{\theta}^2 = \frac{1}{6} m l^2 \cdot \dot{\theta}^2 \quad E_{Kf} = m g \cdot \frac{l \sin \theta_0}{2} \quad \theta_0 = 30^\circ$$

$$\frac{1}{6} m l^2 \dot{\theta}_f^2 = m g \cdot \frac{l}{4} \quad \dot{\theta}_f = \sqrt{\frac{3g}{2l}} \quad \theta_f = 0$$

$$V_A = (0, l \cos \theta \cdot \dot{\theta}) \quad V_f = l \cos \theta_f \cdot \dot{\theta}_f = l \sqrt{\frac{3g}{2l}} = \sqrt{\frac{3gl}{2}}$$

3.

$$\mathbf{q} = (r \cos \theta, r \sin \theta)$$

$$\dot{\mathbf{q}} = (\dot{r} \cos \theta - r \sin \theta \cdot \dot{\theta}, \dot{r} \sin \theta + r \cos \theta \cdot \dot{\theta})$$

$$E_K = \frac{1}{2} (2m) \cdot \dot{r}^2 + \frac{1}{2} (2m) \cdot (r \cdot \dot{\theta})^2 \quad E_{K0} = \frac{1}{2} (2m) (\alpha \cdot \omega_0)^2 = m \alpha^2 \omega_0^2$$

$$J = 2m \cdot r^2 \dot{\theta} \quad \oint_0 = 2m \cdot \alpha^2 \cdot \omega_0$$

动能守恒, 角动量守恒

$$\begin{cases} 2mr^2\dot{\theta} = J_0 & \dot{\theta} = \frac{J_0}{2mr^2} \\ m[\dot{r}^2 + (r\dot{\theta})^2] = E_{k0} \end{cases}$$

$$m\left[\dot{r}^2 + \left(\frac{J_0}{2mr}\right)^2\right] = E_{k0}$$

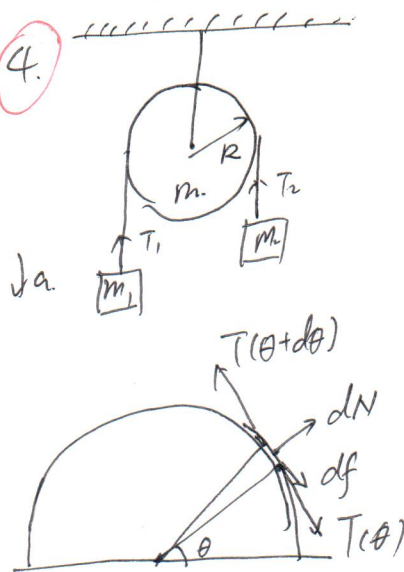
$$\left(\frac{dr}{dt}\right)^2 + \frac{J_0^2}{4m^2 r^2} - \frac{E_{k0}}{m} = 0 \quad \text{令 } C_1 = \frac{E_{k0}}{m} \quad C_2 = \frac{J_0^2}{4m^2}$$

$$\left(\frac{dr}{dt}\right)^2 = C_1 - \frac{C_2}{r^2} \quad \int \frac{r}{\sqrt{C_1 r^2 - C_2}} dr = t$$

$$\therefore t = \frac{1}{2} \int \frac{1}{\sqrt{C_1 r^2 - C_2}} dr^2 = \frac{1}{C_1} \sqrt{C_1 r^2 - C_2} + C_3 \quad \text{代入初始条件 } r(0) = a$$

$$r = a \cdot \sqrt{C_1 t^2 + 1}$$

4.



$$dN = [T(\theta + d\theta) + T(\theta)] \cdot \sin\left(\frac{d\theta}{2}\right) \quad \text{保留一阶}$$

$$= 2T(\theta) \frac{d\theta}{2} = T(\theta) \cdot d\theta$$

$$df = T(\theta + d\theta) - T(\theta) = \mu \cdot dN$$

$$dT = \mu \cdot T \cdot d\theta \quad \frac{dT}{d\theta} = \mu T \quad T = C \cdot e^{\mu\theta}$$

$$\therefore T_1 = e^{\mu\pi} \cdot T_2$$

$$f = T_1 - T_2 = (e^{\mu\pi} - 1) T_2$$

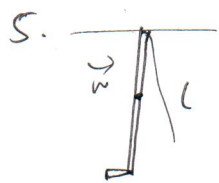
$$\begin{cases} m_1 g - T_1 = m_1 a \\ T_1 = T_2 e^{\mu\pi} \\ T_2 - m_2 g = m_2 a \end{cases}$$

$$(m_1 - m_2 \cdot e^{\mu\pi}) g = (m_1 + m_2 \cdot e^{\mu\pi}) a$$

$$a = \frac{(m_1 - m_2 \cdot e^{\mu\pi}) g}{m_2 e^{\mu\pi} + m_1}$$

$$dm = df \cdot R \quad \therefore m = fR = (e^{\mu\pi} - 1) \cdot T_2 \cdot R = \frac{\mu \pi R^2}{2} \cdot \beta$$

$$\beta = \frac{2}{\mu R} \cdot (e^{\mu\pi} - 1) \cdot (m_2 g + m_2 a)$$



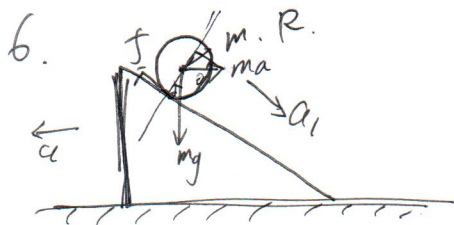
$$L m v_0 = \left( \frac{m_0 l^2}{3} + m l^2 \right) \omega$$

$$\omega = \frac{m v_0}{\frac{m_0 l}{3} + m l}$$

$$E_k = \frac{1}{2} I \cdot \omega^2 = \frac{1}{2} \left( \frac{m_0 l^2}{3} + m l^2 \right) \cdot \omega^2 = \frac{1}{2} \cdot \frac{m^2 v_0^2}{\frac{m_0 l}{3} + m l}$$

$$= m_0 g \cdot \frac{l}{2} (1 - \cos \theta) + m g l (1 - \cos \theta)$$

$$(1 - \cos \theta) = \frac{E_k}{m_0 g \cdot \frac{l}{2} + m g l} = \frac{m^2 v_0^2}{(m_0 g l + 2 m g l) \left( \frac{m_0 l}{3} + m l \right)}$$



设滑块向左加速运动  $a$ , 换算至滑块

$$\begin{cases} m a_1 = m g \sin \theta + m a \cos \theta - f \\ N - m g \cos \theta + m a \sin \theta = 0 \end{cases}$$

$$\frac{1}{2} m r^2 \cdot \frac{a_1}{r} = f r$$

$$N a = N \cdot \sin \theta - f \cos \theta$$

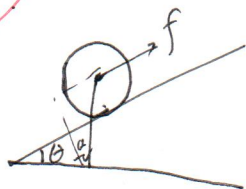
得到  $\begin{cases} M a = m a \sin^2 \theta - m a \cos^2 \theta + m a_1 \cos \theta \\ m g \sin \theta + m a \cos \theta = \frac{3}{2} m a_1 \end{cases}$

$$M a = m a (\sin^2 \theta - \cos^2 \theta) + \cos \theta \cdot \frac{2}{3} (m g \sin \theta + m a \cos \theta)$$

$$M a = -m a \cos(2\theta) + \frac{2}{3} m g \sin 2\theta + \frac{1}{3} m a (1 + \cos 2\theta)$$

$$\therefore a = \frac{\frac{1}{3} m g \sin 2\theta}{M - \frac{1}{3} m + \frac{2}{3} m \cos(2\theta)}$$

7



$$f - mg \sin \theta = ma \quad f = \mu mg \cos \theta$$

$$\frac{2}{5} m R^2 \cdot \frac{d\omega}{dt} = -fR \quad \frac{d\omega}{dt} = - \frac{5 \mu g \cos \theta}{2R}$$

$$\mu g \cos \theta - g \sin \theta = a \quad v = at = (\mu g \cos \theta - g \sin \theta)t$$

$$\omega = \omega_0 - \frac{5 \mu g \cos \theta}{2R} t$$

$$\frac{d}{dt} v = \omega R \quad \omega_0 R - \frac{5 \mu g \cos \theta}{2} t = (\mu g \cos \theta - g \sin \theta)t$$

$$\left( \frac{7}{2} \mu g \cos \theta - g \sin \theta \right) t = \omega_0 R \quad t = \frac{\omega_0 R}{\frac{7}{2} \mu g \cos \theta - g \sin \theta}$$

$$L_0 = \frac{1}{2} a t^2 = \frac{1}{2} (\mu g \cos \theta - g \sin \theta) \cdot \frac{\omega_0^2 R^2}{\left( \frac{7}{2} \mu g \cos \theta - g \sin \theta \right)^2}$$

$$\cancel{R \omega_0} \quad \omega_t = \omega_0 - \frac{5 \mu g \cos \theta \omega_0}{7 \mu g \cos \theta - 2 g \sin \theta} \quad v_t = \frac{\omega_0 R (\mu g \cos \theta - g \sin \theta)}{\frac{7}{2} \mu g \cos \theta - g \sin \theta}$$

$$E_k = \frac{1}{2} m v_t^2 + \frac{1}{2} \cdot \frac{2}{5} m R^2 \cdot \omega_t^2 = m g L_2$$

$$L_{tot} = L_0 + L_2$$

8. case I.

$$R m v_0 = (m R^2 + \frac{1}{2} m_0 R^2) \omega \quad \omega = \frac{2 m v_0}{(2m + m_0) R} \quad E = \frac{m^2 v_0^2}{2m + m_0}$$

$$\text{case II.} \quad \begin{cases} m v_0 = m v_f + m_0 v_{cm} \quad ① \end{cases}$$

$$\bullet \quad \begin{cases} m v_0 R = m (v_f R) + \frac{1}{2} m_0 R^2 \omega \quad ② \end{cases}$$

$$v_f - v_{cm} = \omega R \quad ③$$

$$(m + m_0) v_f = \left[ \frac{m_0}{m + \frac{1}{2} m_0} + 1 \right] m v_0 \quad v_f = \frac{(2m + 3m_0) m v_0}{(2m + m_0)(m + m_0)}$$

$$v_{cm} = \frac{m_0^2 \cdot m}{2m^2 + 3m_0 m + m_0^2} \cdot \frac{1}{m_0} = \frac{m m_0}{2m^2 + 3m m_0 + m_0^2}$$

$$E_k = \frac{1}{2} m v_f^2 + \frac{1}{2} m_0 v_{cm}^2 + \frac{1}{2} \cdot \frac{1}{2} m R^2 \omega^2 = \frac{3}{4} m v_f^2 + \left( \frac{1}{2} m_0 + \frac{1}{4} m \right) v_{cm}^2 + \frac{m}{2} v_f v_{cm}$$



② 代入 ①:  $mV_0 = m(\omega R + V_{cm}) + m_0 V_{cm}$

② - ①:  $\frac{1}{2} m_0 R \omega = m_0 V_{cm}$

解得  $V_{cm} = \frac{mV_0}{m_0 + 3m}$   $\omega = \frac{2V_{cm}}{R} = \frac{2mV_0}{(m_0 + 3m)R}$   $V_f = 3V_{cm}$

$$E_2 = \frac{1}{2} m V_f^2 + \frac{1}{2} m_0 V_{cm}^2 + \frac{1}{2} \cdot \frac{1}{2} m_0 R^2 \cdot \omega^2$$

$$= \frac{1}{2} m (3V_{cm})^2 + \frac{1}{2} m_0 V_{cm}^2 + \frac{1}{4} m_0 R^2 \cdot \frac{4V_{cm}^2}{R^2}$$

$$= \frac{3}{2} m_0 V_{cm}^2 + \frac{9}{2} m V_{cm}^2 = 6 m V_{cm}^2 = 6m \cdot \left( \frac{mV_0}{m_0 + 3m} \right)^2$$

⑨  $fR = \frac{2}{5} m R^2 \frac{d\omega}{dt}$  | 在碰撞中

$f = \mu mg$  |  $\Delta p = 2mV_0 = \int N \cdot dt$

$\therefore \frac{d\omega}{dt} = \frac{5\mu g}{2R}$  |  $\int f dt = \mu R \cdot \int N dt = \int \tau dt = \Delta J = \frac{2}{5} m R^2 \Delta \omega$

$\therefore \Delta \omega = \frac{2mV_0 \mu R}{m R^2} \cdot \frac{5}{2} = \frac{5\mu V_0}{R}$

注意到  $\omega_0 + \Delta \omega < 0$  (摩擦最多使球停转)

其中:  $\omega_0 = -\frac{V_0}{R}$   $\therefore 5\mu - 1 < 0 \quad \mu < \frac{1}{5}$

$$\begin{cases} \omega(t) = -\frac{V_0}{R} + \frac{5\mu g}{2R} t + \frac{5\mu V_0}{R} \\ v(t) = V_0 - \mu g t \end{cases}$$

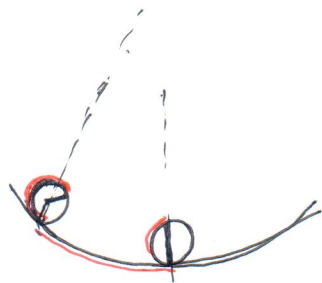
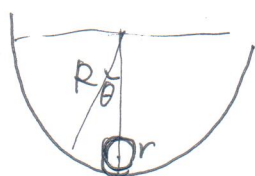
当  $t_0$  时达到纯滚动

$R\omega(t_0) = v(t_0)$  解得  $t_0 = \frac{4 - 10\mu}{7\mu g} V_0$   $v(t_0) = V_0 - \mu g t_0 \quad (\mu < \frac{1}{5})$

⑩  $N = \frac{2mV_0}{\Delta t}$   $f = \mu N < mg$   $\mu < \frac{g\Delta t}{2V_0}$

10. 见 KK 7.6

11.



设小球转过弧度  $\theta$ , 相对于瞬时的法线  
转过  $\varphi$ , 由纯滚动  
 $R\theta = \varphi r$

小球相对于地面转过角度  $\theta + \varphi$ . 设小球自转前速度

$$\omega = \frac{d(\theta + \varphi)}{dt} = \frac{d}{dt} \left( \frac{R}{r} \theta + \theta \right) = \frac{R+r}{r} \frac{d\theta}{dt}$$

$$E_{k1} = \frac{1}{2} m \cdot V_{cm}^2 = \frac{1}{2} m \left[ (R-r) \frac{d\theta}{dt} \right]^2 = \frac{1}{2} m (R-r)^2 \left( \frac{d\theta}{dt} \right)^2$$

$$E_{k2} = \frac{1}{2} I \omega^2 = \frac{1}{2} \cdot \frac{2}{5} m r^2 \left( \frac{R+r}{r} \right)^2 \left( \frac{d\theta}{dt} \right)^2 = \frac{1}{5} m \cdot (R+r)^2 \left( \frac{d\theta}{dt} \right)^2$$

由能量守恒.

$$E_{k1} + E_{k2} + mg(R-r) \cdot (1 - \cos\theta) = E_0.$$

$$\left[ \frac{1}{5} m (R+r)^2 + \frac{1}{2} m (R-r)^2 \right] \left( \frac{d\theta}{dt} \right)^2 + mg(R-r) \cdot \frac{\theta^2}{2} = E_0.$$

$$C_1 = \frac{1}{5} m (R+r)^2 + \frac{1}{2} m (R-r)^2 \quad C_2 = \frac{mg(R-r)}{2} \quad \text{对上式两边同时求导.}$$

$$2C_1 \cdot \left( \frac{d\theta}{dt} \right) \cdot \frac{d^2\theta}{dt^2} + 2C_2 \theta \cdot \frac{d\theta}{dt} = 0$$

$$\frac{d\theta}{dt} \cdot \left( C_1 \frac{d^2\theta}{dt^2} + C_2 \theta \right) = 0.$$

$$\text{得到 } \omega = \sqrt{\frac{C_2}{C_1}} = \sqrt{\frac{\frac{mgR}{2} (1 - \frac{r}{R})}{\frac{1}{5} m R^2 (1 + \frac{r}{R})^2 + \frac{1}{2} m R^2 (1 - \frac{r}{R})^2}}$$

$$\eta = \frac{r}{R}. \text{ 展开至一阶.}$$

$$= \sqrt{\frac{\frac{mgR}{2} (1 - \eta)}{\frac{1}{10} m R^2 (7 - 6\eta)}} = \sqrt{\frac{5g}{R}} \cdot \sqrt{\frac{1 - \eta}{7 - 6\eta}}$$

12. 方程与 1 类似

$$\begin{cases} v \cdot \sin\theta = \frac{L_0}{ml} \\ v^2 \cos\theta = gl \sin^2\theta \end{cases}$$

$$\text{消去 } v \cdot \left\{ \left( \frac{L_0}{ml \sin\theta} \right)^2 \cdot \cos\theta = gl \sin^2\theta \Rightarrow \left( \frac{L_0}{mr} \right)^2 \cos\theta = gl \left( \frac{r}{l} \right)^2 \right. \\ \left. \sin\theta = \frac{r}{l} \right.$$

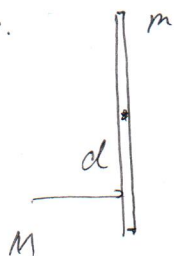
当  $\theta \sim 0$  时  $\sin\theta \sim \theta$   $\cos\theta \sim 1$  (泰勒展开至 1 阶)

$$\left( \frac{L_0}{mr} \right)^2 = g \cdot l \cdot \left( \frac{r}{l} \right)^2 \quad L_0^2 = m^2 g \frac{r^4}{l}$$

$$\text{当 } \theta \sim \frac{\pi}{2} \text{ 时 } \cos\theta = 0 - d\theta \quad \sin\theta = 1 \quad \frac{r}{l} = 1 \quad \frac{gl}{d\theta} = \left( \frac{L_0}{ml} \right)^2 \quad m^2 g l^2 = L_0^2 \cdot d\theta$$

$$l = r \sim 0.$$

13.



$$\begin{cases} MV_0 = MV_f + mV_{cm} \\ MV_0 d = MV_f d + I\omega + mV_{cm} \times 0 \\ \frac{1}{2}MV_0^2 = \frac{1}{2}MV_f^2 + \frac{1}{2}mV_{cm}^2 + \frac{1}{2}I\omega^2 \end{cases}$$

$$\frac{1}{2}MV_0^2 = \frac{1}{2}MV_f^2 + \frac{M^2(V_0 - V_f)^2}{2m} + \frac{M^2(V_0 - V_f)^2 d^2}{2I}$$

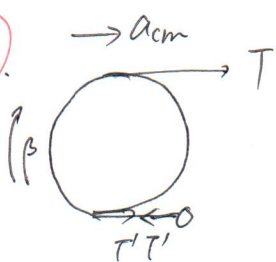
$$\frac{M}{2}(V_0 - V_f)(V_0 + V_f) = \left[ \frac{M^2}{2m} + \frac{M^2 d^2}{2I} \right] (V_0 - V_f)^2 \quad \therefore V_0 \neq V_f \text{ (否则解为碰撞)}$$

$$\frac{M}{2}(V_0 + V_f) = \left( \frac{M^2}{2m} + \frac{M^2 d^2}{2I} \right) (V_0 - V_f)$$

$$\therefore V_f = \frac{\frac{M^2}{2m} + \frac{M^2 d^2}{2I} - \frac{M}{2}}{\frac{M^2}{2m} + \frac{M^2 d^2}{2I} + \frac{M}{2}} V_0 \quad \omega = \frac{M(V_0 - V_f)d}{I}$$

$$V_{cm} = \frac{M(V_0 - V_f)}{m}$$

14.



同轴质点运动:

$$T + T' = m \cdot a_{cm}$$

在质心系下:

$$T' + m \cdot a_{cm} = ma$$

$$a = \beta r$$

$$dT = (T(\theta + d\theta) - T(\theta))r \therefore T = r \int T(\theta + d\theta) - T(\theta) d\theta = r(T - T')$$

$$r(T - T') = \frac{1}{2}mr\beta$$

$$3T = 2ma \quad a = \frac{3T}{2m} \quad a_{cm} = \frac{5T}{4m}$$

$$\text{小球相对地面: } \frac{3T}{2m} - \frac{5T}{4m} = \frac{6T - 5T}{4m} = \frac{T}{4m} \text{ 向左}$$

15.

~~$kx - f = ma$~~

$$\begin{cases} -kx - f = ma \\ f - ma = ma' \\ a' = \beta R \\ I\beta = -Rf \end{cases}$$

$$\beta = -\frac{2f}{mR}$$

$$a' = -\frac{2f}{m}$$

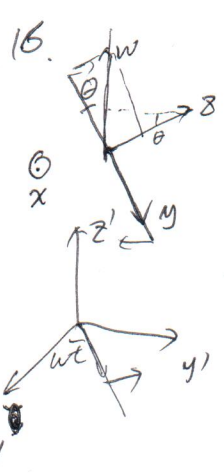
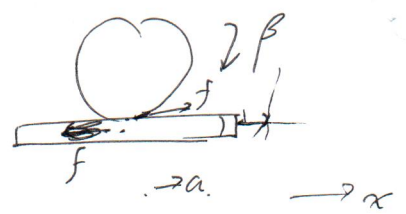
$$f - ma = -2f \quad \Rightarrow \quad f = \frac{ma}{3}$$

$$-kx - \frac{ma}{3} - ma = 0$$

$$\frac{4m}{3} \frac{d^2x}{dt^2} + kx = 0$$

$$\frac{d^2x}{dt^2} + \frac{3k}{4m} x = 0$$

$$\omega = \sqrt{\frac{3k}{4m}}$$



$$\omega = (0, -\omega \cos \theta, \omega \sin \theta)$$

$$J = I\omega = (0, 0, \frac{ml^2}{12} \omega \sin \theta) \quad \text{let } J_0 = \frac{ml^2}{12} \omega \sin \theta$$

换回实验室系  $(J_0 \cos \theta \cos \omega t, J_0 \sin \theta \cos \omega t, J_0 \sin \theta)$

$$\frac{dJ}{dt} = (\omega J_0 \cos \theta \cdot -\sin \omega t, \omega J_0 \cos \theta \cdot \cos \omega t, 0)$$

$$|\tau| = \omega J_0 \cos \theta = 2F \cdot \frac{l}{2} \cos \theta$$

答案不一样?

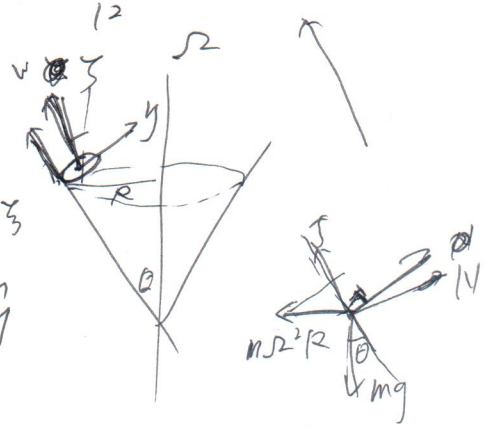
$$\therefore F = \frac{\omega}{l} \cdot J_0 = \frac{\omega}{l} \cdot \frac{ml^2}{12} \omega \sin \theta = \frac{ml^2 \omega^2 \sin \theta}{12}$$

17.

~~$m \Omega^2 R \sin \theta$~~   $mg \sin \theta + m \Omega^2 R \cos \theta = N$  ①

$m \Omega^2 R \sin \theta + f = mg \cos \theta$  ② 无向内摩擦

$$\vec{L} = \frac{d\vec{L}}{dt} \quad \vec{\Omega} = \Omega \cos \theta \hat{z} + \Omega \sin \theta \hat{y}$$



$$\vec{L} = \begin{pmatrix} mr^2 \\ \frac{mr^2}{2} \\ \frac{mr^2}{2} \end{pmatrix} \begin{pmatrix} \omega + \Omega \cos \theta \\ \Omega \sin \theta \\ 0 \end{pmatrix} = mr^2 (\omega + \Omega \cos \theta) \hat{z} + \frac{mr^2}{2} (\Omega \sin \theta) \hat{y}$$

$$\frac{d\vec{L}}{dt} = \vec{\Omega} \times \vec{L} = \Omega \sin \theta \hat{y} \times \hat{z} \quad \frac{d\vec{y}}{dt} = \vec{\Omega} \times \vec{y} = \Omega \cos \theta \hat{z} \times \hat{y}$$



$$\tau = (-r \cdot \hat{y}) \times (f \cdot \hat{z})$$

$$= -rf \cdot \hat{y} \times \hat{z} = \frac{dL}{dt} = mr^2(\omega + \Omega \cos \theta) \Omega \sin \theta \cdot \hat{y} \times \hat{z} + \frac{mr^2}{2} (\Omega \sin \theta) \Omega \cos \theta \cdot \hat{z} \times \hat{y}$$

$$\therefore -rf = mr^2(\omega + \Omega \cos \theta) \Omega \sin \theta - \frac{mr^2}{2} (\Omega \sin \theta) \Omega \cos \theta$$

纯滚动条件  $\vec{\Omega} \times \vec{R} + \vec{\omega} \times \vec{r} = 0$ . (已经做过 - 类似题)

$$\Omega R + \omega r = 0 \quad \therefore \omega = -\frac{R}{r} \Omega$$

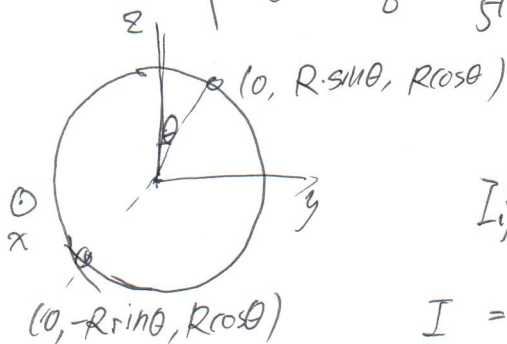
$$-rf = mr^2 \omega \cdot \Omega \sin \theta = m \cdot r^2 \cdot -\frac{R}{r} \Omega^2 \sin \theta$$

答案不一样?

$$f = mR \Omega^2 \sin \theta$$

$$2m \Omega^2 R \sin \theta = mg \cos \theta \quad \Omega = \sqrt{\frac{g \cos \theta}{2R \sin \theta}} = \sqrt{\frac{g}{2R \tan \theta}}$$

$$18. \quad I_0 = \begin{pmatrix} \frac{2}{5}MR^2 & 0 & 0 \\ 0 & \frac{2}{5}MR^2 & 0 \\ 0 & 0 & \frac{2}{5}MR^2 \end{pmatrix} \quad \omega_0 = \begin{pmatrix} 0 \\ 0 \\ \omega \end{pmatrix}$$



$$I'_{ij} = \int (r^2 \delta_{ij} - r_i r_j) dm$$

$$I = 2m \begin{pmatrix} R^2 & 0 & 0 \\ 0 & R^2 & 0 \\ 0 & 0 & R^2 \end{pmatrix} - 2m \begin{pmatrix} 0 \\ R \sin \theta \\ R \cos \theta \end{pmatrix} \begin{pmatrix} 0, R \sin \theta, R \cos \theta \end{pmatrix}$$

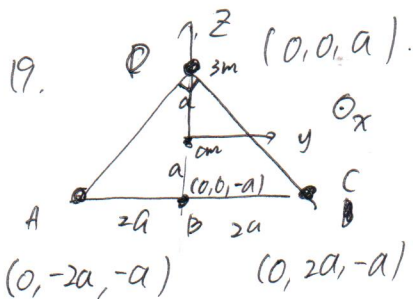
$$= 2m \begin{pmatrix} R^2 & 0 & 0 \\ 0 & R^2 \cos^2 \theta & -R^2 \sin \theta \cos \theta \\ 0 & -R^2 \sin \theta \cos \theta & R^2 \sin^2 \theta \end{pmatrix}$$

$$\therefore I_f = R^2 \begin{pmatrix} \frac{2}{5}M - 2m & 0 & 0 \\ 0 & \frac{2m}{5} - 2m \cos^2 \theta & -2m \sin \theta \cos \theta \\ 0 & -2m \sin \theta \cos \theta & \frac{2m}{5} - 2m \sin^2 \theta \end{pmatrix}$$

$$I_0 \omega_0 = I_f \omega_f$$

$$\frac{2m}{5} - 2m \cos^2 \theta \omega_y - 2m \sin \theta \cos \theta \omega_z = 0$$

$$\frac{\omega_y}{\omega_z} = \frac{\frac{2m}{5} - 2m \cos^2 \theta}{-2m \sin \theta \cos \theta}$$



cm:

$$I_H = m \cdot \begin{pmatrix} 5a^2 & 0 & 0 \\ 0 & 5a^2 & 0 \\ 0 & 0 & 5a^2 \end{pmatrix} - m \cdot \begin{pmatrix} 0 \\ -2a \\ -a \end{pmatrix} \begin{pmatrix} 0 & -2a & -a \end{pmatrix}$$

$$= ma^2 \begin{pmatrix} 5 & 0 & 0 \\ 0 & 1 & -2 \\ 0 & -2 & 4 \end{pmatrix}$$

$$I_B = m \begin{pmatrix} a^2 & 0 & 0 \\ 0 & a^2 & 0 \\ 0 & 0 & a^2 \end{pmatrix} - m \cdot \begin{pmatrix} 0 \\ 0 \\ -a \end{pmatrix} \begin{pmatrix} 0 & 0 & -a \end{pmatrix} = ma^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

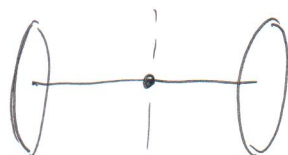
$$I_C = I_H \quad I_Z = 3I_B \quad \therefore I_{tot} = \sum_{i=A,B,C,D} I_i = ma^2 \begin{pmatrix} 14 & 0 & 0 \\ 0 & 6 & -4 \\ 0 & -4 & 8 \end{pmatrix}$$

$$I_{tot} \cdot \vec{\omega} = \int \vec{r} dm \quad \int \vec{r} dm = \begin{pmatrix} \bar{x} & \bar{y} & \bar{z} \\ 0 & 2a & -a \\ -p & 0 & 0 \end{pmatrix} = ap \hat{y} + 2ap \hat{z}$$

$$ma^2 \cdot \begin{pmatrix} 14 & 0 & 0 \\ 0 & 6 & -4 \\ 0 & -4 & 8 \end{pmatrix} \begin{pmatrix} \omega_x \\ \omega_y \\ \omega_z \end{pmatrix} = a \begin{pmatrix} 0 \\ p \\ 2p \end{pmatrix} \quad \omega_x = 0 \quad \begin{pmatrix} 6 & -4 \\ -4 & 8 \end{pmatrix} \begin{pmatrix} \omega_y \\ \omega_z \end{pmatrix} = \frac{p}{ma} \begin{pmatrix} 1 \\ 2 \end{pmatrix}$$

$$\begin{pmatrix} \omega_y \\ \omega_z \end{pmatrix} = \frac{p}{ma} \begin{pmatrix} 0.25 & 0.125 \\ 0.125 & 0.1875 \end{pmatrix} \begin{pmatrix} 1 \\ 2 \end{pmatrix} = \begin{pmatrix} 0.5 \\ 0.5 \end{pmatrix} \quad \therefore \omega_y = \omega_z = \frac{p}{2ma}$$

20.



由KKP298页结论

$$\frac{dL_z}{dt} = 2I \cdot \omega \cdot \Omega$$

叠分开时

这里已经讨论了质心的平动角动量, 但其对时间求导为0.

$$I_{max} = 2mg \cdot \frac{l}{2} = 2I\omega\Omega \quad \therefore \Omega = \frac{mg l}{2I\omega}$$