Density Matrix Renormalization Group for the Heisenberg Model

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I. THE FORMULAS

Some special cases of using the wigner-Eckart theorem:

$$\langle J, J | X_0^{(0)} | J, J \rangle = \frac{1}{\sqrt{2J+1}} \langle J | | X_0^{(0)} | | J \rangle$$

$$\langle J_1', J_2', J' | | X_{1,2}^{(0)} | | J_1, J_2, J \rangle = (-1)^{J_1 + J_2' + J + k} \frac{\sqrt{2J+1}}{\sqrt{2k+1}} \begin{cases} J_1' & J_1 & k \\ J_2 & J_2' & J \end{cases} \langle J_1' | | T_i^{(1)} | | J_1 \rangle \langle J_2' | | T_j^{(1)} | | J_2 \rangle$$

$$\langle J_{sys}', S', J', \alpha' | | T_i^{(1)} | | J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys}' + S' + J + 1} \delta_{S', S} \sqrt{(2J'+1)(2J+1)} \begin{cases} J' & 1 & J \\ J_{sys} & S' & J_{sys}' \end{cases} \langle J_{sys}', \alpha' | | T_i^{(1)} | | J_{sys}, \alpha \rangle$$

$$\langle J_{sys}', S', J', \alpha' | | T_i^{(1)} | | J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J' + 1} \delta_{J_{sys}'}, J_{sys} \sqrt{(2J'+1)(2J+1)} \begin{cases} J' & 1 & J \\ S & J_{sys}' & S' \end{cases} \langle S', \alpha' | | T_i^{(1)} | | S, \alpha \rangle$$

$$\langle J_{sys}', S', J', \alpha' | | X_{ij}^{(1)} | | J_{sys}, S, J, \alpha \rangle =$$

$$\sqrt{(2J'+1)(2K+1)(2J+1)} \begin{cases} J_{sys}' & J_{sys} & 1 \\ S' & S & 1 \\ J' & J & 1 \end{cases} \langle J_{sys}', \alpha' | | T_i^{(1)} | | J_{sys}, \alpha \rangle \langle S' | | T_j^{(1)} | | S \rangle (k=1)$$

II. IRREDUCIBLE TENSOR OPERATORS OF THE HEISENBERG TERMS

The product of two irreducible tensor operators (ITOs) is:

$$\sum_{q_1,q_2} \langle k_1 q_1 k_2 q_2 | kq \rangle T_{q_1}^{(k_1)} T_{q_2}^{(k_2)} = T_q^{(k)} (k = k_1 + k_2, ..., |k_1 - k_2|)$$

$$Q = (U^{(L)}, V^{(L)}) = \sum_{M=-L}^{L} (-1)^M U_M^{(L)} V_{-M}^{(L)}$$
, so:

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$$\sum_{q_1,q_2} \langle k_1 q_1 k_2 q_2 | kq \rangle T_{q_1}^{(k_1)} T_{q_2}^{(k_2)} = T_q^{(k)} \ (k = k_1 + k_2, ..., |k_1 - k_2|)$$
 The scalar product of two irreducible tensor operators (ITOs) is:
$$Q = (U^{(L)}, V^{(L)}) = \sum_{M=-L}^{L} (-1)^M U_M^{(L)} V_{-M}^{(L)}, \text{ so:}$$

$$Q = (U^{(1)}, V^{(1)}) = \sum_{M=-1}^{1} (-1)^M U_M^{(1)} V_{-M}^{(1)} = -U_{-1}^{(1)} V_1^{(1)} - U_1^{(1)} V_{-1}^{(1)} + U_0^{(1)} V_0^{(1)}$$

as the product of two irreducible tensor operators $X_0^{(0)} = [U_L \times V_L]_0 = \frac{(-1)^L}{\sqrt{2L+1}}Q$, so: $X_0^{(0)} = [U_1 \times V_1]_0 = \frac{(-1)}{\sqrt{3}}Q$

For the Heisenberg spin interaction terms, we use the rank-1 ITO $T_{i,q}^{(1)}(q=+1,0,-1)$ to reproduce spin operators, which is: $T_{i,1}^{(1)}=-\frac{1}{\sqrt{2}}S_i^\dagger, T_{i,0}^{(1)}=S_i^z, T_{i,-1}^{(1)}=\frac{1}{\sqrt{2}}S_i^-$ The Hamiltonian of Heisenberg term H_J can be expressed as:

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$$H_J$$
 can be expressed as:
$$H_J = \sum_{i,j} J_{ij} (S_i^x S_j^x + S_i^y S_j^y + S_i^z S_j^z) = \sum_{i,j} J_{ij} (\frac{1}{2} S_i^{\dagger} S_j^{-} + \frac{1}{2} S_i^{-} S_j^{\dagger} + S_i^z S_j^z) = \sum_{i,j} J_{ij} (-T_{i,1}^{(1)} T_{j,-1}^{(1)} - T_{i,-1}^{(1)} T_{j,1}^{(1)} + T_{i,0}^{(1)} T_{j,0}^{(1)}) = \sum_{i,j} J_{ij} Q$$
Thus, the Hamiltonian H_J can be presented as:

$$H_J = -\sqrt{3} \sum_{i,j} J_{ij} X_{0,ij}^{(0)}$$

A. Initialization

There are only one site in system and environment. We need to calculate $\langle J||T^{(1)}||J\rangle$ for initialization.

$$\langle JJ|T_0^{(1)}|JJ\rangle = \frac{1}{\sqrt{2J+1}}\langle JJ10|JJ\rangle\langle J||T^{(1)}||J\rangle = \frac{1}{\sqrt{2J+1}}\frac{J}{\sqrt{J(J+1)}}\langle J||T^{(1)}||J\rangle = J$$

Thus,
$$\langle J||T^{(1)}||J\rangle=\sqrt{J(J+1)(2J+1)}$$

If
$$J = \frac{1}{2}, \langle J || T^{(1)} || J \rangle = \sqrt{\frac{3}{2}}$$

If
$$J = 1$$
, $\langle J || T^{(1)} || J \rangle = \sqrt{6}$.

B. Increasing a new site

In this section, we add a new site in the system block. The irreducible basis is different from before. We need to calculate $\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle$ in the new irreducible basis. There, we need the product of two irreducible tensor operators. There, |J - J'| = 0, 1.

If the site i is in the system block, then

$$\langle J_{sys}', S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys}' + S' + J + 1} \delta_{S', S} \sqrt{(2J' + 1)(2J + 1)} \begin{cases} J' & 1 & J \\ J_{sys} & S' & J_{sys}' \end{cases} \langle J_{sys}', \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle$$
 if $J = J'$
$$\langle J_{sys}', S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys}' + S' + J + 1} \delta_{S', S} (2J + 1) \begin{cases} J & 1 & J \\ J_{sys} & S' & J_{sys}' \end{cases} \langle J_{sys}', \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle$$

$$= (-1)^{J_{sys}' + S + J + 1} (2J + 1) \begin{cases} J & 1 & J \\ J_{sys} & S & J_{sys}' \end{cases} \langle J_{sys}', \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle$$

which would be non-zero only for $J \neq 0$. If J = 0, the 6j – coefficient would be zero.

If the site i is in the added new site, then

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J' + 1} \delta_{J'_{sys}, J_{sys}} \sqrt{(2J' + 1)(2J + 1)} \begin{cases} J' & 1 & J \\ S & J'_{sys} & S' \end{cases} \langle S', \alpha' || T_i^{(1)} || S, \alpha \rangle$$

For the diagonal part with $J=J^{\prime}$, the result can be simplified as

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J + 1} \delta_{J'_{sys}, J_{sys}} (2J + 1) \begin{cases} J & 1 & J \\ S & J'_{sys} & S' \end{cases} \langle S', \alpha' || T_i^{(1)} || S, \alpha \rangle$$

which would be non-zero only for $J \neq 0$. If J = 0, the 6j – coefficient would be zero.

C. Calculating the Hamiltonian

Suppose the site i belong to the system block and the site j is the added new site. Then the Hamiltonian H_J is $H_J = -\sqrt{3} \sum_{i,j} J_{i,j} X_{0,ij}^0$

The matrix element of the Hamiltonian H_J is

$$-\sqrt{3}J_{i,j}\langle J'_{sys},S',J',\alpha'|X^0_{0,ij}|J_{sys},S,J,\alpha\rangle = -\frac{\sqrt{3}}{\sqrt{2J+1}}J_{i,j}\langle J'_{sys},S',J',\alpha'||X^0_{ij}||J_{sys},S,J,\alpha\rangle \\ \text{as } \langle J'_{sys},S',J',\alpha'||X^0_{ij}||J_{sys},S,J,\alpha\rangle = \sqrt{\frac{2J+1}{3}}(-1)^{J_{sys}+S'+1+J} \left\{ \begin{matrix} J'_{sys} & J_{sys} & 1 \\ S & S' & J \end{matrix} \right\} \langle J'_{sys},\alpha'||T^{(1)}_{i}||J_{sys},\alpha\rangle\langle S'||T^{(1)}_{j}||S\rangle$$

Thus, the matrix element of the Hamiltonian H_J is

$$-\sqrt{3}J_{i,j}\langle J'_{sys}, S', J', \alpha' | X^{0}_{0,ij} | J_{sys}, S, J, \alpha \rangle = -J_{i,j}(-1)^{J_{sys}+S'+1+J} \left\{ \begin{matrix} J'_{sys} & J_{sys} & 1 \\ S & S' & J \end{matrix} \right\} \langle J'_{sys}, \alpha' | | T^{(1)}_{i} | | J_{sys}, \alpha \rangle \langle S' | | T^{(1)}_{j} | | S \rangle$$

D. Hamiltonian from the Heisenberg interactions

In this section, we need the matrix-vector multiplication calculations

$$\begin{split} H_{n+1} &= H_{A_n} + H_{B_n} + S_{A_n} \cdot S_{A_n+1} + S_{A_n+1} \cdot S_{B_n+1} + S_{B_n+1} \cdot S_{B_n} \\ H_{n+1} |f\rangle &= |g\rangle, \text{ thus} \\ (H_{A_n} + H_{B_n} + S_{A_n} \cdot S_{A_n+1} + S_{A_n+1} \cdot S_{B_n+1} + S_{B_n+1} \cdot S_{B_n}) |f\rangle &= |g\rangle \\ (H_{A_n}) |f\rangle &= |g_1\rangle, (H_{B_n}) |f\rangle &= |g_2\rangle, (S_{A_n} \cdot S_{A_n+1}) |f\rangle &= |g_3\rangle, (S_{A_n+1} \cdot S_{B_n+1}) |f\rangle &= |g_4\rangle, (S_{B_n+1} \cdot S_{B_n}) |f\rangle &= |g_5\rangle \\ \text{thus, } |g\rangle &= |g_1\rangle + |g_2\rangle + |g_3\rangle + |g_4\rangle + |g_5\rangle \\ \text{generally, } J_{ij} \vec{S}_i \cdot \vec{S}_j |f\rangle &= J_{ij} \vec{S}_i \cdot \vec{S}_j |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= \sum_{J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'} |J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'\rangle \\ &J_{ij} \vec{S}_i \cdot \vec{S}_j |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= \sum_{J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'} \langle J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'\rangle \\ &J_{ij} \vec{S}_i \cdot \vec{S}_j |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle |J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'\rangle \\ &= \sum_{J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'} \langle J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'\rangle |J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'\rangle \\ &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle \\ &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew},$$

1. Hamiltonian from the Heisenberg interactions in the system block

The Hamiltonian contribution is $J_{ij}\vec{S}_i \cdot \vec{S}_j$, where \vec{S}_i and \vec{S}_j belong to the system block. The matrix-vector multiplication is explicitly written as

$$\langle J_{sys}', J_{ns}', J_{sysnew}', J_{env}', J_{ne}', J_{envnew}', J' | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J \rangle \\ f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) \\ \text{As } \vec{S}_i \text{ and } \vec{S}_j \text{ belong to the system block. Thus} \\ \langle J_{sys}', J_{ns}, J_{sysnew}', J_{env}, J_{ne}, J_{envnew}, 0 | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ f(J_{sys}, J_{ns}, J_{sysnew}', J_{env}, J_{ne}, J_{envnew}, 0) &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) \\ \end{cases}$$

We can compute the matrix element as

$$\begin{split} J_{ij}\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -\sqrt{3}J_{ij}\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X^0_{0,ij} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X^0_{ij} | | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys} | | X^0_{ij} | | J_{sys} \rangle \end{split}$$

2. Hamiltonian from the Heisenberg interactions between the system block and the added sites

The Hamiltonian contribution is $J_{ij}\vec{S}_i \cdot \vec{S}_j$, where \vec{S}_i and \vec{S}_j belong to the system block and the added site, respectively. The matrix-vector multiplication is explicitly written as

$$\langle J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J' | J_{ij} \vec{S}_{i} \cdot \vec{S}_{j} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J \rangle$$

$$f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) = g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J)$$

$$As \ \vec{S}_{i} \ and \ \vec{S}_{j} \ belong \ to \ the \ system \ block \ and \ the \ added \ site, \ respectively. \ Thus$$

$$\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | J_{ij} \vec{S}_{i} \cdot \vec{S}_{j} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$f(J_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) = g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0)$$

$$We \ can \ compute \ the \ matrix \ element \ as$$

$$J_{ij} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | \vec{S}_{i} \cdot \vec{S}_{j} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= -\sqrt{3} J_{ij} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X^{0}_{0,ij} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X^{0}_{0,ij} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X^{0}_{0,ij} | J_{sys}, J_{sys}, J_{sys}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | J'_{sys}, J_{sysnew}, J_{env}, J_{sysnew}, J_{env}, J_{sysnew}, J_{env}, J_{sysnew}, J_{env}, J_{sysnew}, J_{env}, J_{sysnew}, J_{sysnew},$$

3. Hamiltonian from the Heisenberg interactions between the environment block and the added sites

The Hamiltonian contribution is $J_{ij}\vec{S}_i\cdot\vec{S}_j$, where \vec{S}_i and \vec{S}_j belong to the environment block and the added site, respectively. The matrix-vector multiplication is explicitly written as

$$\langle J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J' | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J \rangle \\ f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) \\ \text{As } \vec{S}_i \text{ and } \vec{S}_j \text{ belong to the environment block and the added site, respectively. Thus} \\ \langle J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0 | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ f(J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0) &= g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) \\ \text{We can compute the matrix element as} \\ J_{ij} \langle J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0 | \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -\sqrt{3} J_{ij} \langle J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0 | X_{0,ij}^0 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{envnew}+1}} \langle J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0 | X_{0,ij}^0 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{envnew}+1}} \langle J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0 | X_{0,ij}^0 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle \\ &= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{envnew}+1}} \langle J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J_{ne}, J'_{envnew}, 0 | X_{0,ij}^0 | J_{sys}, J_{ns}, J_{sysnew}, J'_{env}, J'_{in}, J'_{in},$$

4. Hamiltonian from the Heisenberg interactions between the added sites

The Hamiltonian contribution is $J_{ij}\vec{S}_i \cdot \vec{S}_j$, where \vec{S}_i and \vec{S}_j belong to the added sites, respectively. The matrix-vector multiplication is explicitly written as

$$\langle J'_{sys}, J'_{ns}, J'_{env}, J'_{ne}, J' | J_{ij}\vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{env}, J_{ne}, J \rangle$$

$$\begin{split} &f(J_{sys},J_{ns},J_{env},J_{ne},J) = g(J_{sys},J_{ns},J_{env},J_{ne},J) \\ &\text{As } \vec{S_i} \text{ and } \vec{S_j} \text{ belong to the added sites, respectively. Thus} \\ &\langle J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne},0|J_{ij}\vec{S_i}\cdot\vec{S_j}|J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne},0\rangle \\ &f(J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne},0) = g(J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne},0) \\ &\text{We can compute the matrix element as} \\ &\langle J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne}|J_{ij}\vec{S_i}\cdot\vec{S_j}|J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne}\rangle \\ &= -\sqrt{3}J_{ij}\langle J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne}|X_{0,ij}^0|J_{sys},J_{ns},J_{env},J_{ne},J_{ns,ne}\rangle \\ &= -J_{ij}\frac{\sqrt{3}}{\sqrt{2J_{ns,ne}+1}}\langle J_{ns},J_{ne},J_{ns,ne}||X_{0,ij}^0||J_{ns},J_{ne},J_{ns,ne}\rangle \\ &= -J_{ij}\frac{\sqrt{3}}{\sqrt{2J_{ns,ne}+1}}\frac{\sqrt{2J_{ns,ne}+1}}{\sqrt{3}}(-1)^{J_{ns}+J_{ne}+1+J_{ns,ne}}\left\{ \begin{matrix} J_{ns}&J_{ns}&1\\J_{ne}&J_{ne}&J_{ns}&1 \end{matrix} \\ J_{ne}&J_{ne}&J_{ns,ne} \end{matrix} \right\}\langle J_{ns},||T_i^{(1)}||J_{ns}\rangle\langle J_{ne}||T_j^{(1)}||J_{ne}\rangle \\ &= -J_{ij}(-1)^{J_{ns}+J_{ne}+1+J_{ns,ne}}\left\{ \begin{matrix} J_{ns}&J_{ns}&1\\J_{ne}&J_{ns,ne} \end{matrix} \right\}\langle J_{ns},||T_i^{(1)}||J_{ns}\rangle\langle J_{ne}||T_j^{(1)}||J_{ne}\rangle \end{split}$$

5. Hamiltonian from the Heisenberg interactions between the system and environment

The Hamiltonian contribution is $J_{ij}\vec{S}_i \cdot \vec{S}_j$, where \vec{S}_i and \vec{S}_j belong to the system and environment, respectively. The matrix-vector multiplication is explicitly written as

$$\langle J'_{sys}, J'_{ns}, J'_{env}, J'_{ne}, J' | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{env}, J_{ne}, J \rangle$$

$$f(J_{sys}, J_{ns}, J_{env}, J_{ne}, J) = g(J_{sys}, J_{ns}, J_{env}, J_{ne}, J)$$

$$As \ \vec{S}_i \ and \ \vec{S}_j \ belong \ to \ the \ system \ and \ environment, \ respectively. \ Thus$$

$$\langle J'_{sys}, J_{ns}, J'_{env}, J_{ne}, J_{sys,env}, 0 | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env}, 0 \rangle$$

$$f(J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env}, 0) = g(J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env}, 0)$$

$$We \ can \ compute \ the \ matrix \ element \ as$$

$$\langle J'_{sys}, J_{ns}, J'_{env}, J_{ne}, J_{sys,env} | J_{ij} \vec{S}_i \cdot \vec{S}_j | J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env} \rangle$$

$$= -\sqrt{3} J_{ij} \langle J'_{sys}, J_{ns}, J'_{env}, J_{ne}, J_{sys,env} | X^0_{0,ij} | J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env} \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2_{sys,env}+1}} \langle J'_{sys}, J_{ns}, J'_{env}, J_{ne}, J_{sys,env} | |X^0_{ij} | J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env} \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sys,env}+1}} \langle J'_{sys}, J_{ns}, J'_{env}, J_{ne}, J_{sys,env} | |X^0_{ij} | |J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env} \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sys,env}+1}} \langle J'_{sys}, J_{ns}, J'_{env}, J_{ne}, J_{sys,env} | |X^0_{ij} | |J_{sys}, J_{ns}, J_{env}, J_{ne}, J_{sys,env} \rangle$$

$$= -J_{ij} \frac{\sqrt{3}}{\sqrt{2J_{sys,env}+1}} \langle J'_{sys}, J_{sys}, J'_{sys}, J'_{sy$$

III. IRREDUCIBLE TENSOR OPERATORS OF THE K TERMS

$$H = \sum_{\langle i,j \rangle} K(\mathbf{S}_{i} \cdot \mathbf{S}_{j})^{2} = \sum_{\langle i,j \rangle} K(1/2(S_{i}^{\dagger} S_{j}^{-} + S_{i}^{-} S_{j}^{\dagger}) + (S_{i}^{z} S_{j}^{z}))(1/2(S_{i}^{\dagger} S_{j}^{-} + S_{i}^{-} S_{j}^{\dagger}) + (S_{i}^{z} S_{j}^{z}))$$

$$= \sum_{\langle i,j \rangle} K(1/4(S_{i}^{\dagger} S_{j}^{-} S_{i}^{\dagger} S_{j}^{-}) + 1/4(S_{i}^{\dagger} S_{j}^{-} S_{i}^{-} S_{j}^{\dagger}) + 1/2(S_{i}^{\dagger} S_{j}^{-} S_{i}^{z} S_{j}^{z})$$

$$+1/4(S_{i}^{-} S_{j}^{\dagger} S_{i}^{\dagger} S_{j}^{-}) + 1/4(S_{i}^{-} S_{j}^{\dagger} S_{i}^{-} S_{j}^{\dagger}) + 1/2(S_{i}^{-} S_{j}^{\dagger} S_{i}^{z} S_{j}^{z})$$

$$+1/2(S_{i}^{z} S_{j}^{z} S_{i}^{\dagger} S_{j}^{-}) + 1/2(S_{i}^{z} S_{j}^{z} S_{i}^{-} S_{j}^{\dagger}) + (S_{i}^{z} S_{j}^{z} S_{i}^{z} S_{j}^{z})$$

$$= \sum_{\langle i,j \rangle} K(1/4(S_{i}^{\dagger} S_{i}^{\dagger} S_{j}^{-} S_{j}^{-}) + 1/4(S_{i}^{-} S_{i}^{-} S_{j}^{\dagger}) + 1/2(S_{i}^{\dagger} S_{i}^{z} S_{j}^{-} S_{j}^{z})$$

$$+1/4(S_{i}^{-} S_{i}^{\dagger} S_{j}^{\dagger} S_{j}^{-}) + 1/4(S_{i}^{-} S_{i}^{-} S_{j}^{\dagger} S_{j}^{\dagger}) + 1/2(S_{i}^{-} S_{i}^{z} S_{j}^{\dagger} S_{j}^{z})$$

$$+1/2(S_{i}^{z} S_{i}^{\dagger} S_{j}^{z} S_{j}^{-}) + 1/2(S_{i}^{z} S_{i}^{-} S_{j}^{z} S_{j}^{\dagger}) + (S_{i}^{z} S_{i}^{z} S_{j}^{z} S_{j}^{z})$$

$$+1/2(S_{i}^{z} S_{i}^{\dagger} S_{j}^{z} S_{j}^{-}) + 1/2(S_{i}^{z} S_{i}^{-} S_{j}^{z} S_{j}^{\dagger}) + (S_{i}^{z} S_{i}^{z} S_{j}^{z} S_{j}^{z})$$

as,
$$T_0^{(0)} = \langle 111 - 1|00\rangle T_1^{(1)} T_{-1}^{(1)} + \langle 1 - 111|00\rangle T_{-1}^{(1)} T_1^{(1)} + \langle 1010|00\rangle T_0^{(1)} T_0^{(1)} = \frac{1}{\sqrt{3}} T_1^{(1)} T_{-1}^{(1)} + \frac{1}{\sqrt{3}} T_{-1}^{(1)} T_1^{(1)} - \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} = \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} + \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} = \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} + \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} = \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} + \frac{1}{\sqrt{3}} T_0^{(1)} T_0^{(1)} = \frac{1}{\sqrt{3}} T_0^{($$

$$\begin{split} T_0^{(1)} &= \langle 111-1|10\rangle T_1^{(1)}T_{-1}^{(1)} + \langle 1-111|10\rangle T_{-1}^{(1)}T_1^{(1)} + \langle 1010|10\rangle T_0^{(1)}T_0^{(1)} = \frac{1}{\sqrt{2}}T_1^{(1)}T_{-1}^{(1)} - \frac{1}{\sqrt{2}}T_{-1}^{(1)}T_1^{(1)} \\ T_1^{(1)} &= \langle 1011|11\rangle T_0^{(1)}T_1^{(1)} + \langle 1110|11\rangle T_1^{(1)}T_0^{(1)} = -\frac{1}{\sqrt{2}}T_0^{(1)}T_1^{(1)} + \frac{1}{\sqrt{2}}T_1^{(1)}T_0^{(1)} \\ T_{-1}^{(1)} &= \langle 101-1|1-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1-110|1-1\rangle T_{-1}^{(1)}T_0^{(1)} = \frac{1}{\sqrt{2}}T_0^{(1)}T_{-1}^{(1)} - \frac{1}{\sqrt{2}}T_{-1}^{(1)}T_0^{(1)} \\ T_0^{(2)} &= \langle 111-1|20\rangle T_1^{(1)}T_{-1}^{(1)} + \langle 1-111|20\rangle T_{-1}^{(1)}T_1^{(1)} + \langle 1010|20\rangle T_0^{(1)}T_0^{(1)} = \frac{1}{\sqrt{6}}T_1^{(1)}T_{-1}^{(1)} + \frac{1}{\sqrt{6}}T_{-1}^{(1)}T_1^{(1)} + \sqrt{\frac{2}{3}}T_0^{(1)}T_0^{(1)} \\ T_1^{(2)} &= \langle 1011|21\rangle T_0^{(1)}T_1^{(1)} + \langle 1110|21\rangle T_1^{(1)}T_0^{(1)} = \frac{1}{\sqrt{2}}T_0^{(1)}T_1^{(1)} + \frac{1}{\sqrt{2}}T_1^{(1)}T_0^{(1)} \\ T_{-1}^{(2)} &= \langle 101-1|2-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1-110|2-1\rangle T_{-1}^{(1)}T_0^{(1)} = \frac{1}{\sqrt{2}}T_0^{(1)}T_{-1}^{(1)} + \frac{1}{\sqrt{2}}T_{-1}^{(1)}T_0^{(1)} \\ T_2^{(2)} &= \langle 1111|22\rangle T_1^{(1)}T_1^{(1)} = T_1^{(1)}T_1^{(1)} \\ T_{-2}^{(2)} &= \langle 1-11-1|2-2\rangle T_{-1}^{(1)}T_{-1}^{(1)} = T_{-1}^{(1)}T_{-1}^{(1)} \end{aligned}$$

so

$$H = \sum_{\langle i,j \rangle} K_{i,j} (T_{i,1}^{(1)} T_{i,1}^{(1)} T_{j,-1}^{(1)} T_{j,-1}^{(1)} + T_{i,1}^{(1)} T_{i,-1}^{(1)} T_{j,-1}^{(1)} T_{j,1}^{(1)} - T_{i,1}^{(1)} T_{i,0}^{(1)} T_{j,-1}^{(1)} T_{j,0}^{(1)}$$

$$+ T_{i,-1}^{(1)} T_{i,1}^{(1)} T_{j,1}^{(1)} T_{j,-1}^{(1)} + T_{i,-1}^{(1)} T_{i,-1}^{(1)} T_{j,1}^{(1)} T_{j,1}^{(1)} - T_{i,-1}^{(1)} T_{i,0}^{(1)} T_{j,1}^{(1)} T_{j,0}^{(1)}$$

$$- T_{i,0}^{(1)} T_{i,1}^{(1)} T_{j,0}^{(1)} T_{j,-1}^{(1)} - T_{i,0}^{(1)} T_{i,-1}^{(1)} T_{j,0}^{(1)} T_{j,1}^{(1)} + T_{i,0}^{(1)} T_{i,0}^{(1)} T_{j,0}^{(1)} T_{j,0}^{(1)}$$

$$(2)$$

$$\begin{split} X_0^{(0)} &= ([T_i^{(0)} \times T_j^{(0)}]_0 + [T_i^{(1)} \times T_j^{(1)}]_0 + [T_i^{(2)} \times T_j^{(2)}]_0) \\ &= (Q_{ij}^{(0)} - \frac{1}{\sqrt{3}}Q_{ij}^{(1)} + \frac{1}{\sqrt{5}}Q_{ij}^{(2)}) \\ &= T_{i,0}^{(0)}T_{j,0}^{(0)} - \frac{1}{\sqrt{3}}(-T_{i,1}^{(1)}T_{j,-1}^{(1)} - T_{i,-1}^{(1)}T_{j,1}^{(1)} + T_{i,0}^{(1)}T_{j,0}^{(1)}) + \frac{1}{\sqrt{5}}(T_{i,2}^{(2)}T_{j,-2}^{(2)} - T_{i,1}^{(2)}T_{j,-1}^{(2)} + T_{i,0}^{(2)}T_{j,0}^{(2)} - T_{i,-1}^{(2)}T_{j,1}^{(2)} + T_{i,-2}^{(2)}T_{j,2}^{(2)}) \\ &= (\frac{1}{\sqrt{3}}T_{i,1}^{(1)}T_{i,-1}^{(1)} + \frac{1}{\sqrt{3}}T_{i,-1}^{(1)}T_{i,1}^{(1)} - \frac{1}{\sqrt{3}}T_{i,0}^{(1)}T_{i,0}^{(1)}) \\ &\qquad \qquad (\frac{1}{\sqrt{3}}T_{j,-1}^{(1)}T_{j,-1}^{(1)} + \frac{1}{\sqrt{3}}T_{j,-1}^{(1)}T_{j,1}^{(1)} - \frac{1}{\sqrt{3}}T_{i,0}^{(1)}T_{i,0}^{(1)}) \\ &\qquad \qquad - \frac{1}{\sqrt{3}}(-(-\frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{i,1}^{(1)} + \frac{1}{\sqrt{2}}T_{i,1}^{(1)}T_{i,0}^{(1)}) \\ &\qquad \qquad (\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,-1}^{(1)} - \frac{1}{\sqrt{2}}T_{j,-1}^{(1)}T_{j,0}^{(1)}) - (\frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{i,-1}^{(1)} - \frac{1}{\sqrt{2}}T_{i,-1}^{(1)}T_{i,0}^{(1)}) (-\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,1}^{(1)} + \frac{1}{\sqrt{2}}T_{j,1}^{(1)}T_{j,0}^{(1)}) \\ &\qquad \qquad + (\frac{1}{\sqrt{2}}T_{i,1}^{(1)}T_{i,-1}^{(1)} - \frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{i,1}^{(1)} + \frac{1}{\sqrt{2}}T_{j,1}^{(1)}T_{j,0}^{(1)}) (\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,-1}^{(1)} - \frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,0}^{(1)}) \\ &\qquad \qquad + (\frac{1}{\sqrt{5}}((T_{i,1}^{(1)}T_{i,1}^{(1)})(T_{j,-1}^{(1)}T_{j,-1}^{(1)}) - (\frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{i,1}^{(1)} + \frac{1}{\sqrt{2}}T_{i,1}^{(1)}T_{i,0}^{(1)}) (\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,-1}^{(1)} + \frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,0}^{(1)}) \\ &\qquad \qquad (\frac{1}{\sqrt{6}}T_{i,1}^{(1)}T_{i,-1}^{(1)} + \frac{1}{\sqrt{6}}T_{i,-1}^{(1)}T_{i,1}^{(1)}) (\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,1}^{(1)} + \frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,0}^{(1)}) \\ &\qquad \qquad \qquad - (\frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{i,-1}^{(1)} + \frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{j,1}^{(1)}) (\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,0}^{(1)}) + T_{i,-1}^{(1)}T_{i,-1}^{(1)}T_{j,1}^{(1)}T_{j,0}^{(1)}) \\ &\qquad \qquad \qquad - (\frac{1}{\sqrt{2}}T_{i,0}^{(1)}T_{i,-1}^{(1)} + \frac{1}{\sqrt{2}}T_{i,-1}^{(1)}T_{i,0}^{(1)}) (\frac{1}{\sqrt{2}}T_{j,0}^{(1)}T_{j,0}^{(1)}) + T_{i,-1}^{(1)}T_{i,-1}^{(1)}T_{j,1}^{(1)}T_{j,0}^{(1$$

$$(Q_{ij}^{(0)} - Q_{ij}^{(1)} + Q_{ij}^{(2)}) = T_{i,0}^{(0)} T_{j,0}^{(0)} - (-T_{i,1}^{(1)} T_{j,-1}^{(1)} - T_{i,-1}^{(1)} T_{j,1}^{(1)}) + (T_{i,2}^{(2)} T_{j,-2}^{(2)} - T_{i,1}^{(2)} T_{j,-1}^{(2)} + T_{i,0}^{(2)} T_{j,0}^{(2)} - T_{i,-1}^{(2)} T_{j,1}^{(2)} + T_{i,-2}^{(2)} T_{j,2}^{(2)}) = (T_{i,1}^{(1)} T_{i,1}^{(1)} T_{j,-1}^{(1)} T_{j,-1}^{(1)} + T_{i,1}^{(1)} T_{i,-1}^{(1)} T_{j,1}^{(1)} T_{j,1}^{(1)} - T_{i,1}^{(1)} T_{i,0}^{(1)} T_{j,1}^{(1)} T_{j,0}^{(1)} + T_{i,-1}^{(1)} T_{i,1}^{(1)} T_{j,1}^{(1)} T_{j,1}^{(1)} T_{j,1}^{(1)} T_{j,0}^{(1)} + T_{i,0}^{(1)} T_{i,1}^{(1)} T_{i,1}^{(1)} T_{j,1}^{(1)} T_{j,1}^{(1)} T_{j,0}^{(1)} T_{j,1}^{(1)} T_{j,0}^{(1)} T_{j,0}^{(1)} + T_{i,0}^{(1)} T_{i,0}^{(1)} T_{j,0}^{(1)} T_{j,0}^{(1)} T_{j,0}^{(1)} + T_{i,0}^{(1)} T_{i,0}^{(1)} T_{j,0}^{(1)} T_{j,0}$$

so,
$$H_K = \sum_{\langle i,j \rangle} K_{i,j} (Q_{ij}^{(0)} - Q_{ij}^{(1)} + Q_{ij}^{(2)}) = \sum_{\langle i,j \rangle} K_{i,j} ([T_i^{(0)} \times T_j^{(0)}]_0 + \sqrt{3} [T_i^{(1)} \times T_j^{(1)}]_0 + \sqrt{5} [T_i^{(2)} \times T_j^{(2)}]_0) = \sum_{\langle i,j \rangle} K_{i,j} (X_{0,ij}^0(0) + \sqrt{3} X_{0,ij}^0(1) + \sqrt{5} X_{0,ij}^0(2))$$

A. Initialization

There are only one site in system and environment. However, there are two operators in one site. The product of two 1 rank irreducible tensor operators can lead to 0 rank, 1 rank and 2 rank irreducible tensor operators. We need to calculate $\langle j||K^{(0)}||j\rangle,\langle j||K^{(1)}||j\rangle,\langle j||K^{(2)}||j\rangle$ for initialization.

$$\begin{split} &\langle j||K^{(0)}||j\rangle, \langle j||K^{(1)}||j\rangle, \langle j||K^{(2)}||j\rangle \text{ for initialization.} \\ &T_{M,i}^{(L)} = \sum_{M} T_{M,i}^{(L)} T_{M,i}^{(L)} T_{M,i}^{(L)} = \sum_{M} T_{M,i}^{(L)} T_{M,i}^{(L)} T_{M,i}^{(L)} = \sum_{M} T_{M,i}^{(L)} T_{M,i}^{(L)} T_{M,i}^{(L$$

B. Increasing a new site

In this section, we add a new site in the system block. The irreducible basis is different from before. We need to calculate $\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle, \langle J'_{sys}, S', J', \alpha' || T_i^{(2)} || J_{sys}, S, J, \alpha \rangle$ in the new irreducible basis. There, we need the product of two irreducible tensor operators.

If the site i is in the system block, then

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(K)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J'_{sys} + S' + J + K} \delta_{S', S} \sqrt{(2J' + 1)(2J + 1)} \left\{ \begin{matrix} J' & K & J \\ J_{sys} & S' & J'_{sys} \end{matrix} \right\} \langle J'_{sys}, \alpha' || T_i^{(K)} || J_{sys}, \alpha \rangle$$

if
$$J = J'$$

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(K)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J'_{sys} + S' + J + K} \delta_{S', S} (2J+1) \left\{ \begin{matrix} J & K & J \\ J_{sys} & S' & J'_{sys} \end{matrix} \right\} \langle J'_{sys}, \alpha' || T_i^{(K)} || J_{sys}, \alpha \rangle$$

$$= (-1)^{J'_{sys} + S + J + K} (2J+1) \left\{ \begin{matrix} J & K & J \\ J_{sys} & S & J'_{sys} \end{matrix} \right\} \langle J'_{sys}, \alpha' || T_i^{(K)} || J_{sys}, \alpha \rangle$$
 which would be non-zero only for $J \neq 0$. If $J = 0$, the $6j-$ coefficient would be zero.

If the site i is in the added new site, then

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(K)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J' + K} \delta_{J'_{sys}, J_{sys}} \sqrt{(2J' + 1)(2J + 1)} \begin{cases} J' & K & J \\ S & J'_{sys} & S' \end{cases} \langle S', \alpha' || T_i^{(K)} || S, \alpha \rangle$$

For the diagonal part with $J=J^{\prime},$ the result can be simplified as

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(K)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J + K} \delta_{J'_{sys}, J_{sys}} (2J + 1) \begin{cases} J & K & J \\ S & J'_{sys} & S' \end{cases} \langle S', \alpha' || T_i^{(K)} || S, \alpha \rangle$$

which would be non-zero only for $J \neq 0$. If J = 0, the 6j – coefficient would be

C. Calculating the Hamiltonian

Suppose the site i belong to the system block and the site j is the added new site. Then the Hamiltonian H_K $\sum_{\langle i,j\rangle} K_{i,j}(Q_{ij}^{(0)} - Q_{ij}^{(1)} + Q_{ij}^{(2)}) = \sum_{\langle i,j\rangle} K_{i,j}([T_i^{(0)} \times T_j^{(0)}]_0 + \sqrt{3}[T_i^{(1)} \times T_j^{(1)}]_0 + \sqrt{5}[T_i^{(2)} \times T_j^{(2)}]_0) = \sum_{\langle i,j\rangle} K_{i,j}(X_{0,ij}^0(0) + X_{0,ij}^0(0)) = \sum_{\langle i,j\rangle} K_{i,j}(X_{0,ij}^0(0) + X_{0,ij}^0(0$ $\sqrt{3}X_{0,ij}^0(1) + \sqrt{5}X_{0,ij}^0(2)$

The matrix element of the Hamiltonian H_K is

for $X_{0,ij}^0(k), k = 0, 1, 2,$

$$K_{i,j}\langle J'_{sys}, S', J', \alpha' | (X^{0}_{0,ij}(k)|J_{sys}, S, J, \alpha) = \frac{1}{\sqrt{2J+1}}K_{i,j}\langle J'_{sys}, S', J', \alpha' | |X^{0}_{ij}(k)||J_{sys}, S, J, \alpha\rangle$$
as $\langle J'_{sys}, S', J', \alpha' | |X^{0}_{ij}(k)||J_{sys}, S, J, \alpha\rangle = \sqrt{\frac{2J+1}{2k+1}}(-1)^{J_{sys}+S'+k+J} \left\{ \begin{matrix} J'_{sys} & J_{sys} & k \\ S & S' & J \end{matrix} \right\} \langle J'_{sys}, \alpha' | |T^{(k)}_{i}||J_{sys}, \alpha\rangle\langle S' | |T^{(k)}_{j}||S\rangle$

Thus, the matrix element of the Hamiltonian H_K is

Thus, the matrix element of the Hamiltonian
$$H_K$$
 is $H_K = K_{i,j}(\langle J'_{sys}, S', J', \alpha' | (X^0_{0,ij}(0)|J_{sys}, S, J, \alpha \rangle + \sqrt{3} \langle J'_{sys}, S', J', \alpha' | (X^0_{0,ij}(1)|J_{sys}, S, J, \alpha \rangle + \sqrt{5} \langle J'_{sys}, S', J', \alpha' | (X^0_{0,ij}(2)|J_{sys}, S, J, \alpha \rangle + \sqrt{5} \langle J'_{sys}, S', J', \alpha' | (X^0_{0,ij}(2)|J_{sys}, S, J, \alpha \rangle)$

$$= (K_{i,j} \frac{1}{\sqrt{2k+1}} (-1)^{J_{sys}+S'+0+J} \begin{cases} J'_{sys} & J_{sys} & 0 \\ S & S' & J \end{cases} \begin{cases} J'_{sys}, \alpha' | |T^{(0)}_i||J_{sys}, \alpha \rangle \langle S' | |T^{(1)}_j| | |S \rangle \end{cases}$$

$$+ (\sqrt{3}K_{i,j} \frac{1}{\sqrt{2k+1}} (-1)^{J_{sys}+S'+2+J} \begin{cases} J'_{sys} & J_{sys} & 1 \\ S & S' & J \end{cases} \langle J'_{sys}, \alpha' | |T^{(2)}_i| | |J_{sys}, \alpha \rangle \langle S' | |T^{(1)}_j| | |S \rangle)$$

$$+ (\sqrt{5}K_{i,j} \frac{1}{\sqrt{2k+1}} (-1)^{J_{sys}+S'+2+J} \begin{cases} J'_{sys} & J_{sys} & 2 \\ S & S' & J \end{cases} \langle J'_{sys}, \alpha' | |T^{(2)}_i| |J_{sys}, \alpha \rangle \langle S' | |T^{(0)}_j| |S \rangle)$$

$$= (K_{i,j}(-1)^{J_{sys}+S'+0+J} (-1)^{J_{sys}+S'+2+J} \begin{cases} J'_{sys} & J_{sys} & 1 \\ S's & S' & J \end{cases} \langle J'_{sys}, \alpha' | |T^{(1)}_i| |J_{sys}, \alpha \rangle \langle S' | |T^{(1)}_j| |S \rangle)$$

$$+ (\sqrt{5}K_{i,j} \frac{1}{\sqrt{3}} (-1)^{J_{sys}+S'+1+J} \begin{cases} J'_{sys} & J_{sys} & 2 \\ S & S' & J \end{cases} \langle J'_{sys}, \alpha' | |T^{(2)}_i| |J_{sys}, \alpha \rangle \langle S' | |T^{(2)}_j| |S \rangle)$$

$$= (K_{i,j} \frac{1}{\sqrt{(2S+1)(2J_{sys}+1)}}} \langle J_{sys}, \alpha' | |T^{(0)}_i| |J_{sys}, \alpha' |S \rangle \langle S' | |T^{(2)}_j| |S \rangle)$$

$$= (K_{i,j} \frac{1}{\sqrt{(2S+1)(2J_{sys}+1)}}} \langle J_{sys}, \alpha' | |T^{(0)}_i| |J_{sys}, \alpha \rangle \langle S | |T^{(1)}_j| |S \rangle)$$

$$+ (K_{i,j}(-1)^{J_{sys}+S+2+J} \begin{cases} J'_{sys} & J_{sys} & 1 \\ S & S & J \end{cases} \langle J'_{sys}, \alpha' | |T^{(1)}_i| |J_{sys}, \alpha \rangle \langle S | |T^{(2)}_j| |S \rangle)$$

$$+ (K_{i,j}(-1)^{J_{sys}+S+2+J} \begin{cases} J'_{sys} & J_{sys} & 1 \\ S & S & J \end{cases} \langle J'_{sys}, \alpha' | |T^{(2)}_i| |J_{sys}, \alpha \rangle \langle S | |T^{(2)}_j| |S \rangle)$$

Hamiltonian from the K interactions

In this section, we need the matrix-vector multiplication calculations

$$\begin{array}{l} H_{n+1} = H_{A_n} + H_{B_n} + (S_{A_n} \cdot S_{A_n+1})^2 + (S_{A_n+1} \cdot S_{B_n+1})^2 + (S_{B_n+1} \cdot S_{B_n})^2 \\ H_{n+1} |f\rangle = |g\rangle, \text{ thus } \\ (H_{A_n} + H_{B_n} + (S_{A_n} \cdot S_{A_n+1})^2 + (S_{A_n+1} \cdot S_{B_n+1})^2 + (S_{B_n+1} \cdot S_{B_n})^2) |f\rangle = |g\rangle \\ (H_{A_n}) |f\rangle = |g_1\rangle, (H_{B_n}) |f\rangle = |g_2\rangle, (S_{A_n} \cdot S_{A_n+1})^2 |f\rangle = |g_3\rangle, (S_{A_n+1} \cdot S_{B_n+1})^2 |f\rangle = |g_4\rangle, (S_{B_n+1} \cdot S_{B_n})^2 |f\rangle = |g_5\rangle \\ \text{thus, } |g\rangle = |g_1\rangle + |g_2\rangle + |g_3\rangle + |g_4\rangle + |g_5\rangle \end{array}$$

$$\begin{aligned} &\text{generally, } J_{ij}(\vec{S}_i \cdot \vec{S}_j)^2 | f \rangle = J_{ij}(\vec{S}_i \cdot \vec{S}_j)^2 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J \rangle f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) \\ &= \sum_{J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'} | J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J' \rangle \langle J'_{sys}, J'_{ns}, J'_{sysnew}, J_{env}, J'_{ne}, J'_{envnew}, J' | J'_{sys}, J'_{ns}, J'_{sysnew}, J_{env}, J'_{ne}, J'_{envnew}, J' | J'_{sys}, J'_{ns}, J'_{sysnew}, J_{env}, J'_{ne}, J'_{envnew}, J' | J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J'_{sysnew}, J'_{envnew}, J'_{sysnew}, J'_{envnew}, J'_{sysnew}, J'_{envnew}, J'_{sysnew}, J'_{sysnew},$$

1. Hamiltonian from the K interactions between the system block and the added sites

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The Hamiltonian contribution is K_{ij}(\vec{S}_i \cdot \vec{S}_j)^2, where \vec{S}_i and \vec{S}_j belong to the system block and the added site, respectively. The matrix-vector multiplication is explicitly written as \langle J'_{sys}, J'_{ns}, J'_{sysnew}, J'_{env}, J'_{ne}, J'_{envnew}, J' | K_{ij}(\vec{S}_i \cdot \vec{S}_j)^2 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J \rangle f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) = g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J) As \vec{S}_i and \vec{S}_j belong to the system block and the added site, respectively. Thus \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | K_{ij}(\vec{S}_i \cdot \vec{S}_j)^2 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) = g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) We can compute the matrix element as K_{ij}\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | (\vec{S}_i \cdot \vec{S}_j)^2 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle = \beta K_{ij} \frac{1}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X_{0,ij}^0(k) | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle = \beta K_{ij} \frac{1}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X_{0,ij}^0(k) | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle = \beta K_{ij} \frac{1}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X_{0,ij}^0(k) | J_{sys}, J_{ns}, J_{sysnew}, X_{env}, J_{ne}, J_{envnew}, 0 \rangle = \beta K_{ij} \frac{1}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X_{0,ij}^0(k) | J_{sys}, J_{ns}, J_{sysnew}, X_{env}, J_{ne}, J_{envnew}, 0 \rangle = \beta K_{ij} \frac{1}{\sqrt{2J_{sysnew}+1}} \langle J'_{sys}, J_{sysnew}, J_{env}, J_{ne}, J_{sysnew}, J_{env}, J_{ne},
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IV. IRREDUCIBLE TENSOR OPERATORS OF THE CHIRAL TERMS

A. J_{χ} term

For i < j < k, the sites follow the clockwise direction.

$$\begin{split} H &= \sum_{\triangle} J_{\chi}(\mathbf{S}_{i} \times \mathbf{S}_{j}) \cdot \mathbf{S}_{k} = \sum_{\triangle} J_{\chi}((S_{i}^{y}S_{j}^{z} - S_{i}^{z}S_{j}^{y})\hat{x}) + (S_{i}^{z}S_{j}^{x} - S_{i}^{x}S_{j}^{z})\hat{y} + (S_{i}^{x}S_{j}^{y} - S_{i}^{y}S_{j}^{x})\hat{z}) \cdot (S_{k}^{x}\hat{x} + S_{y}^{y}\hat{y} + S_{k}^{z}\hat{z}) \\ &= \sum_{\triangle} J_{\chi}((S_{i}^{y}S_{j}^{z} - S_{i}^{z}S_{j}^{y})S_{k}^{x} + (S_{i}^{z}S_{j}^{x} - S_{i}^{x}S_{j}^{z})S_{k}^{y} + (S_{i}^{x}S_{j}^{y} - S_{i}^{y}S_{j}^{x})S_{k}^{z}) \\ &= \sum_{\triangle} J_{\chi}((S_{i}^{y}S_{j}^{z}S_{k}^{x} - S_{i}^{z}S_{j}^{y}S_{k}^{x}) + (S_{i}^{z}S_{j}^{x}S_{k}^{y} - S_{i}^{x}S_{j}^{z}S_{k}^{y}) + (S_{i}^{x}S_{j}^{y}S_{k}^{z} - S_{i}^{y}S_{j}^{x}S_{k}^{z})) \\ &= \sum_{\triangle} J_{\chi}((1/2i(S_{i}^{\dagger} - S_{i}^{-}) \cdot 1/2(S_{k}^{\dagger} + S_{k}^{-}))S_{j}^{z} - (1/2i(S_{j}^{\dagger} - S_{j}^{-}) \cdot 1/2(S_{k}^{\dagger} + S_{k}^{-}))S_{i}^{z}) \\ &+ ((1/2(S_{i}^{\dagger} + S_{j}^{-}) \cdot 1/2i(S_{k}^{\dagger} - S_{k}^{-}))S_{j}^{z} - (1/2i(S_{i}^{\dagger} - S_{i}^{-}) \cdot 1/2i(S_{k}^{\dagger} - S_{k}^{-}))S_{j}^{z}) \\ &+ ((1/2(S_{i}^{\dagger} + S_{i}^{-}) \cdot 1/2i(S_{j}^{\dagger} - S_{j}^{-}))S_{k}^{z} - (1/2i(S_{i}^{\dagger} - S_{i}^{-}) \cdot 1/2i(S_{j}^{\dagger} - S_{k}^{-}))S_{j}^{z}) \\ &+ \sum_{\triangle} J_{\chi}(1/4i(S_{i}^{\dagger}S_{k}^{\dagger} + S_{i}^{\dagger}S_{k}^{-} - S_{i}^{-}S_{k}^{\dagger} - S_{i}^{-}S_{j}^{-}S_{j}^{z})S_{k}^{z} - (S_{j}^{\dagger}S_{k}^{\dagger} + S_{j}^{\dagger}S_{k}^{-} - S_{j}^{-}S_{k}^{\dagger} - S_{j}^{-}S_{k}^{-}S_{j}^{-}S_{k}^{z})S_{i}^{z}) \\ &+ (1/4i(S_{j}^{\dagger}S_{k}^{\dagger} - S_{j}^{\dagger}S_{k}^{-} + S_{j}^{-}S_{k}^{\dagger} - S_{i}^{-}S_{j}^{-}S_{k}^{\dagger} - S_{i}^{\dagger}S_{j}^{-} - S_{i}^{-}S_{j}^{-}S_{k}^{-}S_{j}^{-}S_{k}^{z})S_{j}^{z}) \\ &+ (1/4i(S_{j}^{\dagger}S_{j}^{\dagger} - S_{i}^{\dagger}S_{j}^{-} + S_{i}^{-}S_{j}^{\dagger} - S_{i}^{\dagger}S_{j}^{-} + S_{i}^{-}S_{j}^{\dagger})S_{k}^{z}) \\ &= 1/2i((S_{i}^{\dagger}S_{k}^{-} - S_{i}^{-}S_{k}^{\dagger})S_{j}^{z} + (-S_{j}^{\dagger}S_{k}^{-} + S_{i}^{-}S_{j}^{-}S_{k}^{-} + S_{i}^{-}S_{j}^{-})S_{k}^{z}) \\ &= 1/2i((S_{i}^{\dagger}S_{j}^{-} - S_{i}^{-}S_{j}^{-}S_{k}^{-} + S_{i}^{-}S_{j}^{-}S_{k}^{-} + S_{i}^{-}S_{j}^{-}S_{k}^{-} + S_{i}^{-}S_{j}^{-}S_{k}^{z})S_{j}^{z}) \\ &= 1/2i((S_{i}^{\dagger}S_{j}^{-} - S_{i}^{-}S_{j}^{-}S_{k}^{-} + S_{i}^{-}S_{j}^{-}S_{k}^{-} + S_{i}^{-}S_{j}^{-$$

so, $H = \sum_{\triangle} (J_{\chi,ijk})(i) (T_{1,i}^{(1)} T_{0,j}^{(1)} T_{-1,k}^{(1)} - T_{-1,i}^{(1)} T_{0,j}^{(1)} T_{1,k}^{(1)} - T_{0,i}^{(1)} T_{1,j}^{(1)} T_{-1,k}^{(1)} + T_{0,i}^{(1)} T_{-1,j}^{(1)} T_{1,k}^{(1)} - T_{1,i}^{(1)} T_{0,k}^{(1)} + T_{-1,i}^{(1)} T_{0,k}^{(1)} + T_{-1,$

For i > j > k(k < j < i), the sites follow the clockwise direction(another situation).

$$H = \sum_{\Delta} J_{\chi}(\mathbf{S}_{k} \times \mathbf{S}_{j}) \cdot \mathbf{S}_{i} = \sum_{\Delta} J_{\chi}((S_{k}^{y}S_{j}^{z} - S_{k}^{z}S_{j}^{y})\hat{x}) + (S_{k}^{z}S_{j}^{x} - S_{k}^{x}S_{j}^{z})\hat{y} + (S_{k}^{x}S_{j}^{y} - S_{k}^{y}S_{j}^{x})\hat{z}) \cdot (S_{i}^{x}\hat{x} + S_{i}^{y}\hat{y} + S_{i}^{z}\hat{z})$$

$$= \sum_{\Delta} J_{\chi}((S_{k}^{y}S_{j}^{z} - S_{k}^{z}S_{j}^{y})S_{i}^{x} + (S_{k}^{z}S_{j}^{x} - S_{k}^{x}S_{j}^{z})S_{i}^{y} + (S_{k}^{x}S_{j}^{y} - S_{k}^{x}S_{j}^{z})S_{i}^{y})$$

$$= \sum_{\Delta} J_{\chi}((S_{k}^{y}S_{j}^{z}S_{i}^{x} - S_{k}^{z}S_{j}^{y}S_{i}^{x}) + (S_{k}^{z}S_{j}^{x}S_{i}^{y} - S_{k}^{x}S_{j}^{z}S_{i}^{y}) + (S_{k}^{x}S_{j}^{y}S_{i}^{z} - S_{k}^{y}S_{j}^{x}S_{i}^{z})$$

$$= \sum_{\Delta} (J_{\chi,kji})(i)(T_{1,k}^{(1)}T_{0,j}^{(1)}T_{-1,i}^{(1)} - T_{-1,k}^{(1)}T_{0,j}^{(1)}T_{1,i}^{(1)} - T_{0,k}^{(1)}T_{1,j}^{(1)}T_{-1,i}^{(1)} + T_{0,k}^{(1)}T_{-1,j}^{(1)}T_{1,i}^{(1)} - T_{1,k}^{(1)}T_{0,i}^{(1)} + T_{-1,k}^{(1)}T_{0,i}^{(1)})$$

$$= -\sum_{\Delta} J_{\chi}(\mathbf{S}_{i} \times \mathbf{S}_{j}) \cdot \mathbf{S}_{k}$$

$$(6)$$

The two rank-1 ITO on two sites respectively can be coupled as one rank-1 ITO.

$$\begin{split} T_0^{(1)} &= \langle 1,1,1,-1|1,0\rangle T_1^{(1)}T_{-1}^{(1)} + \langle 1,-1,1,1|1,0\rangle T_{-1}^{(1)}T_1^{(1)} + \langle 1,0,1,0|1,0\rangle T_0^{(1)}T_0^{(1)} = \sqrt{\frac{1}{2}}T_1^{(1)}T_{-1}^{(1)} - \sqrt{\frac{1}{2}}T_{-1}^{(1)}T_1^{(1)} \\ T_1^{(1)} &= \langle 1,1,1,0|1,1\rangle T_1^{(1)}T_0^{(1)} + \langle 1,0,1,1|1,1\rangle T_0^{(1)}T_1^{(1)} + \langle 1,-1,1,2|1,1\rangle T_{-1}^{(1)}T_2^{(1)} = \sqrt{\frac{1}{2}}T_1^{(1)}T_0^{(1)} - \sqrt{\frac{1}{2}}T_0^{(1)}T_1^{(1)} \\ T_{-1}^{(1)} &= \langle 1,1,1,-2|1,-1\rangle T_1^{(1)}T_{-2}^{(1)} + \langle 1,0,1,-1|1,-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1,-1,1,0|1,-1\rangle T_{-1}^{(1)}T_0^{(1)} = \sqrt{\frac{1}{2}}T_0^{(1)}T_{-1}^{(1)} - \sqrt{\frac{1}{2}}T_{-1}^{(1)}T_0^{(1)} \\ T_{-1}^{(1)} &= \langle 1,1,1,-2|1,-1\rangle T_1^{(1)}T_{-2}^{(1)} + \langle 1,0,1,-1|1,-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1,-1,1,0|1,-1\rangle T_{-1}^{(1)}T_0^{(1)} = \sqrt{\frac{1}{2}}T_0^{(1)}T_{-1}^{(1)} - \sqrt{\frac{1}{2}}T_{-1}^{(1)}T_0^{(1)} \\ T_{-1}^{(1)} &= \langle 1,1,1,-2|1,-1\rangle T_1^{(1)}T_{-2}^{(1)} + \langle 1,0,1,-1|1,-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1,-1,1,0|1,-1\rangle T_{-1}^{(1)}T_0^{(1)} = \sqrt{\frac{1}{2}}T_0^{(1)}T_1^{(1)} \\ T_{-1}^{(1)} &= \langle 1,1,1,-2|1,-1\rangle T_1^{(1)}T_{-2}^{(1)} + \langle 1,0,1,-1|1,-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1,-1,1,0|1,-1\rangle T_0^{(1)}T_0^{(1)} = \sqrt{\frac{1}{2}}T_0^{(1)}T_1^{(1)} \\ T_{-1}^{(1)} &= \langle 1,1,1,-2|1,-1\rangle T_1^{(1)}T_{-2}^{(1)} + \langle 1,0,1,-1|1,-1\rangle T_0^{(1)}T_{-1}^{(1)} + \langle 1,-1,1,0|1,-1\rangle T_1^{(1)}T_1^{(1)} = \sqrt{\frac{1}{2}}T_0^{(1)}T_1^{(1)} \\ T_{-1}^{(1)} &= \langle 1,1,1,-2|1,-1\rangle T_1^{(1)}T_1^{(1)} + \langle 1,0,1,-1|1,-1\rangle T_1^{(1)}T_1^{(1)} + \langle 1,-1,1,0|1,-1\rangle T_1^{(1)}T_1^{(1)} + \langle 1,1,1,1,-1|1,-1\rangle T_1^{(1)}T_1^{(1)} + \langle 1,1,1,1,$$

For i < j < k. The new rank-1 ITO can be coupled to one rank-0 ITO with another rank-1 ITO on the third site. $X_0^{(0)} = \langle 1,1,1,-1|0,0\rangle T_1^{(1)}T_{-1}^{(1)} + \langle 1,0,1,0|0,0\rangle T_0^{(1)}T_0^{(1)} + \langle 1,-1,1,1|0,0\rangle T_{-1}^{(1)}T_1^{(1)}$

$$\begin{split} &=\sqrt{\frac{1}{3}}T_{1}^{(1)}T_{-1}^{(1)}-\sqrt{\frac{1}{3}}T_{0}^{(1)}T_{0}^{(1)}+\sqrt{\frac{1}{3}}T_{-1}^{(1)}T_{1}^{(1)}\\ &=\sqrt{\frac{1}{3}}(\sqrt{\frac{1}{2}}T_{1}^{(1)}T_{0}^{(1)}-\sqrt{\frac{1}{2}}T_{0}^{(1)}T_{1}^{(1)})T_{-1}^{(1)}-\sqrt{\frac{1}{3}}(\sqrt{\frac{1}{2}}T_{1}^{(1)}T_{-1}^{(1)}-\sqrt{\frac{1}{2}}T_{-1}^{(1)}T_{1}^{(1)})T_{0}^{(1)}+\sqrt{\frac{1}{3}}(\sqrt{\frac{1}{2}}T_{0}^{(1)}T_{-1}^{(1)}-\sqrt{\frac{1}{2}}T_{-1}^{(1)}T_{0}^{(1)})T_{1}^{(1)}\\ &=\sqrt{\frac{1}{6}}T_{1}^{(1)}T_{0}^{(1)}T_{-1}^{(1)}-\sqrt{\frac{1}{6}}T_{0}^{(1)}T_{1}^{(1)}T_{-1}^{(1)}-\sqrt{\frac{1}{6}}T_{1}^{(1)}T_{-1}^{(1)}T_{0}^{(1)}+\sqrt{\frac{1}{6}}T_{-1}^{(1)}T_{1}^{(1)}T_{0}^{(1)}+\sqrt{\frac{1}{6}}T_{0}^{(1)}T_{-1}^{(1)}T_{1}^{(1)}-\sqrt{\frac{1}{6}}T_{0}^{(1)}T_{1}^{(1)}T_{1}^{(1)}\\ &=\sin(([T_{i}^{(1)}\times T_{j}^{(1)}]_{10}\times T_{0,k}^{(1)}+[T_{i}^{(1)}\times T_{j}^{(1)}]_{11}\times T_{-1,k}^{(1)}+[T_{i}^{(1)}\times T_{j}^{(1)}]_{1-1}\times T_{1,k}^{(1)}+\sqrt{\frac{1}{6}}T_{0,j}^{(1)}T_{1,k}^{(1)}-\sqrt{\frac{1}{6}}T_{0,j}^{(1)}T_{-1,k}^{(1)}+T_{0,j}^{(1)}T_{-1,k}^{(1)}+T_{0,k$$

For
$$i>j>k(k< j< i)$$
. $([T_k^{(1)}\times T_j^{(1)}]_{10}\times T_{0,i}^{(1)}+[T_k^{(1)}\times T_j^{(1)}]_{11}\times T_{(-1),i}^{(1)}+[T_k^{(1)}\times T_j^{(1)}]_{1-1}\times T_{(1),i}^{(1)})=X_{kji}^0=\frac{1}{\sqrt{6}}(T_{1,k}^{(1)}T_{0,j}^{(1)}T_{-1,i}^{(1)}-T_{-1,k}^{(1)}T_{1,j}^{(1)}T_{-1,i}^{(1)}+T_{0,k}^{(1)}T_{-1,j}^{(1)}T_{1,i}^{(1)}-T_{1,k}^{(1)}T_{-1,j}^{(1)}T_{0,i}^{(1)}+T_{-1,k}^{(1)}T_{1,j}^{(1)}T_{0,i}^{(1)})$ so, $H=\sum_{\triangle}J_{\chi}(\mathbf{S}_k\times\mathbf{S}_j)\cdot\mathbf{S}_i=\sum_{\triangle}(J_{\chi,kji})(i)\sqrt{6}X_{kji}^0=-\sum_{\triangle}(J_{\chi,ijk})(i)\sqrt{6}X_{ijk}^0$

B. Initialization

There are only one site in system and environment. As the Heisenberg terms. We need to calculate $\langle J||T^{(1)}||J\rangle$ for initializa-

tion.
$$\langle JJ|T_0^{(1)}|JJ\rangle = \frac{1}{\sqrt{2J+1}}\langle JJ10|JJ\rangle\langle J||T^{(1)}||J\rangle = \frac{1}{\sqrt{2J+1}}\frac{J}{\sqrt{J(J+1)}}\langle J||T^{(1)}||J\rangle = J$$
 Thus,
$$\langle J||T^{(1)}||J\rangle = \sqrt{J(J+1)(2J+1)}.$$
 If $J=\frac{1}{2}, \langle J||T^{(1)}||J\rangle = \sqrt{\frac{3}{2}}.$ If $J=1, \langle J||T^{(1)}||J\rangle = \sqrt{6}.$

C. Increasing a new site

As the J term, we add a new site in the system block. The irreducible basis is different from before. We need to calculate $\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle$ in the new irreducible basis. There, we need the product of two irreducible tensor operators. There, |J - J'| = 0, 1.

If the site i is in the system block, then

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J'_{sys} + S' + J + 1} \delta_{S',S} \sqrt{(2J' + 1)(2J + 1)} \left\{ \begin{matrix} J' & 1 & J \\ J_{sys} & S' & J'_{sys} \end{matrix} \right\} \langle J'_{sys}, \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle$$
 if $J = J'$
$$\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J'_{sys} + S' + J + 1} \delta_{S',S} (2J + 1) \left\{ \begin{matrix} J & 1 & J \\ J_{sys} & S' & J'_{sys} \end{matrix} \right\} \langle J'_{sys}, \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle$$

$$= (-1)^{J'_{sys} + S + J + 1} (2J + 1) \left\{ \begin{matrix} J & 1 & J \\ J_{sys} & S & J'_{sys} \end{matrix} \right\} \langle J'_{sys}, \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle$$
 which would be non-zero only for $J \neq 0$. If $J = 0$, the $6j-$ coefficient would be zero.

If the site i is in the added new site, then

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J' + 1} \delta_{J'_{sys}, J_{sys}} \sqrt{(2J' + 1)(2J + 1)} \begin{cases} J' & 1 & J \\ S & J'_{sys} & S' \end{cases} \langle S', \alpha' || T_i^{(1)} || S, \alpha \rangle$$

For the diagonal part with $J=J^\prime$, the result can be simplified as

$$\langle J'_{sys}, S', J', \alpha' || T_i^{(1)} || J_{sys}, S, J, \alpha \rangle = (-1)^{J_{sys} + S + J + 1} \delta_{J'_{sys}, J_{sys}} (2J + 1) \begin{cases} J & 1 & J \\ S & J'_{sys} & S' \end{cases} \langle S', \alpha' || T_i^{(1)} || S, \alpha \rangle$$

which would be non-zero only for $J \neq 0$. If J = 0, the 6j – coefficient would be zero.

D. Calculating the Hamiltonian

Suppose the site i, j belong to the system block and the site k is the added new site. As i < j < k. Then the Hamiltonian

$$H_{J_\chi} = \sum_{\triangle} J_\chi(\mathbf{S}_i \times \mathbf{S}_j) \cdot \mathbf{S}_k = \sum_{\triangle} (J_{\chi,ijk})(i) \sqrt{6} X^0_{ijk}$$
 The matrix element of the Hamiltonian H_{J_χ} is

$$\begin{split} &(J_{\chi,ijk})(i)\sqrt{6}\langle J_{sys}',S',J',\alpha'|X_{0,ijk}^0|J_{sys},S,J,\alpha\rangle = (J_{\chi,ijk})(i)\sqrt{6}\frac{1}{\sqrt{2J+1}}\langle J_{sys}',S',J',\alpha'||X_{ijk}^0||J_{sys},S,J,\alpha\rangle \\ &\text{as } \langle J_{sys}',S',J',\alpha'||X_{ij}^0||J_{sys},S,J,\alpha\rangle = \sqrt{\frac{2J+1}{3}}(-1)^{J_{sys}+S'+1+J} \left\{ \begin{matrix} J_{sys}' & J_{sys} & 1 \\ S & S' & J \end{matrix} \right\} \langle J_{sys}',\alpha'||T_i^{(1)}||J_{sys},\alpha\rangle\langle S'||T_j^{(1)}||S\rangle \\ &\text{Thus, the matrix element of the Hamiltonian } H_{J_\chi} \text{ is } \\ &(J_{\chi,ijk})(i)\sqrt{6}\langle J_{sys}',S',J',\alpha'|X_{0,ij}^0|J_{sys},S,J,\alpha\rangle = \sqrt{2}(J_{\chi,ijk})(i)(-1)^{J_{sys}+S'+1+J} \left\{ \begin{matrix} J_{sys}' & J_{sys} & 1 \\ S & S' & J \end{matrix} \right\} \\ &(J_{sys}',\alpha'||T_i^{(1)}||J_{sys},\alpha\rangle\langle S'||T_j^{(1)}||S\rangle \end{split}$$

Suppose the site i belong to the system block and the site j is the added new site. As k < j < i. Then $\langle J'_{sus}, S', J', \alpha' | |X_{ij}^{(1)}| |J_{sys}, S, J, \alpha \rangle =$

$$\sqrt{(2J'+1)(2K+1)(2J+1)} \begin{cases} J'_{sys} & J_{sys} & 1 \\ S' & S & 1 \\ J' & J & 1 \end{cases} \langle J'_{sys}, \alpha' || T_i^{(1)} || J_{sys}, \alpha \rangle \langle S' || T_j^{(1)} || S \rangle (k=1)$$

E. Hamiltonian from the chiral interactions

In this section, we need the matrix-vector multiplication calculations $H = \sum_{\triangle} J_{\chi}(\mathbf{S}_{i} \times \mathbf{S}_{j}) \cdot \mathbf{S}_{k} = \sum_{\triangle} (J_{\chi,ijk})(i) \sqrt{6} X_{ijk}^{0}$ $H_{n+1}|f\rangle = |g\rangle \text{, thus}$ $(H_{A_{n},ijk} + H_{B_{n},ijk} + S_{A_{n},ij} \cdot S_{A_{n}+1,k} + S_{B_{n}+1,i} \cdot S_{B_{n},jk})|f\rangle = |g\rangle$ $(H_{A_{n},ijk})|f\rangle = |g_{1}\rangle, (H_{B_{n},ijk})|f\rangle = |g_{2}\rangle, (S_{A_{n},ij} \cdot S_{A_{n}+1,k})|f\rangle = |g_{3}\rangle, (S_{B_{n}+1,i} \cdot S_{B_{n},jk})|f\rangle = |g_{4}\rangle$ thus, $|g\rangle = |g_{1}\rangle + |g_{2}\rangle + |g_{3}\rangle + |g_{4}\rangle$ generally, $X_{ijk}^{0}|f\rangle = X_{ijk}^{0}|J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, J\rangle + |f\rangle = |f\rangle + |f\rangle$

1. Hamiltonian from the chiral interactions between the system block and the added sites

As $T_{ij}^{(1)}$ and $T_k^{(1)}$ belong to the system block and the added site, respectively. Thus $\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0| (J_{\chi,ijk})(i)\sqrt{6}X_{ijk}^0| J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$ $f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) = g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0)$ We can compute the matrix element as $\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0| (J_{\chi,ijk})(i)\sqrt{6}X_{ijk}^0| J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$ $= (J_{\chi,ijk})(i)\sqrt{6}\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0| X_{0,ijk}^0| J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$ $= (J_{\chi,ijk})(i)\sqrt{6}\sqrt{\frac{1}{2J_{sysnew}+1}}\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0| |X_{0,ijk}^0| |J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$ $= (J_{\chi,ijk})(i)\sqrt{6}\sqrt{\frac{1}{2J_{sysnew}+1}}\sqrt{\frac{\sqrt{2J_{sysnew}+1}}{\sqrt{3}}}(-1)^{J_{sys}+J_{ns}+1+J_{sysnew}}\left\{J'_{sys}, J_{sys}, J_{sys}, J_{sysnew}\right\}\langle J'_{sys}, \alpha'||T_{ij}^{(1)}||J_{sys}, \alpha\rangle\langle J_{ns}||T_{k}^{(1)}||J_{ns}\rangle$ $= (J_{\chi,ijk})(i)\sqrt{2}(-1)^{J_{sys}+J_{ns}+1+J_{sysnew}}\left\{J'_{sys}, J_{sys}, J_{sysnew}\right\}\langle J'_{sys}, \alpha'||T_{ij}^{(1)}||J_{sys}, \alpha\rangle\langle J_{ns}||T_{k}^{(1)}||J_{ns}\rangle$ If $J_{sys}=J'_{sys}$, then $J_{sys}\neq 0$

2. Hamiltonian from the chiral interactions between the system block and the added sites (N_s and N_e)

As $T_i^{(1)}$, $T_j^{(1)}$ and $T_k^{(1)}$ belong to the system block and the added sites $(N_s \text{ and } N_e)$, respectively. Thus $\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | (J_{\chi,ijk})(i)\sqrt{6}X^0_{ijk} | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$ $f(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0) = g(J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0)$ We can compute the matrix element as

$$\langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | (J_{\chi,ijk})(i) \sqrt{6} X_{ijk}^0 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= (J_{\chi,ijk})(i) \sqrt{6} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X_{0,ijk}^0 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= (J_{\chi,ijk})(i) \sqrt{6} \sqrt{\frac{1}{2J+1}} \langle J'_{sys}, J_{ns}, J'_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 | X_{ijk}^0 | J_{sys}, J_{ns}, J_{sysnew}, J_{env}, J_{ne}, J_{envnew}, 0 \rangle$$

$$= (J_{\chi,ijk})(i) \sqrt{6} \sqrt{\frac{1}{2J+1}} \frac{\sqrt{2J+1}}{\sqrt{3}} (-1)^{J_{sysnew}+J_{envnew}+1+J}$$

$$\begin{cases} J'_{sysnew}, J_{sysnew}, J_$$